Dipole DM: Model Definitions

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1 Dipole Dark Matter Model

The Dipole DM model extends the SM by adding two Dirac Fermions (χ_1 and χ_0) and a real scalar (ϕ), all singlets under the SM gauge group. The Lagrangian of the model is given by:

$$\mathcal{L} = \mathcal{L}_{SM} + \mathcal{L}_{\phi} + \mathcal{L}_{\chi} + \mathcal{L}_{dipole} + \mathcal{L}_{H\chi}, \tag{1.1}$$

where $\mathcal{L}_{\mathrm{SM}}$ represents the SM Lagrangian and

$$\mathcal{L}_{\phi} = (\partial^{\mu}\phi)^{2} - \mu_{2}^{2}|\phi|^{2} - \lambda_{2}\phi^{4} - \lambda_{3}\phi^{2}|H|^{2}, \qquad (1.2)$$

$$\mathcal{L}_{\chi} = i \overline{\chi}_{i} \partial \chi_{i} - \tilde{M}_{ij} \overline{\chi}_{i} \chi_{j} - (y_{\chi})_{ij} \overline{\chi}_{i} \chi_{j} \phi, \qquad (1.3)$$

$$\mathcal{L}_{\text{dipole}} = \frac{(C_{\gamma\chi\chi})_{ij}}{\Lambda} \overline{\chi}_i \sigma^{\mu\nu} \chi_j F_{\mu\nu} , \qquad (1.4)$$

$$\mathcal{L}_{H\chi} = \frac{(C_{H\chi\chi})_{ij}}{\Lambda} \overline{\chi}_i \chi_j |H|^2 \,. \tag{1.5}$$

In the equations above, $\sigma^{\mu\nu} = \frac{i}{2} \left[\gamma^{\mu}, \gamma^{\nu} \right]$ and H represents the Higgs doublet.

The model Lagrangian given in Eq. (1.1) contains the scalar potential:

$$V_{H,\phi} = \mu_1^2 |H|^2 + \lambda_1 |H|^4 + \mu_2^2 |\phi|^2 + \lambda_2 |\phi|^4 + \lambda_3 |\phi|^2 |H|^2.$$
 (1.6)

Assuming that both H and ϕ develop vevs, $\langle \phi \rangle = v_D/\sqrt{2}$ and $\langle H \rangle = v/\sqrt{2}$, we obtain the mass eigenstates h and S:

$$h = (\sqrt{2}H^0 - v)\cos\alpha - (\sqrt{2}\phi - v_D)\sin\alpha,$$

$$S = (\sqrt{2}\phi - v_D)\cos\alpha + (\sqrt{2}H^0 - v)\sin\alpha.$$
(1.7)

Where the mixing angle (α) is given by

$$\tan(2\alpha) \equiv \frac{\lambda_3 v v_D}{\lambda_1 v^2 - \lambda_2 v_D^2} \,. \tag{1.8}$$

and with masses:

$$m_{S,h}^2 = \lambda_1 v^2 + \lambda_2 v_D^2 \mp (\lambda_1 v^2 - \lambda_2 v_D^2) \sqrt{1 + \tan^2(2\alpha)},$$
 (1.9)

Using the above equations, we can express the quartic couplings (λ_i) and the mass parameters (μ_i) in terms of the physical masses, mixing angle and vevs:

$$\lambda_1 = \frac{1}{2v^2} \left(\cos^2 \alpha \, m_h^2 + m_S^2 \sin^2 \alpha \right), \tag{1.10}$$

$$\lambda_2 = \frac{1}{2v_D^2} \left(\cos^2 \alpha \, m_S^2 + m_h^2 \sin^2 \alpha \right), \tag{1.11}$$

$$\lambda_3 = \frac{1}{v_D v} \left(m_S^2 - m_h^2 \right) \sin \alpha \cos \alpha , \qquad (1.12)$$

$$\mu_1^2 = -\left(\lambda_1 v^2 + \lambda_3 \frac{v_D^2}{2}\right),\tag{1.13}$$

$$\mu_2^2 = -\left(\lambda_2 v_D^2 + \lambda_3 \frac{v^2}{2}\right) \tag{1.14}$$

Therefore, the parameters of the scalar potential $(\mu_{1,2} \text{ and } \lambda_{1,2,3})$ can be replaced by m_h, m_S, v, v_D and $\sin \alpha$.

Finally, the dark fermion mass matrix M_{ij} is defined in the mass eigenstate basis:

$$\tilde{M}_{ij} = M_i \delta_{ij} - \frac{v_D}{\sqrt{2}} (y_\chi)_{ij} - \frac{v^2}{2\Lambda} (C_{H\chi\chi})_{ij}$$

where M_i are the physical masses.

