

Port-Hamiltonian formulation and Symplectic discretization of plate models

Part II : Kirchhoff model for thin plates

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Abstract

The mechanical model of a thin plate with boundary control and observation is presented as a port-Hamiltonian system (PHs¹), both in vectorial and tensorial forms: the Kirchhoff-Love model of a plate is described by using a Stokes-Dirac structure and this represents a novelty with respect to the existing literature. This formulation is carried out both in vectorial and tensorial forms. Thanks to tensorial calculus, this model is found to mimic the interconnection structure of its one-dimensional counterpart, i.e. the Euler-Bernoulli beam.

The Partitioned Finite Element Method (PFEM²) is then extended to obtain a suitable, i.e. structure-preserving, weak form. The discretization procedure, performed on the vectorial formulation, leads to a finite-dimensional port-Hamiltonian system. This part II of the companion paper extends part I, dedicated to the Mindlin model for thick plates. The thin plate model comes along with additional difficulties, because of the higher order of the differential operator under consideration.

Keywords: port-Hamiltonian systems, Kirchhoff plate, Partitioned Finite Element Method, geometric spatial discretization, boundary control.

Introduction

As presented in part I of this companion paper, the port-Hamiltonian (PH) formalism [1, 2, 3] allows the structured modeling and discretization of multi-physics applications involving interconnected finite- and infinite-dimensional systems [4, 5]. Preserving the port-Hamiltonian structure in the discretization process is a keypoint to take benefit of this powerful formalism. This issue was first addressed in [6], with a mixed finite element spatial discretization method, and in [7], with pseudo-spectral methods relying on higher-order global polynomial approximations. All those methods are difficult to implement, especially for those system the spatial dimension of which is bigger than one. Very recently weak formulations

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¹PHs stands for port-Hamiltonian systems.

²PFEM stands for partitioned finite element method.

which lead to Galerkin numerical approximations began to be explored: in [8], a structure preserving finite element method was introduced for the wave equation in a two-dimensional domain; this method exhibits good results, both in the spectral analysis and simulation part, though requiring of a primal and a dual mesh on the geometry of the problem. Another approach is the partitioned finite element method (PFEM) proposed in [9], already largely explored in part I of this companion paper. The advantages of this latter methodology are its simplicity of implementation and its potential to carry over to a wide set of examples, no matter the spatial dimension of the problem. The possible use of open source software like FEniCS [10] or Firedrake [11] is also an appealing feature of this latter method

In part II of this companion paper, the modeling and discretization of thin plates described by the Kirchhoff-Love plate model is carried out within the PH framework, allowing for boundary control and observation. Compared to the existing literature on either the symplectic Hamiltonian formulation of [12, 13], or the port-Hamiltonian framework through jet theory in [14, 15], the main contribution of this paper concerns the PH formalism based on a Stokes-Dirac structure. This formalism is presented both in vectorial and tensorial forms. Moreover, the tensorial formalism [16, Chapter 16] highlights that this model mimics the interconnection structure of its one-dimensional counterpart, i.e. the Euler-Bernoulli beam. Compared to part I dedicated to thick plate Mindlin model in which first-order differential operators are explored in dimension two, and compared to [17] in which second- or higher-order differential operators were explored in dimension one only, the contribution of this paper is the PH formalism of systems of dimension two described with second-order differential operators, such as the Kirchhoff-Love model. The model, once written in a tensorial form, highlights new insights on second-order differential operators: especially the double divergence and the Hessian are proved to be adjoint operators one of another, which represents another important contribution of this paper. Finally, the extension of the PFEM method to the structure-preserving discretization of the Kirchhoff model is also a novelty of the paper. It allows simple implementation of numerical schemes compared to the jet theory formalism, while preserving the structure of PHS at the discrete level. **The last section is dedicated to numerical studies of this model by using Firedrake [11].**

1. Second-order distributed PH systems: Euler-Bernoulli beam

The Euler-Bernoulli beam is the one-dimensional equivalent of the Kirchhoff-Love plate. This model consists of one PDE, describing the vertical displacement along the beam length

$$\rho(x) \frac{\partial^2 w}{\partial t^2}(x, t) + \frac{\partial^2}{\partial x^2} \left(EI(x) \frac{\partial^2 w}{\partial x^2} \right) = 0, \quad x \in (0, L), t \geq 0 \quad (1)$$

where $w(x, t)$ is the transverse displacement of the beam. The coefficients $\rho(x)$, $E(x)$ and $I(x)$ are the mass per unit length, Young's modulus of elasticity and the moment of inertia of a cross section. The energy variables are then chosen as follows

$$\begin{aligned} \alpha_w &= \rho(x) \frac{\partial w}{\partial t}(x, t), & \text{Linear Momentum,} \\ \alpha_\kappa &= \frac{\partial^2 w}{\partial x^2}(x, t), & \text{Curvature.} \end{aligned} \quad (2)$$

Those variables are collected in the vector $\boldsymbol{\alpha} = (\alpha_w, \alpha_\kappa)^T$, so that the Hamiltonian can be written as a quadratic functional in the energy variables

$$H = \frac{1}{2} \int_0^L \boldsymbol{\alpha}^T Q \boldsymbol{\alpha} \, dx, \quad \text{where} \quad Q = \begin{bmatrix} \frac{1}{\rho(x)} & 0 \\ 0 & EI(x) \end{bmatrix}. \quad (3)$$

The co-energy variables are found by computing the variational derivative of the Hamiltonian

$$\begin{aligned} e_w &:= \frac{\delta H}{\delta \alpha_w} = \frac{\partial w}{\partial t}(x, t), & \text{Vertical velocity,} \\ e_\kappa &:= \frac{\delta H}{\delta \alpha_\kappa} = EI(x) \frac{\partial^2 w}{\partial x^2}(x, t), & \text{Flexural momentum.} \end{aligned} \quad (4)$$

Those variables are again collected in vector $\mathbf{e} = (e_w, e_\kappa)^T$, so that the underlying interconnection structure is then found to be

$$\frac{\partial \boldsymbol{\alpha}}{\partial t} = J \mathbf{e}, \quad \text{where} \quad J = \begin{bmatrix} 0 & -\frac{\partial^2}{\partial x^2} \\ \frac{\partial^2}{\partial x^2} & 0 \end{bmatrix}. \quad (5)$$

For an infinite-dimensional system, boundary variables have to be defined as well. Those can be found by evaluating the energy rate flow across the boundary. One possible choice among others (see [18] for a more exhaustive explanation) for this model is the following

$$\mathbf{f}_\partial = \begin{pmatrix} e_w(0) \\ \frac{\partial e_w}{\partial x}(0) \\ \frac{\partial e_\kappa}{\partial x}(L) \\ e_\kappa(L) \end{pmatrix}, \quad \mathbf{e}_\partial = \begin{pmatrix} \frac{\partial e_\kappa}{\partial x}(0) \\ -e_\kappa(0) \\ -e_w(L) \\ \frac{\partial e_w}{\partial x}(L) \end{pmatrix}. \quad (6)$$

The power flow is then easily evaluated as

$$\frac{d}{dt} H(t) = \int_0^L \frac{\partial \boldsymbol{\alpha}}{\partial t} \cdot \mathbf{e} \, dx = \langle \mathbf{e}_\partial, \mathbf{f}_\partial \rangle_{\mathbb{R}^4}. \quad (7)$$

The flow variables can now be defined as $\mathbf{f} = -\frac{\partial \boldsymbol{\alpha}}{\partial t}$, so that the flow space is given by the tuples $(\mathbf{f}, \mathbf{f}_\partial) \in \mathcal{F}$. Equivalently the effort space is given by $(\mathbf{e}, \mathbf{e}_\partial) \in \mathcal{E}$. The bond space is therefore the Cartesian product of the two

$$\mathcal{B} := \{(\mathbf{f}, \mathbf{f}_\partial, \mathbf{e}, \mathbf{e}_\partial) \in \mathcal{F} \times \mathcal{E}\}. \quad (8)$$

The duality pairing between elements of \mathcal{B} is then defined as follows

$$\ll ((\mathbf{f}^a, \mathbf{f}_\partial^a), (\mathbf{e}^a, \mathbf{e}_\partial^a)), ((\mathbf{f}^b, \mathbf{f}_\partial^b), (\mathbf{e}^b, \mathbf{e}_\partial^b)) \gg := \int_0^L \{(\mathbf{f}^a)^T \mathbf{e}^b + (\mathbf{f}^b)^T \mathbf{e}^a\} \, dx + (\mathbf{f}_\partial^a)^T \mathbf{e}_\partial^b + (\mathbf{f}_\partial^b)^T \mathbf{e}_\partial^a. \quad (9)$$

The Stokes-Dirac structure for the Euler-Bernoulli beam is therefore the following

Theorem 1 (From [17], Stokes-Dirac structure for the Bernoulli beam). *Consider the space of power variables \mathcal{B} defined in (8) and the bilinear form (+pairing operator) \ll, \gg given by (9). Define the following linear subspace $\mathcal{D} \subset \mathcal{F} \times \mathcal{E}$*

$$\mathcal{D} = \{(\mathbf{f}, \mathbf{f}_\partial, \mathbf{e}, \mathbf{e}_\partial) \in \mathcal{F} \times \mathcal{E} \mid \mathbf{f} = -J\mathbf{e}\}, \quad (10)$$

where $\mathbf{f}_\partial, \mathbf{e}_\partial$ were defined in (6). Then $\mathcal{D} = \mathcal{D}^\perp$, where \mathcal{D}^\perp is understood in the sense of orthogonality with respect to the bilinear product \ll, \gg , i.e \mathcal{D} is a Stokes-Dirac structure.

Remark 1. *For what concerns the use of this model for control and simulation purposes, the reader can refer to [19] for a stability and stabilization proof of the Euler-Bernoulli beam or to [20] for an illustration of rotating spacecraft with flexible appendages model as PH Bernoulli beams.*

2. Kirchhoff-Love theory for thin plates

In this section the classical variational approach (Hamilton's principle) to derive the equation of motions is first detailed. The physical quantities involved and the different energies, of utmost importance for the PH formalism, are reminded.

2.1. Model and associated variational formulation

The Kirchhoff-Love plate formulation rests on the hypothesis of small thickness compared to the in plane dimensions. The notations and symbols are borrowed from [21] and [22]. The displacement field and the strains are defined by assuming that fibers orthogonal to the middle plane remain orthogonal (see figure 1). This leads to the following relations for the displacement field

$$u(x, y, z) = -z \frac{\partial w}{\partial x}, \quad v(x, y, z) = -z \frac{\partial w}{\partial y}, \quad w(x, y, z) = w(x, y). \quad (11)$$

and for the strains

$$\boldsymbol{\epsilon} = \begin{pmatrix} \epsilon_{xx} \\ \epsilon_{yy} \\ \gamma_{xy} \end{pmatrix} = \begin{pmatrix} \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{pmatrix} \begin{pmatrix} -z \frac{\partial w}{\partial x} \\ -z \frac{\partial w}{\partial y} \end{pmatrix} = -z \begin{pmatrix} \frac{\partial^2 w}{\partial x^2} \\ \frac{\partial^2 w}{\partial y^2} \\ 2 \frac{\partial^2 w}{\partial x \partial y} \end{pmatrix}. \quad (12)$$

The curvature vector is defined as

$$\boldsymbol{\kappa} = \begin{pmatrix} \kappa_{xx} \\ \kappa_{yy} \\ \kappa_{xy} \end{pmatrix} = \begin{pmatrix} \frac{\partial^2 w}{\partial x^2} \\ \frac{\partial^2 w}{\partial y^2} \\ 2 \frac{\partial^2 w}{\partial x \partial y} \end{pmatrix}. \quad (13)$$

The Hooke constitutive law for isotropic material is considered for the constitutive relation

$$\boldsymbol{\sigma} = \mathbf{E} \boldsymbol{\epsilon}, \quad \mathbf{E} := \frac{E}{1 - \nu^2} \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{1-\nu}{2} \end{bmatrix}. \quad (14)$$



