

The Loop-Induced off-shell Higgs process

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Outline

- 1. Loop-induced processes - MadLoop**
- 2. Off-shell Higgs in the SM**
- 3. BSM off-shell Higgs**

Part 1

Loop induced processes and MadLoop

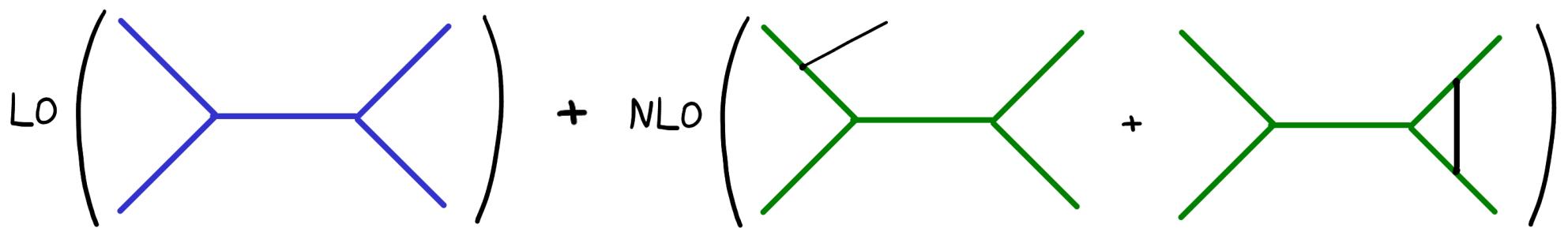
1. NLO calculations

Hadron Collisions:

$$\sigma_{\text{tot}}(\mu_F, \mu_R) = \sum_{i,j} \int_0^1 dx_1 \int_0^1 dx_2 d\Pi_n f_i(x_1, \mu_f) f_j(x_2, \mu_f) \hat{\sigma}_{ij}(\hat{s}, \mu_F, \mu_R)$$

Perturbative Expansion:

$$\hat{\sigma} = \underbrace{\sigma^{\text{Born}}}_{\text{LO}} \left(1 + \underbrace{\frac{\alpha_s}{2\pi} \sigma^{(1)}}_{\text{NLO}} + \underbrace{\left(\frac{\alpha_s}{2\pi}\right)^2 \sigma^{(2)}}_{\text{NNLO}} + \underbrace{\left(\frac{\alpha_s}{2\pi}\right)^3 \sigma^{(3)}}_{\text{N3LO}} + \dots \right)$$



1. NLO calculations

**Automated one-loop
matrix element tools**

- **MadLoop**
- FeynArts, FormCalc, ...
- OpenLoops
- GoSam

3-step program to automate NLO

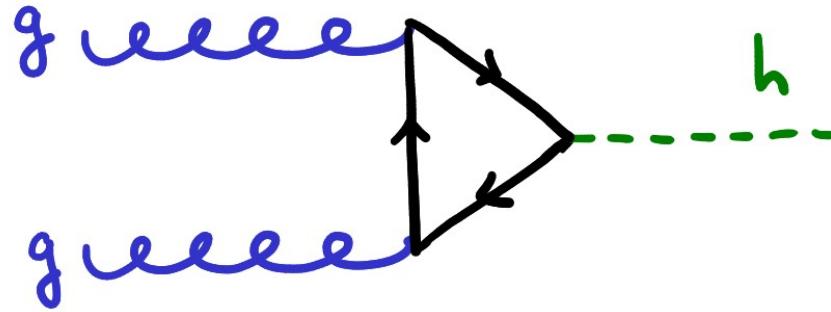
1. Renormalization
2. Integral reduction
3. Rational terms

This seminar: Limit to a particular case

1. Loop-Induced (LI) Processes

LI process → No tree-level contributions

i.e.1 - Gluon Fusion dominates the inclusive Higgs production.



i.e.2 - L.I. $gg \rightarrow ZZ$ is 60% of the full NNLO corrections to had. ZZ production.

i.e.3 - BSM models → Loop suppressed production helps to evade experimental constraints

1. Loop-Induced (LI) Processes

Dealing with LI processes

- Integrate out heavy loops → Effective point-like vertices (i.e. heft)
  **Valid in only a limited kinematic range.**
- Automated implementation of LI processes.
  **Done in 2015! (MG5@NLO). [1507.00020]**
- Dedicated implementations (i.e. $pp \rightarrow h$ in N3LO QCD + NLO EW).
  **Used in state-of-the-art calculations.**

1. Loop-Induced (LI) Processes

LI versus general NLO

- **Advantage:** LI amplitudes are finite!

 **No UV divergences if the model is renormalizable**

- **Disadvantage:** Speed

In usual NLO, bulk contributions come from **born** and **real** emissions.
→ Limited statistics to the evaluation of virtual contributions.

Computation
time

\propto

Execution
speed of the LI
matrix element

1. Loop-Induced (LI) Processes

Loop-integrals reduction

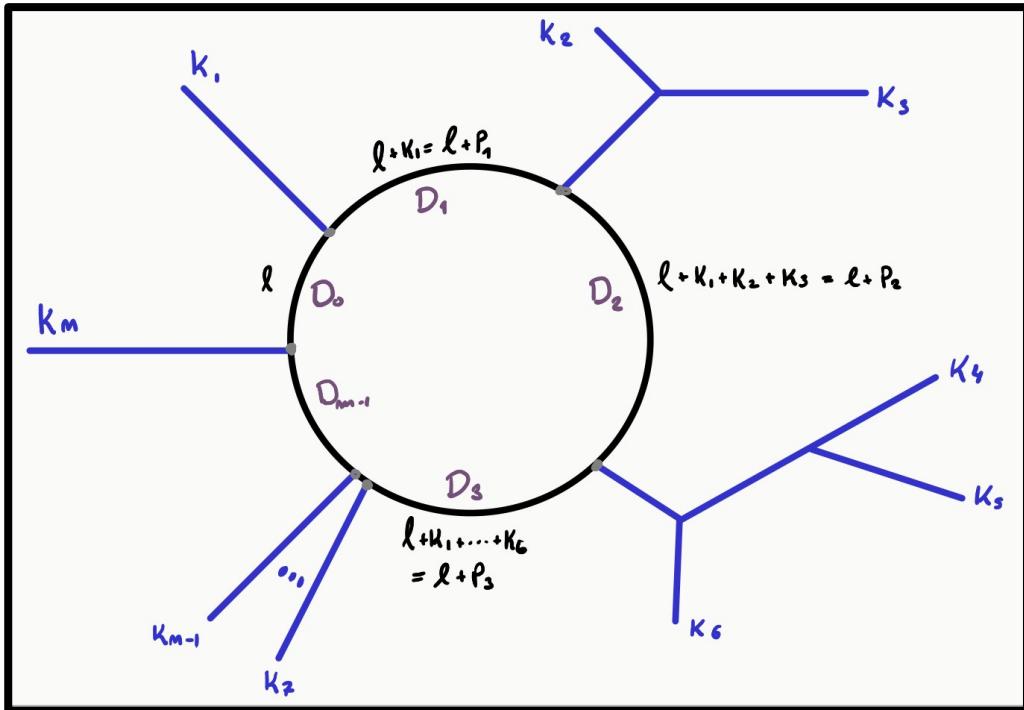
$$I^{(t)} = \sum_i d_i \text{ box}_i + \sum_i c_i \text{ triangle}_i + \sum_i b_i \text{ bubble}_i + \sum_i a_i \text{ tadpole}_i + R$$

Scalar integral basis - available in many libraries i.e. OneLoop, QCDDLoop,...
→ Loop amplitude computation reduces to determining d_i, c_i, b_i, a_i, R

- Methods:**
1. **OPP reduction:** Ossola, Papadopoulos and Pittau.
→ Reduction at the integrand level.
Tools: [CutTools](#), [Ninja](#)
 2. **TIR:** Tensor Integral Reduction.
→ Passarino-Veltman integrals.
Tools: [Collier](#), [PJFry](#), [GOLEM](#)

Question: **which one is better?** (for LI)

1. Loop-Induced (LI) Processes



d-dimensional Amplitude

$$\mathcal{A}_d = \sum_{l_1=1}^L \lambda_{l_1} \int d^d \bar{\ell} \frac{\mathcal{N}_{h,l_1}(\ell)}{\prod_{i=1}^{n_{l_1}} \bar{D}_{i,l_1}}$$

Split the numerator:

$$\overline{\mathcal{N}}(\bar{\ell}) = \underbrace{\mathcal{N}(l)}_{\text{4d part}} + \epsilon \tilde{\mathcal{N}}(\bar{\ell})$$

$$\mathcal{A}_d = \mathcal{A} + \mathcal{A}_{R2}$$

→ Finite contribution

Rational terms: R2 counterterms in UFO (ask me later - Backup slides)

1. Loop-Induced (LI) Processes

$$\mathcal{N}(\ell)_{l,h} = \sum_{r=0}^{r_{max}} C_{\mu_1 \dots \mu_r; h, l}^{(r)} \ell^{\mu_1} \dots \ell^{\mu_r}$$

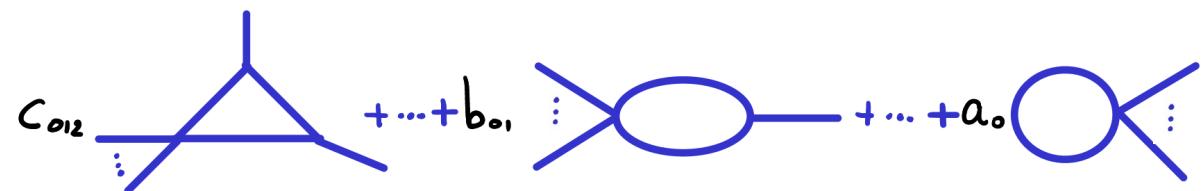
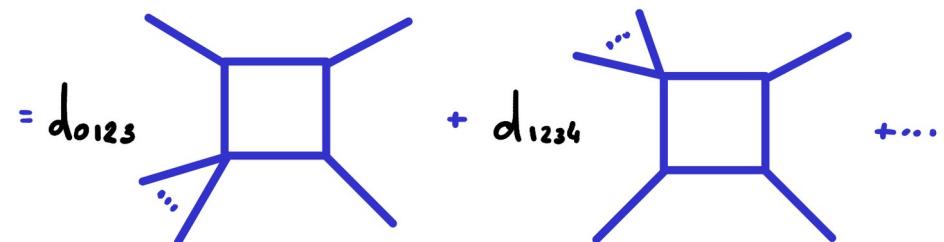
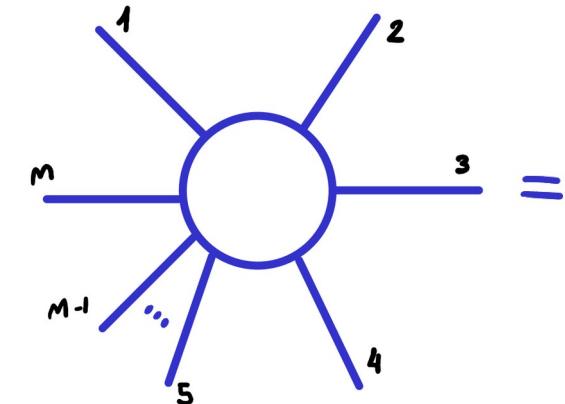
Reduction operation:

$$\text{Red} \left[\frac{\mathcal{N}_{l,h}(\ell)}{\prod_{i=0}^{n_l} \bar{D}_{i,l}} \right] = \begin{cases} \text{OPP} \left[\frac{\sum_{r=0}^{r_{max}} C_{\mu_1 \dots \mu_r; h, l}^{(r)} \ell^{\mu_1} \dots \ell^{\mu_r}}{\prod_{i=1}^{n_l} \bar{D}_{i,l}} \right] \\ \sum_{r=0}^{r_{max}} C_{\mu_1 \dots \mu_r; h, l}^{(r)} \text{TIR} \left[\frac{\ell^{\mu_1} \dots \ell^{\mu_r}}{\prod_{i=1}^{n_l} \bar{D}_{i,l}} \right] \end{cases}$$

1. Brief overview of OPP method

OPP: Decompose the numerator into a **basis of products of the denominators**

$$\begin{aligned}
 N(q) = & \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} [d(i_0 i_1 i_2 i_3) + \tilde{d}(q; i_0 i_1 i_2 i_3)] \prod_{i \neq i_0, i_1, i_2, i_3}^{N_t-1} D_i \\
 & + \sum_{i_0 < i_1 < i_2}^{N_t-1} [c(i_0 i_1 i_2) + \tilde{c}(q; i_0 i_1 i_2)] \prod_{i \neq i_0, i_1, i_2}^{N_t-1} D_i \\
 & + \sum_{i_0 < i_1}^{m-1} [b(i_0 i_1) + \tilde{b}(q; i_0 i_1)] \prod_{i \neq i_0, i_1}^{N_t-1} D_i \\
 & + \sum_{i_0}^{m-1} [a(i_0) + \tilde{a}(q; i_0)] \prod_{i \neq i_0}^{N_t-1} D_i \\
 & + \tilde{P}(q) \prod_i^{N_t-1} D_i.
 \end{aligned}$$



1. Brief overview of OPP method

Algorithm: Sample the numerator by choosing special values for the loop momenta

1. Choose l^\pm such as $D_0(l^\pm) = \dots = D_3(l^\pm) = 0$

→ All terms in the OPP cancel except:

$$N(l^\pm) = [d_{0123} + \tilde{d}_{0123}(l^\pm)] \prod_{i \neq 0,1,2,3}^{n-1} D_i(l^\pm)$$

$$d_{0123} = \frac{1}{2} \left[\frac{N(l^+)}{\prod_{i \neq 0,1,2,3}^{n-1} D_i(l^+)} + \frac{N(l^-)}{\prod_{i \neq 0,1,2,3}^{n-1} D_i(l^-)} \right]$$

2. Compute all box coefficients.
3. Triple cuts to determine triangle coefficients. Use box coefficients from step 2.
4. Proceed to Bubbles, Tadpole and R1.

