

NMPDE/ATSC 2025

Lecture 13

Convection (reaction) dominates
diffusion problem: (Q, Ch. 13)

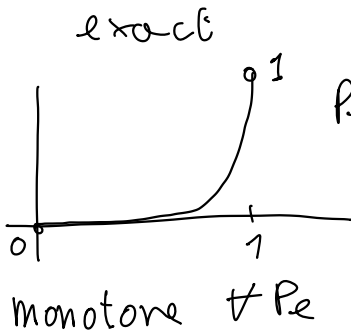
$a, b > 0$ constants

$$\begin{cases} -a u'' + b u' = 0 \\ u(0) = 0 \quad u(1) = 1 \end{cases} \quad \Omega = (0, 1)$$

$$Pe = \frac{b L}{2 a}$$

$$(L = 1 = |\Omega|)$$

} singularly perturbed



$$Pe \gg 1$$

or $a \rightarrow 0 \quad u'(1) \rightarrow \infty$
perturbation of

$$\rightarrow \begin{cases} b u' = 0 \\ u = 0 \end{cases} \rightarrow u \equiv 0$$

Discretisation (by linear FEM or FD) gives
the scheme:

$$\begin{cases} -a \frac{U_{j+1} - 2U_j + U_{j-1}}{h^2} + b \frac{U_{j+1} - U_{j-1}}{2h} = 0 \\ U_0 = 0; U_M = 1 \end{cases}$$

exercise: show that linear FEM gives same scheme

$$P_{eh} := \frac{bh}{2a} \xrightarrow{h \rightarrow 0} 0 \rightarrow P_{eh} < 1$$

if $P_{eh} > 1$ U is not monotone

still the scheme is convergent and monotone for h small enough

example: $\begin{cases} a = 10^{-6} \\ b = 1 \end{cases}$ to get need $h \sim 10^{-6}$ 1D
 $P_{eh} < 1$

2D } h^{-6} requires 10^{12}
 3D } 10^{13} DoF !

→ need to "stabilize" the scheme...

UPWIND method

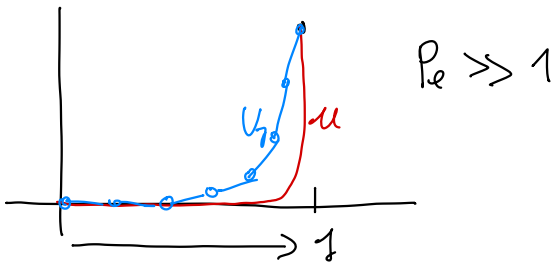
$$\downarrow u'(x) \sim \frac{u(x) - u(x-h)}{h}$$

$$-a \frac{U_{j+1} - 2U_j + U_{j-1}}{h^2} + b \frac{U_j - U_{j-1}}{h} = 0$$

$$U_j = \frac{(1 + 2Pe_h)^j - 1}{(1 + 2Pe_h)^N - 1} \quad \text{monotone !}$$

$\forall h$

(exercise) $|u(x_j) - U_j| = O(h)$



upwind is monotone but overdiffusive
why??

$$b \frac{U_j - U_{j-1}}{h} = b \frac{U_j - \frac{1}{2}U_{j-1} - \frac{1}{2}U_{j-1} + \frac{1}{2}U_{j+1} - \frac{1}{2}U_{j+1}}{h}$$

$\overset{=0}{\phantom{- \frac{1}{2}U_{j+1} + \frac{1}{2}U_{j+1}}}$

$$= b \frac{U_{j+1} - U_{j-1}}{2h} - \frac{b}{2} h \frac{U_{j-1} - 2U_j + U_{j+1}}{h^2}$$

$$\underbrace{b u'} - \frac{b}{2} h u''$$

→ inconsistency

upwind \equiv centered scheme applied to

$$\underbrace{-\alpha(1 + P_{en})}_{=: \alpha_h} u'' + b u' = 0 \quad \alpha_h = \alpha \left(1 + \frac{b h}{2 \alpha}\right)$$

numerical/artificial diffusion

$$P_{en}^{up} = \frac{b h}{2 \alpha_h} = \frac{P_{en}}{1 + P_{en}} < 1$$

In general, consider the perturbed PDE with diffusion $\alpha_h = \alpha (1 + \phi(P_{en}))$

upwind $\phi(t) = t$

or optimize \rightarrow gives

$$\rightarrow \begin{cases} \phi(t) = t - 1 + B(2t) \\ B(t) = \begin{cases} \frac{t}{e^t - 1} & \text{if } t > 0 \\ 1 & t = 0 \end{cases} \end{cases} \quad (\text{Bernoulli})$$

then centred scheme $\mathcal{O}(h^2)$.

↓ Gives the exponential fitting or Scharfetter-Gummel.

