Compliant plugin

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Abstract

1 Constraint forces

This approach unifies soft and hard constraints. Hard constraints are usually implemented using Lagrange multipliers λ in the following equation:

$$\begin{pmatrix} \mathbf{M} & -\mathbf{J}^T \\ \mathbf{J} & \end{pmatrix} \begin{pmatrix} \mathbf{a} \\ \lambda \end{pmatrix} = \begin{pmatrix} \mathbf{f}_{ext} \\ -\phi \end{pmatrix}$$
 (1)

where **M** is the mass matrix, **J** is the Jacobian matrix of the constraint(s), **a** is the acceleration, \mathbf{f}_{ext} is the net external force applied to the system, λ is the constraint force and ϕ is the constraint violation. λ and ϕ are vectors with as many entries as scalar constraints. The equation system is typically solved using a Schur complement to compute the constraint forces:

$$\mathbf{J}\mathbf{M}^{-1}\mathbf{J}^{T}\lambda = -\phi - \mathbf{J}\mathbf{M}^{-1}\mathbf{f}_{ext} \tag{2}$$

and then the acceleration is computed as $\mathbf{a} = \mathbf{M}^{-1}(\mathbf{f}_{ext} + \mathbf{J}^T \lambda)$.

In the generalized approach, the constraint forces are considered proportional to the constraint violations:

$$\lambda = \frac{1}{c} \left(\phi + d\dot{\phi} \right) \tag{3}$$

where c is the compliance of the constraints and d its damping ratio. Combined with a time discretization scheme, this leads to an equation system similar to (2) as shown is the next section.

2 Time integration

Our implicit scheme is:

$$\mathbf{v}_{n+1} = \mathbf{v}_n + h\mathbf{M}^{-1} \left(\alpha \mathbf{f}_{n+1} + (1-\alpha)\mathbf{f}_n\right)$$
 (4)

$$\mathbf{x}_{n+1} = \mathbf{x}_n + h\left(\beta \mathbf{v}_{n+1} + (1-\beta)\mathbf{v}_n\right) \tag{5}$$

where index n denotes current values while index n+1 denotes next values, α is the implicit velocity factor, and β is the implicit position factor. Let

$$\Delta \mathbf{v} = \mathbf{v}_{n+1} - \mathbf{v}_n = h \mathbf{M}^{-1} \left(\alpha \mathbf{f}_{n+1} + (1 - \alpha) \mathbf{f}_n \right)$$
 (6)

$$\Delta \mathbf{x} = \mathbf{x}_{n+1} - \mathbf{x}_n = h(v + \beta \Delta \mathbf{v})$$
 (7)

be the velocity and position changes across the time step. The constraint violation ϕ and its Jacobian **J** are:

$$\mathbf{J} = \frac{\partial \phi}{\partial \mathbf{x}} \tag{8}$$

$$\phi_{n+1} \simeq \phi_n + \mathbf{J}\Delta \mathbf{x} = \phi_n + h\dot{\phi} + \mathbf{J}h\beta\Delta \mathbf{v}$$
 (9)

$$\dot{\phi}_{n+1} \simeq \phi_{n+1} + \mathbf{J} \Delta \mathbf{v} \tag{10}$$

The corresponding forces are:

$$\mathbf{f} = \mathbf{f}_{ext} + \mathbf{J}^T \lambda \tag{11}$$

$$\lambda_i = -\frac{1}{c_i}(\phi_i + d\dot{\phi}_i) \tag{12}$$

where the subscript i denotes a scalar constraint.

The average constraint forces are computed using equations 12, 9 and 10:

$$\bar{\lambda}_{i} = \alpha \lambda_{n+1} + (1 - \alpha) \lambda_{n}$$

$$= -\frac{1}{c_{i}} (\alpha \phi + \alpha h \dot{\phi} + \alpha h \beta \mathbf{J} \Delta \mathbf{v} + \alpha d \dot{\phi} + \alpha d \mathbf{J} \Delta \mathbf{v} + (1 - \alpha) \phi + (1 - \alpha) d \dot{\phi})$$

$$= -\frac{1}{c_{i}} (\phi + d \dot{\phi} + \alpha h \dot{\phi} + \alpha (h \beta + d) \mathbf{J} \Delta \mathbf{v})$$