1. Brief overview of TIR method

TIR Example:

$$\mathcal{I}_{p_1, p_2}^{(1)\mu} = \int d^d \bar{q} \frac{\bar{q}^\mu}{\bar{q}^2(\bar{q} + p_1)^2(\bar{q} + p_2)^2}$$

$$\boxed{\mathcal{I}_{p_1, p_2}^{(1)\mu} = C_1 p_1^\mu + C_2 p_2^\mu}$$

- Contract with external momentum: $[2q \cdot p_1] = \int d^d \bar{q} \frac{2q \cdot p_1}{\bar{q}^2(\bar{q} + p_1)^2(\bar{q} + p_2)^2}$

$$\begin{pmatrix} [2q \cdot p_1] \\ [2q \cdot p_2] \end{pmatrix} = \begin{pmatrix} 2p_1 \cdot p_1 & 2p_1 \cdot p_2 \\ 2p_1 \cdot p_2 & 2p_2 \cdot p_2 \end{pmatrix} \begin{pmatrix} C_1 \\ C_2 \end{pmatrix}$$

- Complete squares of $[2q \cdot p_i]$ integrals: *i.e.* $[2q \cdot p_1] = \mathcal{I}_{p_2}^{(0)} - \mathcal{I}_{p_2-p_1}^{(0)} - p_1^2 \mathcal{I}_{p_1, p_2}^{(0)}$
- Invert the system of equations to obtain C_1 and C_2 .*

1. Loop-Induced (LI) Processes

Which to use?

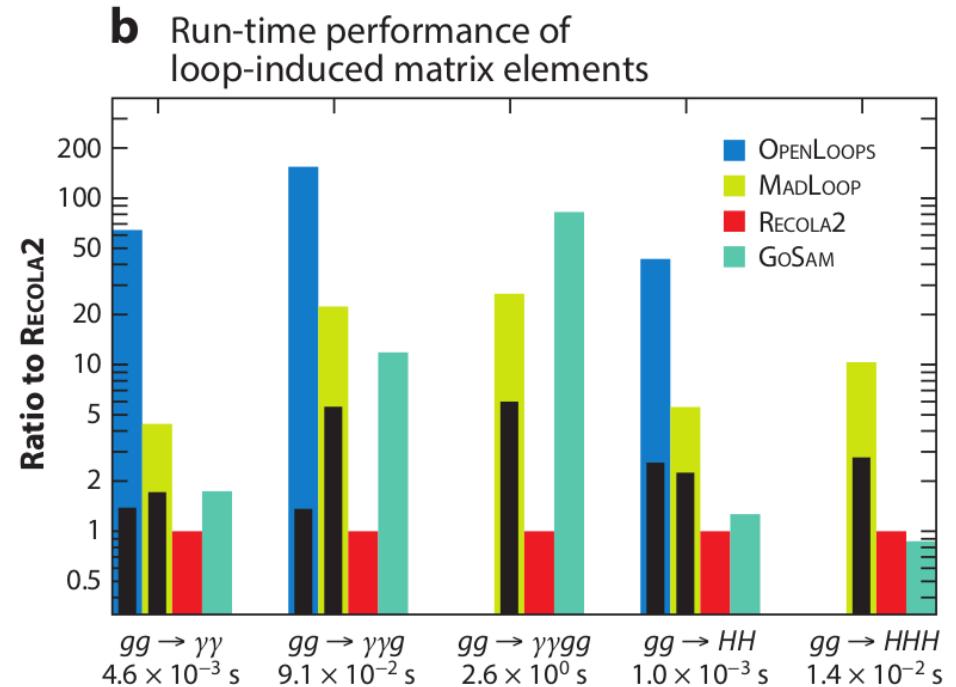
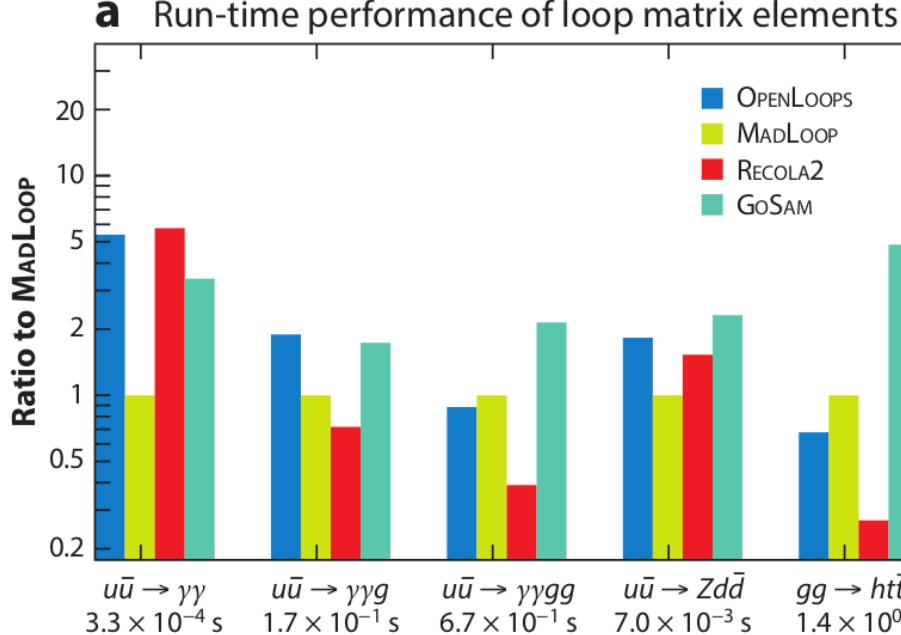
- TIR has a factorial growth in terms of the recursion depth of the algorithm, proportional to both the rank and # of external legs.
- Output of the **OPP reduction depends on both the loop and helicity**, while **TIR depends only on the loop considered**.

OPP recursions = $L \times H$

TIR recursions = L

TIR is usually better for LI processes.

1. Loop-Induced (LI) Processes



Helicity sum	Monte-Carlo	Exact	
Loop Reduction	CUTTOOLS	CUTTOOLS	TIR
Survey			
$pp \rightarrow h j$	13m (125k)	32m (260k)	9m (260k)
$pp \rightarrow h jj$	2d4h (1.2M)	16d10h (5.4M)	9d13h (5.4M)*
$gg \rightarrow zz$	1h06m (34k)	12h50m (255k)	1h44m (255k)
$gg \rightarrow zhg$	11h13m (110k)	1d8h (516k)	1d4h (516k)*

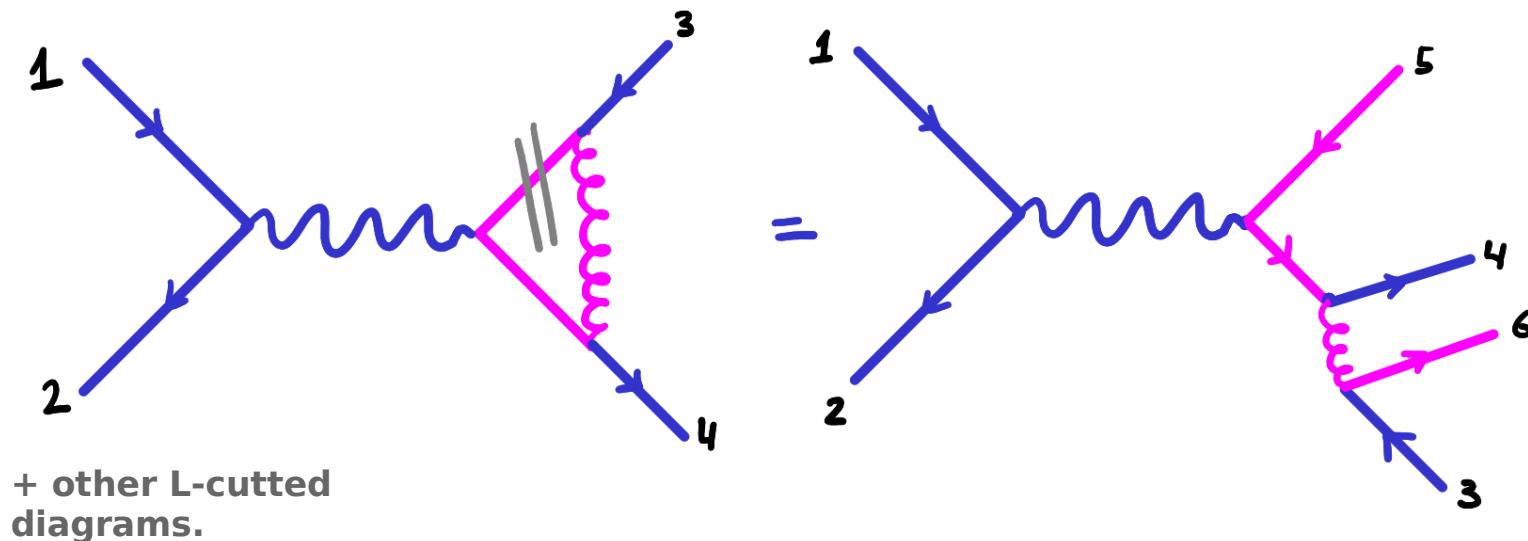
Helicity sum	Monte-Carlo	Exact	
Loop Reduction	CUTTOOLS	CUTTOOLS	TIR
Refine			
$pp \rightarrow h j$	1h43m (385k)	23m (431k)	6m (431k)
$pp \rightarrow h jj$	7d17h (2.18M)	75d1h (20.6M)	51d19h (20.6M)*
$gg \rightarrow zz$	7h20m (407k)	4d13h (4.55M)	23h07m (5.78M)
$gg \rightarrow zhg$	23h03m (277k)	2d22h (1.13M)	3d14h (1.4M)*

1. Loop-Induced (LI) Processes

Event Generation: MadLoop in Madgraph

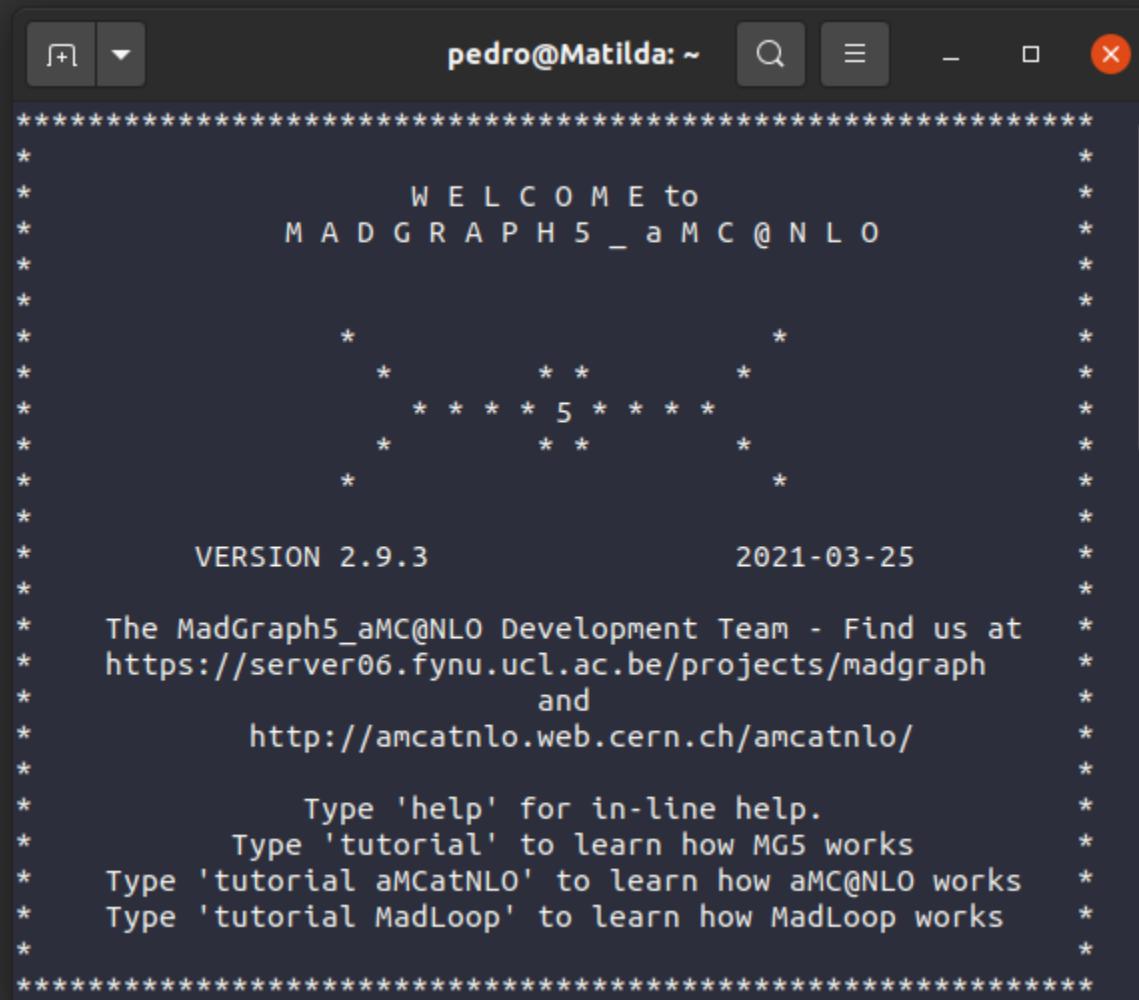
- Madgraph is great at constructing tree-level matrix elements.
- **MadLoop:** generate loop diagrams from tree-level ones.

