

Research Project

Mathematical Modelling of Cell Membrane Dynamics

Isak Hammer



Department of Mathematical Sciences
Norwegian University of Science and Technology

1 Introduction

Cell membrane dynamics has recently had many important applications. For instance, it has been linked to detecting deceases such as Alzheimer's disease, cancer cells and is as well important for development new methods and vaccines [1].

Add more sources.

One of the primary components of the cell membranes are lipids which serve many different functions. A key function is that it is consisting of a bilayer of lipids which controls the structural rigidity and the fluidity of the membrane [2]. It also turns out that the lipids often accumulate into so-called lipid rafts which serves as a rigid platform for proteins with special properties such as intracellular trafficking of lipids and lipid-anchored proteins [3].

Modelling of lipid rafts formation can be modelled as a two-phase separation problem based on minimization of the Ginzburg-Landau energy functional [4]

$$\mathcal{E}_{ch}(\Gamma) = \int_{\Gamma} \Psi(c) + \frac{\gamma}{2} |\nabla c|^2,$$

which is describing the chemical energy for a concentration $c : \Gamma \times [0, T] \mapsto [0, 1]$ over a surface membrane Γ . Several authors have solved this problem often results by deriving variants of Cahn Hilliard Equation or Allen Cahn Equation if the concentration is not conserved both standstill and evolving domains [4, 5, 6, 7, 8, 4].

Assuming that the system is a single-phase system can the elastic bending energy be modelled using the Canham Helrich energy functional [9, 5]

$$\mathcal{E}_e(\Gamma) = \int_{\Gamma} c_b H^2 + c_k K$$

Here is $H = \kappa_1 + \kappa_2$ denoted as the mean curvature and $K = \kappa_1 \kappa_2$ as the gaussian curvature with respectively c_b and c_k as tuning parameters and κ_1 and κ_2 as principal curvatures. Using the Gauss-Bonnet theorem can it be shown that the problem above is equivalent to the so-called Willmore energy functional [10, 11]

$$\mathcal{E}(\Gamma) = \int_{\Gamma} \frac{1}{2} H^2. \quad (1)$$

This is a well known problem in the mathematical community [12, 13, 14]. In fact, it is a mathematical tool used to study the geometry of surfaces because it can be used to study the properties of minimal surfaces, which are surfaces with the least possible area for a given boundary. This is important in many areas of mathematics, including differential geometry, topology and mathematical physics [15, 16, 17].

It has been established many numerical methods for for shape optimization problems [18,

19], evolving surface partial differential equations (PDE) [20, 21, 22] and specific algorithms for the Willmore energy problem (1) [23, 24, 25, 26, 24].

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In this report will we establish a numerical scheme on minimization on this functional. However, we will first establish notation by including a section for definitions and important results from differential geometry. We will then derive the underlying PDE's for this equation. Lastly we will establish the model for the problem and discretize the problem using evolutionary parametric FEM methods.

2 Background Theory

2.1 Differential Calculus

This subsection is inspired by the notation used in [26, 27]. Let some initial surface $\Gamma^0 \subset \mathbb{R}^3$ smooth compact and oriented surface with no boundary where we can assign a unique point $p \in \Gamma^0$. We define the time evolutionary surface to be on the form,

$$\begin{aligned} \Gamma &= \Gamma(t) = \Gamma(\chi(p, t)) \\ &= \{\chi(p, t) : p \in \Gamma^0\} \end{aligned}$$

transformed via the smooth mapping,

$$\chi : \Gamma^0 \times [0, T] \mapsto \mathbb{R}^3.$$

An important regularity result is that if Γ^0 is of class C^∞ , then Γ is also of class C^∞ for $\forall t \in [0, T]$ [18, 27].

Formally in [18, p 48], it might be an idea to formulate normal unit-vector regularity as C^∞

We will define a unique evolutionary point $x \in \Gamma(t)$ based on the smooth mapping $\chi(p, t) = x$. A way to imagine this is to have a initial point in Γ^0 and the mapping χ describes how this point will deform over time. The outer unit normal vector field of $\Gamma(t)$ is defined as the mapping $\nu : \Gamma \mapsto \mathbb{R}^3$.

Using the notation presented in [27] and [26] can we define the basic surface differential operators. Consider a scalar function, $u : \Gamma \mapsto \mathbb{R}$, and a vector-valued function, $\hat{u} : \Gamma \mapsto \mathbb{R}^3$. We can then denote $\nabla_{\Gamma} u : \Gamma \mapsto \mathbb{R}^3$ as the tangential operator,

$$\nabla_{\Gamma} u = \nabla u - \langle \nu, \nabla u \rangle \nu.$$

May be an idea to define a extension $\tilde{u}|_{\Gamma}$ and look into regularity. See definitions in [20].

