

Holonomic Optimal Control for Qudits

Tomas André

Uppsala University

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- 1 Introduction
- 2 Basic Concepts
- 3 Computation from Geometry
- 4 Geometric Quantum Computation
- 5 Results
 - The qutrit
 - Simulation results
 - Generalization
- 6 Conclusions

Outline

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Motivation

"Nature isn't classical, dammit, and if you want to make a simulation of nature, you'd better make it quantum mechanical, and by golly it's a wonderful problem, because it doesn't look so easy."

- Richard Feynman

History (1/2)

- 1980 Paul Benioff proposes a quantum mechanical model of the turing machine. ¹
- 1980s Both Richard Feynman and Yuri Manin suggests that a quantum based computer could perform calculations not possible on a conventional computer. ^{2 3}
- 1994 Peter Shor publishes a quantum algorithm used for finding the prime-factors of an integer. ⁴ (This is big!)

¹Benioff, Paul (1980). Journal of Statistical Physics.22(5):563.

²Feynman, Richard (June 1982). International Journal of Theoretical Physics. 21(6/7):467.

³Manin, Yu. I. (1980). Vychislimoe i nevychislimoe [Computable and Noncomputable]

⁴Shor, Peter (1994). IEEE Comput Soc Press:p.124

History (2/2)

- 1998 The first two-qubit quantum computer capable of running calculations is built by Isaac Chaung, Neil Gershenfeld and Mark Kubinec. ⁵
- 2019 Google claims to have achieved computations on a quantum machine that would be infeasible on a classical computer. ⁶
- Today Both IBM and Google have active QC research and IBM has a quantum computer with 127 qubits. Despite much focus fault tolerant quantum computers are still a far way off.

⁵Chuang, Isaac L.; Gershenfeld, Neil; Kubinec, Markdoi(1998). Phys. Rev. Lett. American Physical Society. 80(15):3408.

⁶Gibney, Elizabeth (October 2019). Nature. 574(7779):461.

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Quantum Computing

What is Quantum Computing?

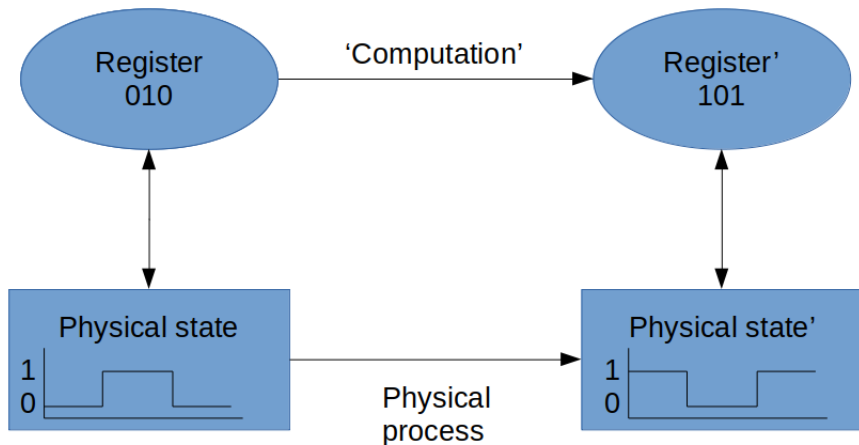
- Qubits
- Quantum Gates
- superposition
- entanglement
- it do be whack yo

Outline

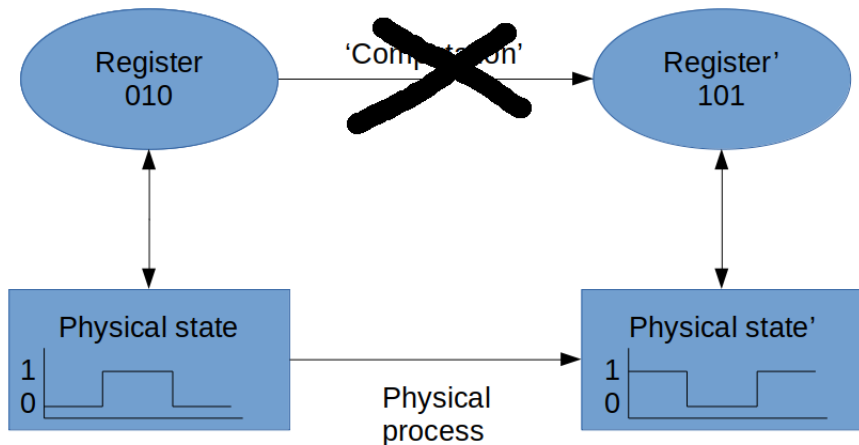
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What is computation?