L-cut diagrams:



- Madgraph generate all L-cutted diagrams.
- **Sew them together to reconstruct loop topology.**

Examples and Madloop options



```
pedro@Matilda: ~
```

```
*****
*          W E L C O M E to
*          M A D G R A P H 5 _ a M C @ N L O
*
*
*          *
*          *           *           *
*          *           * *           *
*          *   * * * 5 * * * * *
*          *           * *           *
*          *           *           *
*          *
*
*          VERSION 2.9.3           2021-03-25
*
*          The MadGraph5_aMC@NLO Development Team - Find us at
*          https://server06.fynu.ucl.ac.be/projects/madgraph
*          and
*          http://amcatnlo.web.cern.ch/amcatnlo/
*
*          Type 'help' for in-line help.
*          Type 'tutorial' to learn how MG5 works
*          Type 'tutorial aMCatNLO' to learn how aMC@NLO works
*          Type 'tutorial MadLoop' to learn how MadLoop works
*
*****
```

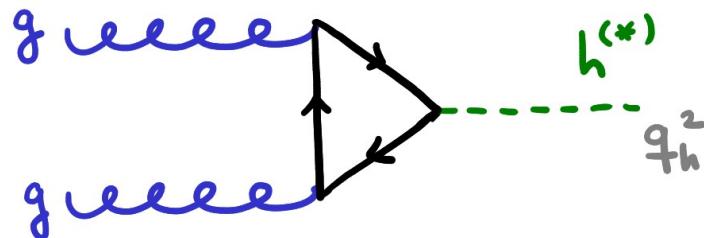
Part 2

Off-Shell Higgs in the Standard Model

2. Off-shell Higgs in the SM

$$\underline{gg \rightarrow h^{(*)} \rightarrow Z^{(*)}Z^{(*)} \rightarrow 4l}$$

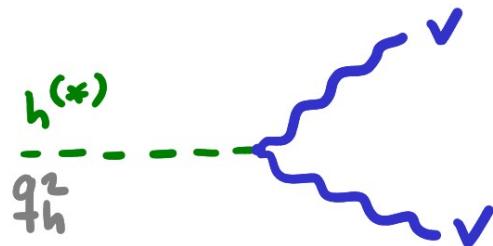
Production



$$\mathcal{L} = \frac{\alpha_s}{4\pi v} g_{ggh}(q_h^2) H G_{\mu\nu} G^{\mu\nu}$$

2 → Triangle loop.

Decay



$$\sum_{\lambda} |\mathcal{M}_{h \rightarrow VV}|^2 = \frac{4M_V^4}{v^2} \left(3 + \frac{q_h^4}{4M_V^4} - \frac{q_h^2}{M_V^2} \right)$$

Expectation: (on-shell production) \times BR should reproduce the total cross-section with good accuracy.

2. Off-shell Higgs in the SM

Expectation: **(on-shell production) x BR** should reproduce the total cross-section with good accuracy.

Reason: We expect the **Narrow Width approximation** (NWA) to be good for the Higgs.

$$m_h = 125 \text{ GeV}$$

$$\Gamma_h = 4.03 \text{ MeV}$$

$$\frac{\Gamma_h}{m_h} \sim 10^{-4}$$

$$D_h^2(q_h^2) = \frac{1}{(q_h^2 - m_h^2)^2 + \Gamma_h^2 m_h^2} \xrightarrow{\text{NWA}} \frac{\pi}{m_h \Gamma_h} \delta(q_h^2 - m_h^2) + \mathcal{O}(\Gamma_h/m_h)$$

2. Off-shell Higgs in the SM

$$\sigma = \frac{1}{2s} \int_{q_{min}^2}^{q_{max}^2} \frac{dq^2}{2\pi} \left(\int d\phi_p |\mathcal{M}_p(q_h^2)|^2 D_h^2(q_h^2) \int d\phi_p |\mathcal{M}_d(q_h^2)|^2 \right)$$

$$\sigma^{\text{NWA}} = \frac{1}{2s} \left(\int d\phi_p |\mathcal{M}_p(q_h^2)|^2 \right) \frac{1}{2m_h \Gamma_h} \left(\int d\phi_p |\mathcal{M}_d(q_h^2)|^2 \right)$$

10^{-4} error... **Or is it?...**

2. Off-shell Higgs in the SM

$$\sigma = \frac{1}{2s} \int_{q_{min}^2}^{q_{max}^2} \frac{dq^2}{2\pi} \left(\int d\phi_p |\mathcal{M}_p(q_h^2)|^2 D_h^2(q_h^2) \int d\phi_p |\mathcal{M}_d(q_h^2)|^2 \right)$$

$$\sigma^{\text{NWA}} = \frac{1}{2s} \left(\int d\phi_p |\mathcal{M}_p(q_h^2)|^2 \right) \frac{1}{2m_h \Gamma_h} \left(\int d\phi_p |\mathcal{M}_d(q_h^2)|^2 \right)$$

10^{-4} error... **Or is it?...**

Assumptions of the NWA approximation

1. Total width is much smaller than the resonance mass. **Ok**
2. There are no relevant thresholds. **False!**
3. There is no significant interference with non-resonant processes. **False!**

2. Off-shell Higgs in the SM

Assumption 2: Thresholds.

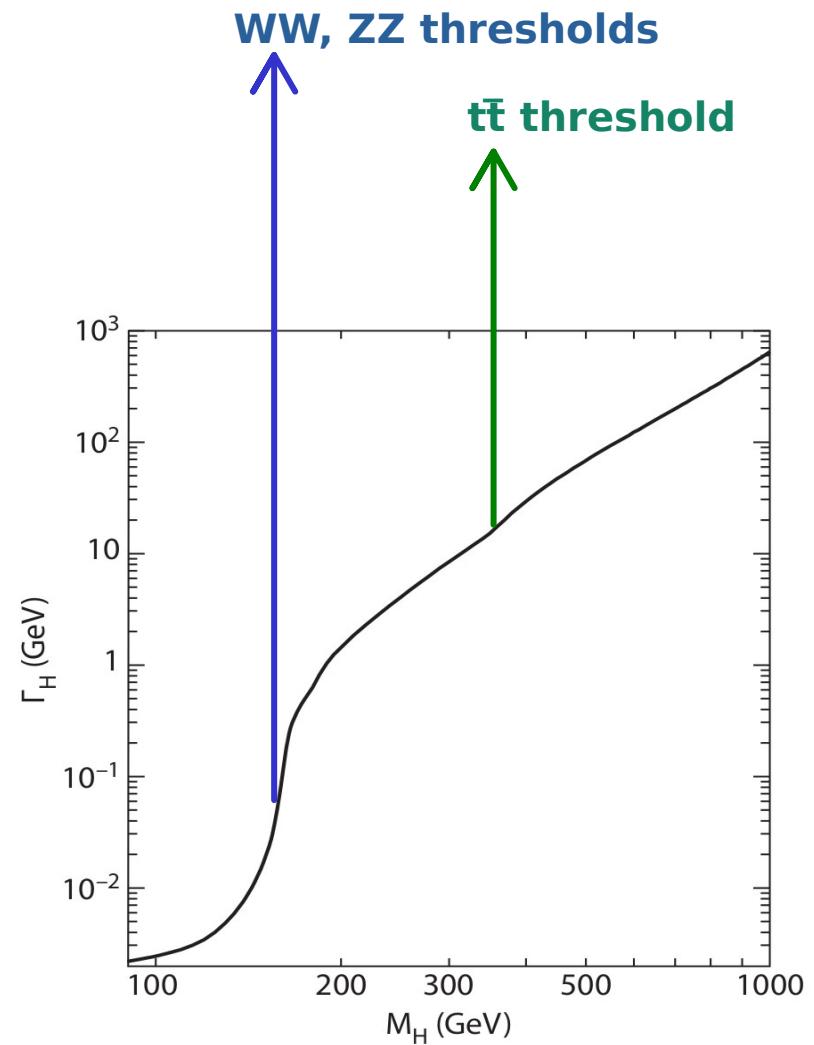
Exact propagator:

$$D_h(q^2) = \frac{i}{q^2 - m_0^2 + \Sigma(q^2)} = \frac{iZ}{q^2 - m_h^2 + iIm\tilde{\Sigma}(q^2)}$$

Use the optical theorem.

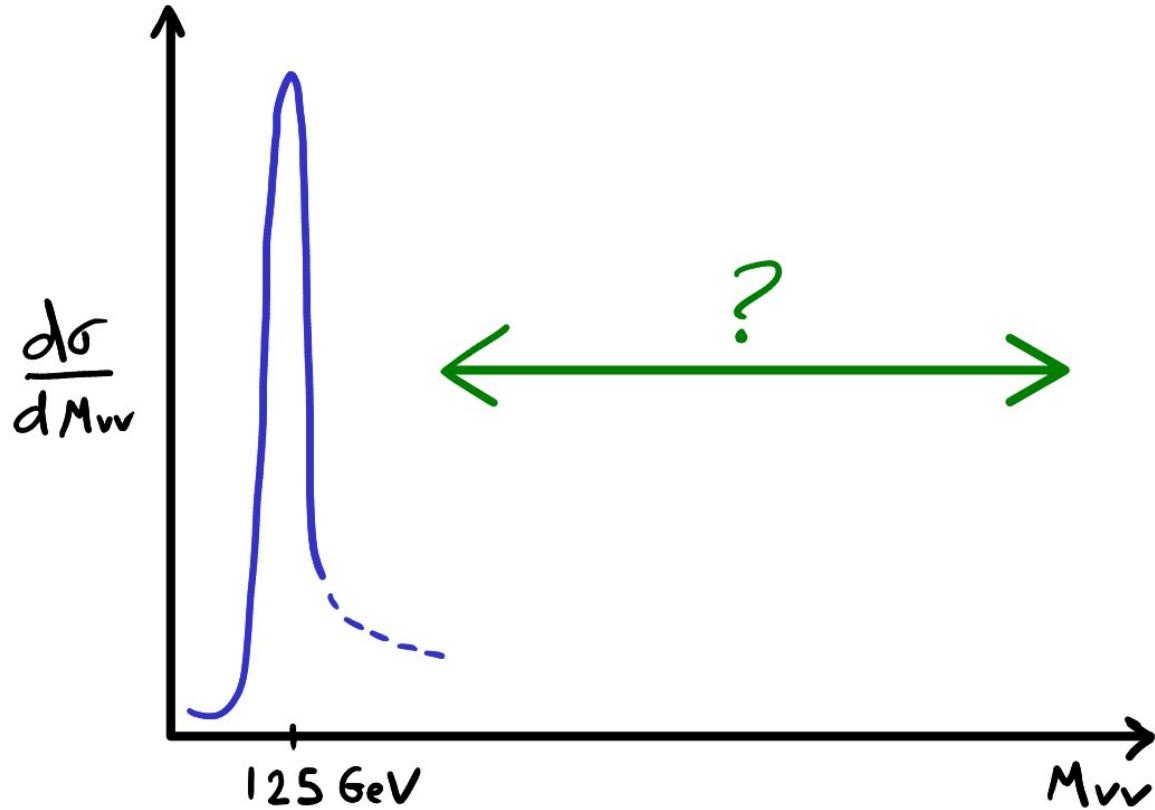
$$\begin{aligned}\Gamma_{X^* \rightarrow 1+2+\dots} &= \frac{1}{2\sqrt{s}} \sum_a |\mathcal{M}|^2 d\Pi_n \\ &= \frac{iIm\tilde{\Sigma}(s)}{\sqrt{s}}\end{aligned}$$

$$D_h(q^2) = \frac{iZ}{q^2 - m_h^2 + i\sqrt{s}\Gamma(s)}$$



2. Off-shell Higgs in the SM

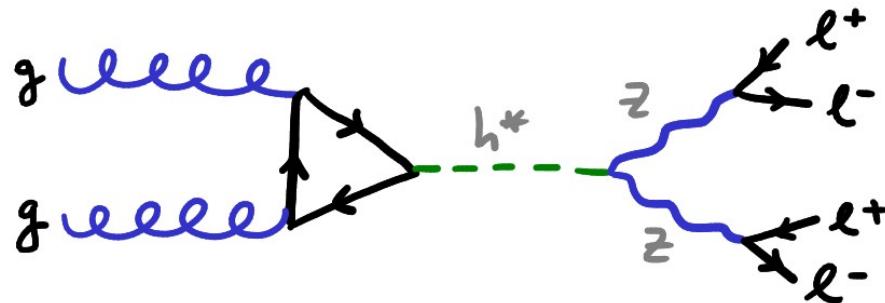
Question: What is the behaviour of the distribution at the high energy regime?



- How small compared to the peak?
- Does it really decreases with M_{vv} ?
- Is there a plateau?
- How does that affects the total cross-section?

