

Improving multiphysics simulation through port-Hamiltonian system theory

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Outline

Multiphysics problems

Port-Hamiltonian systems as a unified language for multiphysics

- Functional analytic structure

- The geometric definition

Mimetic discretization of port-Hamiltonian systems

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Challenges in multiphysics problems

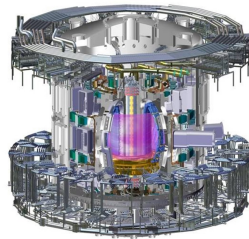
Multiphysics problems are commonly found in industrial applications.



Aeroelasticity



Thermoelasticity



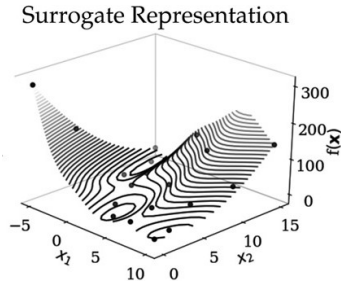
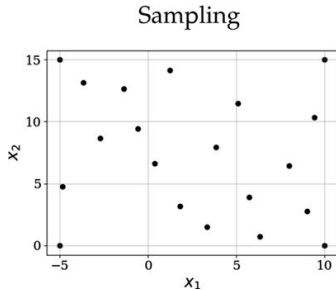
Magnetohydrodynamics

Challenges:

- ▶ Coupling between different models.
- ▶ Huge computational cost due to the large size of the models.
- ▶ Multidisciplinary optimization for dynamical systems.

Typical workflow in industry

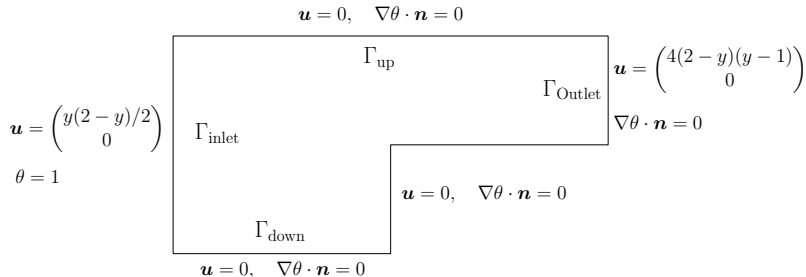
- ▶ Specific modelling and numerical methods for each physical domain.
 - ✗ The open character of systems is not properly considered.
 - ✗ Numerical methods do not preserve the structure required to interconnect systems.
- ▶ Model reduction via statistical methods.
 - ✗ The physical structure of the model is lost and first principles are violated.
 - ✗ This methodology does not generalize to different problems.



Example: convection dominated transport

Convection dominated transport of a passive scalar field in a Stokes flow¹

$$\begin{aligned}\nu \Delta \mathbf{u} + \nabla p &= 0, & \mathbf{u} : \text{Velocity}, \\ \nabla \cdot \mathbf{u} &= 0, & p : \text{Pressure}, \\ -\varepsilon \Delta \theta + \mathbf{u} \cdot \nabla \theta &= 0. & \theta : \text{Temperature}.\end{aligned}$$



Geometry and boundary conditions

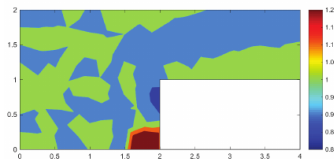
¹Volker John et al. "On the Divergence Constraint in Mixed Finite Element Methods for Incompressible Flows". In: *SIAM Review* 59.3 (2017), pp. 492–544. DOI: 10.1137/15M1047696.

When multiphysics goes wrong

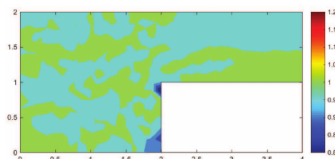
Exact solution for the temperature $\theta_{\text{ex}} = 1$.

- ▶ (\mathbf{u}, p) discretized using the Taylor-Hood element $\mathbb{P}_2/\mathbb{P}_1$;
- ▶ θ discretized via Voronoi finite volume method.

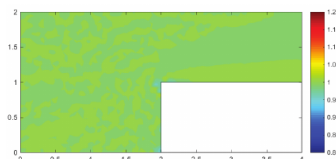
The Taylor-Hood element does not lead to divergence free velocity $\|\nabla \cdot \mathbf{u}\|_{L^2(\Omega)} \neq 0$.