What is computation?



What is computation?



Computation from Geometry

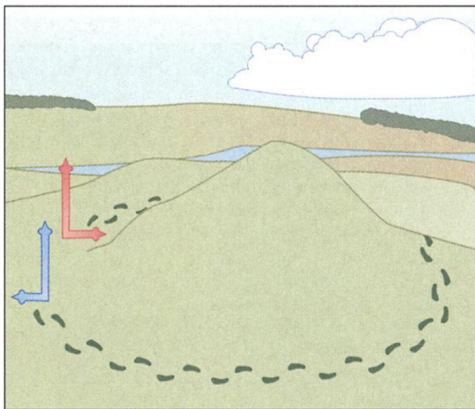


Image credit: Lloyd, Seth (june 2001). "Computation From Geometry". Science vol. 292

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Holonomies (1/2)

The change of a quantum state is called time-evolution, during one period of time T it is written $U(T, 0)$.

The Berry phase is an early example a 1D Holonomy in Quantum mechanics.⁷

$$|\psi(t)\rangle \xrightarrow{\text{time passes}} |\psi(t+T)\rangle = e^{i(\theta+\gamma)} |\psi(t)\rangle \quad (1)$$

in this case

$$U(0, T) = e^{i(\theta+\gamma)} \quad (2)$$

θ - dynamical phase, arises from the Hamiltonian of the system.

γ - 'berry phase', or geometric phase arises from the geometry of the 'landscape'.

⁷Berry, (19..)

Holonomies (2/2)

More generally the holonomy can be higher dimensional and non-abelian (matrix)

$$U(T, 0) = \mathcal{T} \exp \left(i \int_0^T \underbrace{\mathbf{A}(t)}_{\text{Geometric}} - \underbrace{\mathbf{H}(t)}_{\text{Dynamic}} dt \right) \quad (3)$$

where $\mathbf{A}_{mn}(t) = i \langle m(t) | \frac{\partial}{\partial t} | n(t) \rangle$ and $\mathbf{H}_{mn}(t) = \langle m(t) | \mathbf{H}(t) | n(t) \rangle$. In the non-adiabatic case one wishes to remove dynamical phase contribution, then one is left with the Wilczek-Zee holonomy⁸;

$$U(T, 0) = \mathcal{T} \exp \left(i \int_0^T \mathbf{A}(t) dt \right) = \mathcal{P} \exp \left(i \oint_C \mathbf{A}(s) ds \right) = U(C) \quad (4)$$

⁸hej

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The qutrit

Higher dimensional computational elements?

Bit: 0 or 1

Qubit: $\frac{1}{\sqrt{2}}(\alpha|0\rangle + |1\rangle)$

The qutrit

Higher dimensional computational elements?

bit: 0 or 1

$$\text{qubit: } \frac{1}{\sqrt{2}}(\alpha |0\rangle + |1\rangle)$$

trit: 0,1, or 2

$$\text{qutrit: } \frac{1}{\sqrt{3}}(\alpha |0\rangle + |1\rangle + |2\rangle)$$

The qutrit

Higher dimensional computational elements?

bit: 0 or 1

$$\text{qubit: } \frac{1}{\sqrt{2}}(\alpha |0\rangle + \beta |1\rangle)$$

trit: 0,1, or 2

$$\text{qutrit: } \frac{1}{\sqrt{3}}(\alpha |0\rangle + \beta |1\rangle + \gamma |2\rangle)$$

\vdots

dit 0,1,...,or d

$$\text{qudit: } \frac{1}{\sqrt{d}}(\alpha |0\rangle + \beta |1\rangle + \cdots + \Omega |d\rangle)$$