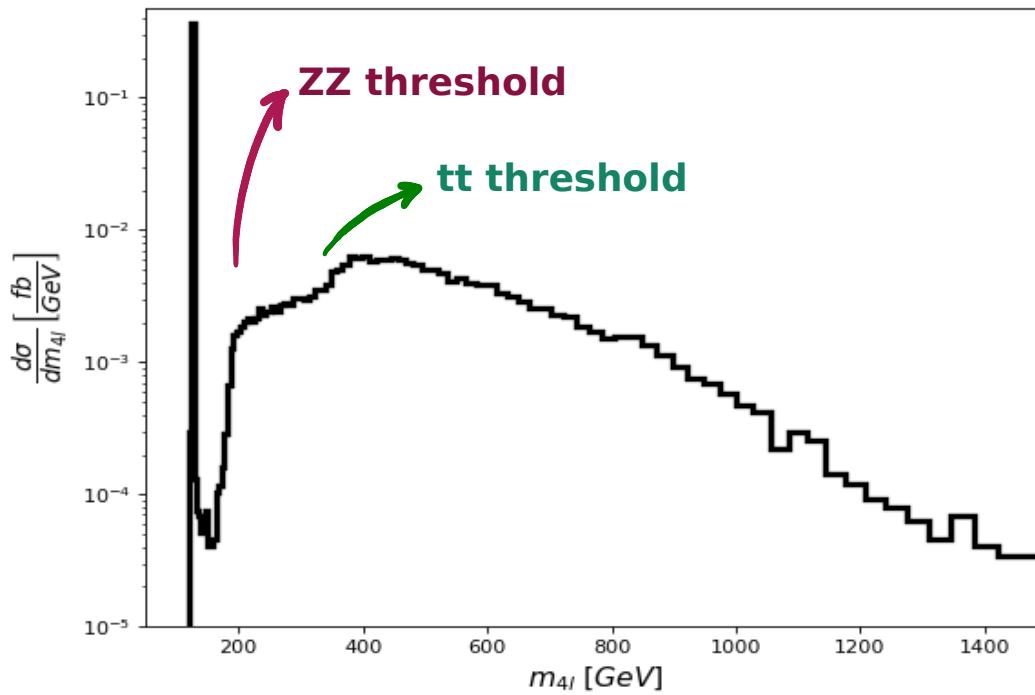
2. Off-shell Higgs in the SM

Signal



```
MG5_aMC>generate g g > l+ l- l+ l- [QCD]
-----
No Born diagrams found. Now switching to the loop-induced
Please cite ref. 'arXiv:1507.00020' when using results fro
-----
3 processes with 20 diagrams generated in 3.430 s
Total: 3 processes with 20 diagrams
```

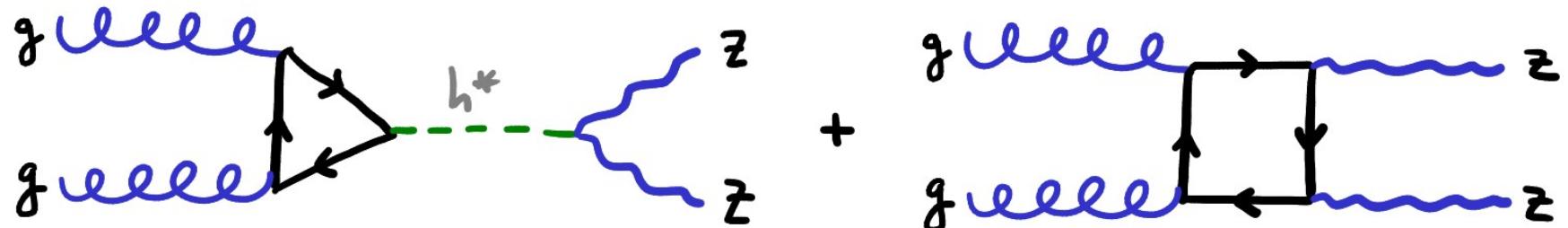
Modify loop_SM UFO → set Zqq coupling to zero to remove irreducible bg (Box diagrams)



~15% net-effect in
the total cross-section
of $M_{zz} > 2M_z$ region.

2. Off-shell Higgs in the SM

Assumption 3: Interference with non-resonant processes.



Box must be included: Irreducible background with important interference.

Unitarizing nature of the Higgs $h^* \rightarrow Z_L Z_L$

$$m_{ZZ} \gg m_t, m_h, m_Z$$

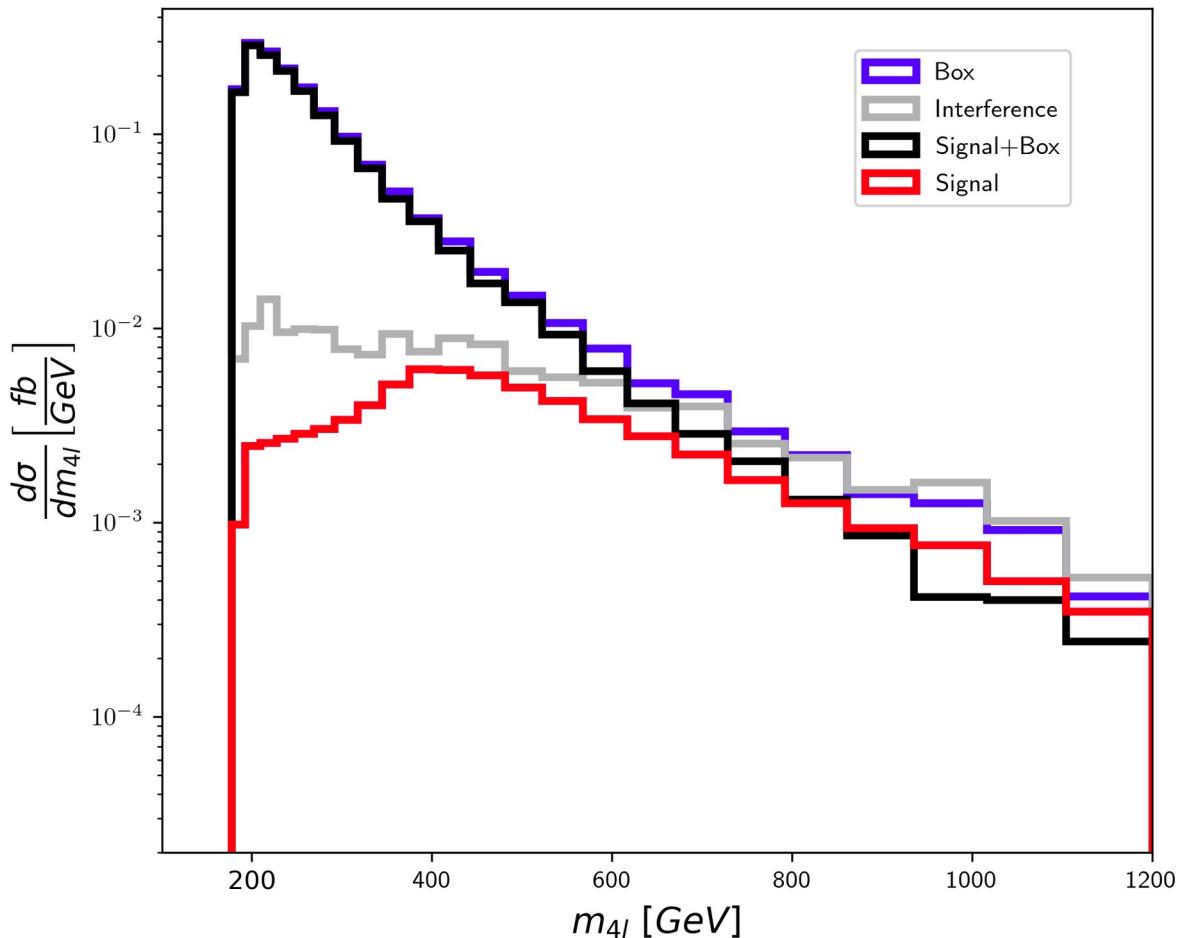
$$\mathcal{M}_{\text{triangle}}^{++00} = + \frac{m_t^2}{2m_Z^2} \log^2 \frac{m_{ZZ}^2}{m_t^2}$$

$$\mathcal{M}_{\text{box}}^{++00} = - \frac{m_t^2}{2m_Z^2} \log^2 \frac{m_{ZZ}^2}{m_t^2}$$

**Destructive
interference**

2. Off-shell Higgs in the SM

$$\frac{d\sigma_{gg \rightarrow ZZ}}{dm_{ZZ}} = \frac{d\sigma_{\text{Sig}}}{dm_{ZZ}} + \frac{d\sigma_{\text{Box}}}{dm_{ZZ}} + \frac{d\sigma_{\text{Int}}}{dm_{ZZ}}$$



MG5@NLOaMC

On-shell Z's

- Signal: $gg \rightarrow h^* \rightarrow ZZ$
- Box: $gg \rightarrow ZZ / h$
- Total: $gg \rightarrow ZZ$
- Interference: Subtraction

$m_{4l} > 180$ GeV

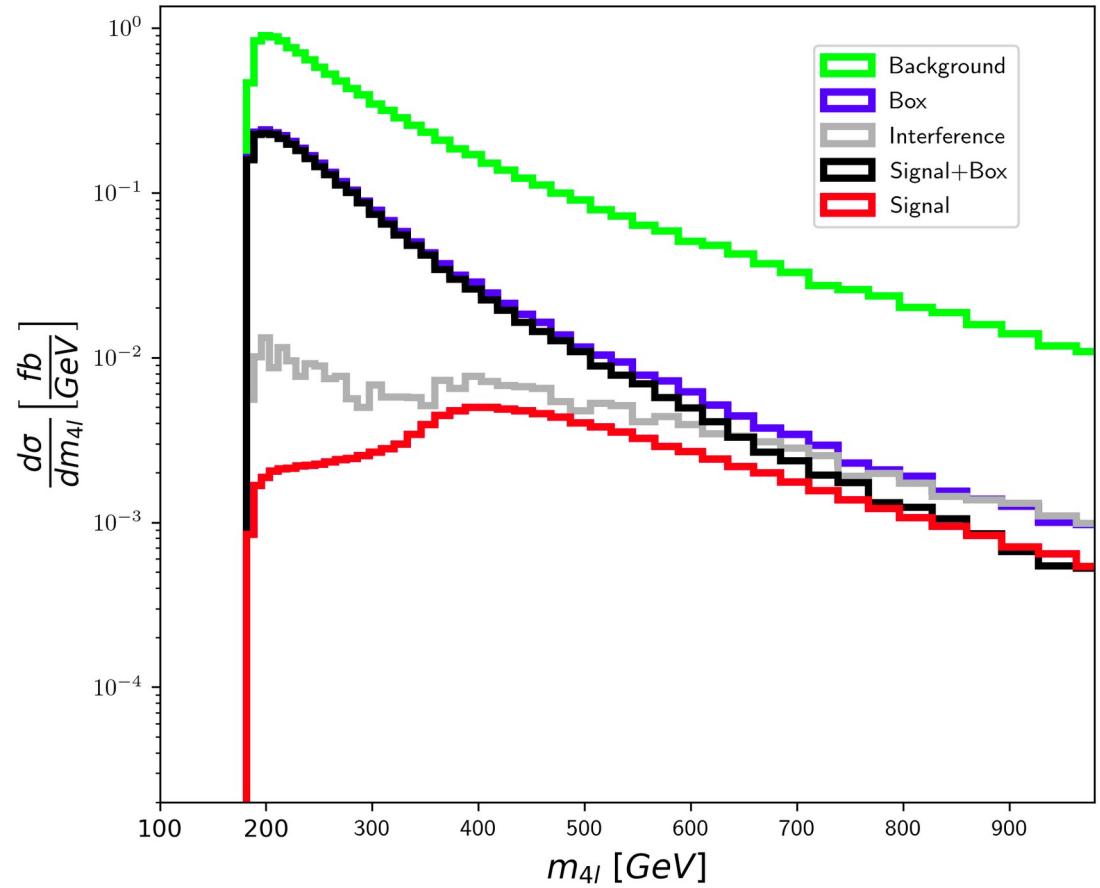
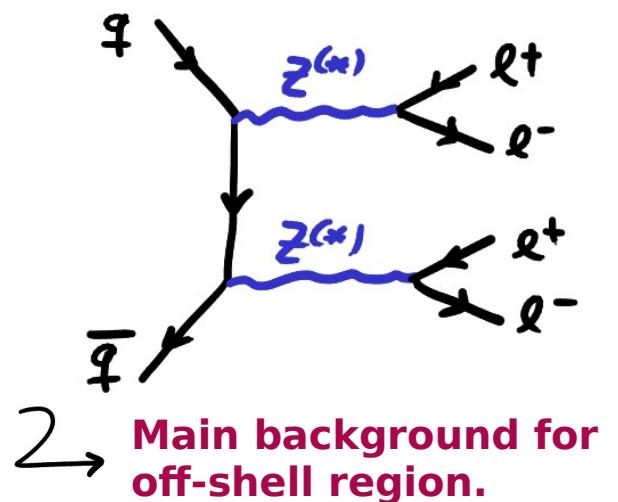
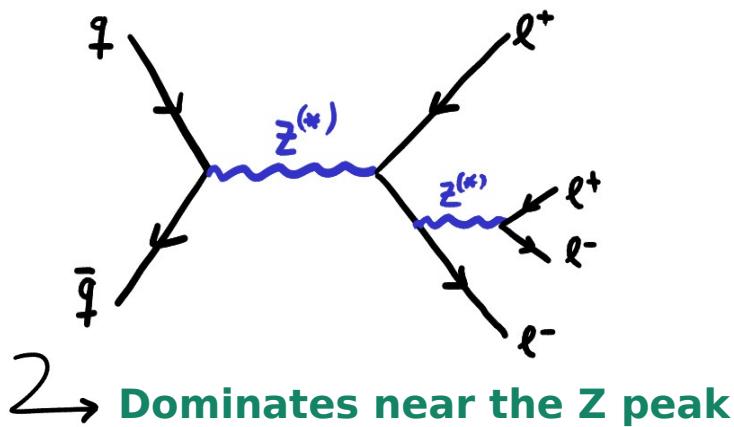
Universal K-factor: 1.76

$$\frac{d\sigma_{gg \rightarrow ZZ}}{dm_{ZZ}} = BR^2(Z \rightarrow l^+ l^-) \frac{d\sigma_{gg \rightarrow 4l}}{dm_{4l}}$$

