Geometric Algebra

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This is a collection of notes from my study of the *Geometric Algebra for Physicists* textbook (Doran and Lasenby 2007).

1 Foundations

The entire structure of geometric algebra can be boiled down to the geometric product, which is a binary operation between two vectors a and b. The usual rules of a vector space apply to these vectors, and the usefulness of geometric algebra is in the structure revealed by the product ab. This product satisfies three fairly natural axioms.

- 1. It is associative: (ab)c = a(bc)
- 2. It distributes over addition: a(b+c) = ab + ac
- 3. The square of any vector is real: $a^2 \in \mathbb{R}$. We do not make the assumption that $a^2 \geq 0$ to allow for the possibility of spaces with mixed signature (e.g., Minkowski space).

It is important not to interpret this product uv as another vector. In fact, uv is a sum of a scalar element $u \in \mathbb{R}$ and a 'bivector' V: ab = u + V – this should be understood as something like a complex number u + iv. In fact, for a vector space of dimension 2, this analogy is exact. In general, the bivector should be pictured as a directed area, or an object like the angular momentum in classical mechanics. This analogy extends further: we can have a trivector, which is an oriented volume, and so on and so forth.

However, this geometric product is not so exotic. We can still understand it in terms of some familiar concepts. For one, we observe that $(a+b)^2$ is always real, so we must have $a^2+ab+ba+b^2 \in \mathbb{R} \implies ab+ba \in \mathbb{R}$. We have found that the symmeterized geometric product is a binary relation between vectors that outputs a real number – that is, a dot product! So we define

$$a \cdot b = \frac{(ab + ba)}{2}.$$

We define the remaining part of the product as the exterior product (which might be thought of as replacing the role of the cross product):

$$a \wedge b = \frac{ab - ba}{2}.$$

This then gives $ab = a \cdot b + a \wedge b$: a sum of a scalar and a bivector, as promised.

Well, technically, not yet: we do not have a good reason to believe we should understand $a \wedge b$ as a bivector (or a directed area) yet. So we define it as such, and show this definition conforms to our intuition later. More precisely, we define exterior product of r vectors as the full anti-symmetrized product over all of them:

$$a_1 \wedge a_2 \wedge \ldots \wedge a_r = \frac{1}{r!} \sum_{\sigma \in \mathbb{S}_r} (-1)^{\sigma} a_{\sigma(1)} a_{\sigma(2)} \ldots a_{\sigma(r)},$$

where S_r is the set of all permutations on r elements and $(-1)^{\sigma}$ is the sign of a permutation. A consequence of this definition is that the product reverses sign under exchange of any two vectors. Then, we see that if any vector is repeated, the exterior product must be zero, which then implies if the vectors are linearly dependent, the exterior product is also zero (this follows simply by distributivity of the geometric product under addition). Therefore, the outer product can be understood to measure the dimensionality of a set of vectors. We say that the outer product of r vectors has grade r (if it does not vanish). A multivector that can be written purely as an outer product is called a blade.

Lemma 1. Any blade $a_1 \wedge ... \wedge a_r$ can be written simply as a product of orthogonal vectors $e_1 ... e_r$. This justifies our idea that a blade with r vectors can be interpreted as a directed area for r = 2, volume for r = 3, and so on.

Proof. Let $A_{i,j} = a_i \cdot a_j$. Since $A_{i,j}$ is a symmetric matrix, it can be diagonalized with P^TDP , where P is orthogonal. Letting $e_i = P_{i,k}a_k$, we see that $e_i \cdot e_j = P_{i,k}P_{j,\ell}a_k \cdot a_\ell = P_{i,k}A_{k,\ell}(P^T)_{\ell,j} = D_{i,j}$. Therefore, these vectors obey $e_ie_j = -e_je_i$ for $i \neq j$. Finally, since $e_i = P_{i,k}a_k \implies a_k = P_{i,k}e_i$:

$$a_1 \wedge \ldots \wedge a_r = \frac{1}{r!} \sum_{\sigma \in \mathbb{S}_r} (-1)^{\sigma} P_{i_1, \sigma(1)} P_{i_2, \sigma(2)} \ldots P_{i_r, \sigma(r)} e_{i_1} e_{i_2} \ldots e_{i_r}$$