Figure 1: Kinematic assumption for the Kirchhoff plate

where ν the Poisson ratio and E the Young modulus. These physical parameters may be inhomogeneous, i.e. $\nu = \nu(x, y, z)$, $E = E(x, y, z)$. The generalized momenta are found by integrating the stresses along the fiber

$$\mathbf{M} = \begin{pmatrix} M_{xx} \\ M_{yy} \\ M_{xy} \end{pmatrix} = \int_{-\frac{h}{2}}^{\frac{h}{2}} \mathbf{E} z^2 \, dz \, \boldsymbol{\kappa}$$

where h is the plate thickness and the relation between momenta and curvatures is expressed by the bending rigidity matrix \mathbf{D}

$$\mathbf{M} = \mathbf{D} \boldsymbol{\kappa} \quad \mathbf{D} := \int_{-\frac{h}{2}}^{\frac{h}{2}} \mathbf{E} z^2 \, dz \quad (15)$$

Now the classical Kirchhoff-Love model for thin can be recalled [23]

$$\mu \frac{\partial^2 w}{\partial t^2} + \frac{\partial^2 M_{xx}}{\partial x^2} + 2 \frac{\partial^2 M_{xy}}{\partial x \partial y} + \frac{\partial^2 M_{yy}}{\partial y^2} = 0, \quad (16)$$

where $\mu = \rho h$ is the surface density and ρ the mass density. If the E, ν coefficients are constant then the ruling EDP becomes

$$\mu \frac{\partial^2 w}{\partial t^2} + D \Delta^2 w = p, \quad (17)$$

where $\Delta^2 = \frac{\partial^4}{\partial x^4} + 2 \frac{\partial^4}{\partial x^2 \partial y^2} + \frac{\partial^4}{\partial y^4}$ is the bilaplacian and $D = \frac{Eh^3}{12(1-\nu^2)}$ is the bending rigidity coefficient. The kinetic and potential energy density per unit area \mathcal{K} and \mathcal{U} , are respectively given by

$$\mathcal{K} = \frac{1}{2} \mu \left(\frac{\partial w}{\partial t} \right)^2, \quad \mathcal{U} = \frac{1}{2} \mathbf{M} \cdot \boldsymbol{\kappa}.$$

The total energy density is split into kinetic and potential energy

$$\mathcal{H} = \mathcal{K} + \mathcal{U}, \quad (18)$$

and the corresponding total energies given by the following relations

$$H = \int_{\Omega} \mathcal{H} \, d\Omega, \quad K = \int_{\Omega} \mathcal{K} \, d\Omega, \quad U = \int_{\Omega} \mathcal{U} \, d\Omega. \quad (19)$$

3. PH formulation of the Kirchhoff plate

In this section the port-Hamiltonian formulation of the Kirchhoff plate is presented first in vectorial form in 3.1 and then in tensorial form in 3.2.

3.1. PH vectorial formulation of the Kirchhoff plate

To obtain a port-Hamiltonian system (PHs) the energy variables as well as the underlying Stokes-Dirac structure, associated with the skew-adjoint operator J , have to be properly defined. Consider the Hamiltonian energy

$$\begin{aligned} H &= \int_{\Omega} \frac{1}{2} \left\{ \mu \left(\frac{\partial w}{\partial t} \right)^2 + \mathbf{M} \cdot \boldsymbol{\kappa} \right\} d\Omega \\ &= \int_{\Omega} \frac{1}{2} \left\{ \mu \left(\frac{\partial w}{\partial t} \right)^2 + \boldsymbol{\kappa}^T \mathbf{D} \boldsymbol{\kappa} \right\} d\Omega \end{aligned} \quad (20)$$

The energy variables are then selected to be the linear momentum $\mu \frac{\partial w}{\partial t}$ and the curvatures $\boldsymbol{\kappa}$, in an analogous fashion with respect to the one-dimensional counterpart of this model, the Euler-Bernoulli beam. The energy variables are collected in vector

$$\boldsymbol{\alpha} := (\mu w_t, \kappa_{xx}, \kappa_{yy}, \kappa_{xy})^T, \quad (21)$$

where $w_t = \frac{\partial w}{\partial t}$. The Hamiltonian density is given by the following expression

$$\mathcal{H} = \frac{1}{2} \boldsymbol{\alpha}^T \begin{bmatrix} \frac{1}{\mu} & 0 \\ 0 & \mathbf{D} \end{bmatrix} \boldsymbol{\alpha}, \quad H = \int_{\Omega} \mathcal{H} d\Omega. \quad (22)$$

So its variational derivative provides as co-energy variables

$$\mathbf{e} := \frac{\delta H}{\delta \boldsymbol{\alpha}} = (w_t, M_{xx}, M_{yy}, M_{xy})^T, \quad (23)$$

The port-Hamiltonian system and the formally skew-adjoint operator relating energy and co-energy variables are found to be

$$\frac{\partial \boldsymbol{\alpha}}{\partial t} = J \mathbf{e} \quad \text{and} \quad J := \begin{bmatrix} 0 & -\frac{\partial^2}{\partial x^2} & -\frac{\partial^2}{\partial y^2} & -\left(\frac{\partial^2}{\partial x \partial y} + \frac{\partial^2}{\partial y \partial x} \right) \\ \frac{\partial^2}{\partial x^2} & 0 & 0 & 0 \\ \frac{\partial^2}{\partial y^2} & 0 & 0 & 0 \\ \frac{\partial^2}{\partial x \partial y} + \frac{\partial^2}{\partial y \partial x} & 0 & 0 & 0 \end{bmatrix}. \quad (24)$$

The first line of the skew-symmetric operator in (24) is found by considering the equation 16. The remaining lines express Clairaut's theorem for the vertical displacement. This theorem states that for smooth function higher order partial derivative commute. To see this consider the second, third and forth line of system (24)

$$\frac{\partial \kappa_{xx}}{\partial t} = \frac{\partial^2 w_t}{\partial x^2}, \quad \frac{\partial \kappa_{yy}}{\partial t} = \frac{\partial^2 w_t}{\partial y^2}, \quad \frac{\partial \kappa_{xy}}{\partial t} = 2 \frac{\partial^2 w_t}{\partial x \partial y}$$

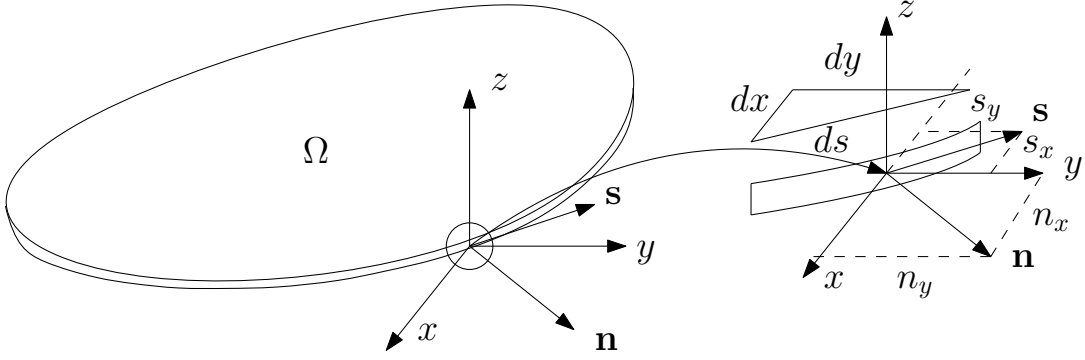


Figure 2: Reference frames and notations

Then, since $w_t = \frac{\partial w}{\partial t}$, $\kappa_{xx} = \frac{\partial^2 w}{\partial x^2}$, $\kappa_{yy} = \frac{\partial^2 w}{\partial y^2}$, $\kappa_{xy} = 2 \frac{\partial^2 w}{\partial x \partial y}$ it holds

$$\frac{\partial^3 w}{\partial t \partial x^2} = \frac{\partial^3 w}{\partial x^2 \partial t}, \quad \frac{\partial^3 w}{\partial t \partial y^2} = \frac{\partial^3 w}{\partial y^2 \partial t}, \quad \frac{\partial^3 w}{\partial t \partial x \partial y} = \frac{\partial^3 w}{\partial x \partial y \partial t},$$

which is exactly Clairaut's theorem for the vertical displacement.

Remark 2. From the Schwarz theorem for C^2 functions the mixed derivative could be expressed as $2 \frac{\partial^2}{\partial x \partial y}$, instead of $\frac{\partial^2}{\partial y \partial x} + \frac{\partial^2}{\partial x \partial y}$. However, in this way the symmetry intrinsically present in $\gamma_{xy} = -z \left(\frac{\partial^2 w}{\partial y \partial x} + \frac{\partial^2 w}{\partial x \partial y} \right)$ would be lost. The mixed derivative is here split to reestablish the symmetric nature of curvatures and momenta (that are of tensorial nature as explained in section 3.2).

The boundary variables are obtained by evaluating the time derivative of the Hamiltonian

$$\begin{aligned} \dot{H} &= \int_{\Omega} \frac{\delta H}{\delta \alpha} \cdot \frac{\partial \alpha}{\partial t} \, d\Omega \\ &= \int_{\Omega} \left\{ w_t \left(-\frac{\partial^2 M_{xx}}{\partial x^2} - \frac{\partial^2 M_{xx}}{\partial y^2} - 2 \frac{\partial^2 M_{xy}}{\partial x \partial y} \right) + M_{xx} \frac{\partial^2 w_t}{\partial x^2} + M_{yy} \frac{\partial^2 w_t}{\partial y^2} + 2 M_{xy} \frac{\partial^2 w_t}{\partial x \partial y} \right\} \, d\Omega \\ &= \int_{\Omega} \left\{ e_1 \left(-\frac{\partial^2 e_2}{\partial x^2} - \frac{\partial^2 e_3}{\partial y^2} - 2 \frac{\partial^2 e_4}{\partial x \partial y} \right) + e_2 \frac{\partial^2 e_1}{\partial x^2} + e_3 \frac{\partial^2 e_1}{\partial y^2} + 2 e_4 \frac{\partial^2 e_1}{\partial x \partial y} \right\} \, d\Omega \end{aligned}$$

In Figure 2 the notations for the different reference frames are introduced. By applying Green theorem, considering the split mixed derivative ($2 \frac{\partial^2}{\partial x \partial y} = \frac{\partial^2}{\partial x \partial y} + \frac{\partial^2}{\partial y \partial x}$)

$$\begin{aligned} \dot{H} &= \int_{\partial \Omega} \left\{ n_x \left(e_2 \frac{\partial e_1}{\partial x} + e_4 \frac{\partial e_1}{\partial y} - e_1 \frac{\partial e_2}{\partial x} - e_1 \frac{\partial e_4}{\partial y} \right) \right. \\ &\quad \left. + n_y \left(e_3 \frac{\partial e_1}{\partial y} + e_4 \frac{\partial e_1}{\partial x} - e_1 \frac{\partial e_3}{\partial y} - e_1 \frac{\partial e_4}{\partial x} \right) \right\} \, ds. \end{aligned} \quad (25)$$

where n_x, n_y are the components along the x - and the y -axis of the normal to the boundary. The variable of integration s is now the curvilinear abscissa which runs along the boundary.

