

ATLAS Experiment
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Particle Physics An Introduction

Module 3: Accelerators and detectors
Part 3.5: Light particles in matter

In this module, we will touch the basics of particle acceleration and detection methods.

In this video we show how light charged particles interact with matter. They are detected via these interactions.

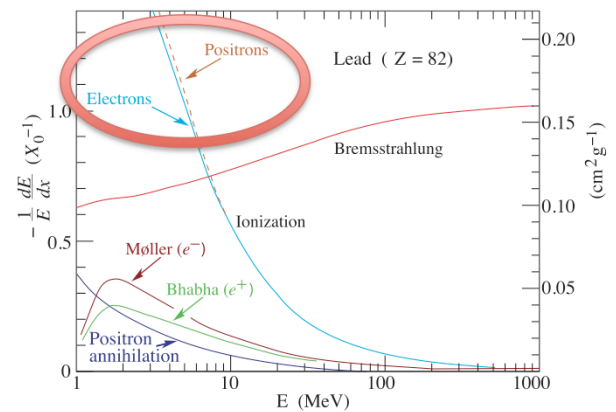
At the end of this video you will know:

- The energy loss by ionization and excitation for electrons
- The energy loss by radiation specific to light particles, called bremsstrahlung.

Bethe-Bloch formula for electrons and positrons:

$$-\frac{dE}{\rho dx}\bigg|_{ion} = K \frac{1}{\beta^2} \frac{Z}{A} \left[\ln \left(\frac{m_e \tau \sqrt{\tau + 2}}{\sqrt{2} I} \right) + \frac{F(\tau)}{2} - \frac{\delta}{2} \right]$$

- τ : kinetic energy of the electron (positron) in units of m_e
- $F(\tau)$: function differentiating between electron and positron



For **light particles** like the electron and positron, things complicate slightly:

- The formula for dE/dx is similar to that for heavy particles, but the special properties of scattering between electrons (**Møller scattering**) and between electrons and positrons (**Bhabha scattering**) come in.
- The difference between the dE/dx for electrons and positrons is, however, not very large.

Energy loss by radiation in terms of radiation length X_0 :

$$-\frac{dE}{dx}\bigg|_{brem} = \frac{E}{X_0}$$

- $E(x) = E_0 \exp(-x/X_0)$: remaining energy
- $x = X_0$: particle energy reduced by a factor e through bremsstrahlung
- $X_0[\text{cm}] = \rho X_0[\text{g/cm}^2]$: radiation length in terms of surface density

But there is an important difference between light and heavy particles:

- Since light projectiles and atomic electrons are of comparable mass, interactions with atomic electrons cause **bremsstrahlung**.
- The energy loss dE/dx by this process is proportional to the energy of the projectile. The material properties that come into play can be combined into a single constant, the **radiation length X_0** .

By integrating the energy loss we can calculate the **remaining energy $E(x)$** of the electron. It is a function which **decreases exponentially** with the penetrated material depth x .

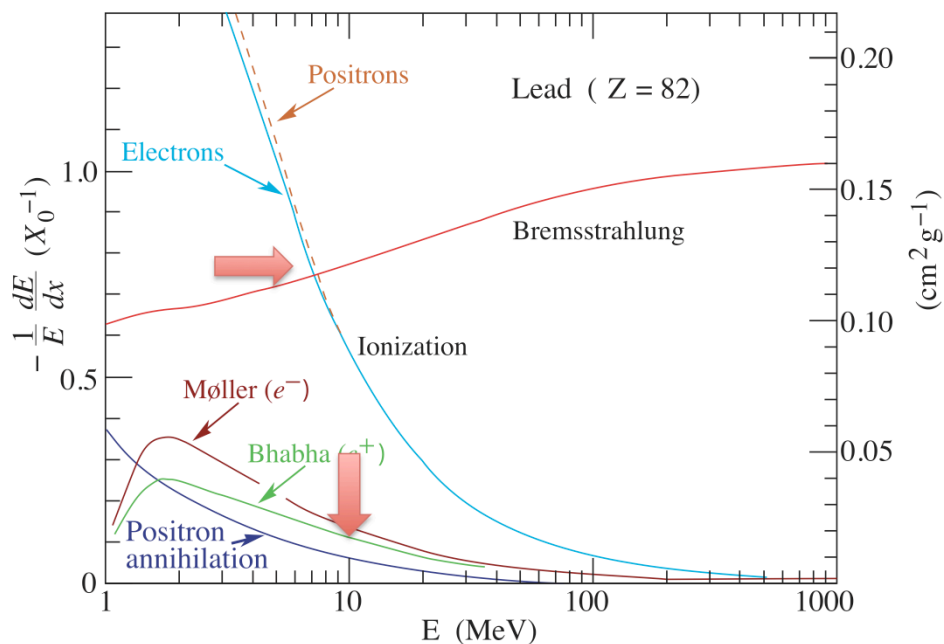
The radiation length is then given as the path length after which **the energy of the projectile is reduced by a factor of e** .