Refinement 1



Refinement 2



Refinement 3

Figure: Discrete temperature field θ obtained

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A unified language for multiphysics in engineering

The port-Hamiltonian (pH) paradigm provides a language to understand multiphysics:

- ▶ The idea of **interconnection** is formalized as **duality pairing**.
- ▶ **Physics** is at the core: port-Hamiltonian systems are **passive** with respect to the **energy storage function**.
- ▶ The **topological** and **metrical** structure of the equation is clearly separated (mimetic discretization).



Finite dimensional pH systems

Still not a well established theory

There is **not a unique definition** of pH systems, even in finite dimension.

Definition (Finite dimensional pH system)

The time-invariant dynamical system

$$\begin{aligned}\dot{\mathbf{x}} &= (\mathbf{J} - \mathbf{R}) \nabla H(\mathbf{x}) + \mathbf{B}\mathbf{u}, \\ \mathbf{y} &= \mathbf{B}^\top \nabla H(\mathbf{x}),\end{aligned}$$

where \mathbf{x} is the state, \mathbf{u} the control input, \mathbf{y} the collocated output and

- ▶ $H(\mathbf{x}) : \mathbb{R}^n \rightarrow \mathbb{R}$, the Hamiltonian, is bounded from below.
- ▶ $\mathbf{J} = -\mathbf{J}^\top$ the interconnection operator.
- ▶ $\mathbf{R} = \mathbf{R}^\top \in \mathbb{R}^{n \times n}$, $\mathbf{R} \geq 0$ the resistive operator.
- ▶ $\mathbf{B} \in \mathbb{R}^{n \times m}$ the control operator.

is a pH system.

The geometric structure of pH systems²

Definition (Finite dimensional Dirac structure)

Given a finite-dimensional vector space F and its dual $E = F'$ with respect to the duality product $\langle \cdot | \cdot \rangle : E \times F \rightarrow \mathbb{R}$, consider the symmetric bilinear form:

$$\langle\langle (\mathbf{f}_1, \mathbf{e}_1), (\mathbf{f}_2, \mathbf{e}_2) \rangle\rangle := \langle \mathbf{e}_1 | \mathbf{f}_2 \rangle + \langle \mathbf{e}_2 | \mathbf{f}_1 \rangle, \text{ where } (\mathbf{f}_i, \mathbf{e}_i) \in B, i = 1, 2.$$

A Dirac structure on $B := F \times E$ is a linear subspace $D \subset B$ which equals its orthogonal companion with respect to $\langle\langle \cdot, \cdot \rangle\rangle$, i.e. $D = D^{[\perp]}$, where:

$$D^{[\perp]} := \left\{ (\mathbf{f}, \mathbf{e}) \in B \mid \langle\langle (\mathbf{f}, \mathbf{e}), (\hat{\mathbf{f}}, \hat{\mathbf{e}}) \rangle\rangle = 0, \forall (\hat{\mathbf{f}}, \hat{\mathbf{e}}) \in D \right\}.$$

Theorem

Given a finite-dimensional vector space F and its dual $E = F'$ a subspace $D \subset F \times E$ is a Dirac structure iff $\langle \mathbf{e} | \mathbf{f} \rangle = 0$ and $\dim D = \dim F$.

²T. J. Courant. "Dirac manifolds". In: *Transactions of the American Mathematical Society* 319.2 (1990), pp. 631–661. ISSN: 0002-9947. DOI: 10.2307/2001258.

Dirac structure and pH systems

From classical matrix factorization $\exists \mathbf{G} \in \mathbb{R}^{k \times n}$ and $\mathbf{K} = \mathbf{K}^\top \in \mathbb{R}^{k \times k}$, $\mathbf{K} \geq 0$ such that $\mathbf{R} = \mathbf{G}^\top \mathbf{K} \mathbf{G}$.

Dirac structure representation

Considering the following **ports**:

- ▶ the **storage ports** $(\mathbf{f}_x, \mathbf{e}_x) := (-\dot{\mathbf{x}}, \nabla H(\mathbf{x})) \in \mathbb{R}^n \times \mathbb{R}^n$;
- ▶ the **resistive ports** $(\mathbf{f}_r, \mathbf{e}_r) \in \mathbb{R}^k \times \mathbb{R}^k$;
- ▶ the **interconnection ports** $(\mathbf{f}_u, \mathbf{e}_u) := (\mathbf{y}, \mathbf{u}) \in \mathbb{R}^m \times \mathbb{R}^m$.