Similarly, for the vector-valued function is the operator defined s.t.

$$\nabla_{\Gamma} \hat{u} = (\nabla_{\Gamma} u_1, \nabla_{\Gamma} u_2, \nabla_{\Gamma} u_3)^T.$$

The surface divergence for a vector-valued function is defined as

$$\nabla_{\Gamma} \cdot \hat{u} = \nabla \cdot \hat{u} - \nu^T D\hat{u} \cdot \nu$$

Here $D\hat{u}$ denotes the Jacobian of \hat{u} . Similarly, the Laplace-Beltrami operator $\Delta_{\Gamma} u : \Gamma \mapsto \mathbb{R}$ is defined s.t.

$$\begin{aligned} \Delta_{\Gamma} u &= \nabla_{\Gamma} \cdot \nabla_{\Gamma} u \\ &= \Delta u - \nu^T D^2 u \cdot \nu - H \partial_{\nu} f. \end{aligned}$$

Here $D^2 u$ denotes as the Hessian of the scalar function u . In the case of a vector valued function is the operator defined as

$$\Delta_{\Gamma} \hat{u} = (\Delta_{\Gamma} u_1, \Delta_{\Gamma} u_2, \Delta_{\Gamma} u_3)^T$$

A method to compute the mean curvature and the so-called Frobenius norm of matrix A involves applying the extended Weingarten map, $A(x) = \nabla_{\Gamma} \nu(x)$, s.t. these identities holds ,

$$\begin{aligned} H &= \text{tr}(A) = k_1 + k_2, \\ |A|^2 &= k_1^2 + k_2^2, \end{aligned}$$

see [26]. We may also want to use these definitions to introduce the following identities ,

$$\begin{aligned} \partial_{\nu} H &= -|A|^2, \\ \nabla_{\Gamma} H &= \Delta_{\Gamma} \nu + |A|^2 \nu. \end{aligned}$$

Again, see Lemma 3.3 and Lemma 3.2 in [27].

2.2 Evolutionary Surface Dynamics

In this section will we develop a framework evolutionary surface dynamics.

First of all, we can denote the velocity $v : \Gamma \mapsto \mathbb{R}^3$ to be

$$\frac{dx}{dt} = v(x) \quad \forall x \in \Gamma(t). \quad (2)$$

Given a model of the velocity v can we solve the ordinary differential equation (ODE) (2) and determine the evolution of a point on the surface $\Gamma(t)$. In this article will we assume that the velocity only has a normal component to the surface, i.e., it exists a scalar function $V : \Gamma \mapsto \mathbb{R}$ s.t. $v = V\nu$.

Recall that the point $x = \chi(p, t)$ is arising from the smooth mapping from the point p in Γ^0 to $\Gamma(t)$. Now, let some arbitrary energy functional have the form,

$$\mathcal{J}(\Gamma(t)) = \int_{\Gamma(t)} \varphi(x).$$

For instance, in the case presented in (1) we define $\varphi = H^2$. Based on [27], the shape derivative of this energy functional at some time t in the direction of the velocity $v(x)$ from (2) is defined as

$$d\mathcal{J}(\Gamma(t); v) = \lim_{\varepsilon \rightarrow 0} \frac{\mathcal{J}(\Gamma(t + \varepsilon)) - \mathcal{J}(\Gamma(t))}{\varepsilon}.$$

Assume we have a scalar function $f : \Gamma(t) \mapsto \mathbb{R}$. Similarly, as for the shape derivative, we can now denote the material derivative at time t as

$$\begin{aligned} \frac{D}{Dt} f(x, t) &= \frac{d}{dt} f(\chi(p, t), t) \\ &= \lim_{\varepsilon \rightarrow 0} \frac{f(\chi(p, t + \varepsilon)) - f(\chi(p, t))}{\varepsilon} \end{aligned}$$

We denote the $L^2(\Gamma)$ as the space of all functions that are square-integrable with respect to the surface measure, i.e.,

$$L^2(\Gamma) = \left\{ u : \Gamma \mapsto \mathbb{R} \mid \int_{\Gamma} |u|^2 < \infty \right\}$$

Let $u, v \in L^2(\Gamma)$, then can we define the norm and the inner-product

$$\begin{aligned} \|u\|_{L^2(\Gamma)}^2 &= \int_{\Gamma} |u|^2 \\ (u, v)_{L^2(\Gamma)} &= \int_{\Gamma} uv \end{aligned}$$

