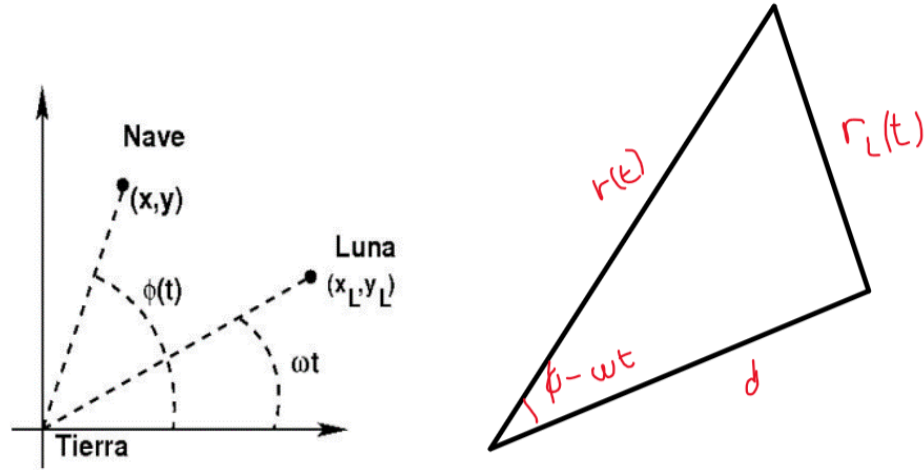


c) Se tiene que



Donde la distancia r_L se puede encontrar mediante la ley de cosenos

$$r_L(r, \phi, t)^2 = r(t)^2 + d^2 - 2r(t)d \cos(\phi - \omega t)$$

$$r_L(r, \phi, t) = \sqrt{r(t)^2 + d^2 - 2r(t)d \cos(\phi - \omega t)}$$

d)

Se tiene que la posición de la nave es

$$\vec{r} = (x, y) = (r \cos \phi, r \sin \phi)$$

$$\dot{\vec{r}} = (\dot{r} \cos \phi - r \dot{\phi} \sin \phi, \dot{r} \sin \phi + r \dot{\phi} \cos \phi)$$

$$|\dot{\vec{r}}|^2 = (\dot{r} \cos \phi - r \dot{\phi} \sin \phi)^2 + (\dot{r} \sin \phi + r \dot{\phi} \cos \phi)^2$$

$$|\dot{\vec{r}}|^2 = \dot{r}^2 \cos^2 \phi - 2\dot{r}r\dot{\phi} \cos \phi \sin \phi + r^2\dot{\phi}^2 \sin^2 \phi + \dot{r}^2 \sin^2 \phi + 2\dot{r}r\dot{\phi} \sin \phi \cos \phi + r^2\dot{\phi}^2 \cos^2 \phi$$

$$|\dot{\vec{r}}|^2 = \dot{r}^2(\cos^2 \phi + \sin^2 \phi) + r^2\dot{\phi}^2(\cos^2 \phi + \sin^2 \phi)$$

$$|\dot{\vec{r}}|^2 = \dot{r}^2 + r^2\dot{\phi}^2$$

Entonces la energía cinética es de la forma

$$T = \frac{1}{2} m |\dot{\vec{r}}|^2 = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\phi}^2)$$

Mientras la energía potencial es

$$U = -G \left(\frac{mm_T}{r} + \frac{mm_L}{r_L} \right)$$

$$\mathcal{L} = \frac{1}{2}m(\dot{r}^2 + r^2\dot{\phi}^2) + G\left(\frac{mm_T}{r} + \frac{mm_L}{r_L}\right)$$

De igual forma los momentos generalizados serán

$$\frac{\partial \mathcal{L}}{\partial \dot{r}} = m\dot{r} = p_r$$

$$\frac{\partial \mathcal{L}}{\partial \dot{\phi}} = m\dot{\phi}r^2 = p_\phi$$

Además, el Hamiltoniano es definido como $H = \sum p_{q_i} \dot{q}_i - \mathcal{L}$

$$H = m\dot{r}^2 + m\dot{\phi}^2 r^2 - \frac{1}{2}m(\dot{r}^2 + r^2\dot{\phi}^2) - G\left(\frac{mm_T}{r} + \frac{mm_L}{r_L}\right)$$

$$H = \frac{1}{2}m\dot{r}^2 + \frac{1}{2}m\dot{\phi}^2 r^2 - G\left(\frac{mm_T}{r} + \frac{mm_L}{r_L}\right)$$