Given this port behavior, the pH system rewrites

$$\begin{pmatrix} \mathbf{f}_x \\ \mathbf{f}_r \\ \mathbf{f}_u \end{pmatrix} = \underbrace{\begin{bmatrix} -\mathbf{J} & \mathbf{G}^\top & -\mathbf{B} \\ \mathbf{G} & \mathbf{0} & \mathbf{0} \\ \mathbf{B}^\top & \mathbf{0} & \mathbf{0} \end{bmatrix}}_{\mathbf{J}_e} \begin{pmatrix} \mathbf{e}_x \\ \mathbf{e}_r \\ \mathbf{e}_u \end{pmatrix}, \quad \mathbf{e}_r = \mathbf{K} \mathbf{f}_r.$$

Since \mathbf{J}_e is skewsymmetric its graph defines a Dirac structure.

A simple definition³

Definition (Port-Hamiltonian system)

Let X_S , X_R , X_P be Banach spaces. A port-Hamiltonian system is a triple $(\mathcal{D}, H, \mathcal{R})$:

- ▶ $\mathcal{D} \subset (X_S, X_R, X_P) \times (X'_S, X'_R, X'_P)$ is a Dirac structure.
- ▶ $\mathcal{H} : U \rightarrow \mathbb{R}$ (with $U \subset X_S$ open) is a Hamiltonian.
- ▶ $\mathcal{R} \subset X_R \times X'_R$ is a resistive relation.

The behavior of the pH system on an interval $\mathbb{I} \subset \mathbb{R}$ consists of all (x, f_R, f_P, e_R, e_P)

- ▶ $x \in W_{\text{loc}}^{1,2}(\mathbb{I}, X_S)$, and $x(t) \in U$, $\forall t \in \mathbb{I}$,
- ▶ $(f_R, e_R) \in L_{\text{loc}}^2(\mathbb{I}; X_R \times X'_R)$ and $(f_P, e_P) \in L_{\text{loc}}^2(\mathbb{I}; X_P \times X'_P)$

that fulfill the differential inclusion

$$\left(-\frac{dx}{dt}, f_R, f_P, D\mathcal{H}(x(t)), e_R, e_P\right) \in \mathcal{D}, \quad (f_R, e_R) \in \mathcal{R}, \quad \text{for almost all } t \in \mathbb{I}.$$

³Timo Reis. "Some notes on port-Hamiltonian systems on Banach spaces". In: *IFAC-PapersOnLine* 54.19 (2021). 7th IFAC Workshop on Lagrangian and Hamiltonian Methods for Nonlinear Control LHMNC 2021, pp. 223–229. DOI: 10.1016/j.ifacol.2021.11.082.

Some mathematical definitions

Dirac structure

Let X be a Banach space. A subspace $\mathcal{D} \subset X \times X'$ is called a Dirac structure, if $\forall f \in X, e \in X'$, it holds

$$(f, e) \in \mathcal{D} \iff \left(\langle \hat{e} | f \rangle + \langle e | \hat{f} \rangle = 0, \quad \forall (\hat{f}, \hat{e}) \in \mathcal{D} \right).$$

Hamiltonian

Let X be a Banach space and $U \subset X$ be open. A mapping $\mathcal{H} : U \rightarrow \mathbb{R}$ is a Hamiltonian if it is locally Lipschitz continuous and Gâteaux differentiable

Resistive relation

Let X be a Banach space. A relation $\mathcal{R} \subset X \times X'$ is called resistive, if

$$\langle e | f \rangle \leq 0, \quad \forall (f, e) \in \mathcal{R}.$$

Operators

If $J \in \mathcal{L}(X', X)$ is a skew-dual operator $\langle w | Jv \rangle = \langle v | -Jw \rangle \forall v, w \in X'$ then $D = \{(Je, e) : e \in X'\}$ is a Dirac structure⁴.

If $K : X \rightarrow X'$ is dissipative $\langle K(x) | x \rangle \leq 0, \forall x \in X$, then $\mathcal{R} = \{(K(f), f) : f \in X\}$ is a resistive relation.

$$\begin{pmatrix} -\partial_t x \\ f_{\mathcal{R}} \\ f_{\mathcal{P}} \end{pmatrix} = J \begin{pmatrix} D_x \mathcal{H} \\ e_{\mathcal{R}} \\ e_{\mathcal{P}} \end{pmatrix}, \quad e_{\mathcal{R}} = K(f_{\mathcal{R}}).$$

⁴T. Reis and T. Stykel. “Passivity, Port-Hamiltonian Formulation and Solution Estimates for a Coupled Magneto-Quasistatic System”. In: *arXiv preprint arXiv:2205.15259* (2022).