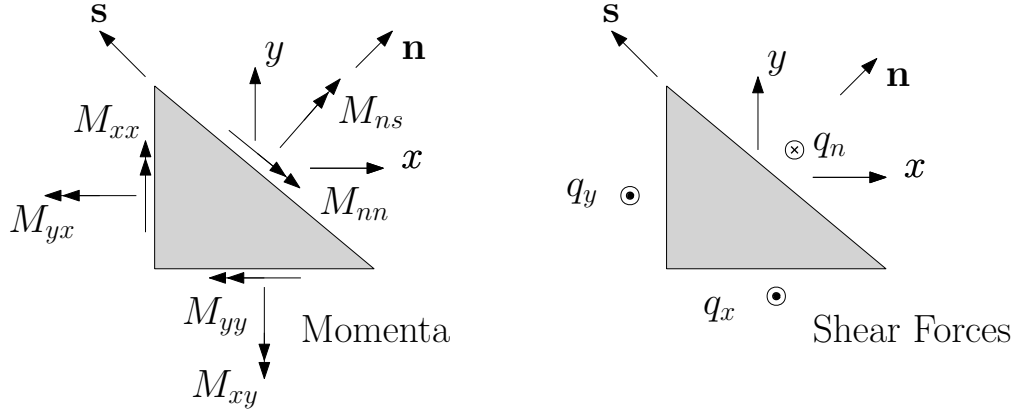


Figure 3: Cauchy law for momenta and forces at the boundary

If the physical variables are introduced

$$\dot{H} = \int_{\partial\Omega} \left\{ n_x \left(M_{xx} \frac{\partial w_t}{\partial x} + M_{xy} \frac{\partial w_t}{\partial y} - w_t \frac{\partial M_{xx}}{\partial x} - w_t \frac{\partial M_{xy}}{\partial y} \right) + n_y \left(M_{yy} \frac{\partial w_t}{\partial y} + M_{xy} \frac{\partial w_t}{\partial x} - w_t \frac{\partial M_{yy}}{\partial y} - w_t \frac{\partial M_{xy}}{\partial x} \right) \right\} ds. \quad (26)$$

Now the following quantities, represented in figure 3, are defined

$$\begin{aligned} \text{Shear Force} & \quad q_n := n_x q_x + n_y q_y, \\ \text{Flexural momentum} & \quad M_{nn} := \mathbf{n}^T \begin{pmatrix} M_{xx} n_x + M_{xy} n_y \\ M_{xy} n_x + M_{yy} n_y \end{pmatrix}, \\ \text{Torsional momentum} & \quad M_{ns} := \mathbf{s}^T \begin{pmatrix} M_{xx} n_x + M_{xy} n_y \\ M_{xy} n_x + M_{yy} n_y \end{pmatrix}, \end{aligned} \quad \begin{aligned} \mathbf{n} &= \begin{pmatrix} n_x \\ n_y \end{pmatrix}, \\ \mathbf{s} &= \begin{pmatrix} -n_y \\ n_x \end{pmatrix}, \end{aligned} \quad (27)$$

where $q_x = -\frac{\partial M_{xx}}{\partial x} - \frac{\partial M_{xy}}{\partial y}$ and $q_y = -\frac{\partial M_{yy}}{\partial y} - \frac{\partial M_{xy}}{\partial x}$. The gradient of the vertical velocity can be projected upon the normal and tangential directions to the boundary

$$\nabla w_t = (\nabla w_t^T \mathbf{n}) \mathbf{n} + (\nabla w_t^T \mathbf{s}) \mathbf{s} = \frac{\partial w_t}{\partial n} \mathbf{n} + \frac{\partial w_t}{\partial s} \mathbf{s}. \quad (28)$$

So the time derivative of the Hamiltonian can be finally written as

$$\dot{H} = \int_{\partial\Omega} \left\{ w_t q_n + \frac{\partial w_t}{\partial s} M_{ns} + \frac{\partial w_t}{\partial n} M_{nn} \right\} ds. \quad (29)$$

The variables w_t and $\frac{\partial w_t}{\partial s}$ are not independent as they are differentially related with respect to derivation along s , the curvilinear abscissa of the boundary domain (see for instance [23]). Another integration by part is needed to highlight appropriate independent power conjugated variables. Let us suppose that the boundary is a closed and regular curve. Then the integration by parts along a closed boundary leads to

$$\int_{\partial\Omega} \frac{\partial w_t}{\partial s} M_{ns} ds = - \int_{\partial\Omega} \frac{\partial M_{ns}}{\partial s} w_t ds. \quad (30)$$

The energy balance can be finally written as

$$\dot{H} = \int_{\partial\Omega} \left\{ w_t \tilde{q}_n + \frac{\partial w_t}{\partial n} M_{nn} \right\} ds \quad (31)$$

where $\tilde{q}_n := q_n - \frac{\partial M_{ns}}{\partial s}$ is the effective shear force. Equation (31) is of utmost importance, since it contains the boundary variables that will be present in the Stokes-Dirac structure defining the port-Hamiltonian system.

3.1.1. Underlying Stokes-Dirac structure

Let \mathcal{F} denote the flow space and let \mathcal{E} denote the effort space. For simplicity we take $\mathcal{F} \equiv \mathcal{E} = \mathcal{C}^2(\Omega, \mathbb{R}^4)$, the space of smooth vector-valued functions in \mathbb{R}^4 . Equation (31) allows identifying the boundary terms of the underlying Stokes-Dirac structures. The space of boundary variables is a vector of four components given by

$$\mathcal{Z} = \{z \mid z = B_{\partial}(e), \forall e \in \mathcal{E}\}, \quad z = \left(\tilde{q}_n, w_t, M_{nn}, \frac{\partial w_t}{\partial n} \right)^T.$$

In the case where the differential J operator of order one, the B_{∂} operator is a linear operator over the trace of the effort variables.. Here, since the differential J is of order two, B_{∂} contains the normal and tangential derivatives at the boundary and so more regularity is required for the boundary variables.

Remark 3. *This fact was already stated for 1-D systems in [17]; here it is the extension to the 2-D system with second order differential operator J .*

This operator reads

$$\begin{aligned} B_{\partial}(e) = & \begin{bmatrix} 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & n_x^2 & n_y^2 & 2n_x n_y \\ 0 & 0 & 0 & 0 \end{bmatrix} e - \begin{bmatrix} 0 & n_x & 0 & n_y \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \frac{\partial e}{\partial x} - \begin{bmatrix} 0 & 0 & n_y & n_x \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \frac{\partial e}{\partial y} \\ & + \frac{\partial}{\partial n} \left(\begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} e \right) - \frac{\partial}{\partial s} \left(\begin{bmatrix} 0 & -n_x n_y & n_x n_y & n_x^2 - n_y^2 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} e \right). \quad (32) \end{aligned}$$

Theorem 2 (Stokes-Dirac structure for the Kirchhoff Plate). *The set*

$$\mathcal{D} := \left\{ (f, e, z) \in \mathcal{F} \times \mathcal{E} \times \mathcal{Z} \mid f = -\frac{\partial \alpha}{\partial t} = -J e, z = B_{\partial}(e) \right\} \quad (33)$$

is a Stokes-Dirac structure with respect to the pairing

$$\ll (f_1, e_1, z_1), (f_2, e_2, z_2) \gg = \int_{\Omega} [e_1^T f_2 + e_2^T f_1] d\Omega + \int_{\partial\Omega} B_J(z_1, z_2) ds, \quad (34)$$

where B_J is a symmetric operator, arising from a double application of the Green theorem. It reads

$$\begin{aligned} B_J(\mathbf{z}_1, \mathbf{z}_2) &= \tilde{q}_{n,2} w_{t,1} + M_{nn,2} \frac{\partial w_{t,1}}{\partial n} \\ &\quad + \tilde{q}_{n,1} w_{t,2} + M_{nn,1} \frac{\partial w_{t,2}}{\partial n}, \quad B_J = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}. \quad (35) \\ &= \mathbf{z}_1^T B_J \mathbf{z}_2 \end{aligned}$$

Proof. A regular boundary will be assumed for the proof.

Step I. The first implication is $\mathbb{D} \subseteq \mathbb{D}^T$. This is true if $\forall \boldsymbol{\omega}_\alpha = (\mathbf{f}_\alpha, \mathbf{e}_\alpha, \mathbf{f}_\alpha)$ and $\boldsymbol{\omega}_\beta = (\mathbf{f}_\beta, \mathbf{e}_\beta, \mathbf{z}_\beta) \in \mathbb{D}$ then $\ll \boldsymbol{\omega}_\alpha, \boldsymbol{\omega}_\beta \gg = 0$. The integral over the domain reads

$$\begin{aligned} \int_{\Omega} [\mathbf{e}_\alpha^T \mathbf{f}_\beta + \mathbf{e}_\beta^T \mathbf{f}_\alpha] d\Omega &= \int_{\Omega} \left\{ e_1^\alpha \left(\frac{\partial^2 e_2^\beta}{\partial x^2} + \frac{\partial^2 e_3^\beta}{\partial y^2} + \frac{\partial^2 e_4^\beta}{\partial x \partial y} \right) - e_2^\alpha \frac{\partial^2 e_1^\beta}{\partial x^2} - e_3^\alpha \frac{\partial^2 e_1^\beta}{\partial y^2} \right. \\ &\quad \left. - 2e_4^\alpha \frac{\partial^2 e_1^\beta}{\partial x \partial y} + e_1^\beta \left(\frac{\partial^2 e_2^\alpha}{\partial x^2} + \frac{\partial^2 e_3^\alpha}{\partial y^2} + \frac{\partial^2 e_4^\alpha}{\partial x \partial y} \right) - e_2^\beta \frac{\partial^2 e_1^\alpha}{\partial x^2} - e_3^\beta \frac{\partial^2 e_1^\alpha}{\partial y^2} - 2e_4^\beta \frac{\partial^2 e_1^\alpha}{\partial x \partial y} \right\} d\Omega. \end{aligned}$$

Once the Green theorem has been applied, the relevant quantities expressed by equations (27) pop up as follows

$$\begin{aligned} \int_{\Omega} [\mathbf{e}_\alpha^T \mathbf{f}_\beta + \mathbf{e}_\beta^T \mathbf{f}_\alpha] d\Omega &= \int_{\partial\Omega} \left\{ e_1^\alpha \left(\frac{\partial e_2^\beta}{\partial x} n_x + \frac{\partial e_3^\beta}{\partial y} n_y + \frac{\partial e_4^\beta}{\partial y} n_x + \frac{\partial e_4^\beta}{\partial x} n_y \right) \right. \\ &\quad \left. + e_1^\beta \left(\frac{\partial e_2^\alpha}{\partial x} n_x + \frac{\partial e_3^\alpha}{\partial y} n_y + \frac{\partial e_4^\alpha}{\partial y} n_x + \frac{\partial e_4^\alpha}{\partial x} n_y \right) - e_2^\alpha \frac{\partial e_1^\beta}{\partial x} n_x - e_2^\beta \frac{\partial e_1^\alpha}{\partial x} n_x \right. \\ &\quad \left. - e_3^\alpha \frac{\partial e_1^\beta}{\partial y} n_y - e_3^\beta \frac{\partial e_1^\alpha}{\partial y} n_y - e_4^\alpha \left(\frac{\partial e_1^\beta}{\partial y} n_x + \frac{\partial e_1^\beta}{\partial x} n_y \right) - e_4^\beta \left(\frac{\partial e_1^\alpha}{\partial y} n_x + \frac{\partial e_1^\alpha}{\partial x} n_y \right) \right\} ds. \quad (36) \end{aligned}$$

Moreover only the kinematically independent quantities have to be considered, leading to the final result

$$- \int_{\partial\Omega} \left(w_t^\alpha \tilde{q}_n^\beta + w_t^\beta \tilde{q}_n^\alpha + \frac{\partial w_t^\alpha}{\partial n} M_{nn}^\beta + \frac{\partial w_t^\beta}{\partial n} M_{nn}^\alpha \right) ds = - \int_{\partial\Omega} B_J(\mathbf{z}_\alpha, \mathbf{z}_\beta) ds. \quad (37)$$

This concludes the first part of the proof.

Step II. For the second implication, i.e. $\mathcal{D}^\perp \subseteq \mathcal{D}$. Let us take $\boldsymbol{\omega}_\alpha \in \mathcal{D}^\perp, \forall \boldsymbol{\omega}_\beta \in \mathcal{D}$. Then the bilinear form, once the Green theorem has been applied, provides the following

$$\begin{aligned} \int_{\Omega} \left\{ e_1^\beta \left(f_1^\alpha - \frac{\partial^2 e_2^\alpha}{\partial x^2} - \frac{\partial^2 e_3^\alpha}{\partial y^2} - 2 \frac{\partial^2 e_4^\alpha}{\partial x \partial y} \right) + e_2^\beta \left(f_2^\alpha + \frac{\partial^2 e_1^\alpha}{\partial x^2} \right) + e_3^\beta \left(f_3^\alpha + \frac{\partial^2 e_1^\alpha}{\partial y^2} \right) + \right. \\ \left. e_4^\beta \left(f_4^\alpha + 2 \frac{\partial^2 e_1^\alpha}{\partial x \partial y} \right) \right\} d\Omega + \int_{\partial\Omega} \left\{ w_t^\beta \left(\frac{\partial e_2^\alpha}{\partial x} n_x + \frac{\partial e_3^\alpha}{\partial y} n_y + \frac{\partial e_4^\alpha}{\partial y} n_x + \frac{\partial e_4^\alpha}{\partial x} n_y \right) \right. \\ \left. - q_n^\beta e_1^\alpha - \frac{\partial w_t^\beta}{\partial x} e_2^\alpha n_x - \frac{\partial w_t^\beta}{\partial y} e_3^\alpha n_y - e_4^\alpha \left(\frac{\partial w_t^\beta}{\partial y} n_x + \frac{\partial w_t^\beta}{\partial x} n_y \right) - \frac{\partial e_1^\alpha}{\partial x} M_{xx}^\beta n_x - \frac{\partial e_1^\alpha}{\partial y} M_{yy}^\beta n_y \right. \\ \left. - M_{xy}^\beta \left(\frac{\partial e_1^\alpha}{\partial y} n_x + \frac{\partial e_1^\alpha}{\partial x} n_y \right) + \tilde{q}_n^\beta z_1^\alpha + w_t^\beta z_2^\alpha + M_{nn}^\beta z_3^\alpha + \frac{\partial w_t^\beta}{\partial n} z_4^\beta \right\} ds = 0. \quad (38) \end{aligned}$$

Since the relation has to be valid for each $\boldsymbol{\omega}_\beta \in \mathcal{D}$ the flux and effort variables are in \mathcal{D} . For the boundary terms the same procedure as before has to be applied by considering the definition of the momenta over the boundary (see equation (27)). Then it can be stated that $\boldsymbol{\omega}_\alpha \in \mathcal{D}$. \square

3.1.2. Including dissipation and external forces in the model

Distributed forces or control and dissipative relations can be easily included in an augmented Stokes-Dirac structure by simply defining the appropriate conjugated variables.