Note that we have the restriction $i_1 \neq i_2 \neq \ldots \neq i_r$ due to the antisymmeterization of the sum. Therefore, we can rewrite as:

$$a_{1} \wedge \ldots \wedge a_{r} = \frac{1}{r!} \sum_{\rho, \sigma \in \mathbb{S}_{r}} (-1)^{\sigma} P_{\rho(1), \sigma(1)} P_{\rho(2), \sigma(2)} \dots P_{\rho(r), \sigma(r)} e_{\rho(1)} e_{\rho(2)} \dots e_{\rho(r)}$$

$$= \frac{1}{r!} \sum_{\rho, \sigma \in \mathbb{S}_{r}} (-1)^{\rho^{-1}} (-1)^{\sigma} P_{1, \sigma(\rho^{-1}(1))} P_{2, \sigma(\rho^{-1}(2))} \dots P_{r, \sigma(\rho^{-1}(r))} e_{1} e_{2} \dots e_{r}$$

$$= \sum_{\alpha \in \mathbb{S}_{r}} (-1)^{\alpha} P_{1, \alpha(1)} P_{2, \alpha(2)} \dots P_{r, \alpha(r)} e_{1} e_{2} \dots e_{r}$$

$$= \det(P) e_{1} \dots e_{r}$$

In the second line, the factor $(-1)^{\rho^{-1}}$ comes from unshuffling $e_{\rho(1)} \dots e_{\rho(r)} \to e_1 \dots e_r$, and in the third line, we use the fact that for any function f that depends only on $\sigma \circ \rho^{-1}$: $\sum_{\rho,\sigma \in \mathbb{S}_r} f(\sigma \circ \rho^{-1}) = r! \sum_{\alpha \in \mathbb{S}_r} f(\alpha)$, where α plays the role of $\sigma \circ \rho^{-1}$. Now, observe that $\det(P) = \pm 1$, and in the case where $\det(P) = -1$, we can simply reorder e_1, \dots, e_r to get rid of the negative sign.

In general, multivectors can be comprised of elements with different grades. For a set of orthogonal vectors, we say that e_i has grade 1, e_ie_j (for $i \neq j$) has grade 2, and so on. We define $\langle \cdot \rangle_r$ to be the component of a multivector with grade r, so that in general a multivector A in a geometric algebra \mathcal{G}_n (whose underlying vector space has dimension n) can be written $A = \sum_{r=0}^{n} \langle A \rangle_r$. A multivector that satisfies $\langle A \rangle_r = A$ (for some r) is called homogenous.

We denote the subspace of \mathcal{G}_n of grade r as \mathcal{G}_n^r . The dimensionality of \mathcal{G}_n^r is $\binom{n}{r}$, because the basis of \mathcal{G}_n^r can be formed by choosing r items from the n basis vectors.¹ The overall dimensionality is therefore $\sum_{r=0}^{n} \binom{n}{r} = 2^n$.

¹It is important to note that not every multivector in \mathcal{G}_n^r is a blade. The simplest nontrivial example is $e_1e_2 + e_3e_4$ in \mathcal{G}_n^4 – there is simply no way to write this as a blade because e_1, e_2 and e_3, e_4 are orthogonal to each other.

1.1 The Dot and Wedge Product

We now study the behavior of aB_r , where a is some vector and B_r is a homogenous multivector of grade r. More specifically, we will show that we can define $a \cdot B_r$ and $a \wedge B_r$ in terms of aB_r and B_r in a way that is similar to the original definition for \cdot and \wedge between two ordinary vectors.