Veamos que con las definiciones de los momentos generalizados de $m\dot{r} = p_r$ y $m\dot{\phi}r^2 = p_\phi$ se puede reescribir como

$$H = \frac{p_r^2}{2m} + \frac{p_\phi^2}{2mr^2} - G\left(\frac{mm_T}{r} + \frac{mm_L}{r_L}\right)$$

e) Las ecuaciones de movimiento entonces son

$$\frac{\partial H}{\partial p_r} = \dot{r} = \frac{\partial}{\partial p_r} \left(\frac{p_r^2}{2m} + \frac{p_\phi^2}{2mr^2} - G\left(\frac{mm_T}{r} + \frac{mm_L}{r_L}\right) \right) = \frac{p_r}{m}$$

$$\frac{\partial H}{\partial p_\phi} = \dot{\phi} = \frac{\partial}{\partial p_\phi} \left(\frac{p_r^2}{2m} + \frac{p_\phi^2}{2mr^2} - G\left(\frac{mm_T}{r} + \frac{mm_L}{r_L}\right) \right) = \frac{p_\phi}{mr^2}$$

$$\dot{p}_r = -\frac{\partial H}{\partial r} = \frac{p_\phi^2}{mr^3} - G\frac{mm_T}{r^2} + G\frac{\partial}{\partial r} \left(\frac{mm_L}{\sqrt{r^2 + d^2 - 2rd \cos(\phi - \omega t)}} \right)$$

$$\begin{aligned} \frac{\partial}{\partial r} \left(\frac{mm_L}{\sqrt{r^2 + d^2 - 2rd \cos(\phi - \omega t)}} \right) &= -\frac{1}{2} \frac{mm_L [2r - 2d \cos(\phi - \omega t)]}{[r^2 + d^2 - 2rd \cos(\phi - \omega t)]^{\frac{3}{2}}} \\ &= -\frac{mm_L}{r_L^3} [r - d \cos(\phi - \omega t)] \end{aligned}$$

$$\dot{p}_r = \frac{p_\phi^2}{mr^3} - G\frac{mm_T}{r^2} - G\frac{mm_L}{r_L^3} [r - d \cos(\phi - \omega t)]$$

$$\begin{aligned}
\dot{p}_\phi &= -\frac{\partial H}{\partial \phi} = G \frac{\partial}{\partial \phi} \left(\frac{mm_L}{\sqrt{r^2 + d^2 - 2rd \cos(\phi - \omega t)}} \right) \\
&= G \left[-\frac{1}{2} \frac{mm_L 2rd \sin(\phi - \omega t)}{[r^2 + d^2 - 2rd \cos(\phi - \omega t)]^3} \right] \\
\dot{p}_\phi &= -G \frac{mm_L}{r_L^3} rd \sin(\phi - \omega t)
\end{aligned}$$

f) definiendo $\tilde{r} = \frac{r}{d}$, $\tilde{p}_r = \frac{p_r}{md}$, $\tilde{p}_\phi = \frac{p_\phi}{md^2}$

Dividiendo la primera ecuación de movimiento por d se obtiene

$$\begin{aligned}
\frac{\dot{r}}{d} &= \frac{p_r}{md} \rightarrow \dot{\tilde{r}} = \tilde{p}_r \\
\dot{\phi} &= \frac{p_\phi}{mr^2} = \tilde{p}_\phi \frac{d^2}{r^2} = \frac{\tilde{p}_\phi}{\tilde{r}^2} \\
\frac{\dot{p}_r}{md} &= \frac{p_\phi^2}{mr^3} \frac{1}{md} - \frac{1}{md} G \frac{mm_T}{r^2} - \frac{1}{md} G \frac{mm_L}{r_L^3} [r - d \cos(\phi - \omega t)] \\
\dot{\tilde{p}}_r &= \frac{\tilde{p}_\phi^2}{r^3} d^3 - \frac{d^2}{d^2} \frac{1}{d} G \frac{m_T}{r^2} - \frac{1}{md} G \frac{mm_L}{[r^2 + d^2 - 2rd \cos(\phi - \omega t)]^{\frac{3}{2}}} [r - d \cos(\phi - \omega t)] \\
\dot{\tilde{p}}_r &= \frac{\tilde{p}_\phi^2}{\tilde{r}^3} - G \frac{m_T}{\tilde{r}^2 d^3} - G \frac{m_L}{[r^2 + d^2 - 2rd \cos(\phi - \omega t)]^{\frac{3}{2}}} [\tilde{r} - \cos(\phi - \omega t)] \\
\dot{\tilde{p}}_r &= \frac{\tilde{p}_\phi^2}{\tilde{r}^3} - G \frac{m_T}{\tilde{r}^2 d^3} - G \frac{\frac{m_L}{d^3}}{\frac{[r^2 + d^2 - 2rd \cos(\phi - \omega t)]^{\frac{3}{2}}}{d^3}} [\tilde{r} - \cos(\phi - \omega t)] \\
\dot{\tilde{p}}_r &= \frac{\tilde{p}_\phi^2}{\tilde{r}^3} - G \frac{m_T}{\tilde{r}^2 d^3} - G \frac{\frac{m_L}{d^3} m_T [\tilde{r} - \cos(\phi - \omega t)]}{\left(\left(\frac{r^2 + d^2 - 2rd \cos(\phi - \omega t)}{d^2} \right)^{\frac{1}{2}} \right)^3 m_T} \\
\dot{\tilde{p}}_r &= \frac{\tilde{p}_\phi^2}{\tilde{r}^3} - G \frac{m_T}{\tilde{r}^2 d^3} - G \frac{\frac{m_L}{d^3} m_T [\tilde{r} - \cos(\phi - \omega t)]}{\left((\tilde{r}^2 + 1 - 2\tilde{r} \cos(\phi - \omega t))^{\frac{1}{2}} \right)^3 m_T}
\end{aligned}$$

