

Research Project

Mathematical Modelling of Cell Membrane Dynamics

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1 Introduction

The value of understanding the basic underlying mechanics of the cell membrane dynamics has quite a lot of application. First can we detect diseases such as Alzheimer's disease, cancer cells and develop new methods and vaccines [1].

One of the primary components of the cell membranes are lipids which serve many different functions. A key function is that it is consisting of a bilayer of lipids which controls the structural rigidity and the fluidity of the membrane [2]. It also turns out that the lipids often accumulate into so-called lipid rafts which serves as a rigid platform for proteins with special properties such as intracellular trafficking of lipids and lipid-anchored proteins [3].

Modelling of lipid rafts formation can be modelled as a two-phase separation problem based on minimization of the Ginzburg-Landau energy functional [4]

$$\mathcal{E}_{ch}(\Gamma) = \int_{\Gamma} \Psi(c) + \frac{\gamma}{2} |\nabla c|^2 dx,$$

which is describing the chemical energy for a concentration $c : \Gamma \times [0, T] \mapsto [0, 1]$ over a surface membrane Γ . Several authors have solved this problem often results by deriving variants of Cahn Hilliard Equation or Allen Cahn Equation if the concentration is not conserved both standstill and evolving domains [4–7].

Assuming that the system is a single-phase system can the elastic bending energy be modelled using the Canham Helrich energy functional [8]

$$\mathcal{E}_e(\Gamma) = \int_{\Gamma} 2c_b H^2 + c_k K dx$$

Here is H denoted as the mean curvature and K as the gaussian curvature with respectively c_b and c_k as tuning parameters.

2 Cahn Hilliard Equation on a Closed Membrane

Let c_0 and c_1 indicate the concentration profile of the substances in a 2-phase system such that $c_0(\mathbf{x}, t) : \Omega \times [0, \infty] \rightarrow [0, 1]$ and similarly $c_1(\mathbf{x}, t) : \Omega \times [0, \infty] \rightarrow [0, 1]$, where \mathbf{x} is a element of some surface Ω and t is time. However, in the 2 phase problem will we will restrict ourself so that $c_0(t, \mathbf{x}) + c_1(t, \mathbf{x}) = 1$ at any \mathbf{x} at time t . A property of the restriction is that we now can express c_0 using c_1 , with no loss of information. Hence, let us now define $c = c_0$ so $c(\mathbf{x}, t) : \Omega \times [0, \infty] \rightarrow [0, 1]$. It has been shown that 2 phase system if thermodynamically unstable can be evolve into a phase separation described by a evolutionary differential equation [9] using a model based on chemical energy of the substances. However, further development has been done [4] to solve this equation on surfaces. Now assume model that we want to describe is a phase-separation on a closed membrane surface Γ , so that $c(\mathbf{x}, t) : \Gamma \times [0, T] \rightarrow [0, 1]$. Then is the surface Cahn Hilliard equation described such that

$$\rho \frac{\partial c}{\partial t} - \nabla_{\Gamma} (M \nabla_{\Gamma} (f'_0 - \varepsilon^2 \nabla_{\Gamma}^2 c)) = 0 \quad \text{on } \Gamma. \quad (1)$$

We define here the tangential gradient operator to be $\nabla_{\Gamma} c = \nabla c - (\mathbf{n} \cdot \nabla c) \mathbf{n}$ applied on the surface Γ restricted to $\mathbf{n} \cdot \nabla_{\Gamma} c = 0$.

Lets define ε to be the size of the layer between the substances c_1 and c_2 . The density ρ is simply defined such that $\rho = \frac{m}{S_{\Gamma}}$ is a constant based on the total mass divided by the total surface area of Γ . Here is the mobility M often derived such that is is dependent on c and is crucial for the result during a possible coarsening event [4]. However, the free energy per unit surface $f_0 = f_0(c)$ is derived based on the thermodynamical model and should according to [4] be non convex and nonlinear.

A important observation is that equation (1) is a fourth order equation which makes it more challenging to solve using conventional FEM methods. This clear when writing the equation on the equivalent weak form and second order equations arise.