Example: the wave equation

Consider the Hamiltonian

$$\mathcal{H} = (p, \kappa p)_{L^2(\Omega)} + (\mathbf{v}, \rho^{-1} \mathbf{v})_{L^2(\Omega, \mathbb{R}^3)}.$$

where κ is the Bulk modulus and ρ is the density.

The wave equation on $\Omega \subset \mathbb{R}^3$ with Dirichlet boundary condition reads:

$$\frac{\partial}{\partial t} \begin{pmatrix} p \\ \mathbf{v} \end{pmatrix} = \begin{bmatrix} 0 & \text{div} \\ \text{grad}_w & 0 \end{bmatrix} \begin{pmatrix} D_p \mathcal{H} \\ D_{\mathbf{v}} \mathcal{H} \end{pmatrix}, \quad \gamma_0(D_p \mathcal{H}) = u \in H^{1/2}(\partial\Omega),$$

where grad_w corresponds to a weak gradient and γ_0 is the Dirichlet trace.

In this case: $X_{\mathcal{S}} = L^2(\Omega) \times H^{\text{div}}(\Omega)'$, $X_{\mathcal{R}} = \emptyset$, $X_{\mathcal{P}} = H^{-1/2}(\partial\Omega)$ and

$$J = \begin{bmatrix} 0 & -\text{div} & 0 \\ -\text{grad}_w & 0 & \text{Id} \\ 0 & \gamma_{\mathbf{n}} & 0 \end{bmatrix}$$

where $\gamma_{\mathbf{n}}$ is the normal trace.

Example: the Maxwell equations

Consider the Hamiltonian:

$$\mathcal{H} = \frac{1}{2}(\mathbf{D}, \varepsilon^{-1} \mathbf{D})_{L^2(\Omega, \mathbb{R}^3)} + \frac{1}{2}(\mathbf{B}, \mu^{-1} \mathbf{B})_{L^2(\Omega, \mathbb{R}^3)}.$$

where ε is the electric permittivity and μ is the magnetic permeability.

The Maxwell equation on $\Omega \subset \mathbb{R}^3$ with conducting boundary condition reads:

$$\frac{\partial}{\partial t} \begin{pmatrix} \mathbf{D} \\ \mathbf{B} \end{pmatrix} = \begin{bmatrix} 0 & \text{curl} \\ -\text{curl}_w & 0 \end{bmatrix} \begin{pmatrix} D_{\mathbf{D}} \mathcal{H} \\ D_{\mathbf{B}} \mathcal{H} \end{pmatrix}, \quad D_{\mathbf{D}} \mathcal{H} \times \mathbf{n}|_{\partial\Omega} = \mathbf{E} \times \mathbf{n} = 0,$$

where curl_w corresponds to a weak curl operator and the field \mathbf{D} , \mathbf{B} satisfy

$$\nabla \cdot \mathbf{D} = 0, \quad \nabla \cdot \mathbf{B} = 0.$$

In this case: $X_{\mathcal{S}} = L^2(\Omega, \mathbb{R}^3 | \text{div} = 0) \times H_0^{\text{curl}}(\Omega | \text{div} = 0)'$, $X_{\mathcal{R}} = \emptyset$, $X_{\mathcal{P}} = \emptyset$ and

$$J = \begin{bmatrix} 0 & \text{curl} \\ \text{curl}_w & 0 \end{bmatrix}.$$

And many more

The same framework applies to

- ▶ Linear and non-linear solid mechanics (beams, plates, shells, etc.).
- ▶ Fluid dynamics.
- ▶ Chemical reactions.

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The canonical geometric port-Hamiltonian system

Distributed port-Hamiltonian were initially defined in a differential geometric setting⁵. In this setting the duality is the Hodge duality, given by the Hodge star \star . Given two fields of smooth differential forms $\alpha^p \in \Lambda^p(\Omega)$ and $\beta \in \Lambda^q(\Omega)$ the following systems

$$\begin{pmatrix} \partial_t \alpha^p \\ \partial_t \beta^q \end{pmatrix} = - \begin{bmatrix} 0 & (-1)^r d \\ d & 0 \end{bmatrix} \begin{pmatrix} \delta_\alpha H^{n-p} \\ \delta_\beta H^{n-q} \end{pmatrix},$$

together with appropriate boundary conditions are port-Hamiltonian distributed systems.

⁵A.J. van der Schaft and B.M. Maschke. "Hamiltonian formulation of distributed-parameter systems with boundary energy flow". In: *Journal of Geometry and Physics* 42.1 (2002), pp. 166–194. DOI: 10.1016/S0393-0440(01)00083-3.






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