

# Model Building for Back-To-The-Future Trigger

- We want event-N triggered + event N-1 in the previous bunch crossing (BC) ( $t=25$  ns,  $ct=7.5$  m) not.
  - We'll see particle  $\chi_1$  in event N-1 and particles  $\chi_0 + Y$  in event N,  $\chi_1 \rightarrow \chi_0 + Y$ .
  - To avoid standard triggers on N-1 event, strong preference for  $\chi_1$  SM neutral (or millicharged).
  - Since MET is involved in the N trigger,  $\chi_0$  should also be neutral.
  - For  $\chi_1$  lifetime largish (to decay it in the next BC), we need small  $\Delta=\chi_1-\chi_0$  and/or  $\chi_1-\chi_0-Y$  coupling
  - Hence the simplest choice, SUSY-reminiscent, is to take  $Y=\gamma$  (SM photon).  $Y=Z, h$  also possible.
  - With only  $\chi_1-\chi_0-\gamma$  vertex very small cross-section ( $\gamma\gamma F$  @LHC with highly off-shell photon).
- We add a new scalar  $s$  produced resonantly in gluon fusion, and decaying into  $\chi_1 \chi_1$ . [2 couplings]

The relevant Lagrangian reads 
$$\mathcal{L} \supset \frac{C_{ggs}}{\Lambda} s G^{\mu\nu} G_{\mu\nu} + \frac{C_{\chi\chi\gamma}}{\Lambda} \bar{\chi}_1 \sigma_{\mu\nu} F^{\mu\nu} \chi_0 + \text{h.c.} + y_\chi^1 \bar{\chi}_1 s \chi_1 + \mathcal{L}_{mass}$$

6 free parameters:  $\frac{C_{ggs}}{\Lambda}, \frac{C_{\chi\chi\gamma}}{\Lambda}, y_\chi^1, m_{\chi_1}, m_{\chi_0}, m_s$   Can we trade them for a more "physical" set? Constraints?

# Parametrization and constraints (I)

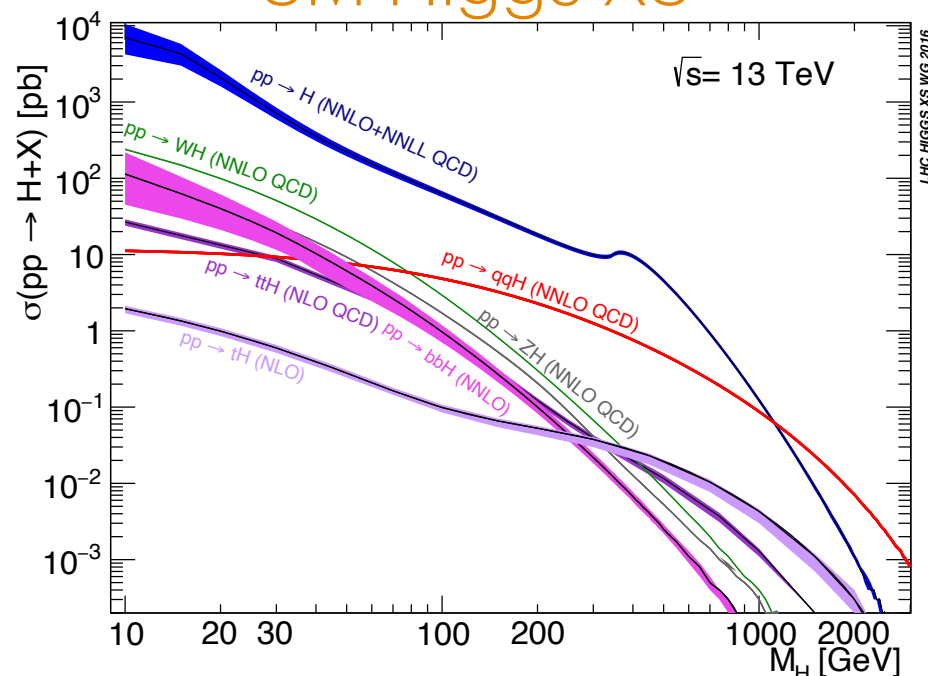
$$\mathcal{L} \supset \frac{C_{ggS}}{\Lambda} s G^{\mu\nu} G_{\mu\nu} + \frac{C_{\chi\chi\gamma}}{\Lambda} \bar{\chi}_1 \sigma_{\mu\nu} F^{\mu\nu} \chi_0 + \text{h.c.} + y_\chi^1 \bar{\chi}_1 s \chi_1 + \mathcal{L}_{mass}$$

Can be traded for  $\sigma_s(m_S) = \sigma(gg \rightarrow s)$  which is only function of  $m_S$ . We can relate this  $\sigma_s$  to  $\sigma_{s,SM}$  by considering either a new heavy quark  $Q$  (or vector-like quark), or via a direct mixing with the SM Higgs ( $h$ ). We will have

