Connection problem for Painlevé tau functions

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Content

1 Airy equation

$$q^{\prime\prime}=tq$$

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$$q(t) = \alpha \mathrm{Ai}(t) + \beta \mathrm{Bi}(t)$$

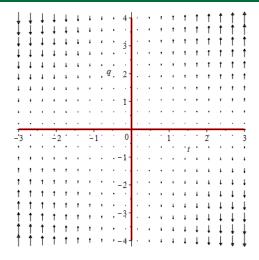
$$q'' = tq$$

$$q(t) = \alpha \operatorname{Ai}(t) + \beta \operatorname{Bi}(t)$$

$$\operatorname{Ai}(t) = \int_{e^{-\frac{i\pi}{3}}\infty}^{e^{\frac{i\pi}{3}}\infty} e^{\frac{z^3}{3} - tz} \frac{dz}{2\pi}$$

$$\operatorname{Bi}(t) = \int_{e^{-\frac{i\pi}{3}}\infty}^{e^{-\frac{i\pi}{3}}\infty} e^{\frac{z^3}{3} - tz} \frac{dz}{2\pi} + \int_{e^{-\frac{i\pi}{3}}\infty}^{e^{-\frac{i\pi}{3}}\infty} e^{\frac{z^3}{3} - tz} \frac{dz}{2\pi}$$

Force field



$$F(q,t)=tq$$

Asymptotics

$$egin{align} ext{Ai}(t) &\simeq rac{e^{-rac{2}{3}t^{rac{3}{2}}}}{2\sqrt{\pi}t^{rac{1}{4}}}, \quad t
ightarrow +\infty \ ext{Bi}(t) &\simeq rac{e^{rac{2}{3}t^{rac{3}{2}}}}{\sqrt{\pi}t^{rac{1}{4}}}, \quad t
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Asymptotics

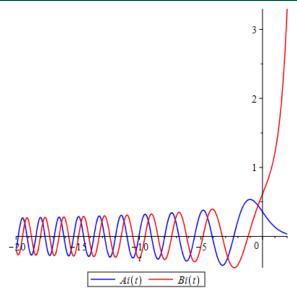
$$\operatorname{Ai}(t) \simeq \frac{e^{-\frac{2}{3}t^{\frac{3}{2}}}}{2\sqrt{\pi}t^{\frac{1}{4}}}, \quad t \to +\infty$$

$$\operatorname{Bi}(t) \simeq \frac{e^{\frac{2}{3}t^{\frac{3}{2}}}}{\sqrt{\pi}t^{\frac{1}{4}}}, \quad t \to +\infty$$

$$\operatorname{Ai}(t) \simeq \frac{1}{\sqrt{\pi}(-t)^{\frac{1}{4}}} \left(\cos\left(\frac{2}{3}(-t)^{\frac{3}{2}} - \frac{\pi}{4}\right)\right), \quad t \to -\infty$$

$$\operatorname{Bi}(t) \simeq -\frac{1}{\sqrt{\pi}(-t)^{\frac{1}{4}}} \left(\sin\left(\frac{2}{3}(-t)^{\frac{3}{2}} - \frac{\pi}{4}\right)\right), \quad t \to -\infty$$

Airy functions graph



$$H(p,q,t) = \frac{p^2}{2} - \frac{tq^2}{2}, \quad \begin{cases} \frac{dq}{dt} = \frac{\partial H}{\partial p}, \\ \frac{dp}{dt} = -\frac{\partial H}{\partial q}, \end{cases}$$

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$$H(t) = \frac{\alpha^2 + \beta^2}{2\pi} \sqrt{-t} + O(t^{-1}), \quad t \to -\infty$$