In this paper will we also the shorthand notation $\|u\|_{L^2(\Gamma)} = \|u\|_{\Gamma}$ and $(u, v)_{L^2(\Gamma)} = (u, v)_{\Gamma}$. The Sobolev space $H^1(\Gamma)$ is defined as the space of all functions and its first weak derivative with a finite L^2 -norm, i.e.,

$$H^1(\Gamma) = \left\{ f : \Gamma \mapsto \mathbb{R} \mid \int_{\Gamma} |f|^2 + |\nabla_{\Gamma} f|^2 < \infty \right\},$$

with the following norm and inner product $u, v \in H^1(\Gamma)$,

$$\begin{aligned} \|u\|_{H^1(\Gamma)} &= \|u\|_{\Gamma} + \|\nabla_{\Gamma} u\|_{\Gamma}, \\ (u, v)_{H^1(\Gamma)} &= (u, v)_{\Gamma} + (\nabla_{\Gamma} u, \nabla_{\Gamma} v)_{\Gamma}. \end{aligned}$$

If we have a vector-valued function that $u : \Gamma \mapsto \mathbb{R}^3$ where each element is in $H^1(\Gamma)$ or $L^2(\Gamma)$, then do we denote is as a member of respectively $[H^1(\Gamma)]^3$ or $[L^2(\Gamma)]^3$.

The method we will use in this paper to minimize the energy functional (1) is to compute the so-called gradient flow. The fundamental idea of the gradient flow is to give rise of evolutionary dynamics to decrease the overall energy functional both in space and time, i.e., $\mathcal{J}(\Gamma(t_2)) < \mathcal{J}(\Gamma(t_1))$ for all $t_2 > t_1$. For information about gradient flows please check out [28, 29]. Now assume we have the velocity defined in (2) to be $v \in [L^2(\Gamma)]^3$, then we define the L^2 gradient flow s.t.

$$(v, \varphi)_{\Gamma} = -d\mathcal{J}(\Gamma; \varphi) \quad \forall \varphi \in [L^2(\Gamma)]^3.$$

It turns out, if $v \neq 0$ then is this equivalent to

$$d\mathcal{J}(\Gamma; v) = -\|v\|_{L^2(\Gamma)}^2 < 0. \quad (3)$$

Hence, we finally have a toolbox which can be used to model evolutionary dynamics for moving surfaces.

3 Evolutionary dynamics of the Willmore energy

Recall that we define the velocity (2) to only have a normal component, i.e., $v(x) = V\nu$. The shape derivative for (1) in the direction of some velocity $v \in [H^1(\Gamma)]^3$ has the form

$$d\mathcal{E}(\Gamma; v) = \int_{\Gamma} \left(-\Delta_{\Gamma} H + \frac{1}{2} H^3 - H |A|^2 \right) V$$

A complete derivation of the shape derivative can be found in [27, Corally 4.7]. Consequently, by applying the gradient flow in (3) and using that $\|v\|_{\Gamma} = \|V\|_{\Gamma}^2$ can we easily see that,

$$\|V\|_{\Gamma}^2 = \int_{\Gamma} V^2 = -d\mathcal{E}(\Gamma; v).$$

Hence, the gradient flow is equivalent to

$$V = \Delta_{\Gamma} H + Q, \quad (4)$$

where we denote the nonlinear term as $Q = -\frac{1}{2} H^3 + H |A|^2$.

Need to define a way to compute $\Delta_{\Gamma} H$ in the background theory.

From [26, Lemma 2.1] it derived that (4) must satisfy the following material derivatives,

$$\begin{aligned} \frac{D}{Dt} H &= -(\Delta_{\Gamma} + |A|^2) V, \\ \frac{D}{Dt} \nu &= (-\Delta_{\Gamma} + (HA - A^2)) z \\ &\quad + |\nabla_{\Gamma} H|^2 \nu - 2(\nabla_{\Gamma} \cdot (A \nabla_{\Gamma} H)) \nu \\ &\quad - A^2 \nabla_{\Gamma} H - \nabla_{\Gamma} Q. \end{aligned}$$

Here is the substitution variable introduced s.t. $z = \Delta_{\Gamma} \nu + |A|^2 \nu$.

It exists methods which do not exploit the material derivatives [25, 30]. However, it turns out that including these material derivatives brings additional computational costs, but provides so-called full-order approximation to the mean curvature, H , and the normal vector, ν , and thus, allows us to construct rigorous convergence proofs for evolving surface FEM, see [26].

An interesting recent used introducing tangential velocity [24] Finally we end up with the following fourth-order time dependent system of equation

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