If distributed forces have to be considered, then the set

$$\mathcal{D}_d := \{(\mathbf{f}, \mathbf{f}_d, \mathbf{e}, \mathbf{e}_d, \mathbf{z}) \in \mathcal{F} \times \mathcal{F}_d \times \mathcal{E} \times \mathcal{E}_d \times \mathcal{Z} \mid \mathbf{f} = -\frac{\partial \boldsymbol{\alpha}}{\partial t} = -J\mathbf{e} - G_d \mathbf{f}_d, \mathbf{e}_d = G_d^* \mathbf{e}, \mathbf{e} = B_\partial(\mathbf{e})\} \quad (39)$$

is a Stokes-Dirac structure with respect to the pairing

$$\begin{aligned} \ll (\mathbf{f}_1, \mathbf{f}_{d,1} \mathbf{e}_1, \mathbf{e}_{d,1}, \mathbf{z}_1), (\mathbf{f}_2, \mathbf{f}_{d,2} \mathbf{e}_2, \mathbf{e}_{d,2}, \mathbf{z}_2) \gg = \\ \int_{\Omega} [\mathbf{e}_1^T \mathbf{f}_2 + \mathbf{e}_2^T \mathbf{f}_1 + \mathbf{e}_{d,1}^T \mathbf{f}_{d,2} + \mathbf{e}_{d,2}^T \mathbf{f}_{d,1}] d\Omega + \int_{\partial\Omega} B_J(\mathbf{z}_1, \mathbf{z}_2) ds. \quad (40) \end{aligned}$$

If gravity has to be included, then $G_d = [1, 0, 0, 0]^T, \mathbf{f}_d = -\mu g$.

Analogously dissipation can be included in an augmented Dirac structure. As an example, the ruling PDE, once a dissipative term of fluid damping type is considered, reads

$$\mu \frac{\partial^2 w}{\partial t^2} + r \frac{\partial w}{\partial t} + \frac{\partial^2 M_{xx}}{\partial x^2} + 2 \frac{\partial^2 M_{xy}}{\partial x \partial y} + \frac{\partial^2 M_{yy}}{\partial y^2} = 0, \quad (41)$$

where $r > 0$ is the damping coefficient. If this equation is rewritten using the port-Hamiltonian formalism then we get

$$\frac{\partial \boldsymbol{\alpha}}{\partial t} = (J - R) \mathbf{e}, \quad R := \begin{bmatrix} r & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}. \quad (42)$$

The R matrix, which is a symmetric, semi-positive definite operator, can be decomposed as

$$R = G_R S G_R^*, \quad (43)$$

where $S = r$ is a coercive operator (in this case simply a positive scalar), $G_R = (1 \ 0 \ 0 \ 0)^T$ and G_R^* denotes the adjoint operator to G_R . The augmented Stokes-Dirac structure reads

$$\mathcal{D}_r := \{(\mathbf{f}, \mathbf{f}_r) \in \mathcal{F}, (\mathbf{e}, \mathbf{e}_r) \in \mathcal{E}, \mathbf{z} \in \mathcal{Z} \mid \begin{pmatrix} \mathbf{f} \\ \mathbf{f}_r \end{pmatrix} = \begin{bmatrix} J & G_R \\ -G_R^* & 0 \end{bmatrix} \begin{pmatrix} \mathbf{e} \\ \mathbf{e}_r \end{pmatrix}, \mathbf{f} = \frac{\partial \alpha}{\partial t}, \mathbf{e}_r = S \mathbf{f}_r, \mathbf{z} = B_\partial(\mathbf{e})\} \quad (44)$$

is a Stokes-Dirac structure with respect to the paring

$$\ll (\mathbf{f}_1, \mathbf{f}_{r,1}, \mathbf{e}_1, \mathbf{e}_{r,1}, \mathbf{z}_1), (\mathbf{f}_2, \mathbf{f}_{r,2}, \mathbf{e}_2, \mathbf{e}_{r,2}, \mathbf{z}_2) \gg = \int_{\Omega} [\mathbf{e}_1^T \mathbf{f}_2 + \mathbf{e}_2^T \mathbf{f}_1] d\Omega - \int_{\Omega} [\mathbf{e}_{r,1}^T \mathbf{f}_{r,2} + \mathbf{e}_{r,2}^T \mathbf{f}_{r,1}] d\Omega - \int_{\partial\Omega} B_J(\mathbf{z}_1, \mathbf{z}_2) ds, \quad (45)$$

where the minus sign for the boundary term arises from the different definition of the flow variables ($\mathbf{f} = \frac{\partial \alpha}{\partial t}$).

Remark 4. *More involved dissipation model can be found in [24] are reference therein. More specifically, for Kirchhoff plate, some specific damping models can be found in [25].*

3.2. PH tensorial formulation of the Kirchhoff plate

In section 3.1 the Stokes-Dirac structure of the Kirchhoff plate was found by using a vectorial notation for the curvatures and momenta. In fact these variables are of tensorial nature and in the following the tensorial formulation takes the place of the vectorial one. First let us rewrite the momenta and curvatures as symmetric matrices (corresponding to the choice of a Cartesian frame for the representation of tensors)

$$\mathbb{K} = \begin{bmatrix} \kappa_{xx} & \kappa_{xy} \\ \kappa_{xy} & \kappa_{yy} \end{bmatrix}, \quad \mathbb{M} = \begin{bmatrix} M_{xx} & M_{xy} \\ M_{xy} & M_{yy} \end{bmatrix}, \quad (46)$$

where now, with a slight abuse of notation, κ_{xy} differs by 1/2 from the definition given in equation (13), i.e. $\kappa_{xy} = \frac{\partial^2 w}{\partial x \partial y}$. All the other quantities stay the same with respect to what stated in section 2.1. The Hamiltonian energy is written as

$$H = \int_{\Omega} \left\{ \frac{1}{2} \mu \left(\frac{\partial w}{\partial t} \right)^2 + \frac{1}{2} \mathbb{M} : \mathbb{K} \right\} d\Omega, \quad (47)$$

where the tensor contraction in Cartesian coordinates is expressed as

$$\mathbb{M} : \mathbb{K} = \sum_{i,j=1}^2 M_{ij} \kappa_{ij} = \text{Tr}(\mathbb{M}^T \mathbb{K})$$

For what concerns the choice of the energy variables a scalar and a tensor variable are grouped together

$$\alpha_w = \mu \frac{\partial w}{\partial t} \quad \mathbb{A}_\kappa = \mathbb{K} \quad (48)$$

The co-energy variables are found by computing the variational derivative of the Hamiltonian

$$e_w := \frac{\delta H}{\delta \alpha_w} = \frac{\partial w}{\partial t} := w_t, \quad \mathbb{E}_\kappa := \frac{\delta H}{\delta \mathbb{A}_\kappa} = \mathbb{M}. \quad (49)$$

Remark 5. For the variational derivative with respect to a tensor, see Propostion 1 in [26].

The port-Hamiltonian system (24) is now rewritten as

$$\begin{cases} \frac{\partial \alpha_w}{\partial t} = -\operatorname{div}(\operatorname{Div}(\mathbb{E}_\kappa)) \\ \frac{\partial \mathbb{A}_\kappa}{\partial t} = \operatorname{Grad}(\operatorname{grad}(e_w)) \end{cases}, \quad (50)$$

where div and Div denote the divergence of a vector and of a tensor respectively. The operator Grad denotes the symmetric gradient

$$\operatorname{Grad}(\mathbf{a}) = \frac{1}{2} \left(\nabla \otimes \mathbf{a} + (\nabla \otimes \mathbf{a})^T \right). \quad (51)$$

The operator $\operatorname{Grad} \circ \operatorname{grad}$ corresponds to the Hessian operator. In Cartesian coordinates it reads

$$\operatorname{Grad} \circ \operatorname{grad} = \begin{bmatrix} \frac{\partial^2}{\partial x^2} & \frac{\partial^2}{\partial x \partial y} \\ \frac{\partial^2}{\partial y \partial x} & \frac{\partial^2}{\partial y^2} \end{bmatrix}. \quad (52)$$

Theorem 3. The operator $\operatorname{Grad} \circ \operatorname{grad}$, corresponding to the Hessian operator, is the adjoint of the double divergence $\operatorname{div} \circ \operatorname{Div}$.

Proof. Let us consider the Hilbert space of the square integrable symmetric square tensors of size $n \times n$ over an open connected set Ω . This space will be denoted by $\mathcal{H}_1 = L^2(\Omega, \mathbb{R}_{\operatorname{sym}}^{n \times n})$. This space is endowed with the integral of the tensor contraction as scalar product

$$\langle \mathbb{E}, \mathbb{F} \rangle_{\mathcal{H}_1} = \int_{\Omega} \mathbb{E} : \mathbb{F} \, d\Omega = \int_{\Omega} \operatorname{Tr}(\mathbb{E}^T \mathbb{F}) \, d\Omega, \quad \forall \mathbb{E}, \mathbb{F} \in [L^2_{\operatorname{sym}}(\Omega)]^{n \times n}.$$

Consider the Hilbert space $\mathcal{H}_2 = L^2(\Omega)$ of scalar square integrable functions, endowed with the inner product

$$\langle e, f \rangle_{\mathcal{H}_2} = \int_{\Omega} e f \, d\Omega. \quad (53)$$

Let us consider the double divergence operator defined as

$$\begin{aligned} A : \mathcal{H}_1 &\rightarrow \mathcal{H}_2, \\ \mathbb{E} &\rightarrow \operatorname{div}(\operatorname{Div}(\mathbb{E})) = e, \end{aligned} \quad \text{with } \mathbf{e} = \operatorname{div}(\operatorname{Div}(\mathbb{E})) = \sum_{i=1}^n \sum_{j=1}^n \frac{\partial^2 \mathbb{E}_{ij}}{\partial x_i \partial x_j}.$$

We shall identify A^*

$$\begin{aligned} A^* : \mathcal{H}_2 &\rightarrow \mathcal{H}_1, \\ f &\rightarrow A^* f = \mathbb{F}, \end{aligned}$$

such that

$$\langle A\mathbb{E}, f \rangle_{\mathcal{H}_2} = \langle \mathbb{E}, A^* f \rangle_{\mathcal{H}_1}, \quad \begin{aligned} \forall \mathbb{E} &\in \operatorname{Domain}(A) \subset \mathcal{H}_1 \\ \forall f &\in \operatorname{Domain}(A^*) \subset \mathcal{H}_2 \end{aligned}$$

The function have to belong to the operator domain, so for instance $f \in \mathcal{C}_0^2(\Omega) \in \operatorname{Domain}(A^*)$ the space of twice differentiable scalar functions with compact support on an open simply connected set Ω and additionally \mathbb{E} can be chosen in the set $\mathcal{C}_0^2(\Omega, \mathbb{R}_{\operatorname{sym}}^{2 \times 2}) \in$

$\text{Domain}(A)$, the space of twice differentiable 2×2 symmetric tensors with compact support on Ω . A classical result is the fact that the adjoint of the vector divergence is $\text{div}^* = -\text{grad}$ as stated in [27]. This may be generalized to the adjoint of the tensor divergence $\text{Div}^* = -\text{Grad}$ (see theorem 4 of [26]). Considering that A is the composition of two different operators $A = \text{div} \circ \text{Div}$ and that the adjoint of a composed operator is the adjoint of each operator in reverse order, i.e. $(B \circ C)^* = C^* \circ B^*$, then it can be stated

$$A^* = (\text{div} \circ \text{Div})^* = \text{Div}^* \circ \text{div}^* = \text{Grad} \circ \text{grad}.$$

Since only formal adjoints are being looked for, this concludes the proof. \square

If the variables in system (50) are grouped together the formally skew-symmetric operator J can be highlighted

$$\frac{\partial}{\partial t} \begin{pmatrix} \alpha_w \\ \mathbb{A}_\kappa \end{pmatrix} = \underbrace{\begin{bmatrix} 0 & -\text{div} \circ \text{Div} \\ \text{Grad} \circ \text{grad} & 0 \end{bmatrix}}_J \begin{pmatrix} e_w \\ \mathbb{E}_\kappa \end{pmatrix}. \quad (54)$$

where all zeros are intended as nullifying operator from the space of input variables to the space of output variables.