$$\begin{split} H(\rho,q,t) &= \frac{\rho^2}{2} - \frac{tq^2}{2}, \quad \left\{ \begin{array}{l} \frac{dq}{dt} &= \frac{\partial H}{\partial \rho}, \\ \frac{d\rho}{dt} &= -\frac{\partial H}{\partial q}, \\ \\ \ln \tau(t_1,t_2,\alpha,\beta) &= \int\limits_{t_1}^{t_2} H dt \\ \\ H(t) &= -\frac{\beta^2}{4\pi t} e^{\frac{4}{3}t^{\frac{3}{2}}} (1 + O(t^{-\frac{3}{2}})) + \frac{\alpha^2}{16\pi t} e^{-\frac{4}{3}t^{\frac{3}{2}}} (1 + O(t^{-\frac{3}{2}})), \quad t \to +\infty \\ \\ H(t) &= \frac{\alpha^2 + \beta^2}{2\pi} \sqrt{-t} + O(t^{-1}), \quad t \to -\infty \\ \\ \ln \tau(t_1,t_2,\alpha,\beta) &= -\frac{\beta^2}{8\pi t_2^{\frac{3}{2}}} e^{\frac{4}{3}t^{\frac{3}{2}}} (1 + O(t_2^{-\frac{3}{2}})) - \frac{\alpha^2}{32\pi t_2^{\frac{3}{2}}} e^{-\frac{4}{3}t^{\frac{3}{2}}} (1 + O(t_2^{-\frac{3}{2}})) \\ \\ &+ \frac{\alpha^2 + \beta^2}{3\pi} (-t_1)^{\frac{3}{2}} + c_0 + O((-t_1)^{-\frac{3}{2}}), \quad t_2 \to +\infty, \quad t_1 \to -\infty. \end{array} \right. \end{split}$$

$$H=\frac{p^2}{2}-\frac{tq^2}{2},$$

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$$S(t_1, t_2, \alpha, \beta) = \int_{t_1}^{t_2} \left(p\frac{dq}{dt} - H\right) dt.$$

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$$\ln \tau(t_1, t_2, \alpha, \beta) = \frac{1}{3}S(t_1, t_2, \alpha, \beta) + \frac{2tH}{3}\Big|_{t_1}^{t_2}$$

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$$\begin{split} S(t_1,t_2,\alpha,\beta) &= \int\limits_{t_1}^{t_2} \left(\rho \frac{dq}{dt} - H \right) dt. \\ \frac{\partial S}{\partial \alpha} &= \rho \frac{\partial q}{\partial \alpha} \bigg|_{t_1}^{t_2} = (\alpha \mathrm{Ai}'(t) + \beta \mathrm{Bi}'(t)) \mathrm{Ai}(t) \bigg|_{t_1}^{t_2} \end{split}$$

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$$\text{Ai}'(t) \text{Bi}(t) - \text{Ai}(t) \text{Bi}'(t)) \Big|_{t_1}^{t_2} = 0$$

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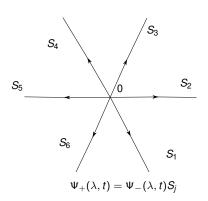
$$S(t_1,t_2,\alpha,\beta) &= \frac{\alpha^2}{2} \mathrm{Ai}'(t) \mathrm{Ai}(t) + \alpha \beta \mathrm{Ai}'(t) \mathrm{Bi}(t) + \frac{\beta^2}{2} \mathrm{Bi}'(t) \mathrm{Bi}(t) \bigg|_{t_1}^{t_2} \end{split}$$

Constant

Theorem

$$c_0 = \frac{\alpha \beta}{6\pi}$$





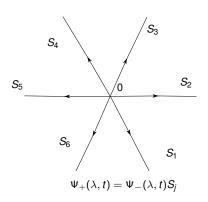
$$S_{2k} = \begin{pmatrix} 1 & 0 \\ s_{2k}e^{-\frac{z^3}{3}+tz} & 1 \end{pmatrix},$$

$$S_{2k+1} = \begin{pmatrix} 1 & s_{2k+1}e^{\frac{z^3}{3}-tz} \\ 0 & 1 \end{pmatrix}$$

$$s_{j+3} = s_j, \quad s_1 - s_2 + s_3 + s_1s_2s_3 = 0$$

$$\lim_{\lambda \to \infty} \Psi(\lambda, t) = I$$





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$$q(t) = -\lim_{\lambda \to \infty} \Psi_{12}(\lambda, t)$$

Asymptotics at $-\infty$

Its, Kapaev (1988), Kapaev (1992)

 \blacksquare special behavior for $1 - s_1 s_3 = 0$

$$q(t) \simeq \sigma \sqrt{\frac{-t}{2}} \sum_{n=0}^{\infty} b_n(-t)^{-\frac{3n}{2}} - \frac{s_1 + s_2}{\sqrt{\pi} 2^{\frac{7}{4}} (-t)^{\frac{1}{4}}} \exp\left(-\frac{2\sqrt{2}}{3} (-t)^{\frac{3}{2}}\right) (1 + O((-t)^{-\frac{1}{4}})), t \to -\infty$$

where $s_1 = -i\sigma$, $\sigma = \pm 1$.