Analysis of (upwind) scheme, FEM style:
instead of $\mathcal{A}(u, v) = \int_0^1 a u' v' + \int_0^1 b u' v$
use $\mathcal{A}_h(u, v) = \int_0^1 a_h u' v' + \int_0^1 b u' v$

Applying FEM results in Generalized Galerkin method
⇒ Analysis via Strang Lemma

$$|u - u_h|_1 \leq C \inf_{v_h} \left\{ \dots + \sup_{w_h \in V_h} \left| \frac{\mathcal{A}(v_h, w_h) - \mathcal{A}_h(v_h, w_h)}{\|w_h\|_1} \right| \right\}$$

$$(Q) \quad |u - u_h|_1 \leq \frac{Ch^k}{2(1 + \phi(P_n))} \|u\|_{k+1} + \frac{\phi(P_n)}{1 + \phi(P_n)} \|u\|_1$$

$$\begin{cases} \mathcal{O}(h) & \text{upwind} \\ \mathcal{O}(h^2) & \text{S.G.} \end{cases}$$

Convection-diffusion in multi-dim.

$$\Omega \subset \mathbb{R}^d \quad a = a(x), \underline{b} = \underline{b}(x)$$

$$\mathcal{L}u := \begin{cases} -a \Delta u + \underline{b} \cdot \nabla u = 0 & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega \end{cases}$$

$$a > 0 \text{ o.e.}$$

conservative term
$\nabla \cdot (\underline{b} u)$
similar to
$\nabla \cdot (\underline{b} u)$
$= \underline{b} \cdot \nabla u + \underbrace{(\nabla \cdot \underline{b}) u}_{\text{extra reaction term}}$

Direct generalization of artificial diffusion method

Apply FE to $\mathcal{L}u \rightarrow \mathcal{L}u - h \overline{b} \Delta u$

where $\overline{b} = \|\underline{b}\|_{\infty, \Omega}$

\uparrow
 $\underbrace{h \overline{b} \Delta u}_{h \overline{b} u''}$

Strongy : $T(u_h) = \overbrace{A_h(u, v_h) - F(v_h)}$

$$= \overline{b} h \int_{\Omega} \nabla u \cdot \nabla v_h$$

$$\left\{ \begin{array}{l} \downarrow \\ \|u - u_h\|_1 \leq O(h^k) + h \overline{b} \|u\|_1 \end{array} \right.$$

$$\rightarrow O(h)$$

We need stronger consistency to retain the order k accuracy.

Streamline diffusion

add instead diffusion only in the direction given by \underline{b}

$$\rightarrow \frac{h}{\underline{b}} \int_{\Omega} (\underbrace{\underline{b} \cdot \nabla u}_{\parallel \frac{\partial u}{\partial \underline{b}}}) (\underbrace{\underline{b} \cdot \nabla v}_{\parallel \frac{\partial v}{\partial \underline{b}}})$$

Hughes-Brooks (1979) fully consistent streamline diffusion term.

$$\begin{aligned} &= \sum_T \frac{h}{\underline{b}} \int_T \left(\alpha \Delta u_h + \underbrace{\underline{b} \cdot \nabla u_h - f}_{\text{RHS}} \right) (\underline{b} \cdot \nabla v_h) \\ &\quad \underbrace{\hspace{10em}}_{=0 \text{ for } u \text{ exact sol}} \\ &= B_h(u_h, v_h) \end{aligned}$$

SD-FEM: Find $u_h \in V_h^k$:

$$\mathcal{J}(u_h, v_h) + B_h(u_h, v_h) = \left(f, v_h \right) + \left(f, \frac{h}{b} (\underline{b} \cdot \nabla u) \right)$$

Petrov-Galerkin interpretation: SD-FEM

\equiv testing with $\underbrace{v_h + \frac{h}{b} (\underline{b} \cdot \nabla u)}_{\text{different space than } V_h^k \text{ or } \underline{b} \cdot \nabla u \notin V_h^k}$

$$\tilde{G}_T := \frac{\delta h_T}{b}$$

$\delta > 0$

SD parameter



$$B_h^z(u_h, v_h) = \sum_T \tilde{G}_T \int_T (a \nabla u + \underline{b} \cdot \nabla u) (\underline{b} \cdot \nabla v)$$

streamline diff.
term

SD norm:

$$\|v\|_{SD} = \left(\underbrace{\bar{\alpha}}_{\bar{\alpha} = \|a\|_{\infty}} |v|_1^2 + \sum_{T \in \mathcal{T}_h} \bar{\alpha}_T \|\underline{b} \cdot \nabla v\|_{0,T}^2 \right)^{1/2}$$

Lemma: If $0 < \bar{\alpha}_T < \frac{h_T^2}{2\bar{\alpha}\mu} \quad \forall T \in \mathcal{T}_h$, then

$$\underbrace{A(v_h, v_h) + B_h(v_h, v_h)}_{=: A_{SD}^{\bar{\alpha}}(v_h, v_h)} \geq \frac{1}{2} \|v_h\|_{SD}^2 \quad \forall v_h \in V_h^k.$$

For $\mu > 0$ some constant (from inverse ineq.).

Theorem: If $\bar{\alpha}_T = \begin{cases} \delta_0 h_T & Pe_h > 1 \\ \delta_1 h_T^2 & Pe_h \leq 1 \end{cases}$

for some $\delta_0, \delta_1 > 0$ large enough, then

$$\|u - u_h\|_{SD} \leq C \left(\bar{\alpha}^{1/2} + h^{1/2} \right) h^k |u|_{k+1}$$

\uparrow
 SD-FEM sol.

$$P_e > 1$$

$\rightarrow h^{k+1/2} |u|_{k+1}$ correct rate
given SD norm dominating
term $\sim h^{1/2} \|b \cdot \nabla (u - u_h)\|_0$

$$P_e \geq 1$$

$$\rightarrow h^k |u|_{k+1}$$

\rightarrow SD-FEM has optimal rate of
convergence ∇