The qutrit

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The qutrit

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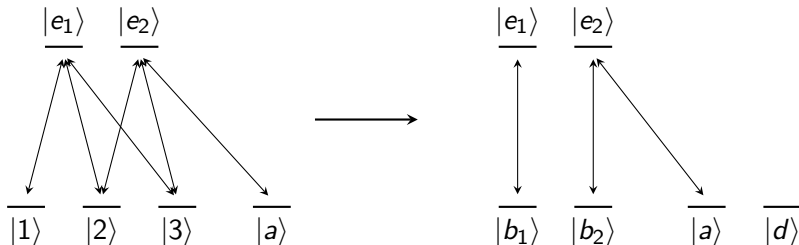
$$\text{qudit: } \frac{1}{\sqrt{d}}(\alpha |0\rangle + |1\rangle + \cdots + \Omega |d\rangle)$$

Pros and Cons

- more information per unit
- operations require fewer gates
- harder to control

Designing the system

Designing the system



$$H = \sum_{j=1}^2 \sum_{i=j}^3 \omega_{ij} |i\rangle \langle e_j| + \frac{\Omega_a(t)}{2} |a\rangle \langle e_2| + \text{h.c} \quad (5)$$

$$H_d = \sum_{j=1}^2 \frac{\Omega_j(t)}{2} e^{-i\phi_j} |b_j\rangle \langle e_j| + \frac{\Omega_a(t)}{2} |a\rangle \langle e_2| + \text{h.c} \quad (6)$$

Dark-Bright basis

The computational basis new computational basis

$\{|1\rangle, |2\rangle, |3\rangle\} \rightarrow \{|d\rangle, |b_1\rangle, |b_2\rangle\}$ parametrized by angles $\theta, \varphi, \chi, \xi$.

$$\begin{aligned}
 |d\rangle &= \cos \theta |1\rangle + e^{i\chi} \sin \theta \cos \varphi |2\rangle + e^{i\xi} \sin \theta \sin \varphi |3\rangle \\
 |b_1\rangle &= \frac{1}{\sqrt{1 - \sin^2 \theta \sin^2 \varphi}} \left(-e^{-i\chi} \sin \theta \cos \varphi |1\rangle + \cos \theta |2\rangle \right) \\
 |b_2\rangle &= \frac{1}{\sqrt{1 - \sin^2 \theta \sin^2 \varphi}} \left(\frac{1}{2} \sin 2\theta \sin \varphi |1\rangle + \frac{e^{i\chi}}{2} \sin^2 \theta \sin 2\varphi |2\rangle \right. \\
 &\quad \left. + e^{i\xi} (\sin^2 \theta \sin^2 \varphi - 1) |3\rangle \right)
 \end{aligned} \tag{7}$$

Unitary gate

Constructing the unitary from the basis

$$\begin{aligned}
 U_1 &= |d\rangle \langle d| - i |e_1\rangle \langle b_1| - i |e_2\rangle \langle b_2|, \quad \phi_1 = \phi_2 = 0 \\
 U_2 &= |d\rangle \langle d| + ie^{i\gamma_1} |b_1\rangle \langle e_1| + ie^{i\gamma_2} |b_2\rangle \langle e_2|, \quad \phi_1 = -\gamma_1, \quad \phi_2 = -\gamma_2
 \end{aligned} \tag{8}$$

$$U(\theta, \varphi, \chi, \xi, \gamma_1, \gamma_2) = U_2 U_1 = |d\rangle \langle d| + e^{i\gamma_1} |b_1\rangle \langle b_1| + e^{i\gamma_2} |b_2\rangle \langle b_2|. \tag{9}$$

Dark path

$$\begin{aligned} |D_1(t)\rangle &= \cos u e^{-i\phi_1} |b_1\rangle + i \sin u |e_1\rangle \\ |D_2(t)\rangle &= \cos u \cos v e^{-i\phi_2} |b_2\rangle - i \sin u |e_2\rangle - \cos u \sin v |a\rangle \end{aligned} \quad (10)$$

such that $u(0) = u(T) = v(0) = v(T) = 0$

$$|\psi(t)\rangle = f_0 |d\rangle + f_1 |D_1(t)\rangle + f_2 |D_2(t)\rangle, t \in [0, T]. \quad (11)$$

so for example

$$\begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} = f_0 |d\rangle + f_1 |D_1(0)\rangle + f_2 |D_2(0)\rangle \xrightarrow{t \rightarrow T} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = f_0 |d\rangle + f_1 |D_1(T)\rangle + f_2 |D_2(T)\rangle \quad (12)$$

or equivalent

$$\begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \rightarrow U \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad (13)$$

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Conclusions