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where $s_1 = -i\sigma$, $\sigma = \pm 1$.

lacksquare singular behavior for $1 - s_1 s_3 < 0$

$$q(t) = \frac{2\sqrt{-t}}{ae^{ig} + be^{-ig} + O((-t)^{-\frac{3}{10}})}$$

where

$$a = \frac{\sqrt{2\pi}e^{\frac{\pi\beta}{2}}}{s_1\Gamma\left(\frac{1}{2} + i\beta\right)}, \quad b = \frac{\sqrt{2\pi}e^{\frac{\pi\beta}{2}}}{s_3\Gamma\left(\frac{1}{2} - i\beta\right)}, \quad ab = 1,$$

$$g = \frac{2}{3}(-t)^{\frac{3}{2}} + \frac{3\beta}{2}\ln(-t) + 3\beta\ln 2 - \frac{\pi}{2}, \quad \beta = \frac{1}{2\pi}\ln(s_1s_3 - 1).$$

Asymptotics at $-\infty$

 \blacksquare generic behavior for $|\arg(1-s_1s_3)| < \pi$

$$q(t) = a_{0,0}^{+} e^{\frac{2i}{3}(-t)^{\frac{3}{2}}} (-t)^{\frac{3\mu}{2} - \frac{1}{4}} + a_{0,0}^{-} e^{-\frac{2i}{3}(-t)^{\frac{3}{2}}} (-t)^{-\frac{3\mu}{2} - \frac{1}{4}} + O\left(|t|^{\frac{9|\operatorname{Re}\mu|}{2} - \frac{7}{4}}\right), \quad t \to -\infty,$$

$$\mu = -\frac{\ln\left(1 - s_{1}s_{3}\right)}{2\pi i}, \quad a_{0,0}^{+} a_{0,0}^{-} = \frac{i\mu}{2},$$

$$a_{0,0}^{+} = \frac{\sqrt{\pi} 2^{3\mu} e^{-\frac{i\pi\mu}{2} - \frac{i\pi}{4}}}{s_{1}\Gamma(\mu)}, \quad a_{0,0}^{-} = \frac{\sqrt{\pi} 2^{-3\mu} e^{-\frac{i\pi\mu}{2} + \frac{i\pi}{4}}}{s_{3}\Gamma(-\mu)},$$

$$(1)$$

Asymptotics at $+\infty$

 \blacksquare special behavior for $s_2 = 0$

$$q(t) \simeq rac{i \mathbf{s}_1}{2\sqrt{\pi}t^{rac{1}{4}}} \exp\left(-rac{2}{3}t^{rac{3}{2}}
ight) (1 + O(t^{-rac{3}{4}})), \ t o -\infty.$$

Asymptotics at $+\infty$

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ight) (1 + O(t^{-rac{3}{4}})), \ t o -\infty.$$

■ singular behavior for $s_2 \in \mathbb{R}, \quad s_2 \neq 0$

$$q(t) = \frac{i\varepsilon}{\sqrt{2}} \left(\frac{ce^{ih} - 1 + O(t^{-\frac{3}{4}})}{ce^{ih} + 1 + O(t^{-\frac{3}{4}})} \right) + O(t^{-\frac{3}{2}})$$

where

$$c = \frac{\sqrt{2\pi}e^{\frac{\sigma}{2}}}{(1+s_2s_3)\Gamma\left(\frac{1}{2}+i\gamma\right)}, \quad h = \frac{2\sqrt{2}}{3}t^{\frac{3}{2}} + \frac{3\gamma}{2}\ln t + \frac{7\gamma}{2}\ln 2, \quad \gamma = \frac{1}{\pi}\ln(\varepsilon s_2), \quad \varepsilon = \mathrm{sign}s_2.$$

(2)