Theorem 2. For any $a \in \mathcal{G}_n^1$ and $B_r \in \mathcal{G}_n^r$, $(aB_r - (-1)^r B_r a)$ is a homogenous multivector with grade r - 1. This motivates the definition

$$a \cdot B_r := \frac{1}{2} (aB_r - (-1)^r B_r a),$$
 (1)

so that \cdot is a grade-lowering operation.

Proof. We assume B_r is a blade, since all of the above statements are linear in B_r . That is, it suffices to show that the above statement is true for $B_r = e_1 e_2 \dots e_r$, for any choice of orthogonal vectors $\{e_1, \dots, e_r\}$. We repeatedly apply $ab = 2a \cdot b - ba$. Observe that $aB_r = 2(a \cdot e_1)e_2 \dots e_r - e_1 ae_2 \dots e_r$. We can repeatedly do this, shifting a further back in the chain, to get:

$$aB_r = 2\sum_{k=1}^r (-1)^{k+1} (a \cdot e_k)(e_1 \dots \check{e}_k \dots e_r) + (-1)^r B_r a,$$

where \check{e}_k denotes the fact that e_k is omitted from the product (i.e. $e_1\check{e}_2e_3=e_1e_3$). A simple rearrangement gives:

$$\frac{1}{2}(aB_r - (-1)^r B_r a) = \sum_{k=1}^r (-1)^{k+1} (a \cdot e_k)(e_1 \dots \check{e}_k \dots e_r)$$
 (2)

Note that in the sum, $e_1 ldots \check{e}_k ldots e_r$ is grade r-1. So, $\frac{1}{2}(aB_r-(-1)^rB_ra)$ is a linear combination of blades, each of which are grade r-1 so it is on the whole a homogenous multivector of grade r-1.

A similar result holds for the wedge product \wedge .

Theorem 3. For any blade $B_r = b_1 \wedge ... \wedge b_r$ and any vector a:

$$a \wedge b_1 \wedge \ldots \wedge b_r = \frac{1}{2}(aB_r + (-1)^r B_r a)$$

Proof. Let $e_1, \ldots e_r$ be an orthogonalization of $b_1, \ldots b_r$ using the same technique as Lemma 1, such that $b_1 \wedge \ldots \wedge b_r = e_1 \ldots e_r$. Then, let $a_{\perp} = a - \sum_{k=1}^r \beta_k e_k$, where:

$$\beta_k = \begin{cases} \frac{a \cdot e_k}{(e_k \cdot e_k)^2} & \text{if } e_k \cdot e_k \neq 0\\ 0 & \text{otherwise} \end{cases}$$

Letting $a_{\parallel} = \sum_{k=1}^{r} \beta_k e_k$, we can write $a = a_{\parallel} + a_{\perp}$, and it is now simple to see that $a_{\perp} \cdot e_k = 0$ for all $k = 1, \ldots, r$. Furthermore, since each of the e_k are merely linear combinations of b_k , $a_{\perp} \cdot b_k = 0$ for all $k = 1, \ldots, r$ as well. Therefore, $a \wedge b_1 \wedge \ldots \wedge b_r = a_{\perp} \wedge b_1 \wedge \ldots \wedge b_r$.

Now, observe that:

$$a_{\parallel}B_r + (-1)^r B_r a_{\parallel} = \sum_{k=1}^r \beta_k \left((-1)^{k+1} (e_k)^2 e_1 \dots \check{e}_k \dots e_r + (-1)^r (-1)^{r-k} e_1 \dots \check{e}_k \dots e_r (e_k)^2 \right)$$

$$= 0$$

Therefore, using the crucial fact that $a_{\perp}b_k = -b_k a_{\perp}$:

$$aB_{r} + (-1)^{r}B_{r}a = a_{\perp}B_{r} + (-1)^{r}B_{r}a_{\perp}$$

$$= \frac{1}{r!} \sum_{\sigma \in \mathbb{S}_{r}} (-1)^{\sigma} (a_{\perp}b_{\sigma(1)} \dots b_{\sigma(r)} + (-1)^{r}b_{\sigma(1)} \dots b_{\sigma(r)}a_{\perp})$$

$$= \frac{2}{r!} \sum_{\sigma \in \mathbb{S}_{r}} (-1)^{\sigma} a_{\perp}b_{\sigma(1)} \dots b_{\sigma(r)}$$

$$= \frac{2}{(r+1)!} \sum_{k=1}^{r+1} \sum_{\sigma \in \mathbb{S}} (-1)^{\sigma} a_{\perp}b_{\sigma(1)} \dots b_{\sigma(r)}$$

