OPTICAL PUMPING

Jung Lin (Doris) Lee

ABSTRACT

In this lab, we examined whether

1. INTRODUCTION

Optical pumping is a experimental technique that enable us to measure the —- splittling levels of Rb atoms in a magnetic field and —- with degenerate perturbation theory. energetically excited to a higher energy state and spontaneous emmision back down to the — state. Due to the energy splitting of —-

2. THEORY

In degenerate perturbation theory of quantum mechanics, we can find the first order corrections to spin orbit coupling and relativistic correction, which together yields the fine structure correction to the eigenstate energies. In addition, there is hyperfine splitting due to interactions of the nucleus of the Rb atom which generates its own electromagnetic fields that gives rise to the perturbation Wikipedia (2016). Additionally, in this experiment, we also apply an external magnetic field (in addition to the Earth's magnetic field) which results in Zeeman splitting in the case of the weak field. In Fig. 1, we can see that the fine structure, hyperfine, and Zeeman interactions splits the energy levels into j , f and m_f levels respectively.
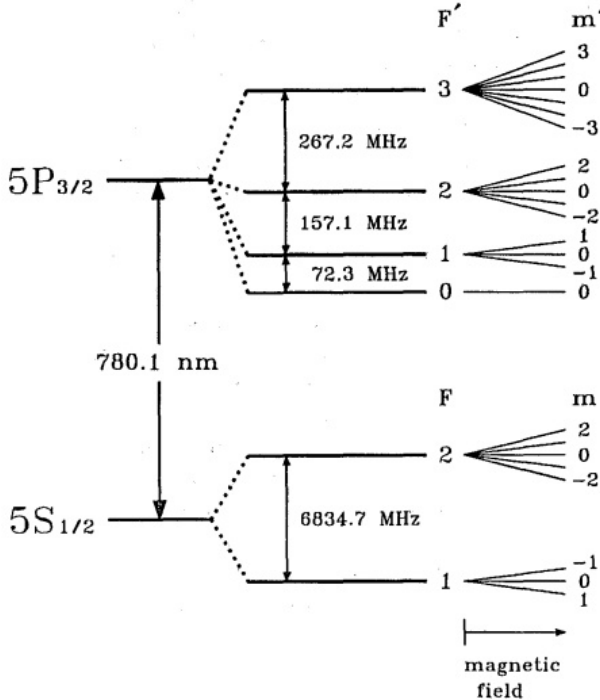
Figure 1. Energy level splitting of a Rubidium atom. (Source: The Optical Society)

$$B_H = 0.9 \times 10^{-6} \frac{T \cdot m}{A} \frac{Ni}{a} \qquad (1)$$

3. APPARATUS AND PROCEDURE

We measure the minimum and maximum current value that still yields a symmetric Lissajous curve and make this the error range of our current measurement, since within this range, we can not determine what is the best value corresponding to the resonant frequency. This error is propagated through our subsequent analysis in Sec. 4 and to get an error estimate on the final computed value of Earth's magnetic field.

4. ANALYSIS

Determining Nuclear Moments of Observed Species— By taking the ratio of the Breit-Rabi equation for the two species:

$$\frac{\frac{\nu_1}{B_1} = \frac{2.799}{2I_1+1}}{\frac{\nu_2}{B_2} = \frac{2.799}{2I_2+1}} = \frac{2I_2+1}{2I_1+1}$$

$$\frac{\nu_1}{\nu_2} = \frac{B_1}{B_2}\left(\frac{2I_2+1}{2I_1+1}\right)$$

We chosing the "best" value of $\frac{\nu_1}{\nu_2}$ as one and this yielded a 2I+1 ratio of about 2/3. We deduced that if Species 1 was Rb-85 (I=5/2) and Species 2 was Rb-87 (I = 3/2) (Millman and Fox 1936), then we would get the appropriate 2I+1 ratio. The errors on these deduced values are accurate to 3 significant figures, because we know that nuclear moments must be half-integer values, so the only uncertainty is on the numerical coefficient 2.799.

Knowing the I_1 and I_2 values, we can use the Breit-Rabi equation to compute the magnetic field and compare this with the one that we obtain from Eq.1. For species 1, the mean square difference between the two measures of the magnetic field is 0.0915 and for species 2 the mean squared difference is 0.0973. This value gives us an estimate of the systematic error of the possible non-uniformity of the Helmholtz coil number and radius. Using this technique, we were able to obtain a more accurate estimate of the numerical coefficient in Eq.1 of 9.963×10^{-7}.

Data Fitting and Transformation— We perform a linear regression on the data using the model obtained from rearranging the Breit-Rabi equation into the linear form (y = ax+b), where the frequency is the dependent variable and the magnetic field strength is the independent variable.