Asymptotics at $+\infty$

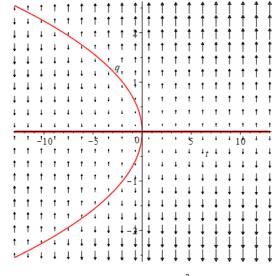
generic behavior for

$$|\arg(i\sigma s_{2})| < \frac{\pi}{2}, \quad \sigma = \operatorname{sign} \operatorname{Re}(is_{2}) = \pm 1$$

$$\sigma q(t) = i\sqrt{\frac{t}{2}} + b_{1,1}^{+} e^{\frac{2i\sqrt{2}}{3}t^{\frac{3}{2}}} t^{\frac{3\nu}{2}} - \frac{1}{4} + b_{1,1}^{-} e^{-\frac{2i\sqrt{2}}{3}t^{\frac{3}{2}}} t^{\frac{3\nu}{2}} - \frac{1}{4} + O\left(t^{3|\operatorname{Re}\nu|-1}\right), \qquad t \to + \frac{\nu}{2} + \frac{\ln\left(i\sigma s_{2}\right)}{\pi i}, \qquad b_{1,1}^{+} b_{1,1}^{-} = \frac{i\nu}{4\sqrt{2}},$$

$$b_{1,1}^{+} = \frac{\sqrt{\pi} 2^{-\frac{7\nu}{2} - \frac{3}{4}} e^{\frac{i\pi\nu}{2} - \frac{i\pi}{4}}}{(1 + \sin\sin\left(-\nu\right))}, \qquad b_{1,1}^{-} = -\frac{\sqrt{\pi} 2^{\frac{7\nu}{2} - \frac{3}{4}} e^{\frac{i\pi\nu}{2} + \frac{i\pi}{4}}}{(1 + \sin\sin\left(-\nu\right))}.$$
(3)

Painlevé II: force field for real solutions



Painlevé II: asymptotics of real solutions

The real nonsingular solutions are parametrized by number $s_1 \in i\mathbb{R}, \quad |s_1| \leq 1$. (Kapaev, 1992) The asymptotic at $+\infty$ is given by

$$q(t) = \frac{is_1}{2\sqrt{\pi}t^{\frac{1}{4}}}e^{-\frac{2}{3}t^{\frac{3}{2}}}\left(1 + O\left(\frac{1}{t^{\frac{3}{4}}}\right)\right), \quad t \to +\infty.$$

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If $|s_1| < 1$, then the asymptotics at $-\infty$ is given by Ablowitz-Segur solution

$$q(t) = \frac{d}{(-t)^{\frac{1}{4}}} \sin\left(\frac{2}{3}(-t)^{\frac{3}{2}} + \frac{3}{4}d^2\ln(-t) + \phi\right) + O\left(\frac{1}{|t|}\right), \quad t \to -\infty,$$

where

$$d = \sqrt{\frac{1}{\pi}\ln\left(1-|s_1|^2\right)}, \quad \phi = -\frac{\pi}{4} + \frac{3}{2}d^2\ln 2 - \arg\left(\Gamma\left(i\frac{d^2}{2}\right)\right) - \arg(s_1).$$

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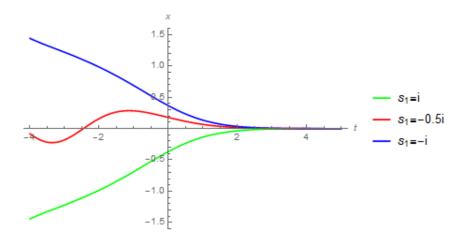
where

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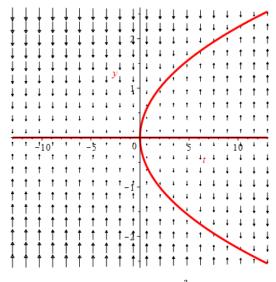
If $s_1 = \pm i$, then the asymptotics at $-\infty$ is given by Hastings-Mcleod solution

$$q(t)=is_1\sqrt{\frac{-t}{2}}+O(t^{-\frac{5}{2}}),\quad t\to-\infty.$$

Painlevé II: real solutions



Painlevé II: force field for imaginary solutions



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Painlevé II: asymptotics of imaginary solutions