MadAnalysis+MatPlotLib

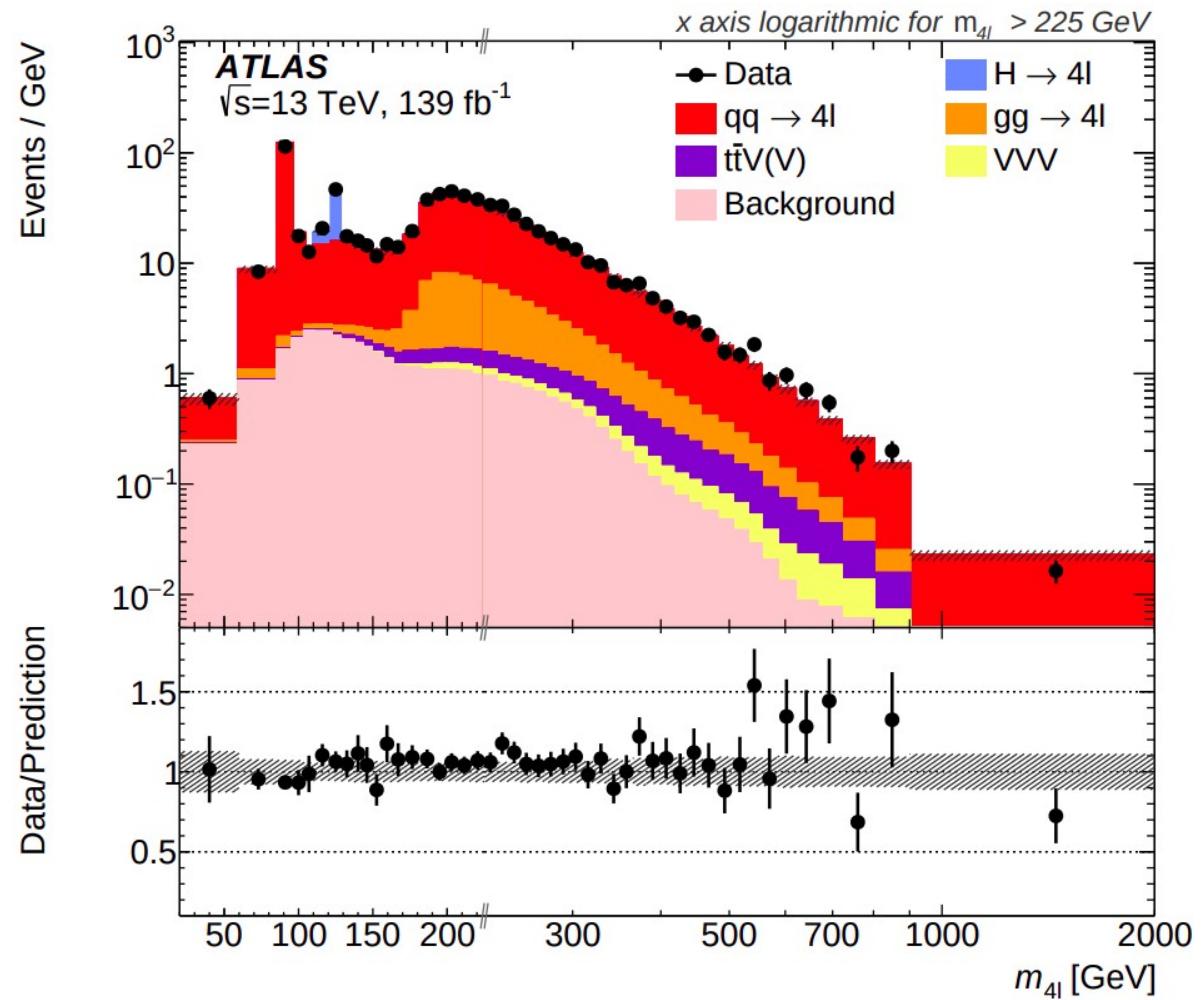
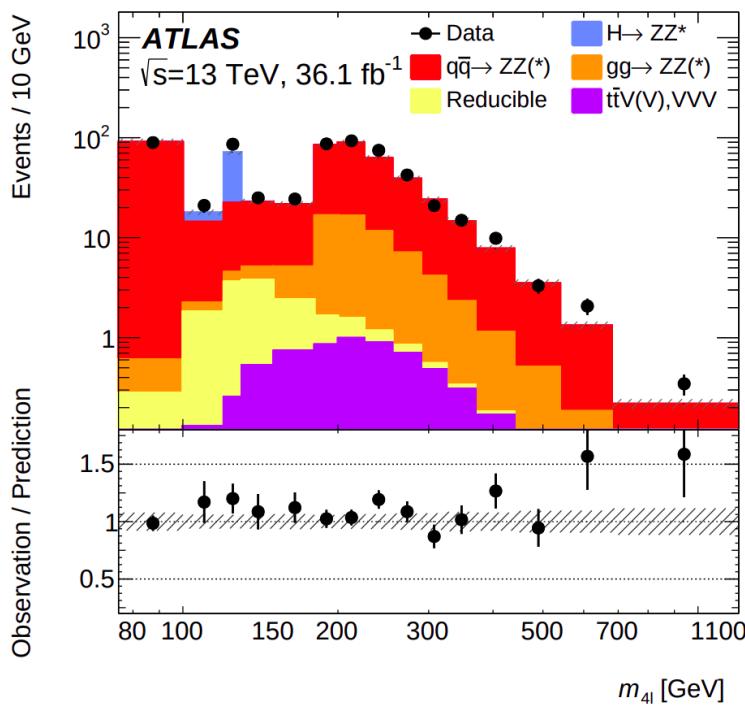
2. Off-shell Higgs in the SM

$q\bar{q} \rightarrow 4l$ Background: reducible



2. Off-shell Higgs in the SM

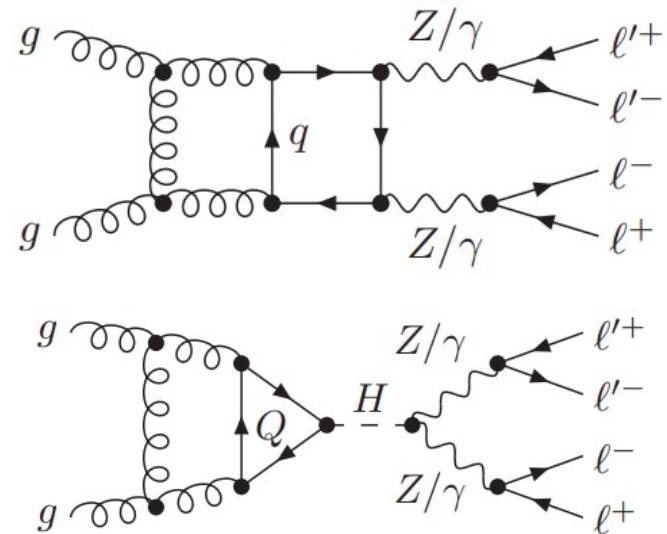
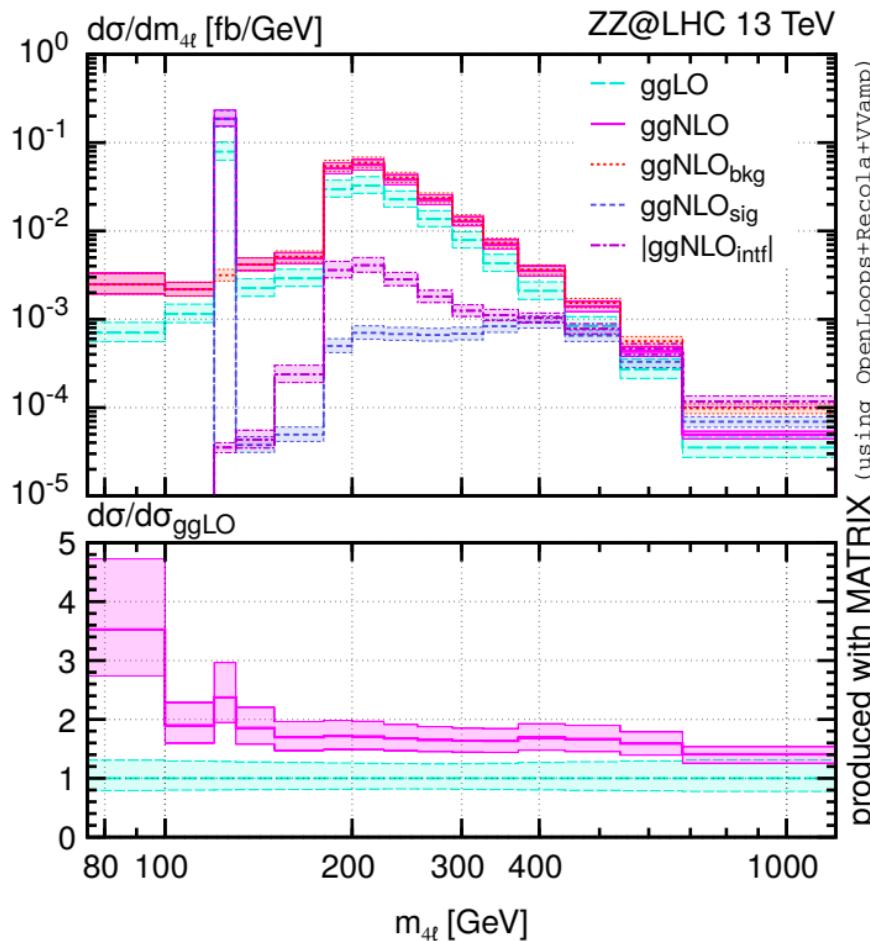
Experimental status



2. Off-shell Higgs in the SM

State-of-the-art simulations

- NLO (two-loop) $gg \rightarrow ZZ \rightarrow 4l$
- NNLO QCD && NLO EW $qq \rightarrow 4l$



Large K-factors: +67.3%
in $m_{4l} > 200$ GeV region

Part 3

Off-Shell Higgs Beyond the Standard Model

3. BSM off-shell Higgs

So far, all the LHC Higgs measurements are in agreement with the SM

2 → Focus on on-shell Higgs production.

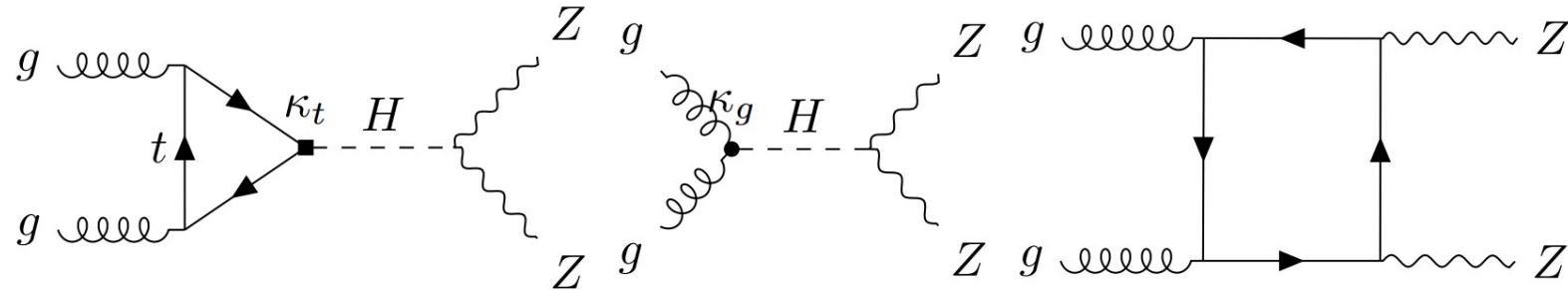
BSM effects should be of order v^2/Λ^2 for the on-shell Higgs.

The off-shell Higgs allow us to probe higher scales ($p \gg v$) with an enhanced p^2/Λ^2 sensitivity over the usual v^2/Λ^2 .

- Higgs couplings at higher scales.
- Additional decays and the Higgs width.
- EFT anomalous contributions.
- Momentum-dependent effects.

