• The Dipole DM model extends the SM by adding three Dirac Fermions and a real scalar, all <u>singlets</u> under the SM gauge group:

$$\chi_2, \chi_1, \chi_0 \text{ and } \phi$$

• The Lagrangian of the model is given by:

$$\mathcal{L} = \mathcal{L}_{\mathrm{SM}} + \mathcal{L}_{\phi} + \mathcal{L}_{\chi} + \mathcal{L}_{H\chi} + \mathcal{L}_{\mathrm{dipole}},$$

$$\mathcal{L}_{\phi} = (\partial^{\mu}\phi)^{2} - \mu_{2}^{2}|\phi|^{2} - \lambda_{2}\phi^{4} - \lambda_{3}\phi^{2}|H|^{2},$$

$$\mathcal{L}_{\chi} = i\overline{\chi}_{i}\partial\chi_{i} - \tilde{M}_{ij}\overline{\chi}_{i}\chi_{j} - (y_{\chi})_{ij}\overline{\chi}_{i}\chi_{j}\phi,$$

$$\mathcal{L}_{H\chi} = \frac{(C_{H\chi\chi})_{ij}}{\Lambda}\overline{\chi}_{i}\chi_{j}|H|^{2}.$$

$$\mathcal{L}_{\text{dipole}} = \frac{(C_{\gamma\chi\chi})_{ij}}{\Lambda} \overline{\chi}_i \sigma^{\mu\nu} \chi_j F_{\mu\nu}$$

• Mass eigenstates: $\chi_2,\,\chi_1,\,\chi_0,\,S$

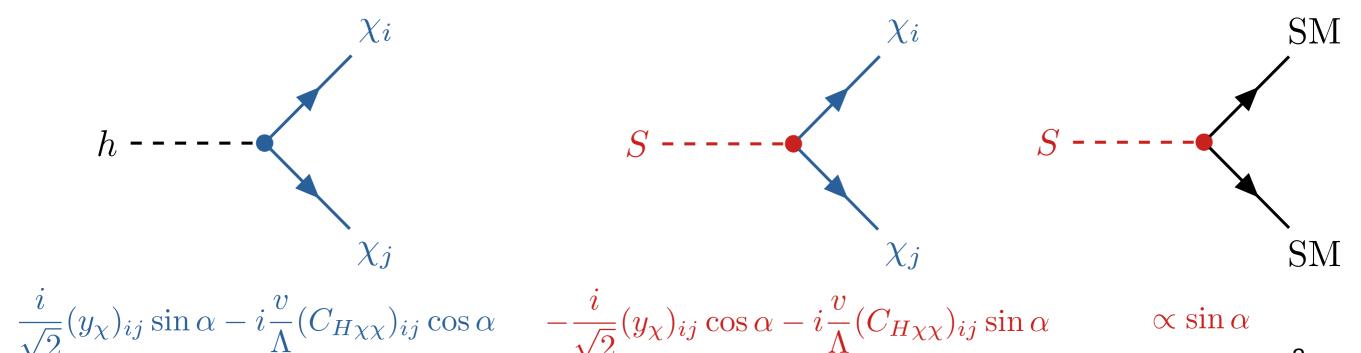
$$h = (\sqrt{2}H^0 - v)\cos\alpha - (\sqrt{2}\phi - v_D)\sin\alpha,$$

$$S = (\sqrt{2}\phi - v_D)\cos\alpha + (\sqrt{2}H^0 - v)\sin\alpha.$$

• We assume the following mass hierarchy:

$$m_S > m_{\chi_2} > m_{\chi_1} > m_{\chi_0} > m_h/2$$

• Vertices:

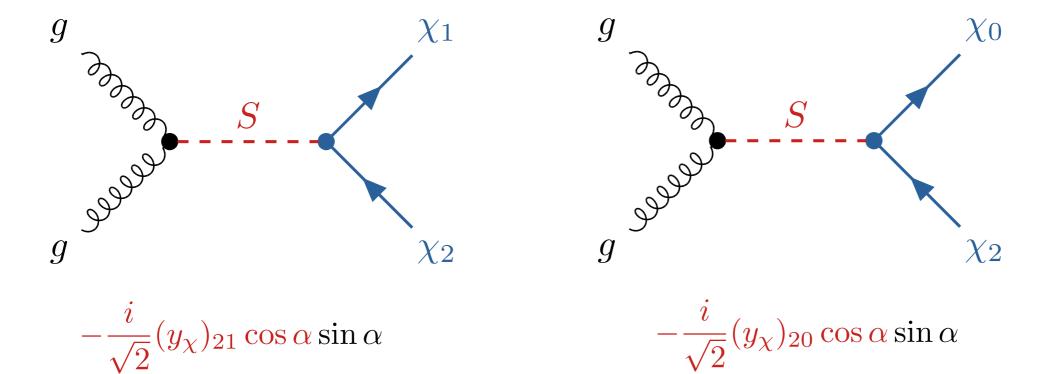


• Minimal "12"-Model:

$$(C_{H\chi\chi})_{ij} = 0$$
, if $i = j$ or $i = 2$ or $j = 2$
 $(y_{\chi})_{ij} = 0$, if $i = j$ or $i = 1, j = 0$

$$\mathcal{L}_{\chi} = \overline{\chi}_i \left(i\partial - M_i \right) \chi_i - \left[\left(y_{\chi} \right)_{21} \overline{\chi}_2 \chi_1 \phi + \left(y_{\chi} \right)_{20} \overline{\chi}_2 \chi_0 \phi + \frac{\left(C_{H\chi\chi} \right)_{01}}{\Lambda} \overline{\chi}_0 \chi_1 |H|^2 + h.c. \right]$$

• "Leading" Production: $(v/\Lambda \ll 1, \alpha < 1)$

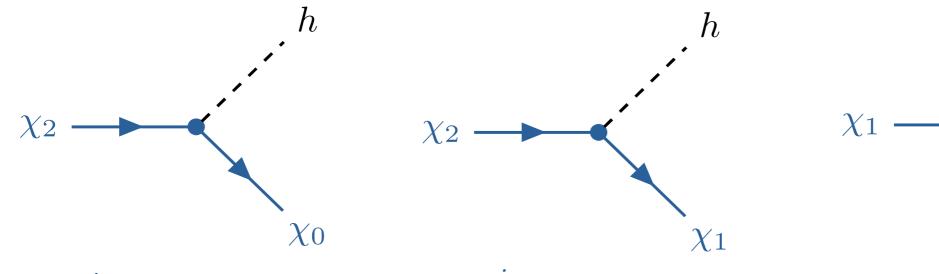


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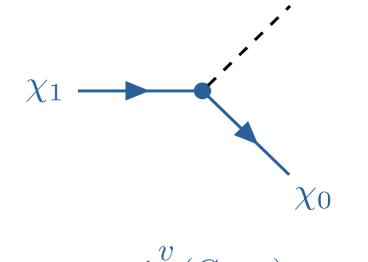
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• "Leading" Decays: $(v/\Lambda \ll 1, \alpha < 1)$



$$\simeq \frac{\imath}{\sqrt{2}}(y_{\chi})_{20}\sin\alpha$$

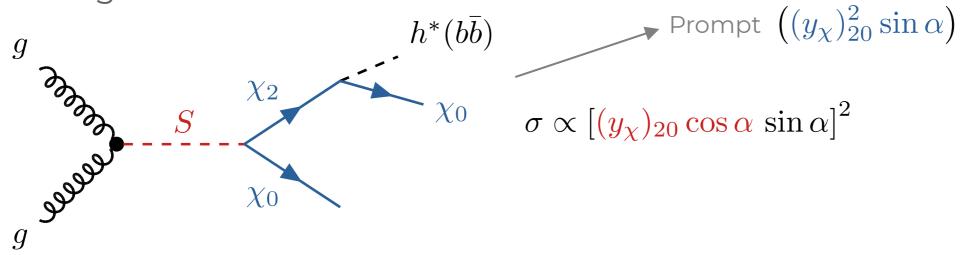
$$\simeq \frac{\imath}{\sqrt{2}}(y_{\chi})_{21} \sin \alpha$$

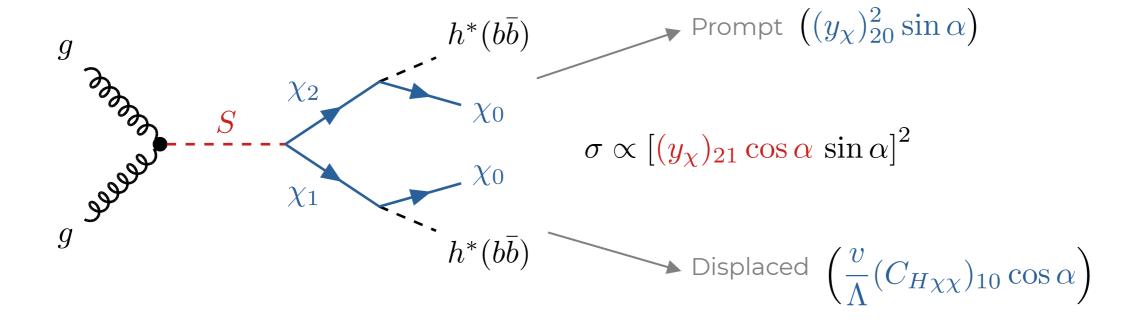


 $\simeq -i \frac{v}{\Lambda} (C_{H\chi\chi})_{10} \cos \alpha$

Note that chi2 → chi0+h dominates over chi2→chi1+h, if mchi2 ~ mchi1

• Leading channels:

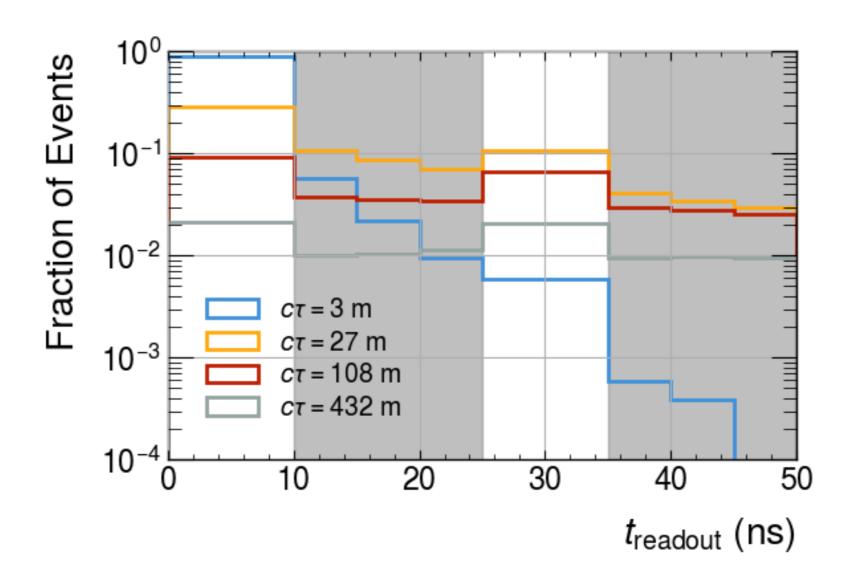


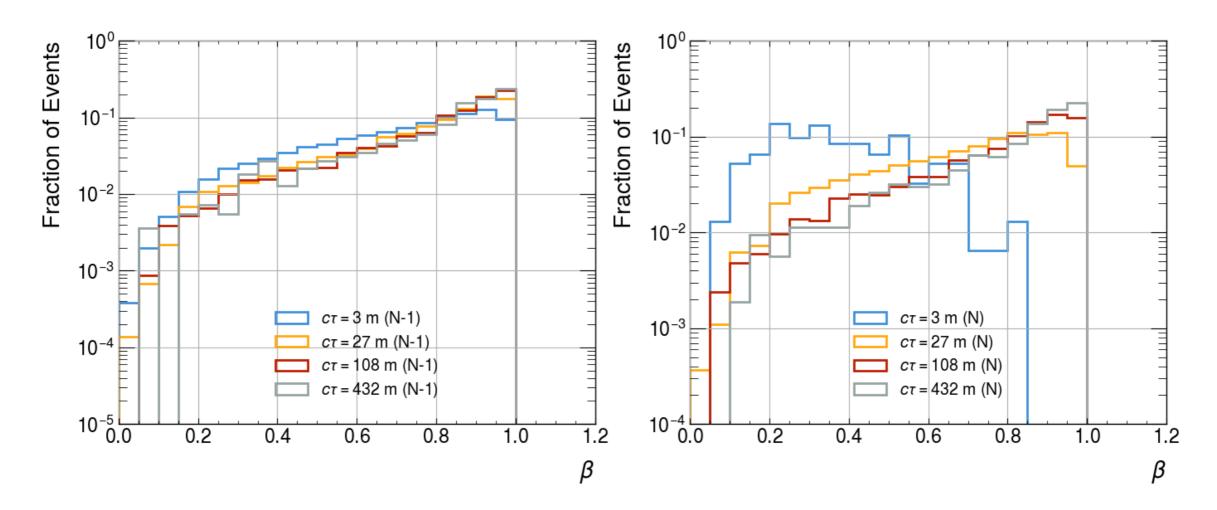


• Benchmark: $pp \to \chi_1 \chi_2, \ \chi_1 \to b\bar{b}\chi_0, \ \chi_2 \to b\bar{b}\chi_0,$

$$m_S = 500 \,\text{GeV}, \, m_2 = 250 \,\text{GeV}, \, m_1 = 244 \,\text{GeV}, \, m_0 = 154 \,\text{GeV}$$

 $(y_\chi)_{21} = (y_\chi)_{20} = 1, \, (C_{H\chi\chi})_{10} = 0.1, \, \sin\alpha = 0.2 \qquad c\tau = 3, 27, 108, 432 \,\text{m}$





The above behavior for the small lifetime in *N* events is due to the requirement that the readout time is given by:

$$t_r = t_d - L_d/c \sim t_d - c\beta t_d/c = t_d(1-\beta) \sim \gamma \tau_0(1-\beta) = \tau_0 \sqrt{\frac{1-\beta}{1+\beta}}$$

Hence for small proper lifetime we need a small boost to achieve large readout times.

As a result, decays taking place in the *N* event record will correspond to small beta and therefore result in smaller traveled distance (when compare to scenarios with larger lifetimes). This enhances the fraction of *N*-decays within the HCAL.

- f(N-1) = # Events with the LLP decaying in N/(Total # Events)
- f(HCAL) = # Events with the LLP decaying in the HCAL/(f(N) or f(N-1))