Now, observe that we are free to scramble a_{\perp} somewhere into the kth position of $b_{\sigma(1)} \dots b_{\sigma(r)}$ – that is, $a_{\perp}b_{\sigma(1)}\dots b_{\sigma(r)}=(-1)^kb_{\sigma(1)}\dots b_{\sigma(k-1)}a_{\perp}b_{\sigma(k)}\dots b_{\sigma(r)}$. But then, viewing a_{\perp} as the first element of the set $S=[a_{\perp},b_1,\dots,b_r]$, so that we identify $b_0=a_{\perp}$ and $S\cong[0,\dots,r]$, we see that the permutation that sends S to $[b_{\sigma(1)},\dots,b_{\sigma(k-1)},a_{\perp},b_{\sigma(k)},\dots,b_{\sigma(r)}]\cong[\sigma(1),\dots,\sigma(k-1),0,\sigma(k),\dots,\sigma(r)]$ has parity $(-1)^k(-1)^{\sigma}$! Therefore, the double sum over $\sum_{k=1}^{r+1}\sum_{\sigma\in\mathbb{S}_r}$ can be rewritten as a sum over permutations mapping $[0,\dots,r]$ onto itself – that is, \mathbb{S}_{r+1} . We are then left with:

$$aB_r + (-1)^r B_r a = \frac{2}{(r+1)!} \sum_{\sigma \in \mathbb{S}_{r+1}} b_{\sigma(0)} \dots b_{\sigma(r)}$$
$$= 2a \wedge b_1 \wedge \dots \wedge b_r,$$

as required.

Corollary 1. For any multivector $B_r \in \mathcal{G}_n^r$, let $a \wedge B_r$ be defined by:

$$a \wedge B_r := \frac{1}{2} (aB_r + (-1)^r B_r a),$$
 (3)

so that $a \wedge B_r$ is always a homogenous multivector of grade r+1. That is, \wedge raises grade by 1.

 $A_s \cdot B_r$, we mean the lowest grade component of $A_s B_r$, and the highest for $A_s \wedge B_r$:

Proof. This follows simply by decomposing B_r into a linear combination of blades and applying Theorem 3 to each blade (observing that $a \wedge B_r$ is bilinear)

It follows trivially from the definitions in Equation (1) and Equation (3) that $aB_r = a \cdot B_r + a \wedge B_r$, which says that the product aB_r is a sum of homogenous multivectors with grade $r \pm 1$. This can be generalized for a general multivector A_s . Since A_s can be written as a linear combination of blades, and each blade can be written as a product of anticommuting (orthogonal) vectors, A_sB_r is composed of multivectors with grades $|r-s|, |r-s|+2, \ldots, r+s$. To further generalize the dot and wedge product notation, when we write

$$A_s \cdot B_r = \langle A_s B_r \rangle_{|s-r|} \tag{4}$$

$$A_s \wedge B_r = \langle A_s B_r \rangle_{s+r} \tag{5}$$

We also define a scalar product $\langle A_1 A_2 \dots A_k \rangle$ to be the grade 0 (i.e., scalar) component of the product $A_1 A_2 \dots A_k$.

Lemma 4. The wedge product is associative:

$$(A_r \wedge B_s) \wedge C_t = A_r \wedge (B_s \wedge C_t),$$

so we can write $A_r \wedge B_s \wedge C_t$ unambiguously.