3. BSM off-shell Higgs

EFT anomalous couplings



$$\begin{aligned} \mathcal{L} &\supset c_g \frac{\alpha_s}{12\pi v^2} |H|^2 G_{\mu\nu} G^{\mu\nu} + c_t \frac{y_t}{v^2} |H| \bar{Q}_L \tilde{H} t_r + h.c. \\ &= \kappa_g \frac{\alpha_s}{12\pi v^2} h G_{\mu\nu} G^{\mu\nu} - \kappa_t \frac{y_t}{v^2} h (\bar{t}_R t_L + h.c.) \end{aligned}$$

$\kappa_t = 1 - Re(c_t)$
 $\kappa_g = c_g$

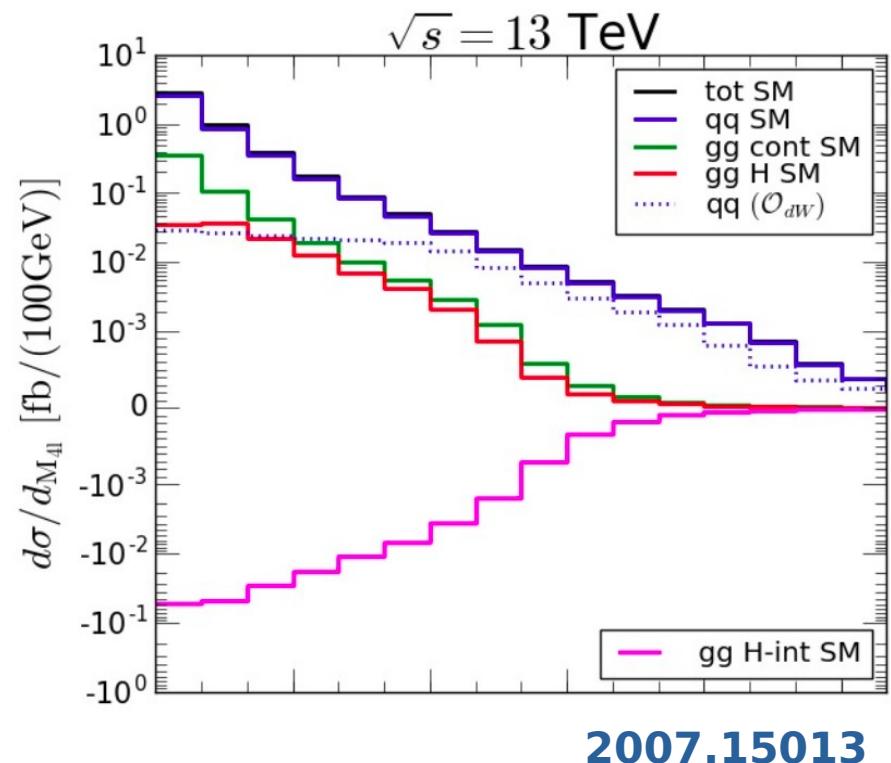
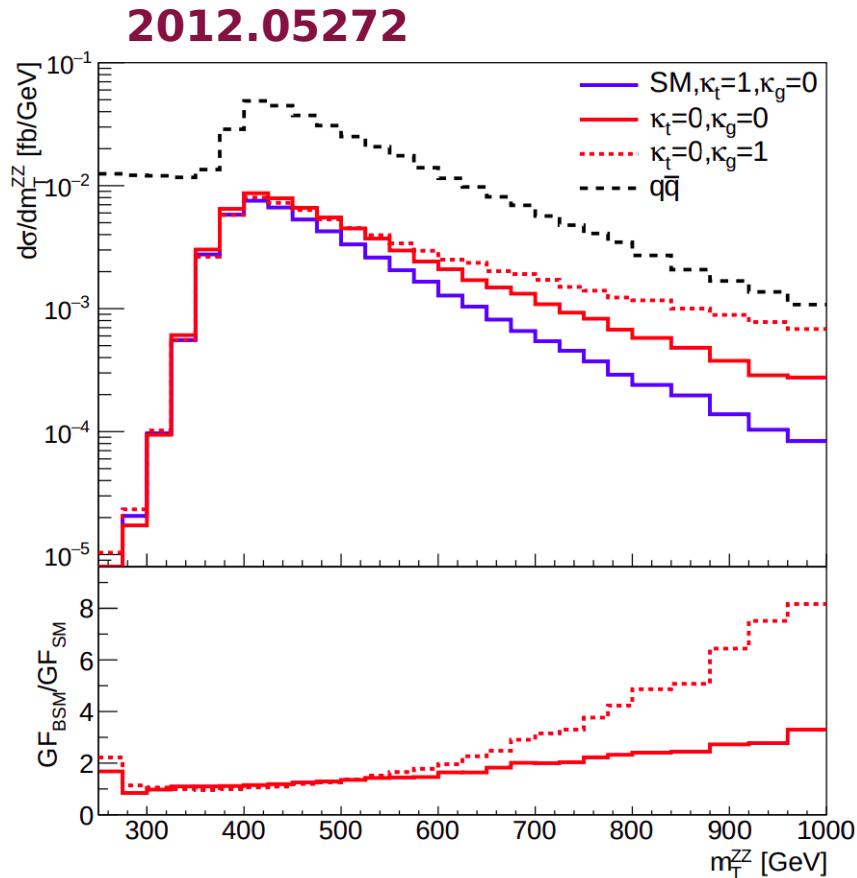
$$\left\{ \begin{array}{l} \mathcal{M}_{\text{triangle}}^{++00} = +\kappa_t \frac{m_t^2}{2m_Z^2} \log^2 \frac{m_{ZZ}^2}{m_t^2} \\ \mathcal{M}_{\text{box}}^{++00} = -\frac{m_t^2}{2m_Z^2} \log^2 \frac{m_{ZZ}^2}{m_t^2} \\ \mathcal{M}_g^{++00} = -\kappa_g \frac{m_{ZZ}^2}{2m_z^2} \end{array} \right.$$

← **Destructive interference**
← **Constructive interference**

3. BSM off-shell Higgs

Other possibilities:

- Anomalous hVV couplings
- Anomalous Zqq couplings



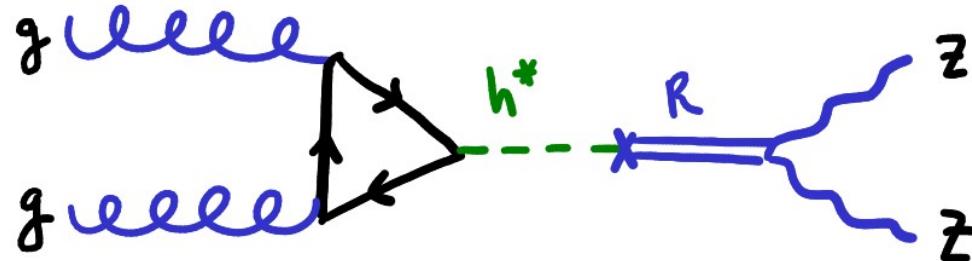
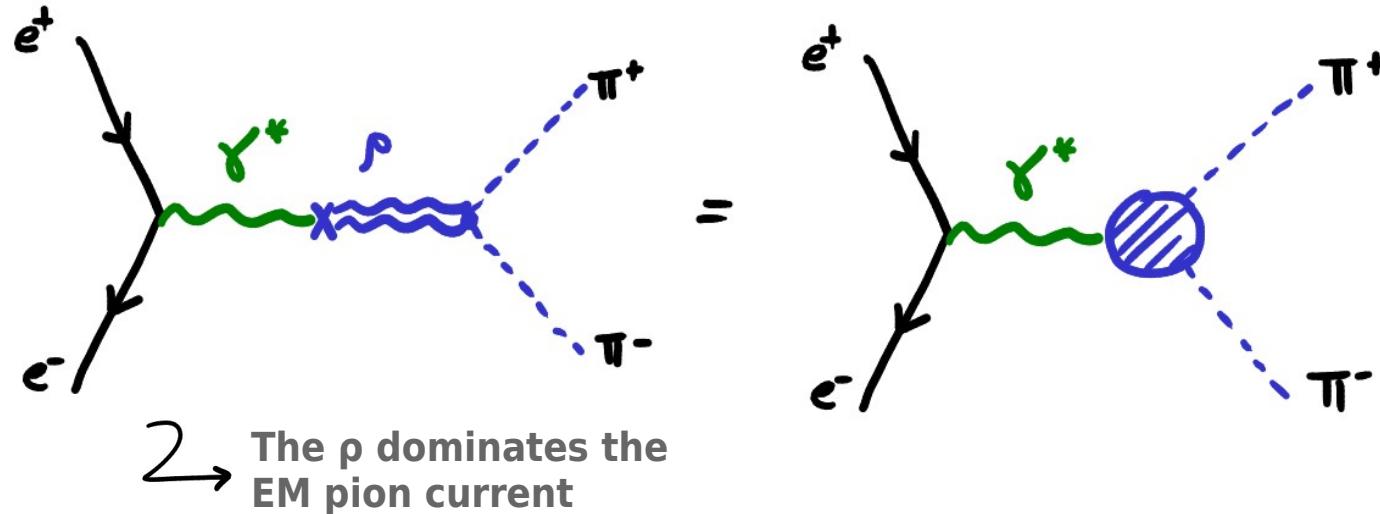
2007.15013

3. BSM off-shell Higgs

Momentum-dependent effects: Form Factors

Explore compositeness → Non-local momentum dependent Higgs couplings.

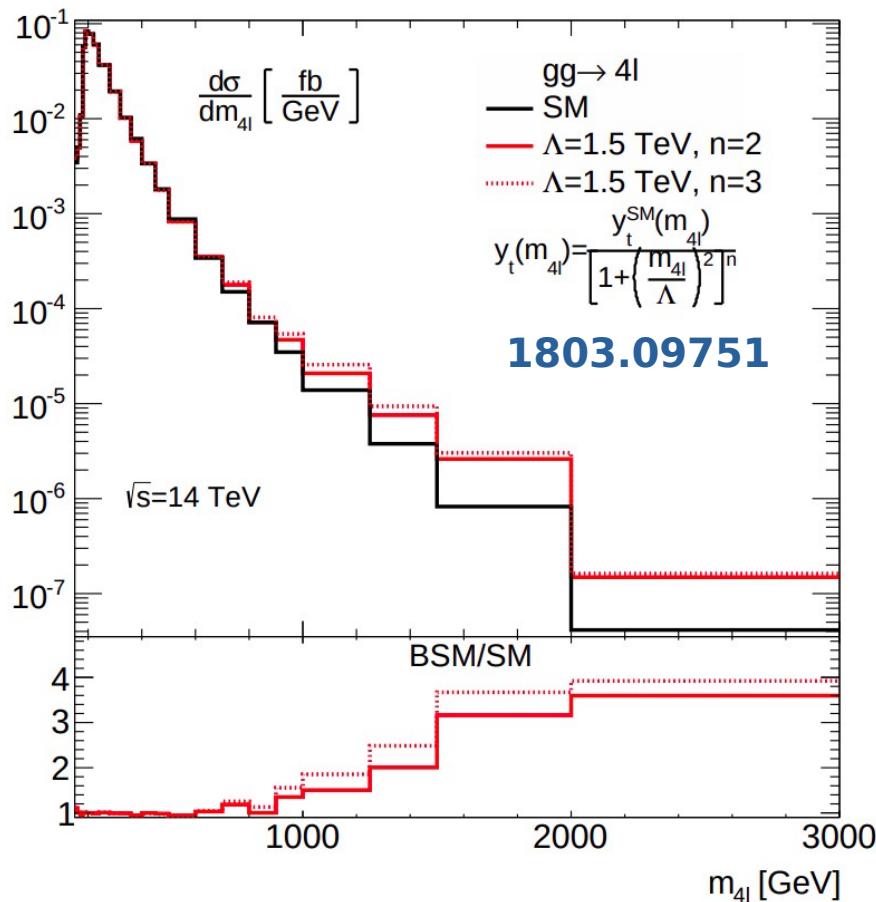
Example: Sakurai's Vector Meson Dominance.