3 Energy Functionals

Let $c(x, t) : \Gamma \times [0, T] \mapsto [0, 1]$. From [4] can we observe the energy functionals

$$E_1(c) = \int_{\Gamma} f(c).$$

where

$$f(c) = f_0(c) + \frac{1}{2}\varepsilon^2 |\nabla_{\Gamma} c|^2$$

and the conservation law $\rho \frac{\partial c}{\partial t} + \text{div}_{\Gamma} \mathbf{j} = 0$ for the evolution of c , derived from the Ficks Law $\mathbf{j} = -M \nabla_{\Gamma} \mu$ for the chemical potential derived by the functional derivative $\mu = \frac{\delta f}{\delta c}$. The double well function is denoted as

$$f_0(c) = \frac{\zeta}{4} c^2 (1 - c)^2$$

References

- [1] Scott A Small and Sam Gandy. “Sorting through the cell biology of Alzheimer’s disease: intracellular pathways to pathogenesis”. In: *Neuron* 52.1 (2006), pp. 15–31.
- [2] Saul L. Neidleman. “Effects of Temperature on Lipid Unsaturation”. In: *Biotechnology and Genetic Engineering Reviews* 5.1 (1987). PMID: 3314900, pp. 245–268. DOI: [10.1080/02648725.1987.10647839](https://doi.org/10.1080/02648725.1987.10647839). eprint: <https://doi.org/10.1080/02648725.1987.10647839>. URL: <https://doi.org/10.1080/02648725.1987.10647839>.
- [3] Michael Edidin. “The State of Lipid Rafts: From Model Membranes to Cells”. In: *Annual Review of Biophysics and Biomolecular Structure* 32.1 (2003). PMID: 12543707, pp. 257–283. DOI: [10.1146/annurev.biophys.32.110601.142439](https://doi.org/10.1146/annurev.biophys.32.110601.142439). eprint: <https://doi.org/10.1146/annurev.biophys.32.110601.142439>. URL: <https://doi.org/10.1146/annurev.biophys.32.110601.142439>.
- [4] Vladimir Yushutin et al. “A computational study of lateral phase separation in biological membranes”. In: *International Journal for Numerical Methods in Biomedical Engineering* 35.3 (2019). e3181 cnm.3181, e3181. DOI: <https://doi.org/10.1002/cnm.3181>. eprint: <https://onlinelibrary.wiley.com/doi/pdf/10.1002/cnm.3181>. URL: <https://onlinelibrary.wiley.com/doi/abs/10.1002/cnm.3181>.
- [5] Andreas Rätz. “A benchmark for the surface Cahn–Hilliard equation”. In: *Applied Mathematics Letters* 56 (2016), pp. 65–71. ISSN: 0893-9659. DOI: <https://doi.org/10.1016/j.aml.2015.12.008>. URL: <https://www.sciencedirect.com/science/article/pii/S0893965915300045>.
- [6] P. Gera and David Salac. “Cahn-Hilliard on surfaces: A numerical study”. In: *Appl. Math. Lett.* 73 (2017), pp. 56–61.
- [7] D. CAETANO and C. M. ELLIOTT. “Cahn–Hilliard equations on an evolving surface”. In: *European Journal of Applied Mathematics* 32.5 (2021), 937–1000. DOI: [10.1017/S0956792521000176](https://doi.org/10.1017/S0956792521000176).
- [8] Xiaolang Wang and Qiang Du. “Modelling and simulations of multi-component lipid membranes and open membranes via diffuse interface approaches”. In: *Journal of mathematical biology* 56 (Apr. 2008), pp. 347–71. DOI: [10.1007/s00285-007-0118-2](https://doi.org/10.1007/s00285-007-0118-2).

- [9] John W. Cahn and John E. Hilliard. “Free Energy of a Nonuniform System. I. Interfacial Free Energy”. In: *The Journal of Chemical Physics* 28.2 (1958), pp. 258–267. DOI: [10.1063/1.1744102](https://doi.org/10.1063/1.1744102). eprint: <https://doi.org/10.1063/1.1744102>. URL: <https://doi.org/10.1063/1.1744102>.