All pure imaginary solutions q=iy are parametrized by number $s_1\in\mathbb{C}$. (Its, Kapaev, 1988) The asymptotic at $-\infty$ is given by

$$y(t) = \frac{d}{(-t)^{\frac{1}{4}}} \sin\left(\frac{2}{3}(-t)^{\frac{3}{2}} + \frac{3}{4}d^2\ln(-t) + \phi\right) + O\left(\frac{1}{|t|}\right), \quad t \to -\infty,$$

$$d = \sqrt{\frac{1}{\pi}\ln\left(1+|s_1|^2\right)}, \quad \phi = -\frac{\pi}{4} + \frac{3}{2}d^2\ln 2 - \arg\left(\Gamma\left(i\frac{d^2}{2}\right)\right) - \arg(s_1).$$

Painlevé II: asymptotics of imaginary solutions

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If $\operatorname{Im} s_1 \neq 0$ then the asymptotic at $+\infty$ is given by

$$\begin{split} y(t) &= \sigma \sqrt{\frac{t}{2}} + \frac{\sigma \rho}{(2t)^{\frac{1}{4}}} \cos\left(\frac{2\sqrt{2}}{3}t^{\frac{3}{2}} - \frac{3}{2}\rho^2 \ln x + \theta\right) + O\left(\frac{1}{t}\right), \quad t \to +\infty. \\ \rho &= \sqrt{\frac{1}{\pi} \ln\left(\frac{1 + |s_1|^2}{2|\mathrm{Im}(s_1)|}\right)}, \quad \sigma &= -\mathrm{sign}(\mathrm{Im}(s_1)), \\ \theta &= -\frac{3\pi}{4} - \frac{7}{2}\rho^2 \ln 2 + \mathrm{arg}(\Gamma(i\rho^2)) + \mathrm{arg}(1 + s_1^2). \end{split}$$

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$$\begin{split} y(t) &= \frac{d}{(-t)^{\frac{1}{4}}} \sin\left(\frac{2}{3}(-t)^{\frac{3}{2}} + \frac{3}{4} d^2 \ln(-t) + \phi\right) + O\left(\frac{1}{|t|}\right), \quad t \to -\infty, \\ d &= \sqrt{\frac{1}{\pi} \ln\left(1 + |s_1|^2\right)}, \quad \phi = -\frac{\pi}{4} + \frac{3}{2} d^2 \ln 2 - \arg\left(\Gamma\left(i\frac{d^2}{2}\right)\right) - \arg(s_1). \end{split}$$

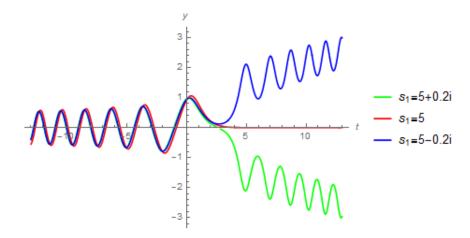
If $\operatorname{Im} s_1 \neq 0$ then the asymptotic at $+\infty$ is given by

$$\begin{split} y(t) &= \sigma \sqrt{\frac{t}{2}} + \frac{\sigma \rho}{(2t)^{\frac{1}{4}}} \cos\left(\frac{2\sqrt{2}}{3}t^{\frac{3}{2}} - \frac{3}{2}\rho^2 \ln x + \theta\right) + O\left(\frac{1}{t}\right), \quad t \to +\infty. \\ \rho &= \sqrt{\frac{1}{\pi} \ln\left(\frac{1 + |\mathbf{s}_1|^2}{2|\mathrm{Im}(\mathbf{s}_1)|}\right)}, \quad \sigma = -\mathrm{sign}(\mathrm{Im}(\mathbf{s}_1)), \\ \theta &= -\frac{3\pi}{4} - \frac{7}{2}\rho^2 \ln 2 + \mathrm{arg}(\Gamma(i\rho^2)) + \mathrm{arg}(1 + \mathbf{s}_1^2). \end{split}$$

If $\operatorname{Im} s_1 = 0$ the asymptotic at $+\infty$ is given by the Ablowitz-Segur solution,

$$y(t) = \frac{s_1}{2\sqrt{\pi}t^{\frac{1}{4}}}e^{-\frac{2}{3}t^{\frac{3}{2}}}\left(1 + O\left(\frac{1}{t^{\frac{3}{4}}}\right)\right), \quad t \to +\infty.$$