Remark 6. *The interconnection structure J now resembles that of the Bernoulli beam. The double divergence and the double gradient coincide, in dimension one, with the second derivative.*

Again the boundary **port variables** can be found by evaluating the time derivative of the Hamiltonian

$$\begin{aligned} \dot{H} &= \int_{\Omega} \left\{ \frac{\partial \alpha_w}{\partial t} e_w + \frac{\partial \mathbb{A}_\kappa}{\partial t} : \mathbb{E}_\kappa \right\} d\Omega, \\ &= \int_{\Omega} \{ -\text{div}(\text{Div}(\mathbb{E}_\kappa)) e_w + \text{Grad}(\text{grad}(e_w)) : \mathbb{E}_\kappa \} d\Omega, && \text{Integration by parts} \\ &= \int_{\partial\Omega} \left\{ \underbrace{-\mathbf{n} \cdot \text{Div}(\mathbb{E}_\kappa)}_{q_n} e_w + [\mathbf{n} \otimes \text{grad}(e_w)] : \mathbb{E}_\kappa \right\} ds, && \text{See (27) and (28)} \\ &= \int_{\partial\Omega} \left\{ q_n e_w + \frac{\partial e_w}{\partial n} \underbrace{(\mathbf{n} \otimes \mathbf{n}) : \mathbb{E}_\kappa}_{M_{nn}} + \frac{\partial e_w}{\partial s} \underbrace{(\mathbf{n} \otimes \mathbf{s}) : \mathbb{E}_\kappa}_{M_{ns}} \right\} ds, && \text{Dyadic properties} \\ &= \int_{\partial\Omega} \left\{ q_n w_t + \frac{\partial w_t}{\partial n} M_{nn} + \frac{\partial w_t}{\partial s} M_{ns} \right\} ds. \end{aligned} \quad (55)$$

Remark 7. *The definitions*

$$q_n = -\mathbf{n} \cdot \text{Div}(\mathbb{E}_\kappa), \quad M_{nn} = (\mathbf{n} \otimes \mathbf{n}) : \mathbb{E}_\kappa, \quad M_{ns} = (\mathbf{n} \otimes \mathbf{s}) : \mathbb{E}_\kappa$$

are exactly the same given in (27). The tensorial formalism allows a more compact writing.

And now only the kinematically independent variables must be highlighted. The tangential derivative has to be moved on the torsional momentum. In order to do that, the boundary needs to be split in a collection of regular subsets Γ_i , such that $\partial\Omega = \bigcup_{\Gamma_i \subset \partial\Omega} \Gamma_i$

$$\begin{aligned} \int_{\partial\Omega} \frac{\partial w_t}{\partial s} M_{ns} \, ds &= \sum_{\Gamma_i \subset \partial\Omega} \int_{\Gamma_i} \frac{\partial w_t}{\partial s} M_{ns} \, ds \\ &= \sum_{\Gamma_i \subset \partial\Omega} [M_{ns} w_t]_{\partial\Gamma_i} - \int_{\partial\Omega} \frac{\partial M_{ns}}{\partial s} w_t \, ds. \end{aligned} \quad (56)$$

If a regular boundary is considered the final energy balance is exactly the same obtained with the vectorial notation, namely

$$\dot{H} = \int_{\partial\Omega} \left\{ w_t \tilde{q}_n + \frac{\partial w_t}{\partial n} M_{nn} \right\} \, ds, \quad \text{where } \tilde{q}_n := q_n - \frac{\partial M_{ns}}{\partial s}. \quad (57)$$

The tensorial formulation allows highlighting the intrinsic differential operators. Furthermore the symmetric nature of the variables is explicitly expressed by the usage of symmetric tensors. Now that the energy balance has been established in terms of the boundary variables the Stokes-Dirac structure for the Kirchhoff plate in tensorial form can be defined. Consider now the bond space

$$\mathcal{B} := \{(\mathbf{f}, \mathbf{e}, \mathbf{z}) \in \mathcal{F} \times \mathcal{E} \times \mathcal{Z}\}, \quad (58)$$

where $\mathcal{F} = \mathcal{L}^2(\Omega) := L^2(\Omega) \times L^2(\Omega, \mathbb{R}_{\text{sym}}^{2 \times 2})$ and $\mathcal{E} = \mathcal{H}^2(\Omega) = H^2(\Omega) \times H^{\text{div Div}}(\Omega, \mathbb{R}_{\text{sym}}^{2 \times 2})$. The space $H^{\text{div Div}}(\Omega, \mathbb{R}_{\text{sym}}^{2 \times 2})$ is such that

$$H^{\text{div Div}}(\Omega, \mathbb{R}_{\text{sym}}^{2 \times 2}) = \{\mathbb{A} \in L^2(\Omega, \mathbb{R}_{\text{sym}}^{2 \times 2}) \mid \text{div}(\text{Div}(\mathbb{A})) \in L^2(\Omega)\}.$$

Consider the space of boundary port variables

$$\mathcal{Z} := \left\{ \mathbf{z} \mid \mathbf{z} = \begin{pmatrix} \mathbf{f}_\partial \\ \mathbf{e}_\partial \end{pmatrix} \right\}, \quad \text{with } \mathbf{f}_\partial = \begin{pmatrix} w_t \\ \frac{\partial w_t}{\partial n} \end{pmatrix}, \quad \mathbf{e}_\partial = \begin{pmatrix} \tilde{q}_n \\ M_{nn} \end{pmatrix} \quad (59)$$

The duality pairing between elements of \mathcal{B} is then defined as follows

$$\langle\langle (\mathbf{f}_1, \mathbf{e}_1, \mathbf{z}_1), (\mathbf{f}_2, \mathbf{e}_2, \mathbf{z}_2) \rangle\rangle := \langle \mathbf{e}_1, \mathbf{f}_2 \rangle_{\mathcal{L}^2(\Omega)} + \langle \mathbf{e}_2, \mathbf{f}_1 \rangle_{\mathcal{L}^2(\Omega)} + \int_{\partial\Omega} B_J(\mathbf{z}_1, \mathbf{z}_2) \, ds, \quad (60)$$

where the pairing $\langle \cdot, \cdot \rangle_{\mathcal{L}^2(\Omega)}$ is the L^2 inner product on space $\mathcal{L}^2(\Omega)$ and $B_J(\mathbf{z}_1, \mathbf{z}_2) = (\mathbf{f}_{\partial,1})^T \mathbf{e}_{\partial,2} + (\mathbf{f}_{\partial,2})^T \mathbf{e}_{\partial,1}$.

Theorem 4 (Stokes-Dirac Structure for the Kirchhoff plate in tensorial form). *Consider the space of power variables \mathcal{B} defined in (58) and the matrix differential operator J defined in (54). By theorem 2 in [26] the linear subspace $\mathcal{D} \subset \mathcal{B}$*

$$\mathcal{D} = \left\{ (\mathbf{f}, \mathbf{e}, \mathbf{z}) \in \mathcal{B} \mid \mathbf{f} = -\frac{\partial \boldsymbol{\alpha}}{\partial t} = -J\mathbf{e}, \mathbf{z} = \begin{pmatrix} \mathbf{f}_\partial \\ \mathbf{e}_\partial \end{pmatrix} \right\}, \quad (61)$$

is a Stokes-Dirac structure with respect to the pairing $\langle\langle \cdot, \cdot \rangle\rangle$ given by (60)

4. Discretization of the Kirchhoff plate using a Partioned Finite Element Method

Following the procedure illustrated in [9] the Kirchhoff plate written as a port-Hamiltonian system can be discretized by using a Partitioned Finite Element Method (PFEM). This method is an extension of the Mixed Finite Element Method to the case of pH systems and requires the integration by parts to be performed so that the symplectic structure is preserved. It consist in three different steps:

1. the system is first put into weak form;
2. once the boundary control of interest is selected, the corresponding subsystem is integrated by parts;
3. the problem is discretized by using a Mixed Finite Element method.

The weak form is illustrated using the tensorial formulation. Two different kind of boundary control will be shown:

1. boundary control through forces and momenta, in this case the first line of (54) is integrated by parts;
2. boundary control through kinematic variables, the second line of (54) is integrated by parts.

4.1. Weak form

The same procedure detailed above can be used on system (54). In this case the test functions are of scalar or tensorial nature. Keeping the same notation of section 3.2 the scalar test function is denoted by v_w , the tensorial one by \mathbb{V}_κ .

4.1.1. Boundary control through forces and momenta

The fist line of (54) is multiplied by v_w (scalar multiplication), the second line by \mathbb{V}_κ (tensor contraction).

$$\int_{\Omega} v_w \frac{\partial \alpha_w}{\partial t} \, d\Omega = \int_{\Omega} -v_w \operatorname{div}(\operatorname{Div}(\mathbb{E}_\kappa)) \, d\Omega, \quad (62)$$

$$\int_{\Omega} \mathbb{V}_\kappa : \frac{\partial \mathbb{A}_\kappa}{\partial t} \, d\Omega = \int_{\Omega} \mathbb{V}_\kappa : \operatorname{Grad}(\operatorname{grad}(e_w)) \, d\Omega. \quad (63)$$

The right hand side of equation (62) has to be integrated by parts twice

$$\int_{\Omega} -v_w \operatorname{div}(\operatorname{Div}(\mathbb{E}_\kappa)) \, d\Omega = \int_{\partial\Omega} \underbrace{-\mathbf{n} \cdot \operatorname{Div}(\mathbb{E}_\kappa)}_{q_n} v_w \, ds + \int_{\Omega} \operatorname{grad}(v_w) \cdot \operatorname{Div}(\mathbb{E}_\kappa) \, d\Omega \quad (64)$$

Applying again the integration by parts

$$\int_{\Omega} \operatorname{grad}(v_w) \cdot \operatorname{Div}(\mathbb{E}_\kappa) \, d\Omega = \int_{\partial\Omega} \operatorname{grad}(v_w) \cdot (\mathbf{n} \cdot \mathbb{E}_\kappa) \, ds - \int_{\Omega} \operatorname{Grad}(\operatorname{grad}(v_w)) : \mathbb{E}_\kappa \, d\Omega \quad (65)$$

The usual additional manipulation is performed on the boundary term containing the momenta, so that the proper boundary values arise

$$\begin{aligned}
\int_{\partial\Omega} \text{grad}(v_w) \cdot (\mathbf{n} \cdot \mathbb{E}_\kappa) \, ds &= \int_{\partial\Omega} \left(\frac{\partial v_w}{\partial n} \mathbf{n} + \frac{\partial v_w}{\partial s} \mathbf{s} \right) \cdot (\mathbf{n} \cdot \mathbb{E}_\kappa) \, ds \\
&= \int_{\partial\Omega} \left\{ \frac{\partial v_w}{\partial n} \underbrace{(\mathbf{n} \otimes \mathbf{n}) : \mathbb{E}_\kappa}_{M_{nn}} + \frac{\partial v_w}{\partial s} \underbrace{(\mathbf{n} \otimes \mathbf{s}) : \mathbb{E}_\kappa}_{M_{ns}} \right\} \, ds \\
&= \sum_{\Gamma_i \subset \partial\Omega} [M_{ns} v_w]_{\partial\Gamma_i} + \int_{\partial\Omega} \left\{ \frac{\partial v_w}{\partial n} M_{nn} - v_w \frac{\partial M_{ns}}{\partial s} \right\} \, ds
\end{aligned} \tag{66}$$