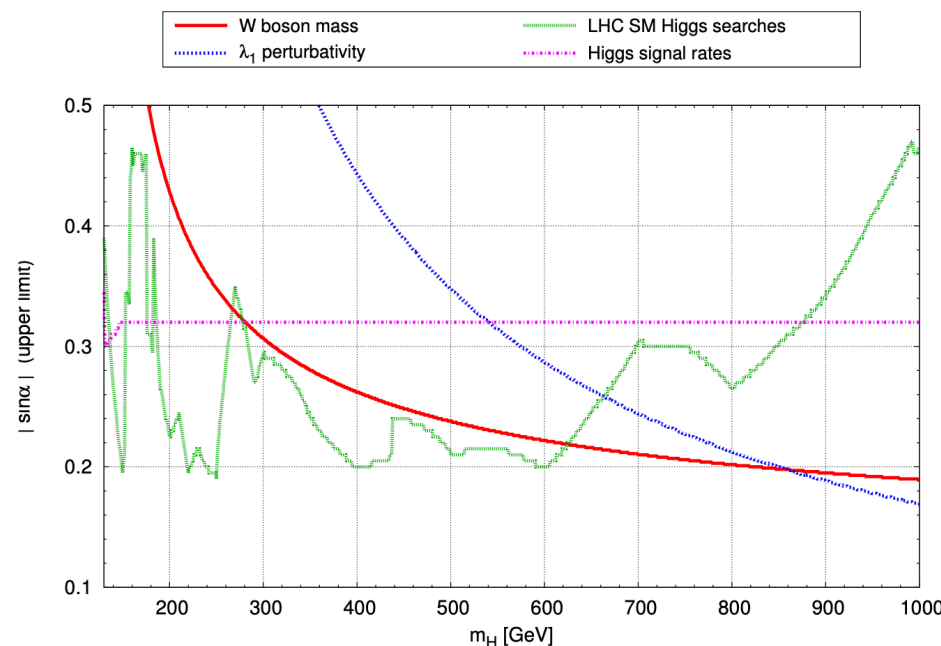
$$\sigma(pp \rightarrow s) = \left[ \frac{C_{ggS}^{(Q)}}{\Lambda} \right]^2 f(m_S) = \left[ \frac{C_{ggS}^{(Q)}}{C_{ggS}^{(t)}} \frac{C_{ggh}^{(t)}}{\Lambda} \right]^2 f(m_S) = \left( \frac{C_{ggS}^{(Q)}}{C_{ggh}^{(t)}} \right)^2 \times \sigma_{s,SM}(m_S)$$

For a heavy quark  $Q$ , and the parenthesis would be traded by  $\sin^2(\alpha)$  if we consider that  $s$  mixes with  $h$ .

## SM Higgs XS



## Bounds on mixing angle



$m_S = 600 \text{ GeV}$

$\sigma_{s,SM}(600) \sim 3 \text{ pb}$ ,  $\sin^2 \alpha \sim 0.04$ ,  
 $\sigma_s \leq 120 \text{ fb}$

$m_S = 1000 \text{ GeV}$

$\sigma_{s,SM}(1000) \sim 0.2 \text{ pb}$ ,  $\sin^2 \alpha \sim 0.02$ ,  
 $\sigma_s \leq 5 \text{ fb}$

$m_S = 1500 \text{ GeV}$

$\sigma_{s,SM}(1500) \sim 0.02 \text{ pb}$ ,  $\sin^2 \alpha \sim 0.02$ ,  
 $\sigma_s \leq 0.5 \text{ fb}$

Robens, 2205.06295

[https://twiki.cern.ch/twiki/pub/LHCPhysics/HiggsXSBR/plotAll\\_13tev\\_BSM\\_sqrt.pdf](https://twiki.cern.ch/twiki/pub/LHCPhysics/HiggsXSBR/plotAll_13tev_BSM_sqrt.pdf)

# Parametrization and constraints (II)

$$\mathcal{L} \supset \frac{C_{ggs}}{\Lambda} s G^{\mu\nu} G_{\mu\nu} + \frac{C_{\chi\chi\gamma}}{\Lambda} \bar{\chi}_1 \sigma_{\mu\nu} F^{\mu\nu} \chi_0 + \text{h.c.} + y_\chi^1 \bar{\chi}_1 s \chi_1 + \mathcal{L}_{mass}$$

Other searches that might apply: 1) di-jet Z' searches 2) monojet

1)  $\sigma_{jj}(m_s) = \sigma_s(m_s) BR(s \rightarrow gg)$ , but the latter is given by  $1/[1+(y_\chi \Lambda / C_{ggs})^2]$ , small for large enough  $y_\chi$ .

2)  $\sigma(pp \rightarrow sj, s \rightarrow \chi_1 \chi_1)$  constrained by monojet, if we assume  $\chi_1$  decays are undetected. This should be only a function of  $m_s$ , but a non-trivial one, since the shape of  $pT(j)$  would depend on  $m_s$ .