To obtain the energy loss dE/dx or the radiation length in terms of the target **surface density**, one must take into account the volume density ρ of the material.

Material	Z	A	$\langle Z/A \rangle$	Nucl.coll. length λ_T {g cm ⁻² }	Nucl.inte length λ {g cm ⁻² }	Rad.len. X_0 {g cm ⁻² }	$dE/dx _{\min}$ { MeV g ⁻¹ cm ² }	Density {g cm ⁻³ } {gℓ ⁻¹ }	Melting point (K)	Boiling point (K)	Refract. index @ Na D
H ₂	1	1.008(7)	0.99212	42.8	52.0	63.04	(4.103)	0.071(0.084)	3.81	20.28	1.11[132.]
D ₂	1	2.01410177803(8)	0.49650	51.3	71.8	125.97	(2.053)	0.169(0.168)	18.7	23.65	1.11[138.]
He	2	4.002602(2)	0.49967	51.8	71.0	94.32	(1.937)	0.125(0.166)		4.220	1.02[35.0]
Li	3	6.94(2)	0.43221	52.2	71.3	82.78	1.639	0.534	453.6	1615.	
Be	4	9.0121831(5)	0.44384	55.3	77.8	65.19	1.595	1.848	560.	2744.	
C diamond	6	12.0107(8)	0.49955	59.2	85.8	42.70	1.725	3.520			2.42
C graphite	6	12.0107(8)	0.49955	59.2	85.8	42.70	1.742	2.210			
N ₂	7	14.007(2)	0.49976	61.1	89.7	37.99	(1.825)	0.807(1.165)	63.15	77.29	1.20[298.]
O ₂	8	15.999(3)	0.50002	61.3	90.2	34.24	(1.801)	1.141(1.332)	54.36	90.20	1.22[271.]
F ₂	9	18.998403163(6)	0.47372	65.0	97.4	32.93	(1.676)	1.507(1.580)	53.53	85.03	[195.]
Ne	10	20.1797(6)	0.49555	65.7	99.0	28.93	(1.724)	1.204(0.839)	54.56	27.07	1.09[67.1]
Al	13	26.9815385(7)	0.48181	69.7	107.2	24.01	1.615	2.699	933.5	2792.	
Si	14	28.0855(3)	0.49848	70.2	108.4	21.82	1.664	2.329	687.	3538.	3.95
Cl ₂	17	35.453(2)	0.47951	73.8	115.7	19.28	(1.630)	1.574(2.980)	71.6	239.1	[773.]
Ar	18	39.948(1)	0.45059	75.7	119.7	19.55	(1.519)	1.396(1.662)	83.81	87.26	1.23[281.]
Ti	22	47.867(1)	0.45961	78.8	126.2	16.16	1.477	4.540	941.	3560.	
Fe	26	55.845(2)	0.46557	81.7	132.1	13.84	1.451	7.874	811.	3134.	
Cu	29	63.546(3)	0.45636	84.2	137.3	12.86	1.403	8.960	1358.	2835.	
Ge	32	72.630(1)	0.44053	86.9	143.0	12.25	1.370	5.323	211.	3106.	
Sn	50	118.710(7)	0.42119	98.2	166.7	8.82	1.263	7.310	405.1	2875.	
Xe	54	131.293(6)	0.41129	100.8	172.1	8.48	(1.255)	2.953(5.483)	61.4	165.1	1.39[701.]
W	74	183.84(1)	0.40252	110.4	191.9	6.76	1.145	19.300	3695.	5828.	
Pt	78	195.084(9)	0.39983	112.2	195.7	6.54	1.128	21.450	2042.	4098.	
Au	79	196.966569(5)	0.40108	112.5	196.3	6.46	1.134	19.320	337.	3129.	
Pb	82	207.2(1)	0.39575	114.1	199.6	6.37	1.122	11.350	600.6	2022.	
U	92	[238.02891(3)]	0.38651	118.6	209.0	6.00	1.081	18.950	908.	4404.	
Air (dry, 1 atm)			0.49919	61.3	90.1	36.62	(1.815)	(1.205)		78.80	[289]
Shielding concrete			0.50274	65.1	97.5	26.57	1.711	2.300			
Borosilicate glass (Pyrex)			0.49707	64.6	96.5	28.17	1.696	2.230			
Lead glass			0.42101	95.9	158.0	7.87	1.255	6.220			
Standard rock			0.50000	66.8	101.3	26.54	1.688	2.650			