We can rewrite the previous equation as:

$$\mathbf{J}\Delta\mathbf{v} + \frac{1}{\alpha(h\beta+d)}\mathbf{C}\bar{\lambda} = -\frac{1}{\alpha(h\beta+d)}(\phi + (d+\alpha h)\dot{\phi})$$
 (13)

where values without indices denote current values. The complete equation system is:

$$\begin{pmatrix} \frac{1}{h} \mathbf{P} \mathbf{M} & -\mathbf{P} \mathbf{J}^T \\ \mathbf{J} & \frac{1}{l} \mathbf{C} \end{pmatrix} \begin{pmatrix} \mathbf{\Delta} \mathbf{v} \\ \bar{\lambda} \end{pmatrix} = \begin{pmatrix} \mathbf{P} \mathbf{f}_{ext} \\ -\frac{1}{l} (\phi + (d + \alpha h) \dot{\phi}) \end{pmatrix}$$
(14)

where $l = \alpha(h\beta + d)$ The system is singular due to matrix **P**, however we can use $\mathbf{P}\mathbf{M}^{-1}\mathbf{P}$ as inverse mass matrix to compute a Schur complement:

$$\begin{pmatrix} h\mathbf{J}\mathbf{P}\mathbf{M}^{-1}\mathbf{P}\mathbf{J}^{T} + \frac{1}{l}\mathbf{C} \end{pmatrix} \bar{\lambda} = -\frac{1}{l} \left(\phi + (d + h\alpha)\dot{\phi} \right) - h\mathbf{J}\mathbf{M}^{-1}\mathbf{f}_{ext}$$

$$\Delta \mathbf{v} = h\mathbf{P}\mathbf{M}^{-1}(\mathbf{f}_{ext} + \mathbf{J}^{T}\bar{\lambda})$$

$$\Delta \mathbf{x} = h(\mathbf{v} + \beta\Delta\mathbf{v})$$

3 Matrix assembly

The equation system, in its most general form, can be written as:

$$\begin{pmatrix} \mathbf{M} & -\mathbf{J}^T \\ \mathbf{J} & \mathbf{C} \end{pmatrix} \begin{pmatrix} \mathbf{x} \\ \lambda \end{pmatrix} = \begin{pmatrix} \mathbf{f} \\ \phi \end{pmatrix}$$
 (15)

We assemble the 7 terms of equation 15 separately. Figure 3 shows an example of mechanical system and matrix. The independent DOFs are X_a and X_d . State X_b is attached to X_a using a simple mapping, and a mass matrix M_{bb} is defined at this level. State X_c is attached to X_b using a simple mapping, and a compliance matrix $C_{\alpha\alpha}$ (possibly a deformation force) is applied to these DOFs. State X_c is attached to X_d and X_d at the same time, using a MultiMapping. A compliance matrix $C_{\beta\beta}$, possibly an interaction force, is applied to these DOFs, while a mass M_{dd} is applied to X_d . The corresponding equation system has the block structure shown in the right of Figure 3.

The 8 parts of the equation system are shown in Figure 3. The J matrices are the mapping matrices. The bottom row has two mappings, since the state X_e impacted by compliance β depends on two parent states. Offset matrices J_{*0} are shown in the bottom of the figure. They are composed of identity matrices (represented with a diagonal in a block) and null blocks. These matrices are used to shift the indices of local vector and matrices to global numbering.

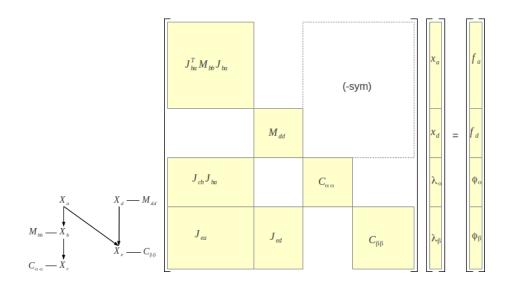


Figure 1: A mechanical system and its matrix equation. Left: scene graph. Right: block view of the corresponding equation system, with non-null blocks highlighted in yellow.

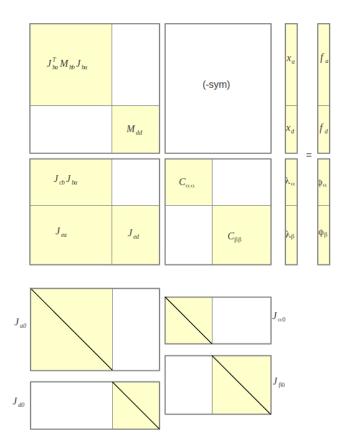


Figure 2: Top: the 8 parts of the equation system. Bottom: the offset matrices used in the assembly.