Exclusive Signal Region	EM2	EM4	EM6	EM8	EM9	EM11
Data events (139 fb <sup>-1</sup> )	313 912	102 888	10 203	1663	738	187
SM prediction	314 000 ± 3500	101 600 ± 1200	10 000 ± 180	1640 ± 40	754 ± 20	182 ± 6
$W \rightarrow e\nu$	16 000 ± 1000	3980 ± 250	280 ± 19	35.8 ± 2.6	13.4 ± 1.0	3.01 ± 0.24
$W \rightarrow \mu\nu$	23 600 ± 500	5940 ± 120	481 ± 12	66.8 ± 2.3	31.2 ± 1.2	7.8 ± 0.4
$W \rightarrow \tau\nu$	54 900 ± 800	15 430 ± 260	1243 ± 29	167 ± 6	77.4 ± 2.9	15.5 ± 0.8
VBF $W$ + jets	2340 ± 300	1010 ± 150	140 ± 27	29 ± 7	16 ± 5	5.0 ± 1.9
$Z \rightarrow ee$	—	—	—	—	—	—
$Z \rightarrow \mu\mu$	597 ± 15	97.4 ± 2.7	4.51 ± 0.15	1.49 ± 0.05	0.60 ± 0.02	—
$Z \rightarrow \tau\tau$	530 ± 14	115.0 ± 3.3	8.31 ± 0.28	0.90 ± 0.04	0.40 ± 0.02	2.10 ± 0.08
$Z \rightarrow \nu\nu$	192 800 ± 2100	67 400 ± 1000	7000 ± 170	1180 ± 40	534 ± 20	126 ± 6
VBF $Z$ + jets	3900 ± 500	2170 ± 310	370 ± 60	86 ± 17	45 ± 10	13.7 ± 3.3
single- $t$	2800 ± 700	550 ± 180	15 ± 8	—	—	—
$t\bar{t}$	8900 ± 700	2000 ± 150	100 ± 8	8.2 ± 1.0	2.4 ± 0.4	0.30 ± 0.05
Diboson	6200 ± 1000	2700 ± 500	350 ± 70	71 ± 15	33 ± 8	—
Multijet	1100 ± 1100	57 ± 57	0.6 ± 0.6	0.1 ± 0.1	—	—
Noncollision background	240 ± 240	46 ± 46	8 ± 8	6 ± 6	—	—
SUSY, $m(\tilde{t}, \tilde{\chi}^0) = (600, 580)$ GeV	562 ± 70	516 ± 60	159 ± 19	44 ± 6	28 ± 4	8.2 ± 1.1
DMA, $m(\chi, Z_A) = (1, 2000)$ GeV	770 ± 30	684 ± 27	212 ± 9	79 ± 4	47.9 ± 2.3	18.7 ± 1.1
Dark energy, $M_2 = 1486$ GeV	286 ± 7	320 ± 11	125 ± 7	52 ± 5	33.6 ± 3.2	14.6 ± 1.8

From the paper:

“Values of  $\sigma \times A \times \epsilon$  above 736 fb (for IM0) and above 0.3 fb (for IM12) are excluded at 95% CL.”

We’d need to compute  $A, \epsilon$  for our model. But they should be  $O(10^{-2,-3})$  and hence comparable with those from single  $s$  production from last slide.

ATLAS, 2102.10874, 139 fb<sup>-1</sup>

## Parametrization and constraints (III)

$$\mathcal{L} \supset \frac{C_{ggS}}{\Lambda} S G^{\mu\nu} G_{\mu\nu} + \frac{C_{\chi\chi\gamma}}{\Lambda} \bar{\chi}_1 \sigma_{\mu\nu} F^{\mu\nu} \chi_0 + \text{h.c.} + y_\chi^1 \bar{\chi}_1 S \chi_1 + \mathcal{L}_{mass}$$

I can't fathom a relevant constraint here, except the coupling not being too large (unitarity)

This coefficient enters only in the  $\chi_1$  decay, then  $\tau(\chi_1) = \tau = f(C_{\chi\chi\gamma}/\Lambda, m_{\chi_1}, m_{\chi_0})$ . So use  $\tau$  instead of  $\chi\chi\gamma$  coupling!

Note: in the current Pythia simulation we force the  $\chi_1 \rightarrow \chi_0 \gamma$  decay using oneChannel. This might not give the required tensor structure (need to read and think about it!).

-We can always make a MG5 model and then shower the parton level events with Pythia.

-I recall Gareth saying Pythia lifetime does not agree with the estimation of 2301.05252, this ought to be checked.

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! 5) Process selection
Higgs:useBSM = on
HiggsBSM:allH2 = on
35:m0 = 600.
1000023:m0 = 300.0
1000022:m0 = 247.5
1000023:tau0 = 7500
1000022:mayDecay = off ! this line is probably not necessary. Better safe than sorry

! Force H0 > N2 N2, N2 > N1 gamma decays
35:oneChannel = 1 1.000 100 1000023 1000023
1000023:oneChannel 1 1.000 100 1000022 22
```