$$F(q^2) = \frac{g_R \mu^2}{q^2 - M_R^2}$$

3. BSM off-shell Higgs

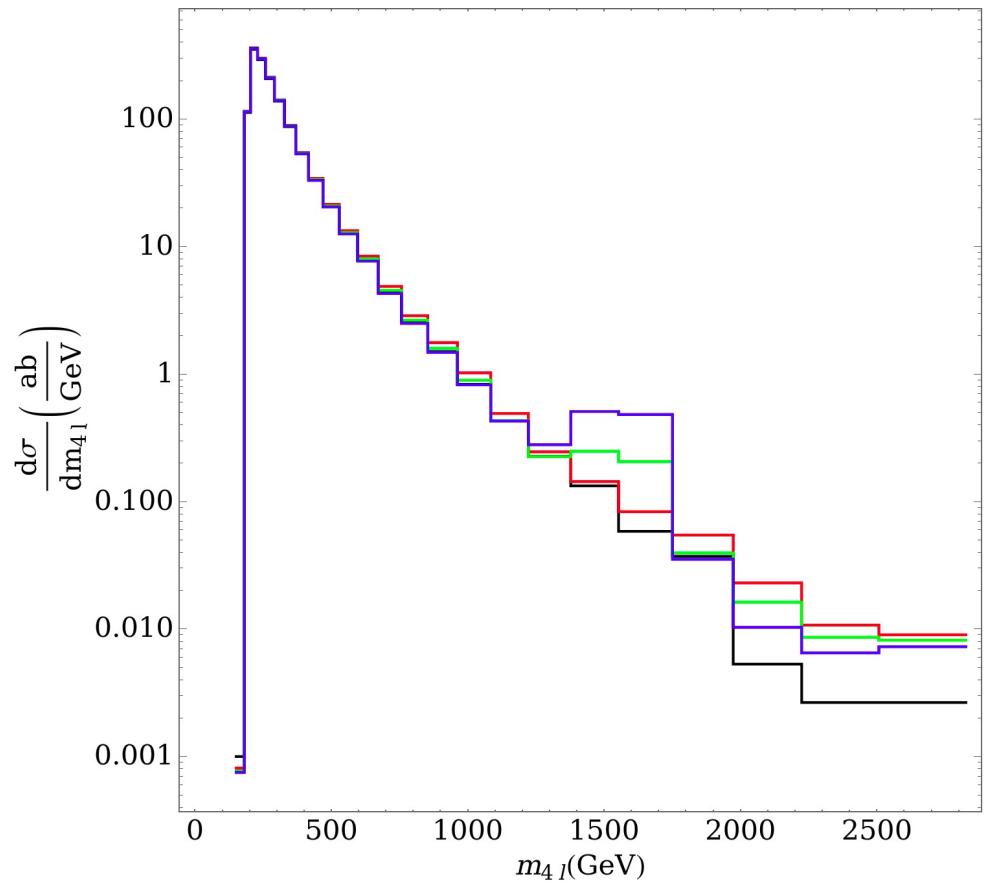
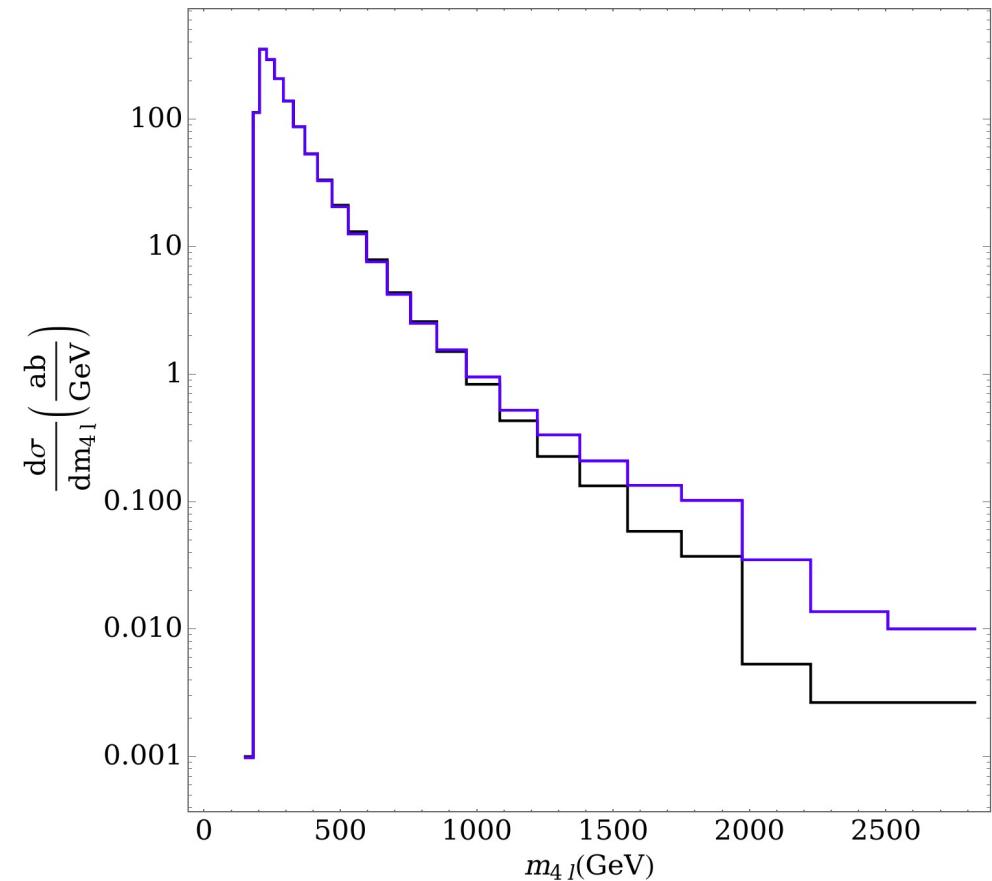
Momentum-dependent effects: Form Factors



$$F(q^2) = \frac{1}{(1 + q^2/\Lambda^2)^n}$$

3. BSM off-shell Higgs

Momentum-dependent effects: Form Factors



$$F(q^2) = \frac{1}{(1 + q^2/\Lambda^2)^n}$$

$$F(q^2) = \frac{\mu^2}{q^2 - M_R^2 - i\Gamma M_R}$$

3. BSM off-shell Higgs

Momentum-dependent effects in the EFT

$$F(q^2) = c_1 + c_{q^2} \frac{q^2}{\Lambda^2} + c_{q^4} \frac{q^4}{\Lambda^4} + \dots$$

Induce effective operators like $\frac{c_{q^2}}{\Lambda^2} v \partial^2 h V_\mu V^\mu$

$$\partial^2 h V_\mu V^\mu \rightarrow -m_h^2 h V_\mu V^\mu + \frac{3m_h^2}{2v} h^2 V_\mu V^\mu + y_d \bar{d} q V_\mu V^\mu + \dots$$



Not q^2 dependent

Off-shell sensitivity enhancement is lost in SMEFT since the reduction of the basis eliminates the derivative operators via EoM.

Missing off-shell effects in SMEFT

Conclusions

- Automatization of NLO is broadly available and very useful.
- Choosing between OPP v.s. TIR in MC simulations can save some time.
- Sizable contributions in off-shell Higgs region → Inadequacy of the NWA.
- BSM physics in high energy tails of the off-shell Higgs distributions.

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References - Part 3

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- Azatov, Aleksandr, et al. "**Taming the off-shell Higgs boson.**" Journal of Experimental and Theoretical Physics 120.3 (2015): 354-368.

Form Factor tutorial: <https://cp3.irmp.ucl.ac.be/projects/madgraph/wiki/FormFactors>

Backup

Rational terms

$\mathbf{R} = \mathbf{R}_1 + \mathbf{R}_2$ {

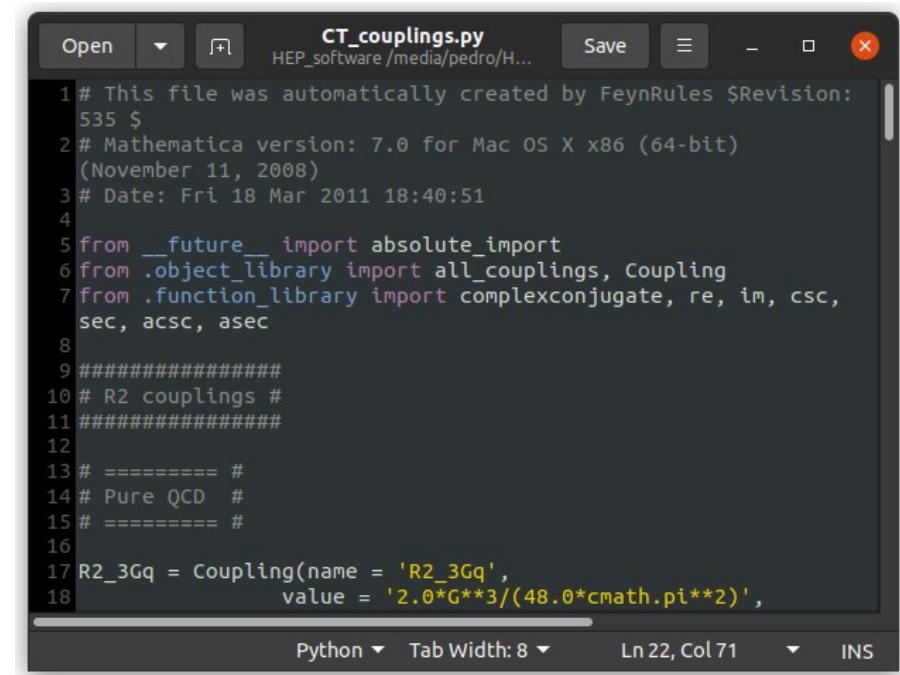
\mathbf{R}_1 : Originates from the ϵ part of the integral denominator.
→ **Computed using the reduction methods.**

\mathbf{R}_2 : Originates from the ϵ part of the integral denominator.

$$R_2 \equiv \lim_{\epsilon \rightarrow 0} \sum_c^{N_c} \mathcal{C}^{(c)} \sum_t^{N_t} \int d^d \bar{q} \frac{\tilde{\mathcal{N}}^{(c,t)}(\bar{q})}{\bar{D}_0^{(t)} \dots \bar{D}_{N_t}^{(t)}}$$

R2 Counterterms must be included in UFO model.

→ **FeynRules level.**



```
CT_couplings.py
# This file was automatically created by FeynRules $Revision:
# 535 $
# Mathematica version: 7.0 for Mac OS X x86 (64-bit)
# (November 11, 2008)
# Date: Fri 18 Mar 2011 18:40:51
#
# ===== #
# R2 couplings #
#####
#
# ===== #
# Pure QCD #
# ===== #
#
#2.0*G**3/(48.0*cmath.pi**2)',
```

Bounding the Higgs width

Constraints set by considering the relationship between the on-shell ($105 < m_H < 140 \text{ GeV}$) and off-shell ($m_H > 220 \text{ GeV}$) regions.

$$\sigma_{\text{vv} \rightarrow H \rightarrow 4\ell}^{\text{on-shell}} \propto \mu_{\text{vv}H} \quad \text{and} \quad \sigma_{\text{vv} \rightarrow H \rightarrow 4\ell}^{\text{off-shell}} \propto \mu_{\text{vv}H} \Gamma_H$$

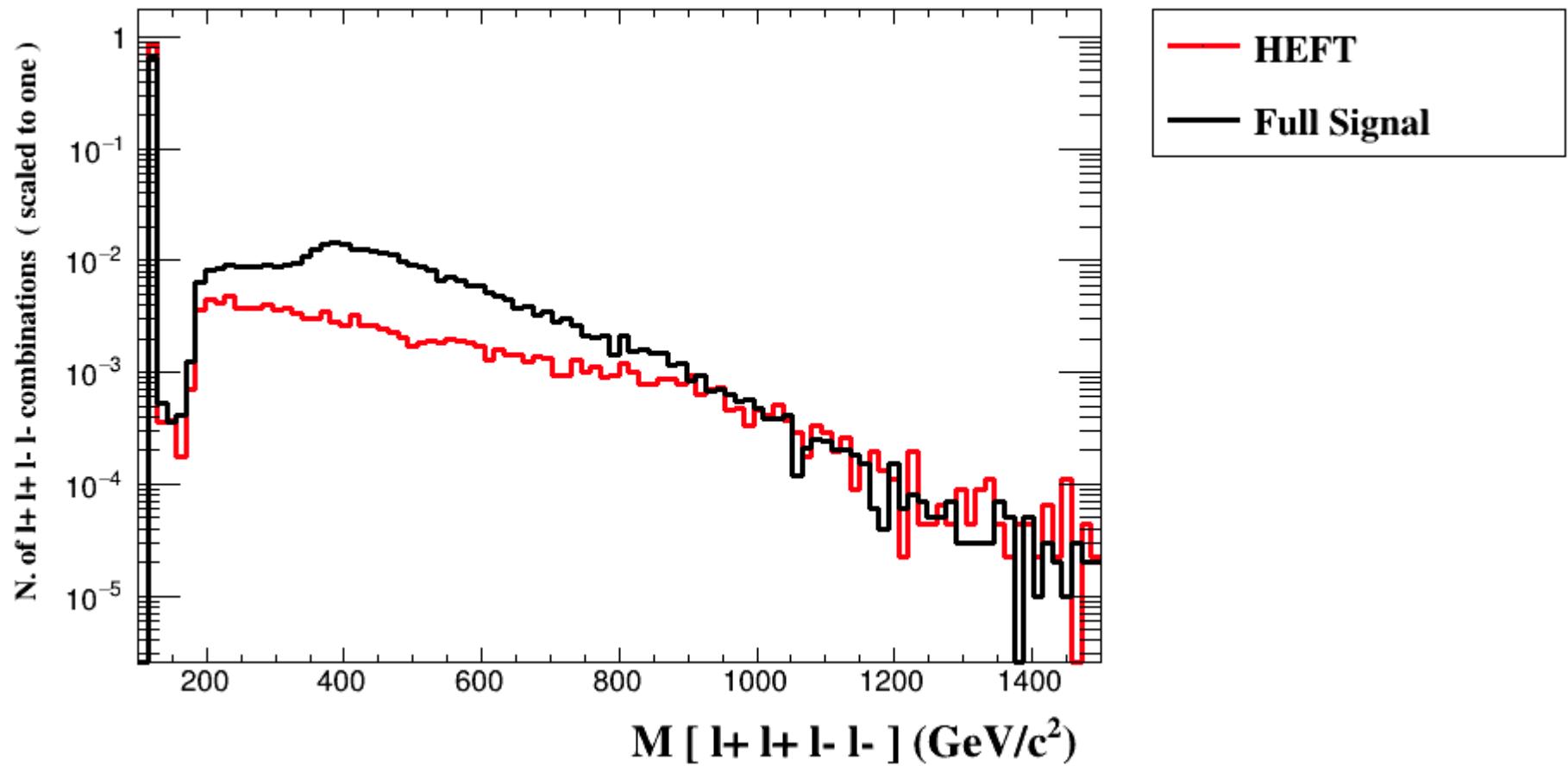
→ $\sigma_{i \rightarrow H \rightarrow f}^{\text{on-shell}} \propto \frac{g_i^2(m_H) g_f^2(m_H)}{\Gamma_H}$

On-shell degeneracy
 $g_{i,f}(m_H) \rightarrow \xi g_{i,f}(m_H)$
 $\Gamma_H \rightarrow \xi^4 \Gamma_H$