Con

$$\Delta \equiv \frac{Gm_T}{d^3}; \quad \mu \equiv \frac{m_L}{m_T}; \quad \tilde{r}' = \sqrt{1 + \tilde{r}^2 - 2\tilde{r} \cos(\phi - \omega t)}$$

$$\dot{\tilde{p}}_r = \frac{\tilde{p}_\phi^2}{\tilde{r}^3} - \frac{\Delta}{\tilde{r}^2} - \frac{\Delta\mu[\tilde{r} - \cos(\phi - \omega t)]}{\tilde{r}'^3}$$

$$\dot{\tilde{p}}_r = \frac{\tilde{p}_\phi^2}{\tilde{r}^3} - \Delta \left\{ \frac{1}{\tilde{r}^2} + \frac{\mu}{\tilde{r}'^3} [\tilde{r} - \cos(\phi - \omega t)] \right\}$$

$$\frac{\dot{p}_\phi}{md^2} = -G \frac{m_L}{r_L^3 d^2} r d \sin(\phi - \omega t)$$

$$\dot{\tilde{p}}_\phi = -\left(\frac{m_T}{m_T}\right) G \frac{m_L}{r_L^3} \tilde{r} \sin(\phi - \omega t) = -G \frac{\mu m_T}{r_L^3} \tilde{r} \sin(\phi - \omega t)$$

$$\dot{\tilde{p}}_\phi = -G \frac{\frac{\mu m_T}{d^3}}{\left(\left(\frac{r^2 + d^2 - 2rd \cos(\phi - \omega t)}{d^2} \right)^{\frac{1}{2}} \right)^3} \tilde{r} \sin(\phi - \omega t)$$

$$\dot{\tilde{p}}_\phi = -\frac{\Delta\mu\tilde{r}}{\left((\tilde{r}^2 + 1 - 2\tilde{r} \cos(\phi - \omega t))^{\frac{1}{2}} \right)^3} \sin(\phi - \omega t)$$

$$\dot{\tilde{p}}_\phi = -\frac{\Delta\mu\tilde{r}}{\tilde{r}'^3} \sin(\phi - \omega t)$$

g) Finalmente para asignar los momentos canónicos originales se demuestra que

Como se sabe p_r no es mas que el momento lineal, entonces

$$r_0 = (r_0 \cos \phi, r_0 \sin \phi); \dot{r}_0 = (v_0 \cos \theta, v_0 \sin \theta)$$

$$p_r = m \frac{dr}{dt} = m \frac{d}{dt} \sqrt{x^2 + y^2}$$

$$\tilde{p}_r^0 = \frac{p_r^0}{md} = \frac{m}{md} \frac{dr_0}{dt} = \frac{1}{d} \frac{d}{dt} \left(\sqrt{x_0^2 + y_0^2} \right) = \frac{2x_0\dot{x}_0 + 2y_0\dot{y}_0}{2d\sqrt{x_0^2 + y_0^2}} = \frac{x_0\dot{x}_0 + y_0\dot{y}_0}{dr_0}$$

$$\tilde{p}_r^0 = \frac{r_0 \cos \phi v_0 \cos \theta + r_0 \sin \phi v_0 \sin \theta}{dr_0} = \frac{\cos \phi v_0 \cos \theta + v_0 \sin \phi \sin \theta}{d}$$

$$= \tilde{v}_0 (\cos \phi v_0 \cos \theta + \sin \phi \sin \theta)$$

$$\tilde{p}_r^0 = \tilde{v}_0 \cos(\theta - \phi)$$

Como $\arctan\left(\frac{y}{x}\right) = \phi$ y p_ϕ no es más que el momento angular, que sigue entonces la forma $p_\phi = \frac{d\phi}{dt} I$. Para una partícula $I = mr^2$

$$\tilde{p}_\phi^0 = \frac{p_\phi^0}{md^2} = \frac{1}{md^2} mr^2 \frac{d\phi}{dt} = \tilde{r}^2 \frac{d}{dt} \left(\arctan \left(\frac{y}{x} \right) \right) = \tilde{r}^2 \frac{1}{\left(\frac{y}{x} \right)^2 + 1} \frac{d}{dt} \left(\frac{y}{x} \right)$$

$$= \tilde{r}^2 \frac{1}{\left(\frac{y}{x} \right)^2 + 1} \frac{\dot{y}x - \dot{x}y}{x^2} = \tilde{r}^2 \frac{1}{y^2 + x^2} (\dot{y}x - \dot{x}y) = \frac{\tilde{r}^2}{r^2} (\dot{y}x - \dot{x}y)$$

$$\frac{1}{d^2} (v_0 r_0 \sin \theta \cos \phi - v_0 r_0 \cos \theta \sin \phi)$$

$$\tilde{p}_\phi^0 = \tilde{v}_0 \tilde{r}_0 (\sin \theta \cos \phi - \cos \theta \sin \phi) = \tilde{v}_0 \tilde{r}_0 \sin(\theta - \phi)$$