

The 1d wave equation looks like

$$u_{tt} = c^2 u_{xx}$$

a)

We use a finite difference scheme in the time dimension

$$\frac{u^{n-1} - 2u^n + u^{n+1}}{\Delta t^2} = c^2 u_{xx}$$

This can be solved for the next time step

$$u^{n+1} = \Delta t^2 c^2 u_{xx} - u^{n-1} + 2u^n$$

c)

We rearrange the last equation

$$u^{n+1} - \Delta t^2 c^2 u_{xx} + u^{n-1} - 2u^n = \mathcal{L}(u) = R$$

We want  $R$  to be orthogonal to the vector space  $\mathcal{V}$ , so we multiply by a basis function  $\varphi_i$  and integrate

$$\begin{aligned} (R, \varphi_i) &= \int_0^L dx R \varphi_i \\ &= \int_0^L dx \left[ u^{n+1} - 2u^n + u^{n-1} - c^2 \Delta t^2 u'' \right] \varphi_i \\ &= \int_0^L dx u^{n+1} \varphi_i - 2 \int_0^L dx u^n \varphi_i + \int_0^L dx u^{n-1} \varphi_i - c^2 \Delta t^2 \int_0^L dx u'' \varphi_i, \end{aligned}$$

where we use  $u''$  instead of  $u_{xx}$ . We look at the second derivative term and apply integration by parts

$$\int_0^L dx u'' \varphi_i = u' \varphi_i \Big|_0^L - \int_0^L dx u' \varphi_i' = - \int_0^L dx u' \varphi_i' = -(u', \varphi_i'),$$

because the boundary terms are zero. We can then write the residual as

$$(R, \varphi_i) = (u^{n+1}, \varphi_i) - 2(u^n, \varphi_i) + (u^{n-1}, \varphi_i) + C^2 (u', \varphi_i') = 0$$

where we have defined  $C^2 = \Delta t^2 c^2$ , using that the residual should be orthogonal to all basis vectors and evaluating each  $u$  at timestep  $n$ . By inserting  $u = \sum_j c_j \varphi_j$ , we get

$$\left( \sum_j c_j^{n+1} \varphi_j, \varphi_i \right) = 2 \left( \sum_j c_j^n \varphi_j, \varphi_i \right) - \left( \sum_j c_j^{n-1} \varphi_j, \varphi_i \right) - C^2 \left( \sum_j c_j^n \varphi_j', \varphi_i' \right).$$

Since the P1 elements have the property that  $u(x_i) = \sum_j c_j \varphi_j(x_i) = c_j \delta_{ij}$ , we have

$$c_j^{n+1}(\varphi_j, \varphi_i) = 2c_j^n(\varphi_j, \varphi_i) - c_j^{n-1}(\varphi_j, \varphi_i) + C^2 c_j^n(\varphi_j', \varphi_i').$$

We recognize this as equation  $i$  in the matrix multiplication

$$M c^{n+1} = 2M c^n - M c^{n-1} + C^2 K c^n,$$

where  $M_{ij} = (\varphi_i, \varphi_j)$  and  $K_{ij} = (\varphi'_i, \varphi'_j)$ . With the P1 elements, we then have the  $5 \times 5$   $M$  matrix given as

$$M = \frac{h}{6} \begin{pmatrix} 2 & 1 & 0 & 0 & 0 \\ 1 & 4 & 1 & 0 & 0 \\ 0 & 1 & 4 & 1 & 0 \\ 0 & 0 & 1 & 4 & 1 \\ 0 & 0 & 0 & 1 & 2 \end{pmatrix}$$

and the  $5 \times 5$   $K$  matrix looks like

$$K = \frac{1}{h} \begin{pmatrix} 2 & -1 & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 \\ 0 & -1 & 2 & -1 & 0 \\ 0 & 0 & -1 & 2 & -1 \\ 0 & 0 & 0 & -1 & 2 \end{pmatrix}.$$

$$\left[ D_t D_t \left( u + \frac{1}{6} \Delta x^2 D_x D_x u \right) = c^2 D_x D_x u \right]_i^n \quad (1)$$