In addition to the Lagragian in Eq. (1.1), the effective h - G - G and S - G - G couplings induced by a top quark loop were also included as effective operators:

$$\mathcal{L}_{GG\phi} = -\frac{1}{4}\cos\alpha F(m_h^2/m_t^2)G^{\mu\nu}G_{\mu\nu}h - \frac{1}{4}\sin\alpha F(m_S^2/m_t^2)G^{\mu\nu}G_{\mu\nu}S$$

where F(x) is the effective loop function.

Model Parameters

The model parameters are given in Table Table 1 as well as their naming convention in the UFO model and their default values. The $C_{\gamma\chi\chi}$ and $C_{H\chi\chi}$ matrices are assumed to be real and symmetric.

Feynman rules

The relevant Feynman rules for the BSM particles are given in Table 2 below. Here, f represents any of the SM fermions and q any SM quark.

2 Minimal Dipole Model

The minimal scenario assumes that χ_0 only couples through the effective operators and these have zero entries in the diagonal:

$$(C_{\gamma\chi\chi})_{ij} = (C_{H\chi\chi})_{ij} = 0 , \text{ if } i = j$$

$$(2.1)$$

$$(y_x)_{ij} = 0$$
, if $i = 0$ or $j = 0$ (2.2)

Table 1. Model parameters, their respective names in the UFO model and default values.

Parameter	UFO name	Default Value
Λ	LambdaUV	5 TeV
$(C_{\gamma\chi\chi})_{11}$	Caxx1	0
$(C_{\gamma\chi\chi})_{00}$	Caxx0	0
$(C_{\gamma\chi\chi})_{10}$	Caxx10	0.1
$(C_{H\chi\chi})_{11}$	Chxx1	0
$(C_{H\chi\chi})_{00}$	Chxx0	0
$(C_{H\chi\chi})_{10}$	Chxx10	0.1
$(y_{\chi})_{11}$	ychi1	1.0
$(y_\chi)_{00}$	ychi0	0
$(y_{\chi})_{10}$	ychi10	0
$\sin \alpha$	sina	0.2
v_D	vevD	1 TeV

Table 2. Feynman rules for the relevant interactions.

Interaction	Vertex term	
$\chi_i \chi_j h$	$\frac{i}{\sqrt{2}}(y_{\chi})_{ij}\sin\alpha - i\frac{v}{\Lambda}(C_{H\chi\chi})_{ij}\cos\alpha$	
$\chi_i \chi_j A^\mu$	$-irac{1}{\Lambda}(C_{\gamma\chi\chi})_{ij}(\gamma^{\mu} ot\!\!/- ot\!\!/\gamma^{\mu})$	
$S \chi_i \chi_j$	$-\frac{i}{\sqrt{2}}(y_{\chi})_{ij}\cos\alpha - i\frac{v}{\Lambda}(C_{H\chi\chi})_{ij}\sin\alpha$	
Shh	$-i\frac{m_S^2}{2v}\left(1+2\frac{m_h^2}{m_S^2}\right)\left(\cos\alpha+2\frac{v}{v_D}\sin\alpha\right)\sin(2\alpha)$	
$S f \bar{f}$	$-i\frac{m_f}{v}\sin\alpha$	
$SW_{\mu}^{-}W_{\nu}^{+}$	$2ig^{\mu\nu}rac{m_W^2}{v}\sinlpha$	
$S Z_{\mu} Z_{\nu}$	$2ig^{\mu\nu} rac{m_Z^2}{v} \sin \alpha$	

With the above assumptions the Lagrangian simplifies to:

$$\mathcal{L}_{\chi} = \overline{\chi}_i \left(i \partial - \tilde{M}_i \right) \chi_i - (y_{\chi})_{11} \overline{\chi}_1 \chi_1 \phi, \qquad (2.3)$$

$$\mathcal{L}_{\text{dipole}} = \frac{(C_{\gamma\chi\chi})_{01}}{\Lambda} \overline{\chi}_0 \sigma^{\mu\nu} \chi_1 F_{\mu\nu} + h.c., \qquad (2.4)$$

$$\mathcal{L}_{H\chi} = \frac{(C_{H\chi\chi})_{01}}{\Lambda} \overline{\chi}_0 \chi_1 |H|^2 + h.c..$$
 (2.5)

resulting in the following vertices:

$$\chi_1 \chi_1 h : \frac{i}{\sqrt{2}} (y_\chi)_{11} \sin \alpha, \quad S \chi_1 \chi_1 : -\frac{i}{\sqrt{2}} (y_\chi)_{11} \cos \alpha$$
(2.6)

$$\chi_1 \chi_0 h : -i \frac{v}{\Lambda} (C_{H\chi\chi})_{01} \cos \alpha, \quad S \chi_1 \chi_0 : -i \frac{v}{\Lambda} (C_{H\chi\chi})_{01} \sin \alpha$$
 (2.7)

$$\chi_1 \, \chi_0 \, A^{\mu} : -i \frac{1}{\Lambda} (C_{\gamma \chi \chi})_{01} (\gamma^{\mu} \not p - \not p \gamma^{\mu}) \tag{2.8}$$