Proof. Recall that we can expand

$$A_r B_s = \langle A_r B_s \rangle_{r+s} + \langle A_r B_s \rangle_{r+s-2} + \ldots + \langle A_r B_s \rangle_{|r-s|}.$$

Since $(A_r \wedge B_s) \wedge C_t = \langle A_r B_s \rangle_{r+s} \wedge C_t$:

$$\langle A_r B_s \rangle_{r+s} C_t = \left(\langle A_r B_s \rangle_{r+s} C_t \right)_{r+s+t} + \left(\langle A_r B_s \rangle_{r+s} C_t \right)_{r+s+t-2} + \dots$$
$$= \langle (A_r B_s) C_t \rangle_{r+s+t} + \dots,$$

where the ... means all remaining terms have lower grade. Therefore, $(A_r \wedge B_s) \wedge C_t = \langle (A_r B_s) C_t \rangle_{r+s+t}$. Finally, by associativity of the geometric product (recall, this is one of the axioms), $\langle (A_r B_s) C_t \rangle_{r+s+t} = \langle A_r (B_s C_t) \rangle_{r+s+t} = A_r \wedge (B_s \wedge C_t)$.

Lemma 5. For any vectors a, b_1, \ldots, b_r :

$$a \cdot (b_1 \wedge \ldots \wedge b_r) = \sum_{k=1}^r (-1)^{k+1} (a \cdot b_k) b_1 \ldots \check{b}_k \ldots b_r.$$

Proof. Since

$$a \cdot (b_1 \dots b_r) = a \cdot \langle b_1 \dots b_r \rangle_r + a \cdot \langle b_1 \dots b_r \rangle_{r-2} + \dots,$$

by taking the r-1th component of each side, we find:

$$\langle a \cdot (b_1 \dots b_r) \rangle_{r-1} = a \cdot \langle b_1 \dots b_r \rangle_r$$

and since $\langle b_1 \dots b_r \rangle_r = b_1 \wedge \dots \wedge b_r$:

$$a \cdot (b_1 \wedge \ldots \wedge b_r) = \frac{1}{2} \langle ab_1 \ldots b_r - (-1)^r b_1 \ldots b_r a \rangle_{r-1}$$

Repeatedly applying the identity $ab = 2(a \cdot b) - ba$:

$$a \cdot (b_1 \wedge \ldots \wedge b_r) = \left\langle \sum_{k=1}^r (-1)^{k+1} (a \cdot b_k) b_1 \ldots \check{b}_k \ldots b_r \right\rangle_{r-1}$$
$$= \sum_{k=1}^r (-1)^{k+1} (a \cdot b_k) b_1 \ldots \check{b}_k \ldots b_r$$

Remark 1. The above result (Lemma 5) is a generalization of the vector triple product identity in \mathbb{R}^3 which says $a \times (b \times c) = (a \cdot c)b - (a \cdot b)c$. The analogy is not with $a \cdot (b \times c)$, because in three dimensions, cross product between a vector and a pseudovector $b \times c$ is actually a dot product the language of geometric algebra. This fact originates in the fact that the pseudovector $b \times c$ is really a bivector, and ordinary vector algebra lacks the language the make this fact clear. We discuss this further below.

1.2 A Few Additional Operations

We define the reverse of a product of vectors to be

$$(a_1 \dots a_r)^{\dagger} = a_r \dots a_1.$$

Therefore, for an arbitrary blade A_r , by writing $A_r = e_1 \dots e_r$, we see that $A_r^{\dagger} = (-1)^{r(r-1)/2} A_r$, since r(r-1)/2 swaps occur in order for a complete reversion $e_1 \dots e_r \to e_r \dots e_1$. If we enforce $(A+B)^{\dagger} = A^{\dagger} + B^{\dagger}$, we can easily see that for any homogenous multivector A of grade r (not just blades), $A^{\dagger} = (-1)^{r(r-1)/2} A$.