In order to decrease the error on the parameters for linear regression , we flipped the negative data along the x axis so that we could perform two linear fits so that we have double the size of the sample. The error estimates on the parameters can be obtained from the derivation of the maximum log-likelihood estimator on a linear regression model (Press 1988):

$$\sigma_a^2 = S_{xx}/\triangle \quad \sigma_b^2 = S/\triangle \qquad (2)$$

where $\triangle = SS_{xx} - (S_x)^2$, $S_{xx} = \sum_{i=1}^{N} \frac{x_i^2}{\sigma_i^2}$, $S_x = \sum_{i=1}^{N} \frac{x_i}{\sigma_i^2}$, $S = \sum_{i=1}^{N} \frac{1}{\sigma_i}$. With this error estimate, we summarize
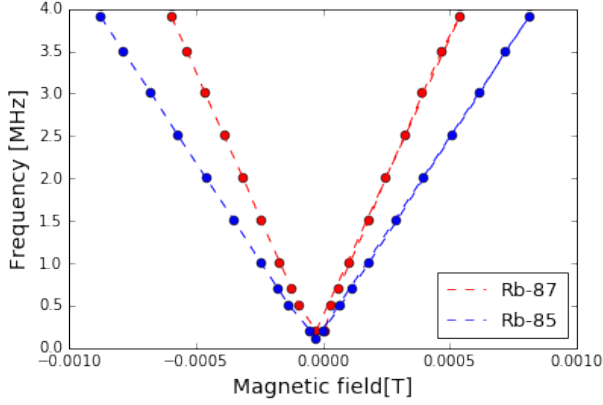
Figure 2. Experimental data of magnetic field strength of Helmholtz coil versus resonant frequency before data transformation.

Table 1
Fitting coefficients on the experimental data, where a is the slope and b is the y-intercept.

Species	a	b	σ_a	σ_b
Rb-85	4644.29	0.1342	1.091×10^{-6}	0.231
Rb-87	6983.93	0.2069	1.201×10^{-6}	0.573

the fitting coefficients and their respective errors in Table. 1. We obtained a chi squared value of 3.571 for Rb-85 and 3.724 for Rb-87. The chi squared goodness-to-fit test shows that the experimental values are very close to the modelled values (p¡0.01) and with the values of nuclear spins found in our analysis.

Estimating Earth's magnetic field— At the zero current point, the magnetic field is zero, so there is no magnetic field contribution from the Helmholtz coil so $B_{total} =$ B_{earth}. This could also be thought of as the y intercept of the linear regression:

$$B_H = \frac{(2I+1)}{2.799}\nu - B_e$$

. From the two species, we obtained two estimates of the total magnetic field as 59.28μ T for species 1 and 91.428μ T. The actual magnetic field strength in Berkeley is $48.6\mu T$, according to Wolfram Alpha. This discrepancy can not be completely accounted for by the error on our recorded measurement. Possible sources error may include instrumentation systematics, non-uniformity of the magnetic field and electronics reading noise.
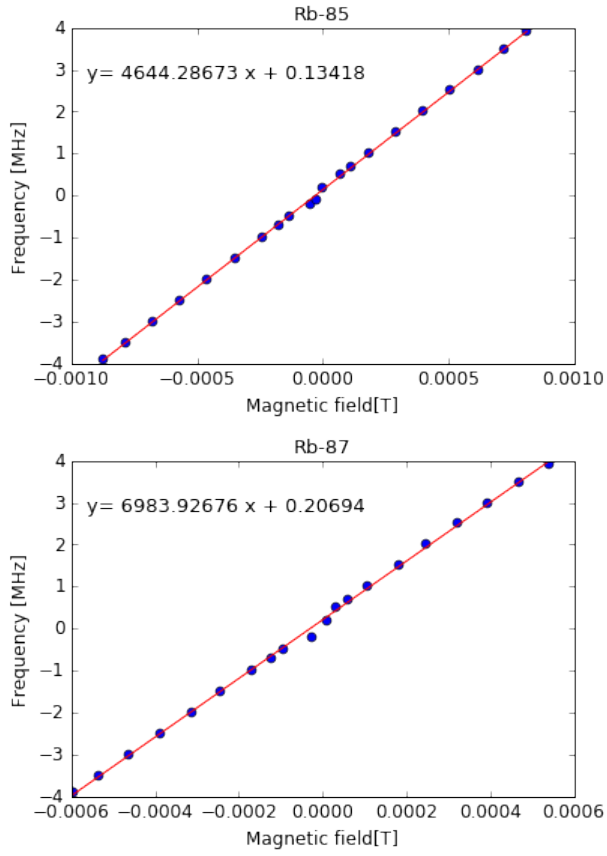
5. CONCLUSION

REFERENCES

W. H. Press, *Numerical recipes in C* (Cambridge University Press, 1988).
S. Millman and M. Fox, Phys. Rev. **50**, 220 (1936), URL http://link.aps.org/doi/10.1103/PhysRev.50.220.
Wikipedia, *Hyperfine structure* (2016).

Figure 3. This is the linear regression on the experimental data for Rb-85 and Rb-87. The error bar is too small compared to the pixels spanned by the datapoint, so it can not be seen on this plot.