## Stability analysis

We expand  $u$  in a discrete Fourier basis so that

$$u_p^n = \exp \left[ i(kp\Delta x - \tilde{\omega}n\Delta t) \right] \quad (2)$$

Writing equation (1) on discretized form

$$\begin{aligned} 6C^2 \left[ u_{i-1}^n - 2u_i^n + u_{i+1}^n \right] &= 6(u_i^{n-1} - 2u_i^n + u_i^{n+1}) + (u_{i-1}^{n-1} - 2u_{i-1}^n + u_{i-1}^{n+1}) \\ &\quad - 2(u_i^{n-1} - 2u_i^n + u_i^{n+1}) + (u_{i+1}^{n-1} - 2u_{i+1}^n + u_{i+1}^{n+1}) \\ &= 4(u_i^{n-1} - 2u_i^n + u_i^{n+1}) + (u_{i-1}^{n-1} - 2u_{i-1}^n + u_{i-1}^{n+1}) \\ &\quad + (u_{i+1}^{n-1} - 2u_{i+1}^n + u_{i+1}^{n+1}) \end{aligned}$$

and inserting (2) into the LHS

$$\begin{aligned} L.H.S. &= 6C^2 e^{-i\tilde{\omega}n\Delta t} \left[ e^{ik(p-1)\Delta x} - 2e^{ikp\Delta x} + e^{ik(p+1)\Delta x} \right] \\ &= 6C^2 e^{i(kp\Delta x - \tilde{\omega}n\Delta t)} \left[ e^{-ik\Delta x} - 2 + e^{ik\Delta x} \right] \\ &= -24C^2 e^{i(kp\Delta x - \tilde{\omega}n\Delta t)} \sin^2 \left( \frac{k\Delta x}{2} \right) \end{aligned}$$

and into RHS

$$\begin{aligned} R.H.S. &= -16e^{i(kp\Delta x - \tilde{\omega}n\Delta t)} \sin^2 \left( \frac{\tilde{\omega}\Delta t}{2} \right) - 4e^{i(k(p-1)\Delta x - \tilde{\omega}n\Delta t)} \sin^2 \left( \frac{\tilde{\omega}\Delta t}{2} \right) \\ &\quad - 4e^{i(k(p+1)\Delta x - \tilde{\omega}n\Delta t)} \sin^2 \left( \frac{\tilde{\omega}\Delta t}{2} \right) \\ &= -4e^{i(kp\Delta x - \tilde{\omega}n\Delta t)} \sin^2 \left( \frac{\tilde{\omega}\Delta t}{2} \right) \left[ e^{-ik\Delta x} + 4 + e^{ik\Delta x} \right] \\ &= -8e^{i(kp\Delta x - \tilde{\omega}n\Delta t)} \sin^2 \left( \frac{\tilde{\omega}\Delta t}{2} \right) \left[ \cos(k\Delta x) + 2 \right]. \end{aligned}$$

Combining both sides gives

$$\sin^2 \left( \frac{\tilde{\omega}\Delta t}{2} \right) = \frac{3C^2}{\cos(k\Delta x) + 2} \sin^2 \left( \frac{k\Delta x}{2} \right),$$

and by square rooting, we get

$$\sin\left(\frac{\tilde{\omega}\Delta t}{2}\right) = \frac{\sqrt{3}C}{\left[\cos(k\Delta x) + 2\right]^{\frac{1}{2}}} \sin\left(\frac{k\Delta x}{2}\right).$$

The sine on the right hand side has its maximum values  $\pm 1$  when  $k\Delta x = (n+1)\pi$ . The cosine in the denominator is then  $\pm 1$ . The worst case scenario is if  $\cos(k\Delta x) = -1$ , where we have the stability criterion  $C\sqrt{3} \leq 1$  giving

$$C \leq \frac{1}{\sqrt{3}},$$

because the sine on the left hand side must evaluate to less than 1. The solution -1 will not appear since we actually have a sine squared.