Lemma 6. The scalar product is invariant under cyclic permutation:

$$\langle A_1 A_2 \dots A_k \rangle = \langle A_2 A_3 \dots A_k A_1 \rangle = \dots = \langle A_k A_1 \dots A_{k-1} \rangle$$

We sometimes write $A * B = \langle AB \rangle$.

Proof. For any A, B, we have $\langle AB \rangle = \langle AB \rangle^{\dagger}$ since the grade of $\langle AB \rangle$ is r=0, so $(-1)^{r(r-1)/2}=1$. By the linearity of the reversion operator, $\langle AB \rangle^{\dagger} = \langle (AB)^{\dagger} \rangle$. This can be seen by writing AB as a sum of homogenous multivectors, and applying † to both sides. Since $(AB)^{\dagger} = B^{\dagger}A^{\dagger}$, we have $\langle AB \rangle = \langle B^{\dagger}A^{\dagger} \rangle$. Finally, we observe that if the grade of A and B is r and s respectively, the lowest grade component of AB is |r-s|. Therefore, if $r \neq s$, $\langle AB \rangle = 0 = \langle BA \rangle$. If r=s, then $\langle AB \rangle = \langle B^{\dagger}A^{\dagger} \rangle = (-1)^{r(r-1)} \langle BA \rangle = \langle BA \rangle$. So, it holds in general that $\langle AB \rangle = \langle BA \rangle$.

Now, for any m, we can choose $A = A_1 \dots A_m$ and $B = A_{m+1}A_{m+2} \dots A_k$ to find $\langle A_1 \dots A_k \rangle = \langle A_{m+1}A_{m+2} \dots A_k A_1 \dots A_m \rangle$ for any $m = 1, \dots, k$. This gives the desired result.

Another product is the commutator

$$A \times B = \frac{1}{2}(AB - BA).$$

Note that for a bivector B and a vector a, the commutator is equal to the dot product:

$$B\times a=\frac{1}{2}(Ba-aB)=-a\cdot B=-\left\langle aB\right\rangle _{1}=-\left\langle B^{\dagger}a^{\dagger}\right\rangle _{1}=\left\langle Ba\right\rangle _{1}=B\cdot a$$

It turns out that this is a specific case of a more general property, which we state below.

Lemma 7. For any homogenous multivector A_r , $B \times A_r$ is also a homogenous multivector of grade r. Specifically, $B \times A_r = \langle BA_r \rangle_r$.

Proof. We decompose $BA_r = \langle BA_r \rangle_{r-2} + \langle BA_r \rangle_r + \langle BA_r \rangle_{r+2}$. Now, we reverse both sides.

$$-(-1)^{r(r-1)/2}A_rB = (-1)^{(r-2)(r-3)/2} \langle BA_r \rangle_{r-2} + (-1)^{r(r-1)/2} \langle BA_r \rangle_r + (-1)^{(r+1)(r+2)/2} \langle BA_r \rangle_{r+2} - A_rB = (-1)^{-2r+3} \langle BA_r \rangle_{r-2} + \langle BA_r \rangle_r + (-1)^{4r+1} \langle BA_r \rangle_{r+2}$$

$$A_rB = \langle BA_r \rangle_{r-2} - \langle BA_r \rangle_r + \langle BA_r \rangle_{r+2}$$

Therefore, it is simple to see that $\frac{1}{2}(BA_r - A_rB) = \langle BA_r \rangle_r$.

Lemma 7 gives a useful identity:

$$BA_r = B \cdot A_r + B \times A_r + B \wedge A_r$$

Finally, in spaces of a finite dimension n, all multivectors with grade n are unique up to a scalar multiplier. We define the pseudoscalar I to be the unique multivector of grade n such that $|I^2| = 1$. The pseudoscalar defines an orientation for the space – that is, any blade will have either the same or different sign as I, which corresponds to a positive or negative orientation, respectively. The pseudoscalar is useful because it defines a duality transformation mapping blades to their orthogonal component $A_r \to IA_r$. Note that IA_r has grade n-r because the r basis vectors from A_r combine with r components of I, leaving n-r basis vectors. Because of this duality transformation, as one might expect, it can also be used to transform between the dot and wedge product.