Painlevé II: imaginary solutions



$$H = \frac{p^2}{2} - \frac{tq^2}{2} - \frac{q^4}{2}$$

$$\ln \tau(t_1,t_2,s_1,s_2) \simeq \frac{t_2^3}{24} + \frac{i\sqrt{2}}{3}\nu t_2^{\frac{3}{2}} - \frac{(6\nu^2+1)}{16}\ln t_2 + \frac{2i\mu}{3}\left(-t_1\right)^{\frac{3}{2}} + \frac{3\mu^2}{4}\ln \left(-t_1\right) + \ln \Upsilon$$

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$$H(c^2p, cq, c^2t) = c^4H(q, p, t),$$

$$H = \frac{p^2}{2} - \frac{tq^2}{2} - \frac{q^4}{2}$$

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$$H(c^2p, cq, c^2t) = c^4H(q, p, t),$$

$$H = p\frac{dq}{dt} - H + \frac{d}{dt}\left(\frac{2tH}{3} - \frac{pq}{3}\right)$$

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$$H(c^2p, cq, c^2t) = c^4H(q, p, t),$$

$$H = p\frac{dq}{dt} - H + \frac{d}{dt}\left(\frac{2tH}{3} - \frac{pq}{3}\right)$$

$$\left(2tH - pq\right)^{t_2} + 2(t + t)$$

$$\ln \tau(t_1, t_2, s_1, s_2) = \left. \left(\frac{2tH}{3} - \frac{pq}{3} \right) \right|_{t_1}^{t_2} + S(t_1, t_2, s_1, s_2) \tag{4}$$

Classical action

$$\begin{split} \frac{\partial S}{\partial \mathbf{s}_1} &= p \frac{\partial q}{\partial s_1} \bigg|_{t_1}^{t_2} \\ &\frac{\partial S}{\partial \mathbf{s}_2} = p \frac{\partial q}{\partial s_2} \bigg|_{t_2}^{t_2} \\ &\ln \tau(t_1, t_2, s_1, s_2) = \left. \left(\frac{2tH}{3} - \frac{pq}{3} \right) \right|_{t_1}^{t_2} + S(t_1, t_2, s_1, s_2) \end{split}$$

Classical action

$$\begin{split} \frac{\partial S}{\partial s_1} &= \rho \frac{\partial q}{\partial s_1} \bigg|_{t_1}^{t_2} \\ \frac{\partial S}{\partial s_2} &= \rho \frac{\partial q}{\partial s_2} \bigg|_{t_2}^{t_2} \\ \ln \tau(t_1, t_2, s_1, s_2) &= \left(\frac{2tH}{3} - \frac{\rho q}{3} \right) \bigg|_{t_1}^{t_2} + S(t_1, t_2, s_1, s_2) \\ \frac{\partial}{\partial s_j} \ln \tau(t_1, t_2, s_1, s_2) &= \frac{\partial}{\partial s_j} \left(\frac{2tH}{3} - \frac{\rho q}{3} \right) \bigg|_{t_1}^{t_2} + \rho \frac{\partial q}{\partial s_j} \bigg|_{t_1}^{t_2} \,. \\ \ln \tau(t_1, t_2, s_1, s_2) &= \ln \tau(t_1, t_2, 0, -i) + \left(\left(\frac{2tH}{3} - \frac{\rho q}{3} \right) \bigg|_{t_1}^{t_2} \right) \bigg|_{(0, -i)}^{(s_1, s_2)} \\ &+ \int_{(0, -i)}^{(s_1, s_2)} \rho \frac{\partial q}{\partial s_1} \bigg|_{t_1}^{t_2} \, ds_1 + \rho \frac{\partial q}{\partial s_2} \bigg|_{t_1}^{t_2} \, ds_2 \end{split}$$