Combining equations (64), (65) and (66) the final expression which make appear the dynamic boundary terms (forces and momenta) is found

$$\int_{\Omega} v_w \frac{\partial \alpha_w}{\partial t} \, d\Omega = - \int_{\Omega} \text{Grad}(\text{grad}(v_w)) : \mathbb{E}_\kappa \, d\Omega + \int_{\partial\Omega} \left\{ \frac{\partial v_w}{\partial n} M_{nn} + v_w \tilde{q}_n \right\} \, ds + \sum_{\Gamma_i \subset \partial\Omega} [M_{ns} v_w]_{\partial\Gamma_i}. \tag{67}$$

If the boundary is regular the final expression simplifies

$$\int_{\Omega} v_w \frac{\partial \alpha_w}{\partial t} \, d\Omega = - \int_{\Omega} \text{Grad}(\text{grad}(v_w)) : \mathbb{E}_\kappa \, d\Omega + \int_{\partial\Omega} \left\{ \frac{\partial v_w}{\partial n} M_{nn} + v_w \tilde{q}_n \right\} \, ds. \tag{68}$$

So the final weak form obtained from system (54) reads

$$\begin{cases} \int_{\Omega} v_w \frac{\partial \alpha_w}{\partial t} \, d\Omega &= - \int_{\Omega} \text{Grad}(\text{grad}(v_w)) : \mathbb{E}_\kappa \, d\Omega + \int_{\partial\Omega} \left\{ \frac{\partial v_w}{\partial n} M_{nn} + v_w \tilde{q}_n \right\} \, ds, \\ \int_{\Omega} \mathbb{V}_\kappa : \frac{\partial \mathbb{A}_\kappa}{\partial t} \, d\Omega &= \int_{\Omega} \mathbb{V}_\kappa : \text{Grad}(\text{grad}(e_w)) \, d\Omega. \end{cases} \tag{69}$$

The control input \mathbf{u}_∂ and the corresponding conjugate outputs \mathbf{y}_∂ are

$$\mathbf{u}_\partial = \begin{pmatrix} \tilde{q}_n \\ M_{nn} \end{pmatrix}_{\partial\Omega}, \quad \mathbf{y}_\partial = \begin{pmatrix} w_t \\ \frac{\partial w_t}{\partial n} \end{pmatrix}_{\partial\Omega}.$$

4.1.2. Boundary control through kinematic variables

Alternatively, the same procedure can be performed on the second line of the system to make appear the kinematic boundary conditions, i.e. the value of the vertical velocity and its normal derivative along the border. Once the necessary calculations are carried out, the following result is found

$$\begin{cases} \int_{\Omega} v_w \frac{\partial \alpha_w}{\partial t} \, d\Omega &= \int_{\Omega} -v_w \text{div}(\text{Div}(\mathbb{E}_\kappa)) \, d\Omega, \\ \int_{\Omega} \mathbb{V}_\kappa : \frac{\partial \mathbb{A}_\kappa}{\partial t} \, d\Omega &= \int_{\Omega} \text{div}(\text{Div}(\mathbb{V}_\kappa)) e_w \, d\Omega + \int_{\partial\Omega} \left\{ v_{M_{nn}} \frac{\partial w_t}{\partial n} + v_{\tilde{q}_n} w_t \right\} \, ds. \end{cases} \tag{70}$$

where $v_{M_{nn}} = (\mathbf{n} \otimes \mathbf{n}) : \mathbb{V}_\kappa$ and $v_{\tilde{q}_n} = -\text{Div}(\mathbb{V}_\kappa) \cdot \mathbf{n} - \frac{\partial v_{M_{ns}}}{\partial s}$ with $v_{M_{ns}} = (\mathbf{n} \otimes \mathbf{s}) : \mathbb{E}_\kappa$. The control input \mathbf{u}_∂ and the corresponding conjugate outputs \mathbf{y}_∂ are

$$\mathbf{u}_\partial = \begin{pmatrix} w_t \\ \frac{\partial w_t}{\partial n} \end{pmatrix}_{\partial\Omega}, \quad \mathbf{y}_\partial = \begin{pmatrix} \tilde{q}_n \\ M_{nn} \end{pmatrix}_{\partial\Omega}.$$

4.2. Finite-dimensional port-Hamiltonian system

In this section the formulation (69) is used in order to explain the discretization procedure. The same procedure may be performed using the formulation given in (70) and in 5.2 both strategies will be used to compute the eigenvalues of a square plate. Test and co-energy variables are discretized using the same basis function (Galerkin Method)

$$\begin{aligned} v_w &= \sum_{i=1}^{N_w} \phi_w^i(x, y) v_w^i, & e_w &= \sum_{i=1}^{N_w} \phi_w^i(x, y) e_w^i(t), \\ \mathbb{V}_\kappa &= \sum_{i=1}^{N_\kappa} \Phi_\kappa^i(x, y) v_\kappa^i, & \mathbb{E}_\kappa &= \sum_{i=1}^{N_\kappa} \Phi_\kappa^i(x, y) e_\kappa^i(t), \end{aligned} \quad (71)$$

The basis function ϕ_w^i, Φ_κ^i , have to be chosen in a suitable function space \mathcal{V}^h in the domain of operator J , i.e. $\mathcal{V}^h \subset \mathcal{V} \in \mathcal{D}(J)$. This will be discussed in section 5. The discretized skew-symmetric bilinear form on the right side of (69) then reads

$$\mathbf{J}_d = \begin{bmatrix} 0 & -\mathbf{D}_H^T \\ \mathbf{D}_H & 0 \end{bmatrix}, \quad (72)$$

where the matrix \mathbf{D}_H is computed in the following way

$$\mathbf{D}_H(i, j) = \int_{\Omega} \Phi_\kappa^i : \text{Grad}(\text{grad}(\phi_w^j)) \, d\Omega, \quad \in \mathbb{R}^{N_\kappa \times N_w}, \quad (73)$$

where the notation $A(i, j)$ indicates the entry in the matrix corresponding to the i th row and j th column. The energy variables are deduced from the co-energy variables

$$\alpha_w = \mu e_w, \quad \mathbb{A}_\kappa = \mathbb{D}^{-1} \mathbb{E}_\kappa, \quad (74)$$

where \mathbb{D}_{ijkl} is the symmetric bending rigidity tensor, the tensorial analogous of matrix \mathbf{D} defined in (15). The symmetric bilinear form on the left side of (69) is discretized as

$$\begin{aligned} \mathbf{M} &= \text{diag}[\mathbf{M}_w, \mathbf{M}_\kappa], \quad \text{with} \\ \mathbf{M}_w(i, j) &= \int_{\Omega} \rho h \phi_w^i \phi_w^j \, d\Omega \in \mathbb{R}^{N_w \times N_w}, \\ \mathbf{M}_\kappa(i, j) &= \int_{\Omega} (\mathbb{D}^{-1} \Phi_\kappa^i) : \Phi_\kappa^j \, d\Omega \in \mathbb{R}^{N_\kappa \times N_\kappa}, \end{aligned} \quad (75)$$

The boundary variables are then discretized as

$$\tilde{q}_n = \sum_{i=1}^{N_{\tilde{q}_n}} \phi_{\tilde{q}_n}^i(s) \tilde{q}_n^i, \quad M_{nn} = \sum_{i=1}^{N_{M_{nn}}} \phi_{M_{nn}}^i(s) M_{nn}^i. \quad (76)$$

The variables are defined only over the boundary $\partial\Omega$. Consequently, the input matrix reads

$$\mathbf{B} = \begin{bmatrix} \mathbf{B}_{\tilde{q}_n} & \mathbf{B}_{M_{nn}} \\ 0 & 0 \end{bmatrix} \quad (77)$$

where the inner components are computed as

$$\begin{aligned} \mathbf{B}_{\tilde{q}_n}(i, j) &= \int_{\partial\Omega} \phi_w^i \phi_{q_n}^j \, ds \quad \in \mathbb{R}^{N_w \times N_{q_n}}, \\ \mathbf{B}_{M_{nn}}(i, j) &= \int_{\partial\Omega} \frac{\partial \phi_w^i}{\partial n} \phi_{M_{nn}}^j \, ds \quad \in \mathbb{R}^{N_w \times N_{M_{nn}}}, \end{aligned} \quad (78)$$

The final port-Hamiltonian system, as defined in [28] is written as

$$\begin{aligned} \mathbf{M} \dot{\mathbf{e}} &= \mathbf{J}_d \mathbf{e} + \mathbf{B} \mathbf{u}_\partial \\ \mathbf{y}_\partial &= \mathbf{B}^T \mathbf{e} \end{aligned} \quad (79)$$

where $\mathbf{e} = (e_w^1, \dots, e_\kappa^{N_\kappa})^T$ and $\mathbf{u}_\partial = (\tilde{q}_n^1, \dots, M_{nn}^{N_{M_{nn}}})^T$ are the concatenations of the degrees of freedom for the different variables. The discrete Hamiltonian is then found as

$$\begin{aligned} H_d &= \frac{1}{2} \int_{\Omega} \{ \alpha_w e_w + \mathbb{A}_\kappa : \mathbb{E}_\kappa \} \, d\Omega \\ &= \frac{1}{2} \{ \mathbf{e}_w^T \mathbf{M}_w \mathbf{e}_w + \mathbf{e}_\kappa^T \mathbf{M}_\kappa \mathbf{e}_\kappa \} \\ &= \frac{1}{2} \mathbf{e}^T \mathbf{M} \mathbf{e} \end{aligned} \quad (80)$$

Using equations (79), (80) the time derivative of the Hamiltonian is given by the scalar product of the boundary flows

$$\dot{H}_d = \mathbf{y}_\partial^T \mathbf{u}_\partial. \quad (81)$$

The above Equation is equivalent to the energy balance of the continuous system, expressed by (31). Definition (80), together with system (79) are the finite-dimensional equivalent of (47) and (54). The discretized system obtained via PFEM shares the [port-Hamiltonian structure](#) of the original infinite-dimensional system, the discretization method is therefore structure preserving.

5. Numerical studies

In this section we illustrate numerically the consistency of the proposed model. For this purpose the computation of the eigenvalues of a square plate and temporal simulations for several boundary conditions are presented.

5.1. Finite Element Choice

The domain of the operator J in (54) is $\mathcal{D}(J) = H^2(\Omega) \times H^{\text{div Div}}(\Omega, \mathbb{R}_{\text{sym}}^{2 \times 2}) + \text{boundary conditions}$. For this reason a suitable choice for the functional space is

$$(v_w, \mathbb{V}_\kappa) \in H^2(\Omega) \times H^2(\Omega, \mathbb{R}_{\text{sym}}^{2 \times 2}) \equiv \mathcal{H}, \quad (82)$$

since $\mathcal{H} \subset \mathcal{D}(J)$.

Remark 8. It has to be appointed that the space $H^{\text{div Div}}(\Omega, \mathbb{R}_{\text{sym}}^{2 \times 2})$ was never addressed in the mathematical literature. A theoretical analysis of this space is needed. For this reason the only way to deal with this problem numerically is to use $H^2(\Omega)$ conforming finite elements.

The H^2 conforming finite elements (like the Hermite, Bell or Argyris finite elements) do not satisfy the proper equivalence properties to give a simple relationship between the reference basis and nodal basis on a general cell [29]. For this reason the Firedrake library [11] was used to implement the numerical analysis as it provides functionalities to automate the generalized mappings for these elements.

Then for the Finite Element choice denote by

$$H_r^k(\mathbb{P}_l, \Omega) = \{v \in H^k(\Omega) \mid v|_T \in \mathbb{P}_l \ \forall T \in \mathbb{T}_r\}$$

the finite element space which is a subspace of $H^k(\Omega)$, based on the shape function space of piecewise polynomials of degree l . The shape function space is defined over the mesh $\mathbb{T}_r = \bigcup_i T_i$, where the cells T_i are triangles. These spaces can be scalar-valued or symmetric matrix valued, depending on the variables to be discretized. The parameter r is the average size of a mesh element. All the variables, i.e. the velocity e_w and the momenta tensor \mathbb{V}_κ as well as the corresponding test functions, are discretized by the same finite element space, the Bell finite element [30], denoted by $H_r^2(\mathbb{P}_5, \Omega)$. For this element the field is computed using quintic polynomials whose degrees of freedom are the values of the function, its gradient and its Hessian at the vertex of each triangular element. In the following mixed boundary conditions will be considered. To deal with those Lagrange multipliers have to be introduced (the reader can refer to [26], section 4.3 for an explanation). The multipliers are therefore discretized by using second degree Lagrange polynomials defined over the boundary $H_r^1(\mathbb{P}_2, \partial\Omega)$.