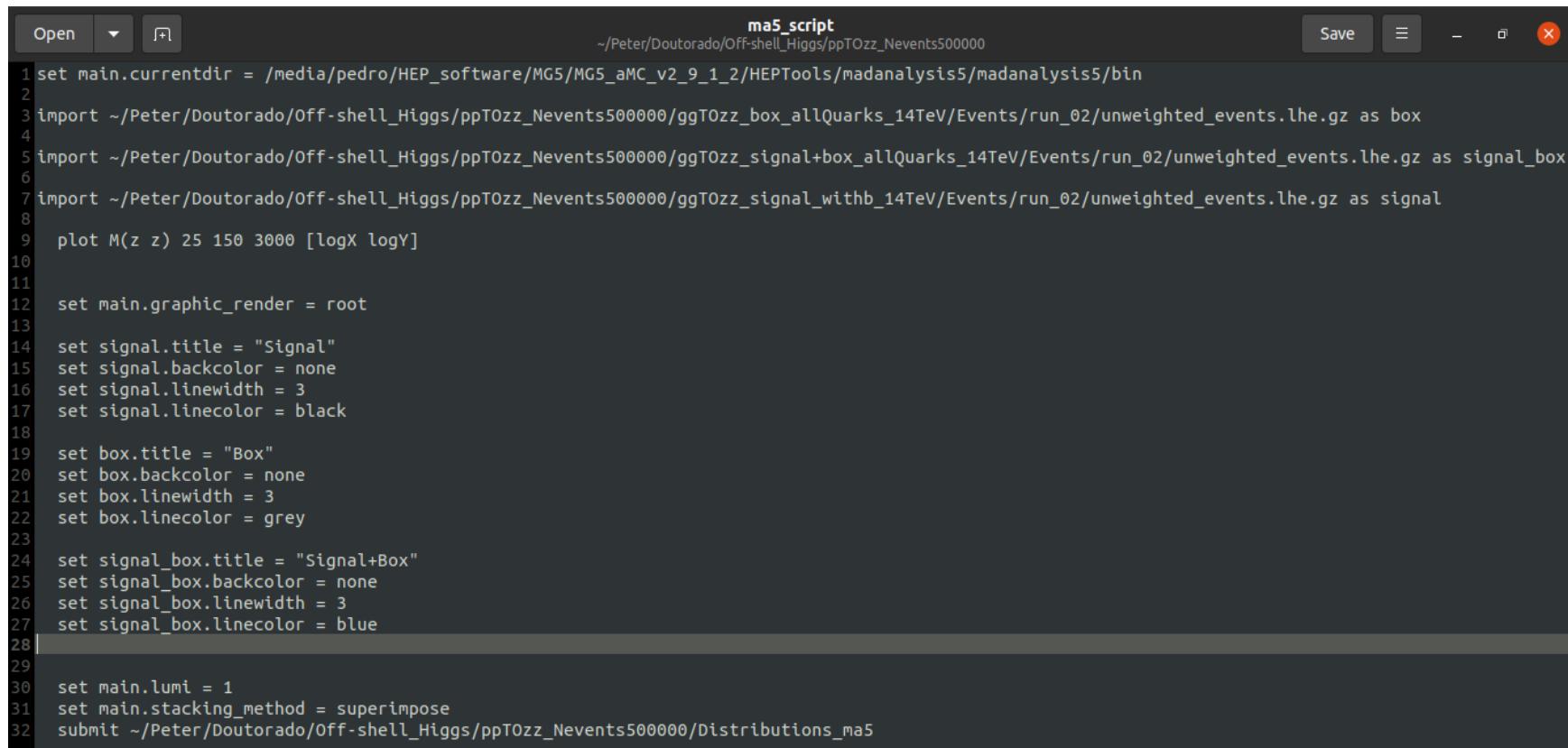
→ $\sigma_{i \rightarrow H^* \rightarrow f}^{\text{off-shell}} \propto g_i^2(\sqrt{\hat{s}}) g_f^2(\sqrt{\hat{s}})$

$$\boxed{\mu_{\text{off-shell}} / \mu_{\text{on-shell}} = \Gamma_H / \Gamma_H^{SM}}$$

Full Off-shell Higgs v.s. heft

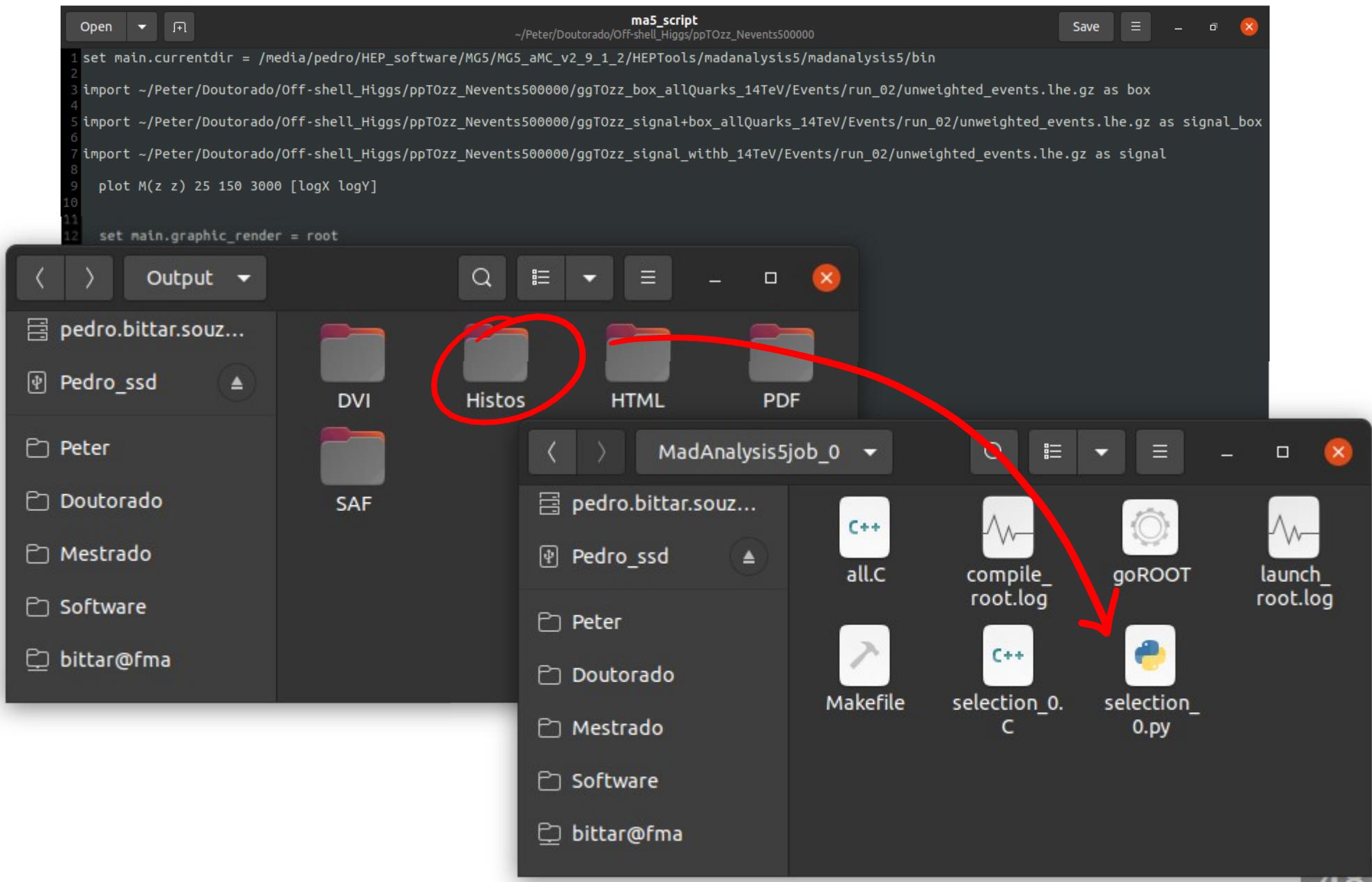


MadAnalysis + Matplotlib



```
ma5_script
~/Peter/Doutorado/Off-shell_Higgs/ppTozz_Nevents500000
1 set main.currentdir = /media/pedro/HEP_software/MG5/MG5_aMC_v2_9_1_2/HEPTools/madanalysis5/madanalysis5/bin
2
3 import ~/Peter/Doutorado/Off-shell_Higgs/ppTozz_Nevents500000/ggTozz_box_allQuarks_14TeV/Events/run_02/unweighted_events.lhe.gz as box
4
5 import ~/Peter/Doutorado/Off-shell_Higgs/ppTozz_Nevents500000/ggTozz_signal+box_allQuarks_14TeV/Events/run_02/unweighted_events.lhe.gz as signal_box
6
7 import ~/Peter/Doutorado/Off-shell_Higgs/ppTozz_Nevents500000/ggTozz_signal_withb_14TeV/Events/run_02/unweighted_events.lhe.gz as signal
8
9 plot M(z z) 25 150 3000 [logX logY]
10
11 set main.graphics_render = root
12
13 set signal.title = "Signal"
14 set signal.backcolor = none
15 set signal.linewidth = 3
16 set signal.linecolor = black
17
18 set box.title = "Box"
19 set box.backcolor = none
20 set box.linewidth = 3
21 set box.linecolor = grey
22
23 set signal_box.title = "Signal+Box"
24 set signal_box.backcolor = none
25 set signal_box.linewidth = 3
26 set signal_box.linecolor = blue
27
28
29
30 set main.lumi = 1
31 set main.stackning_method = superimpose
32 submit ~/Peter/Doutorado/Off-shell_Higgs/ppTozz_Nevents500000/Distributions_ma5
```

MadAnalysis + Matplotlib



MadAnalysis + MatPlotLib

```
selection_wlth_int.py ~ /Peter/Doutorado/Off-shell_Higgs/ppTOzz_Nevents500000/Distributions_ma5_25bins/Output... - □ ×

File Edit View Selection Find Packages Help
selection_0.py selection_wlth_int.py

1 def selection_0():
2
3     # Library import
4     import numpy
5     import matplotlib
6     import matplotlib.pyplot as plt
7     import matplotlib.gridspec as gridspec
8
9     # Library version
10    matplotlib_version = matplotlib.__version__
11    numpy_version       = numpy.__version__
12
13    # Histo binning
14    xBinning =
15    [161.825742626, 174.58380651, 188.347694261, 203.196703305, 219.216382743, 236.499026222, 255.144203659, 275.
16    259334892, 296.960308559, 320.372149754, 345.629740338, 372.878596026, 402.275704743, 433.99043107, 468.20549
17    2005, 505.118009636, 544.940646822, 587.902832395, 634.252082963, 684.255428919, 738.200952872, 796.399449374
18    , 859.186215481, 926.922982497, 1000.0, 1000.0]
19
20
21    # Creating data sequence: middle of each bin
22    xData =
23    numpy.array([150.0, 168.083771195, 181.335207314, 195.631363919, 211.054604981, 227.693787905, 245.644775461
24    , 265.010988075, 285.903999663, 308.444180435, 332.761390385, 358.995727443, 387.298334621, 417.832270786, 450
25    .773450088, 486.311655445, 524.651631935, 566.014266387, 610.637859968, 658.779501101, 710.716546618, 766.748
26    219688, 827.197333723, 892.412152189, 962.768395045])
27
28    # Creating weights for histo: y1_M_0
29    y1_M_0_weights =
30    numpy.array([0.0, 0.0, 0.80.5657393517, 233.543998121, 224.321998195, 192.58893845, 157.448778733, 123.48049900
31    , 694.8706952366, 72.0145194205, 53.4200595701, 39.513199682, 29.695439761, 22.5777798183, 17.0634678627, 12.4
32    62177897, 9.72589792173, 7.4182939403, 5.57276795516, 4.03621996752, 3.12876597482, 2.15450598266, 1.6979959
33    8634, 1.30272478952, 0.971476392182])
34
35    # Creating weights for histo: y1_M_1
36    y1_M_1_weights =
37    numpy.array([0.0, 0.0, 0.78.2236955773, 223.820387345, 216.415787764, 186.219369471, 150.735491477, 118.3370733
38    , 09.92.1015547926, 68.5296361254, 50.9243371207, 37.5213778786, 27.4877384458, 20.5660588372, 15.1539351432, 1
39    1.5697333458, 8.2349115344, 6.25312164645, 4.457039748, 2.93156583425, 2.12239988, 1.50955691465, 1.130177936
40    1, 0.7481459577, 0.52794697015])
41
42    # Creating weights for histo: y1_M_2
43    y1_M_2_weights =
44    numpy.array([0.0, 0.0, 0.404.479407862, 1507.19682929, 1493.49062903, 1274.12842476, 1053.65482048, 868.2702168
45
46    ~ /Peter/Doutorado/Off-shell_Higgs/ppTOzz_Nevents500000/Distributions_ma5_25bins/Output/Hists/MadAnalysis5Job, LF, UTF-8, Python, GitHub, Git (0)
```

Modify weights, add K-factor ...

```

# Creating a new Stack
BR = 0.034
Kfactor=1.76 #arxiv 1311.3589
y1_M_0_weights = Kfactor*BR**2*pbtotfb*y1_M_0_weights
y1_M_1_weights = Kfactor*BR**2*pbtotfb*y1_M_1_weights
y1_M_2_weights = BR**2*pbtotfb*y1_M_2_weights
y1_M_3_weights = Kfactor*BR**2*pbtotfb*y1_M_3_weights
interf = -y1_M_1_weights+y1_M_3_weights+y1_M_0_weights

```