²The sign of I^2 depends on the dimension n and the signature of the space.

Lemma 8. For any multivectors A_r, B_s such that $r + s \le n$:

$$A_r \cdot (B_s I) = A_r \wedge B_s I.$$

Proof.

$$A_r \cdot (B_s I) = \langle A_r B_s I \rangle_{|r-(n-s)|}$$

$$= \langle A_r B_s I \rangle_{n-(r+s)}$$

$$= \langle A_r B_s \rangle_{r+s} I$$

$$= A_r \wedge B_s I$$

1.3 Rotations and Reflections

Vectors can be reflected with respect to some direction (i.e. unit vector) n. This reflection ought to take $n \to -n$, and leave all vectors perpendicular to n unchanged. Finally, it ought to be linear. The map

$$f(a) = -nan$$

does exactly this. We see that f(n) = -nnn = -n, and for vectors n_{\perp} orthogonal to n, $f(n_{\perp}) = -nn_{\perp}n = nnn_{\perp} = nnn_{\perp} = n_{\perp}$. This also ensures that for any vector a, f(a) is still a vector.

Vectors can also be rotated with respect to some plane. This plane is described by two unit basis vectors n and m, or by the bivector $n \wedge m$. Two successive reflections about vectors m then n results in a rotation. That is,

$$f(a) = nmamn$$

is a rotation. It is trivial to see that this leaves vectors perpendicular to both n and m unchanged. We also verify that $f(a) \cdot f(b) = a \cdot b$, as required by definition of a rotation.

$$f(a) \cdot f(b) = \frac{1}{2}(nmamnnmbmn + nmbmnnmamn)$$
$$= \frac{1}{2}(nmabmn + nmbamn)$$
$$= nm(a \cdot b)mn$$
$$= a \cdot b$$

We call the product R = nm a rotor, and therefore rotors generate rotations via RaR^{\dagger} . Since rotations form a group, composing rotations $a \to R_1 R_2 a R_2 R_1$ results in another rotation. Therefore, we say that $R_1 R_2$ is also a rotor. That is, rotors also form a group. In general, rotors are simply geometric products of an even number of vectors with the normalization condition $RR^{\dagger} = 1$.

This description of rotation leads to interesting behavior on multivectors. For instance, take some geometric product $a_1
ldots a_k$. If we rotate each of the component vectors, we get $(Ra_1R^{\dagger})(Ra_2R^{\dagger})\dots(Ra_kR^{\dagger})=Ra_1\dots a_kR^{\dagger}$. We get similar behavior for a blade (since the blade is just a linear combination of geometric products). Therefore, the rotation transformation for arbitrary multivectors A is still RAR.

1.4 Frames

A set of linearly independent vectors $\{e_1, \dots, e_n\}$ is called a frame in geometric algebra. We do not necessarily assume these are orthonormal vectors, so they do not (in general) commute. We define the volume element

$$E_n = \mathbf{e}_1 \wedge \ldots \wedge \mathbf{e}_n$$

and the reciprocal frame $\left\{\mathbf{e}^{1},\ldots,\mathbf{e}^{n}\right\}$ by

$$\mathbf{e}_j \cdot \mathbf{e}^k = \delta_j^k. \tag{6}$$

We can write an explicit form for the reciprocal vectors:

$$\mathbf{e}^k = (-1)^{k+1} (\mathbf{e}_1 \wedge \ldots \wedge \check{\mathbf{e}}_k \wedge \ldots \mathbf{e}_n) E_n^{-1}.$$

Recall the $\check{\mathbf{e}}_k$ means that this term is omitted. Intuitively, what this says is that \mathbf{e}^k is orthogonal to each basis vector except \mathbf{e}_k , and we add the inverse volume element E_n^{-1} at the end to normalize. More rigorously, we can verify that this definition satisfies Equation (6) by making use of the duality transformation in Lemma 8.