5.2. Eigenvalues Computation

The test case for this analysis is a simple square plate of side L , a benchmark problem which has been studied in [31, 32] for different boundary conditions on each plate side. The possible cases are the following:

- clamped side (C), for which $w_t = 0$, $\frac{\partial w_t}{\partial n} = 0$;
- simply supported side (S), $w_t = 0$, $M_{nn} = 0$;
- free side (F), $\tilde{q}_n = 0$, $M_{nn} = 0$.

In order to compare our results the computed frequencies ω_n^h in the following non-dimensional form:

$$\hat{\omega}_n^h = 4L^2\omega_n^h \left(\frac{\rho h}{D}\right)^{1/2}, \quad (83)$$

The only parameters which influence the results are the Poisson modulus $\nu = 0.3$, but only for cases containing free sides [12]. The reported non-dimensional frequencies are independent of the remaining geometrical and physical parameters. The error is computed as

$$\varepsilon = \frac{\text{abs}(\hat{\omega}_n^h - \omega_n^{DR})}{\omega_n^L}, \quad (84)$$

where ω_n^L are the eigenvalues computed in [31]. The results are reported in table 1 and 2, together with the results from Leissa [31]. The boundary condition are denoted by string of four letters, one for each side of the square. The results are reported either by using the

n	CSCS			SSCS			SSSS		
	Leissa	Hessian	divDiv	Leissa	Hessian	divDiv	Leissa	Hessian	divDiv
$\hat{\omega}_1$	28.946	28.950	28.951	23.646	23.640	23.646	19.739	19.730	19.739
$\hat{\omega}_2$	54.743	54.747	54.744	51.674	51.666	51.675	49.348	49.333	49.348
$\hat{\omega}_3$	69.32	69.331	69.330	58.641	58.641	58.647	49.348	49.336	49.348
$\hat{\omega}_4$	94.584	94.602	94.593	86.126	86.121	86.138	78.957	78.920	78.958
$\hat{\omega}_5$	102.213	102.245	102.221	100.259	100.284	100.275	98.696	98.692	98.700
$\hat{\omega}_6$	129.086	129.141	129.110	113.217	113.250	113.237	98.696	98.709	98.703

Table 1: Eigenvalues obtained with 5 Bell element per side for $\nu = 0.3$, considering either the Grad \circ grad formulation (69), either the Div \circ div formulation (70).

■ reference, ■ $\varepsilon < 0.1\%$.

forces and momenta as control (69) (column Hessian in tables 1, 2) or the vertical linear and angular velocity (70) (column divDiv in tables 1, 2). The results are obtaining using a regular mesh composed by 5 Bell element on each side. Hence, the state vector has a total dimension of 864. The dimension of the Lagrange multiplier vector depends on the boundary condition upon consideration. When using $H_r^2(\mathbb{P}_2, \partial\Omega)$ on the considered mesh, this number can vary from 0 depending on the bcs. The results obtained by using (69) are in perfect agreement with the reference. This formulation was also used to compute the eigenvectors corresponding to the vertical velocity, for the different cases under examination: CSCS in figure 4, SSCS in 5, SSSS in 6, CSFS in 7, SSFS in 8 and FSFS 9. For what concerns the weak formulation (70) the results deteriorate when a free condition (table 2) is present.

n	CSFS			SSFS			FSFS		
	Leissa	Hessian	divDiv	Leissa	Hessian	divDiv	Leissa	Hessian	divDiv
$\hat{\omega}_1$	12.69	12.681	14.336	11.68	11.679	13.136	9.631	9.630	11.110
$\hat{\omega}_2$	33.06	33.041	33.895	27.76	27.732	28.681	16.13	16.117	18.421
$\hat{\omega}_3$	41.7	41.692	43.791	41.2	41.185	43.073	36.72	36.683	38.058
$\hat{\omega}_4$	63.01	62.982	64.753	59.07	59.027	60.764	38.94	38.939	40.808
$\hat{\omega}_5$	72.4	72.371	72.756	61.86	61.825	62.298	46.74	46.709	50.252
$\hat{\omega}_6$	90.61	90.602	92.960	90.29	90.283	92.436	70.75	70.666	73.550

Table 2: Eigenvalues obtained with 5 Bell element per side for $\nu = 0.3$, considering either the Grad \circ grad formulation (69), either the Div \circ div formulation (70).

■ reference, ■ $\varepsilon < 0.1\%$, ■ $\varepsilon < 1\%$, ■ $\varepsilon < 5\%$, ■ $\varepsilon < 15\%$.

5.3. Temporal Simulations

In this analysis we consider a square plate, subject either to a non null shear force on the boundaries either to a distributed force over the domain. The physical parameters and simulation setting are reported in table 3. The energy variables and Lagrange multipliers are discretized using Bell shape functions (five elements for each side) and second order Lagrange polynomials respectively. The symplectic Störmer-Verlet integrator is employed, so that when no solicitation is applied to the system, the Hamiltonian is preserved. Two different simulation with different BC are considered. The initial condition are set to zero

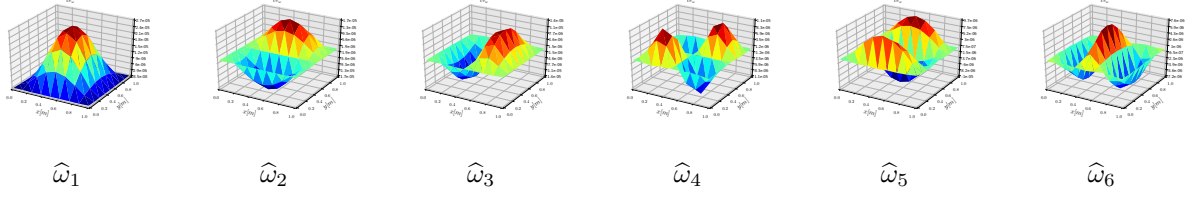


Figure 4: Eigenvectors for the CSCS case

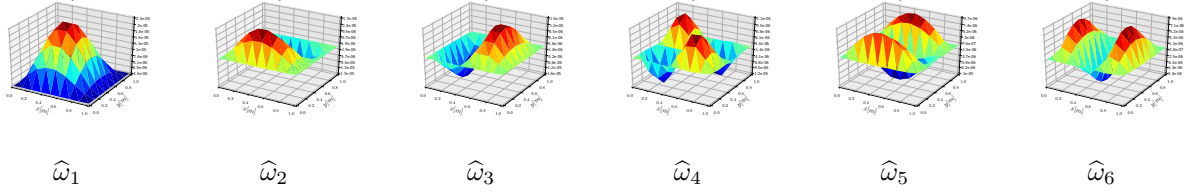


Figure 5: Eigenvectors for the SSCS case

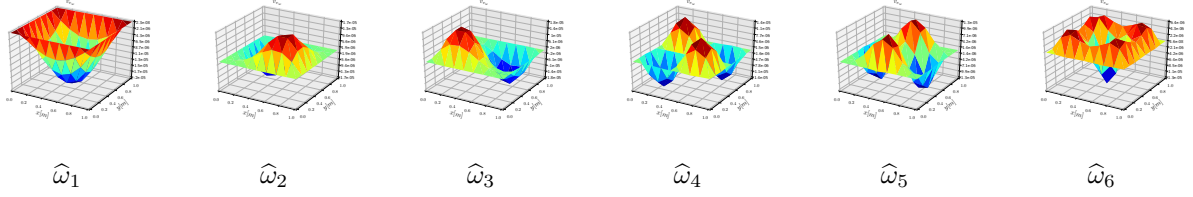


Figure 6: Eigenvectors for the SSSS case

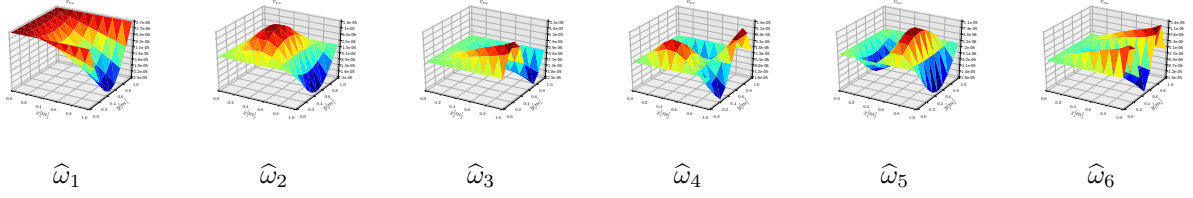


Figure 7: Eigenvectors for the CSFS case

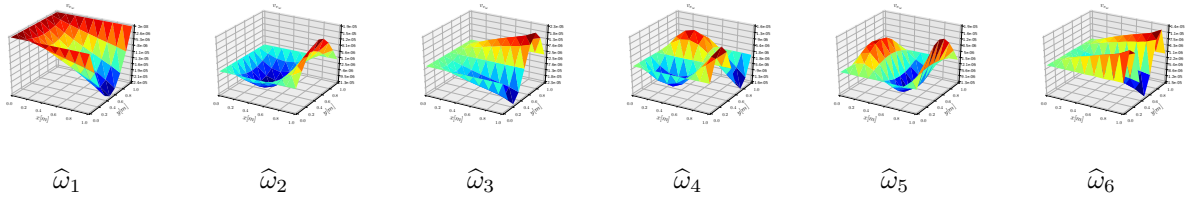


Figure 8: Eigenvectors for the SSFS case

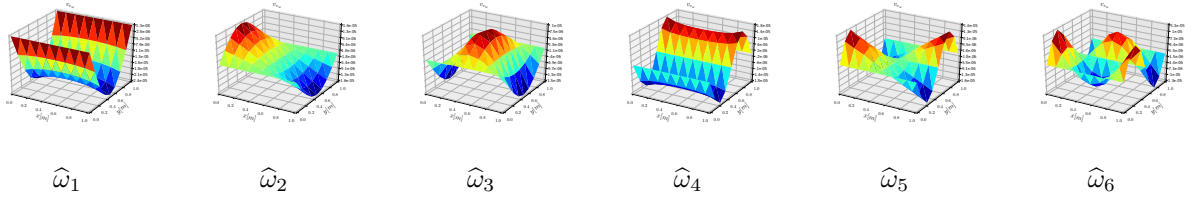


Figure 9: Eigenvectors for the FSFS case

Plate Parameters		Simulation Parameters	
E	70 [GPa]	Δt	0.001 [ms]
ν	0.35	t_{end}	1 [ms]
h/L	0.05	Integrator	Störmer-Verlet
L	1 [m]	N° Elements	5
		FE space	$H_{r=L/5}^2(\mathbb{P}_5, \Omega)$ for $\mathbf{e} \times H_{r=L/5}^1(\mathbb{P}_2, \partial\Omega)$ for $\boldsymbol{\lambda}$

Table 3: Simulations settings

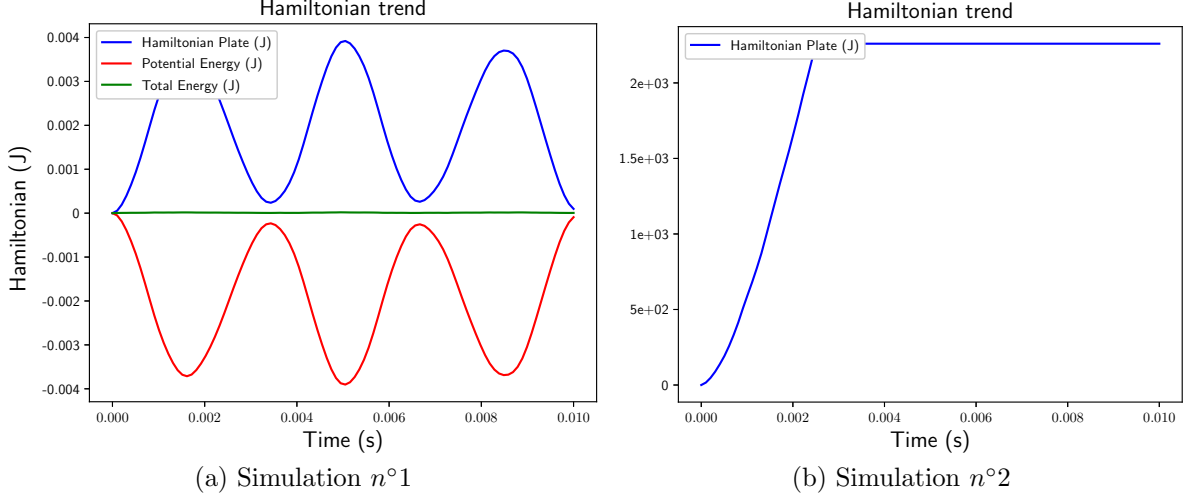


Figure 10: Hamiltonian trend for the two simulations

for each variables. For the first simulation a plate subject to gravity is considered. For this simulation the set of BC reads:

$$\text{Simulation } n^\circ 1 \quad \begin{cases} w_t = 0, \frac{\partial w_t}{\partial n} = 0, & \text{for } x = 0, y = 0 \text{ and } y = 1 \\ \tilde{q}_n = 0, M_{nn} = 0, & \text{for } x = 1 \end{cases} \quad (85)$$