<http://www-pdg.lbl.gov/2016/reviews/rpp2016-rev-atomic-nuclear-prop.pdf>

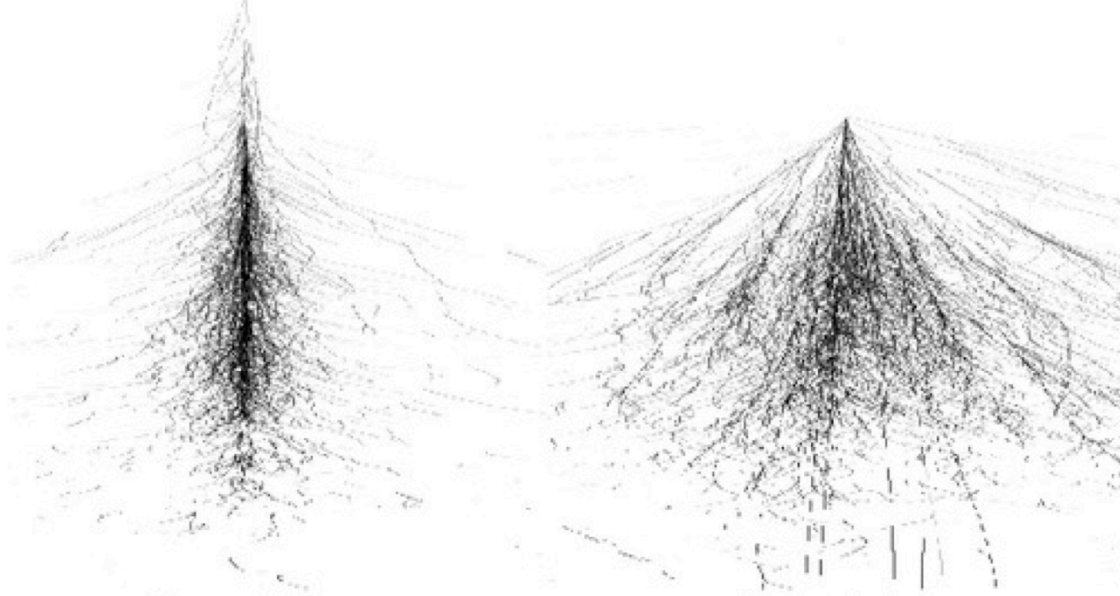
Many material properties are tabulated by the Particle Data Group, particularly the **volume density** and the **radiation length**. You will find them on the Web at the indicated address.



<http://www-pdg.lbl.gov/2016/reviews/rpp2016-rev-passage-particles-matter.pdf>

A property specific to positrons is that they can annihilate with atomic electrons of the material. This process is important at low energy, i.e. at the **end of the path** just before the positron stops. From annihilation at rest, two photons are emitted back-to-back, each with an energy of $E = m_e = 511 \text{ keV}$.

At energies above about 10 MeV, **electrons and positrons** lose energy mainly by bremsstrahlung, which dominates over ionization and excitation. This means that many photons are produced when a high energy electrons goes through matter. Those are mostly converted into electron-positron pairs, as discussed in the next video. These pairs will emit photons in turn.



Electron/photon shower

Hadronic shower

In this way, high energy electrons generate **electron/photon showers** in matter, which are populated by a large number of particles.

The strong interaction of hadrons with the nuclei of the material can also cause showers, called **hadronic showers**. These are typically less populated at comparable energy, because the lightest hadron, the pion, has a mass of about hundred MeV. For the same reason, they are wider and less regular than electromagnetic showers.

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- Again, a table provided by the **Particle Data Group** summarizes the properties of materials often used in particle detectors.
- The characteristic lengths for nuclear interactions are typically **comparable to the radiation length for lighter materials, but much larger for heavy materials.**

In the next video, we will discuss in more detail how photons interact with matter.