$$\mathbf{e}_{j} \cdot \mathbf{e}^{k} = \mathbf{e}_{j} \cdot \left((-1)^{k+1} (\mathbf{e}_{1} \wedge \ldots \wedge \check{\mathbf{e}}_{k} \wedge \ldots \mathbf{e}_{n}) E_{n}^{-1} \right)$$
$$= (-1)^{k+1} \mathbf{e}_{j} \wedge \mathbf{e}_{1} \wedge \ldots \wedge \check{\mathbf{e}}_{k} \wedge \ldots \mathbf{e}_{n} E_{n}^{-1}$$

Note that if $j \neq k$, the wedge product vanishes. Otherwise, $(-1)^{k+1} \mathbf{e_j} \wedge \mathbf{e_1} \wedge \ldots \wedge \check{\mathbf{e}_k} \wedge \ldots \mathbf{e_n} = E_n$. Therefore, we recover our desired result $\mathbf{e_j} \cdot \mathbf{e^k} = \delta_i^k$.

Once we have chosen a frame for our space, we can write vectors as linear combinations of the frame basis vectors:

$$a = a^i \mathbf{e}_i = a_i \mathbf{e}^i,$$

where we are using Einstein summation notation. By the reciprocal property in Equation (6), we can see that $a^i = a \cdot \mathbf{e}^i$ and $a_i = a \cdot \mathbf{e}_i$. Therefore, $a \cdot \mathbf{e}_i \mathbf{e}^i = a \cdot \mathbf{e}^i \mathbf{e}_i = a$ (recall that the order of operations is \cdot and \wedge before the geometric product). This is an instance of a more general fact.

Lemma 9. For any frame $\{e_i\}$ and any homogenous multivector A_r ,

$$\mathbf{e}_i \mathbf{e}^i \cdot A_r = r A_r.$$

Proof. We apply Equation (2). We also assume A_r is a blade that can be written as a product of orthogonal vectors $\mathbf{a}_1 \mathbf{a}_2 \dots \mathbf{a}_r$ (the general case follows trivially because the desired identity is linear in A_r).

$$\mathbf{e}_{i}\mathbf{e}^{i} \cdot A_{r} = \mathbf{e}_{i} \sum_{k=1}^{r} (-1)^{k+1} (\mathbf{e}^{i} \cdot \mathbf{a}_{k}) (\mathbf{a}_{1} \dots \check{\mathbf{a}}_{k} \dots \mathbf{a}_{r})$$

$$= \sum_{k=1}^{r} (-1)^{k+1} \mathbf{a}_{k} \mathbf{a}_{1} \dots \check{\mathbf{a}}_{k} \dots \mathbf{a}_{r}$$

$$= \sum_{k=1}^{r} \mathbf{a}_{1} \dots \mathbf{a}_{r}$$

$$= r A_{r}$$

Since $\{\mathbf{e}_1, \dots, \mathbf{e}_n\}$ spans \mathbb{R}^n , any blade of degree r can be written as a linear combination of wedge products $\mathbf{e}_i \wedge \dots \wedge \mathbf{e}_j \wedge \mathbf{e}_k$. This holds similarly for the reciprocal frame $\{\mathbf{e}^1, \dots, \mathbf{e}^n\}$. Therefore, any homogenous multivector A of grade r can be expanded:

$$A = \sum_{i_1 < i_2 < \dots < i_r} A_{i_1 \dots i_r} \mathbf{e}^{i_1} \wedge \dots \mathbf{e}^{i_2} \wedge \mathbf{e}^{i_r}$$

An explicit formula for the components $A_{i_1...i_r}$ can be given:

$$A_{i_1...i_r} = \langle (e_{i_r} \wedge \ldots \wedge e_{i_1})A \rangle$$

References

Doran, Chris, and A. N. Lasenby. 2007. *Geometric Algebra for Physicists*. 1st ed. Cambridge; New York: Cambridge University Press.