Since the solicitation admits a potential the Hamiltonian does not represent the total energy, that now includes the potential energy

$$E_p = \int_{\Omega} \rho h g w \, d\Omega, \quad (86)$$

where w is the vertical displacement and $g = 10[m/s^2]$ is the gravity acceleration. For the second simulation the following boundary conditions, corresponding to the case $CFFF$ are considered

$$\text{Simulation } n^\circ 2 \quad \begin{cases} w_t = 0, \frac{\partial w_t}{\partial n} = 0, & \text{for } x = 0, \\ q_n = +f(1, t), M_{nn} = M_{ns} = 0, & \text{for } x = 1, \\ q_n = +f(x, t), M_{nn} = M_{ns} = 0, & \text{for } y = 0, \\ q_n = +f(x, t), M_{nn} = M_{ns} = 0, & \text{for } y = 1, \end{cases} \quad (87)$$

where

$$f(x, t) = \begin{cases} 10^5 x [Pa \cdot m], & \forall t < 0.2 t_{\text{end}}, \\ 0, & \forall t \geq 0.2 t_{\text{end}}. \end{cases} \quad (88)$$

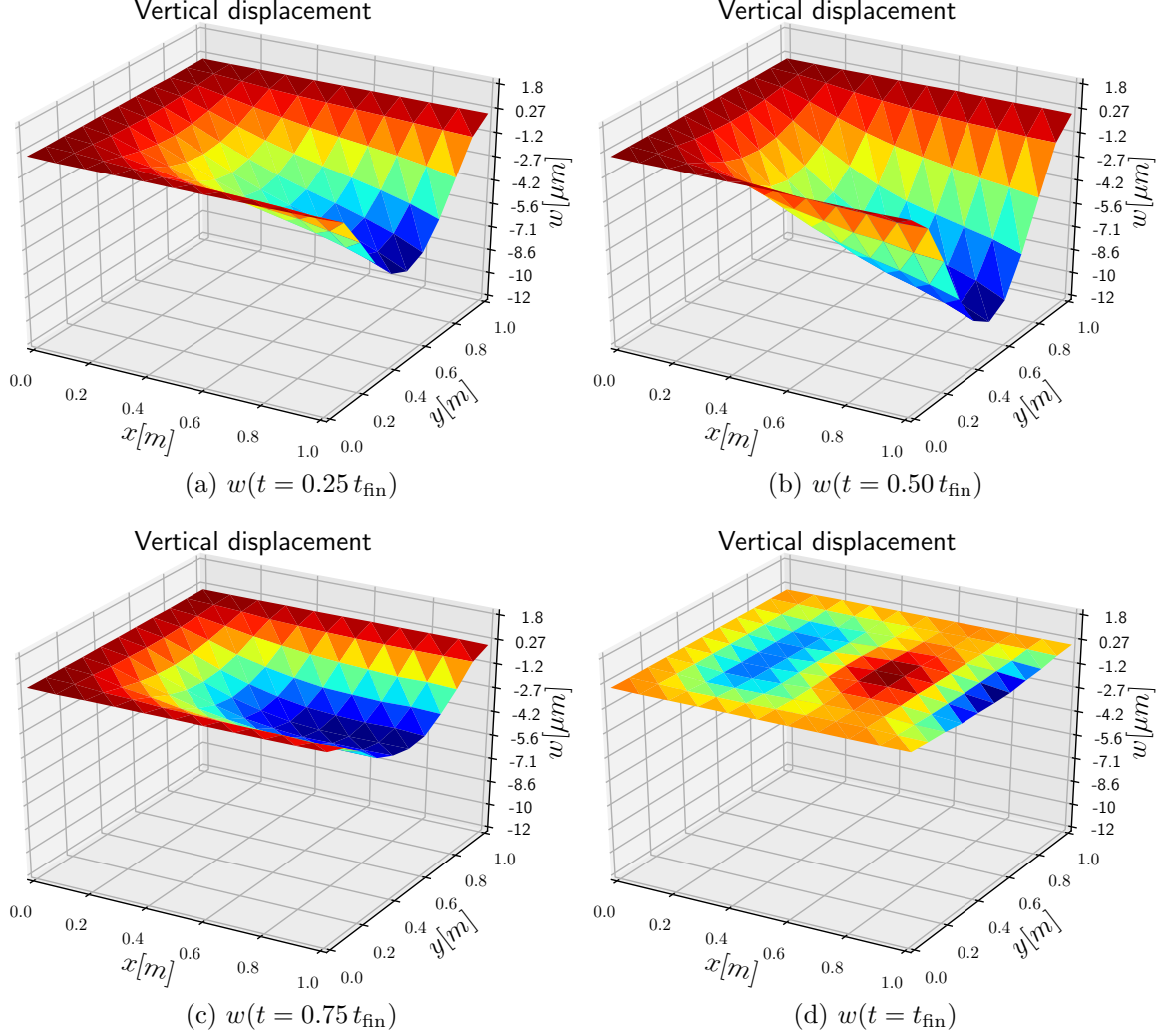


Figure 11: Snapshots for Simulation $n^\circ 1$.

In this case inhomogeneous boundary conditions are considered. Snapshots of the vertical displacement are reported in figure 11, 12. This field is obtained from the velocity field $e_w = \frac{\partial w}{\partial t}$ by applying the trapezoidal rule integration. For both simulations the output is consistent with the imposed BC and with the physical intuition of the observed phenomenon. The symplectic integration has been used to demonstrate numerically the conservation of total energy. The Hamiltonian for both simulation is reported in figure 10. For the first one the total energy E_p is the conserved quantity and it stays constant equal to zero as expected. For the second simulation the Hamiltonian is conserved quantity once the excitation is set back to zero.

Conclusions and Future Perspectives

In this paper the port-Hamiltonian formulation of the Kirchhoff plate was detailed with the equivalent vectorial and a tensorial representation. The tensorial formalism allowed showing the adjointness relation between the double divergence of a symmetric tensor and

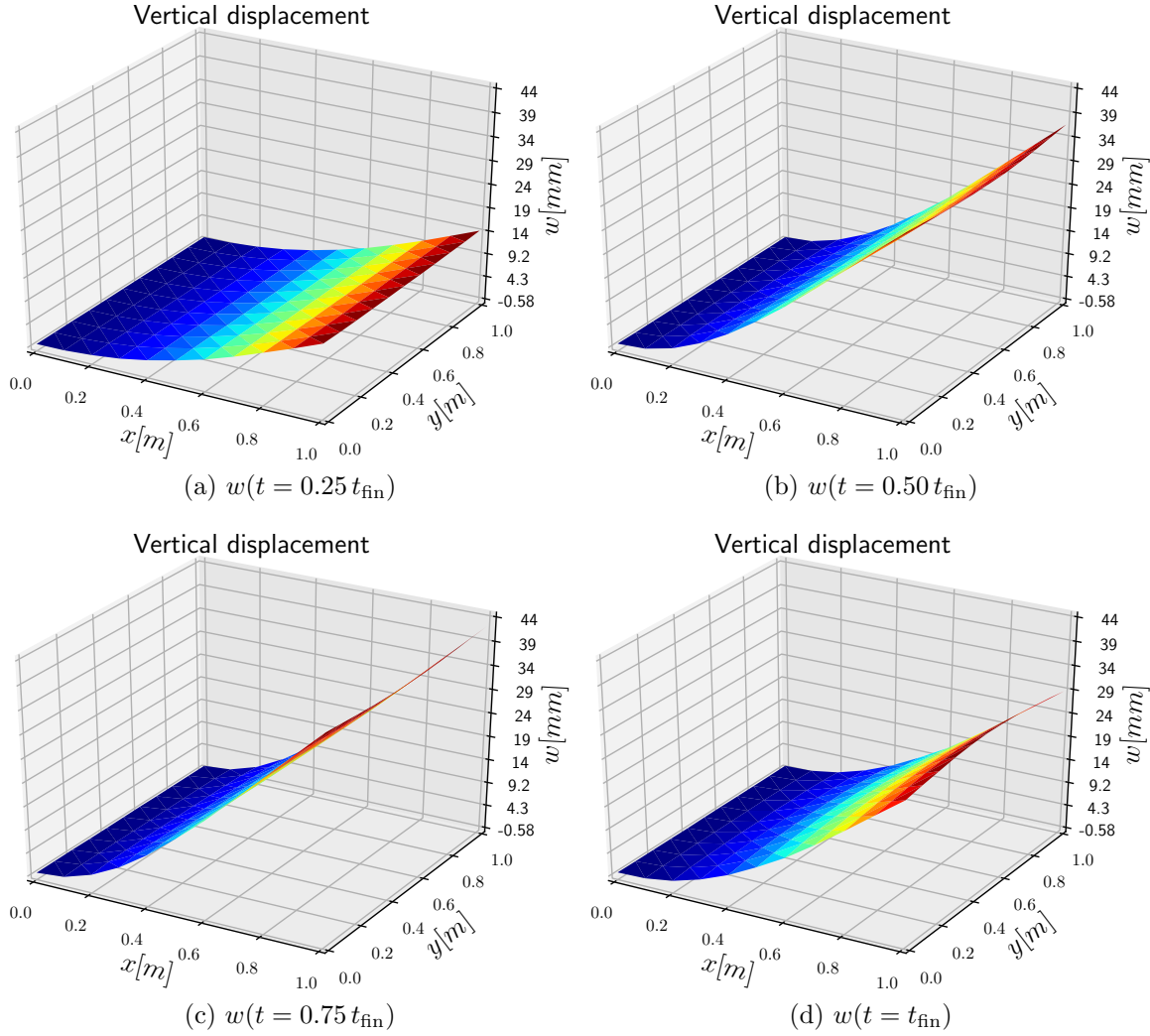


Figure 12: Snapshots for Simulation $n^\circ 2$.

the Hessian of a scalar field. This result represents an appealing novelty for mathematical working on functional analysis. Moreover, many features of the PFEM are of interest:

- its capability of preserving the port-Hamiltonian structure;
- the natural derivation of boundary port variables as inputs;
- the possibility of dealing with mixed boundary conditions inside the framework of port-Hamiltonian descriptor systems PHDAEs detailed in [28];
- the easy implementability of the method using standard Finite Element libraries (Fire-drake [11] in this case);

The computation of eigenvalues with different boundary conditions and the numerical simulations demonstrate the validity of the proposed model.

The model presented in this paper should be completed with a precise analysis of the well-posedness, in the input-output sense. It must be appointed that a complication arises in this

formulation. The differential operator of the PH model of the Kirchhoff plate requires the momenta to belong to the space $H^{\text{div Div}}(\Omega, \mathbb{R}_{\text{sym}}^{d \times d})$. To the best of our knowledge this space was never analyzed in the mathematical literature and a precise study of its peculiarities is needed. The results obtained in [33] for the wave equation in \mathbb{R}^d could be generalized to the second order differential operator presented herein. A numerical analysis focusing on the convergence of appropriate finite elements should be carried out.

The discretization procedure details in the paper open new scenarios on the interconnection of PH systems. Starting from the results stated in [5] this system may be interconnected over its boundary to other finite or infinite dimensional PH systems, such as rigid bodies or other flexible appendages. This may find useful applications in simulating a multi-body environment for spatial applications, like the attitude motion of a satellite with flexible solar panels [20]. Since no causality is imposed on the boundary ports, the discretization method herein proposed allows the construction of arbitrarily complex connections among different modules. This feature is particularly appealing for complex applications.

Acknowledgments

This work is supported by the project ANR-16-CE92-0028, entitled *Interconnected Infinite-Dimensional systems for Heterogeneous Media*, INFIDHEM, financed by the French National Research Agency (ANR) and the Deutsche Forschungsgemeinschaft (DFG). Further information is available at <https://websites.isae-supaero.fr/infidhem/the-project>. Moreover the author would like to thank Michel Salaün for the fruitful and insightful discussions

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