Classical action

$$\frac{\partial S}{\partial s_{1}} = \rho \frac{\partial q}{\partial s_{1}} \Big|_{t_{1}}^{t_{2}}$$

$$\frac{\partial S}{\partial s_{2}} = \rho \frac{\partial q}{\partial s_{2}} \Big|_{t_{2}}^{t_{2}}$$

$$\ln \tau(t_{1}, t_{2}, s_{1}, s_{2}) = \left(\frac{2tH}{3} - \frac{pq}{3}\right) \Big|_{t_{1}}^{t_{2}} + S(t_{1}, t_{2}, s_{1}, s_{2})$$

$$\frac{\partial}{\partial s_{j}} \ln \tau(t_{1}, t_{2}, s_{1}, s_{2}) = \frac{\partial}{\partial s_{j}} \left(\frac{2tH}{3} - \frac{pq}{3}\right) \Big|_{t_{1}}^{t_{2}} + \rho \frac{\partial q}{\partial s_{j}} \Big|_{t_{1}}^{t_{2}}.$$

$$\ln \tau(t_{1}, t_{2}, s_{1}, s_{2}) = \ln \tau(t_{1}, t_{2}, 0, -i) + \left(\left(\frac{2tH}{3} - \frac{pq}{3}\right) \Big|_{t_{1}}^{t_{2}}\right) \Big|_{(0, -i)}^{(s_{1}, s_{2})}$$

$$+ \int_{(0, -i)}^{(s_{1}, s_{2})} \rho \frac{\partial q}{\partial s_{1}} \Big|_{t_{1}}^{t_{2}} ds_{1} + \rho \frac{\partial q}{\partial s_{2}} \Big|_{t_{1}}^{t_{2}} ds_{2}$$

$$\ln \tau(t_{1}, t_{2}, 0, -i) \sim \frac{t_{2}^{3}}{24} - \frac{1}{16} \ln t_{2} + \ln \Upsilon_{0}$$
(5)

Numerical constant

$$\ln au(t_1,t_2,-i,0) \sim -rac{t_1^3}{24} - rac{1}{16} \ln(-t_1) + \ln \Upsilon_{HM}, \quad t_1 o -\infty, \quad t_2 o +\infty$$

Numerical constant

$$\ln \tau(t_1, t_2, -i, 0) \sim -\frac{t_1^3}{24} - \frac{1}{16} \ln(-t_1) + \ln \Upsilon_{HM}, \quad t_1 \to -\infty, \quad t_2 \to +\infty$$

$$\Upsilon_{HM} = 2^{-\frac{1}{48}} e^{-\zeta'(-1)}.$$

Baik, Buckingham, DiFranco (2008) and Deift, Its, Krasovsky (2008)

Numerical constant

$$\ln \tau(t_1, t_2, -i, 0) \sim -\frac{t_1^3}{24} - \frac{1}{16} \ln(-t_1) + \ln \Upsilon_{HM}, \quad t_1 \to -\infty, \quad t_2 \to +\infty$$

$$\Upsilon_{HM} = 2^{-\frac{1}{48}} e^{-\zeta'(-1)}.$$

Baik, Buckingham, DiFranco (2008) and Deift, Its, Krasovsky (2008)

$$q(t;0,-i) = e^{\frac{2\pi i}{3}} q\left(te^{\frac{2\pi i}{3}};-i,0\right).$$

Results

Theorem

Its, Lisovyv, P., 2018

$$\Upsilon = \Upsilon_0 2^{\frac{3}{2}\mu^2 - \frac{7\nu^2}{8}} (2\pi)^{-\frac{\mu}{2} - \frac{\nu}{4}} e^{\frac{\pi i}{8} \left(\eta^2 + 2\mu^2 + 2\eta\nu - 8\mu\eta\right)} \frac{\sqrt{G(1-\nu)\,\hat{G}(\eta)}}{G(1-\mu)\,\hat{G}\left(\frac{\eta-\nu}{2}\right)} \ ,$$

$$\Upsilon_0 = 2^{\frac{1}{48}} e^{\frac{\zeta'(-1)}{2} + \frac{i\pi}{48}},$$

where $\zeta(z)$ - Riemann Zeta function, G(z) - Barnes G-function, $\hat{G}(z) = \frac{G(1+z)}{G(1-z)}$, and $\mu, \nu, \sigma, \eta, \Upsilon, \Upsilon_0$ are described by (1), (2), (3),(4),(5) and (6)

$$s_3^{-1} = e^{i\pi\eta} e^{i\pi\frac{\sigma}{2}}.$$

(6)

Last slide

THANK YOU