

# MEK 4300

## Mandatory Assignment

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### Taylor-Green Vortex

#### Physical problem

The Taylor-Green vortex is a unsteady flow where we observe the flow of a decaying vortex. This flow has an exact solution for the incompressible Navier-Stokes equation in 2D, while for the 3D case there are several numerical results for comparison. Assuming the fluid is incompressible, the exact solution for velocity and pressure are given as ( $t = 0$ )

$$u(x, y, t) = (V_0 \cos(\pi x) \sin(\pi y) e^{-2t\nu\pi^2}, \quad V_0 \cos(\pi y) \sin(\pi x) e^{-2t\nu\pi^2}) \quad (1)$$

$$p(x, y, t) = -0.25(\cos(2\pi x) + \cos(2\pi y)) e^{-4t\nu\pi^2} \quad (2)$$

For the 3D case we will consider the kinetic energy  $E_k$  and kinetic energy dissipation rate  $\epsilon$  for the system. These quantities are explored thoroughly by other authors, and will be used as comparison. The initial field set up in the 3D Taylor-Green vortex is defined as

$$u(x, y, z) = (V_0 \sin(\frac{x}{L}) \cos(\frac{y}{L}) \cos(\frac{z}{L}), \quad -V_0 \cos(\frac{x}{L}) \sin(\frac{y}{L}) \cos(\frac{z}{L}), \quad 0) \quad (3)$$

$$p(x, y, z) = \rho_0 + \frac{\rho_0 V_0^2}{16} (\cos(\frac{2x}{L}) + \cos(\frac{2y}{L})) (\cos(\frac{2z}{L}) + 2) \quad (4)$$

Exploring the incompressible flow condition we define  $\rho_0 = \rho$ , and for simplicity we let  $V_0 = 1$

#### Governing Equation and Computations

The incompressible Navier-Stokes equation describes the flow motion, from the principles of conservation of momentum and continuum.

$$\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} = -\nabla p + \nu \nabla^2 \mathbf{v} \quad (5)$$

$$\nabla \cdot \mathbf{v} = 0 \quad (6)$$

There is a open sea full of different approaches to solve this non-linear equation. We will for this time explore Chorin's method, and the incremental pressure correction scheme (*IPCS*).

We will use the FEniCS project, a open-source PDE solver using the finite element method approach. As verification of our schemes, the known analytical solution for the 2D Taylor-Green Vortex will be used as comparison. Finally we will move on to the 3D Taylor-Green vortex, using ... as reference.

The Reynolds number, discovered Osborne Reynolds as the relation between inertial and viscous forces, is defined as

$$Re = \frac{\rho V_0 L}{\mu} = \frac{V_0 L}{\nu} \quad (7)$$

Where  $\nu$  denotes the kinematic viscosity, while  $U_0$  and  $D$  is some characteristic velocity and length. We define the kinetic energy as  $E_k = \frac{1}{2} \|u\|_{L^2}^2$ , and the kinetic energy dissipation rate  $\epsilon = \frac{-dE_k}{dt}$

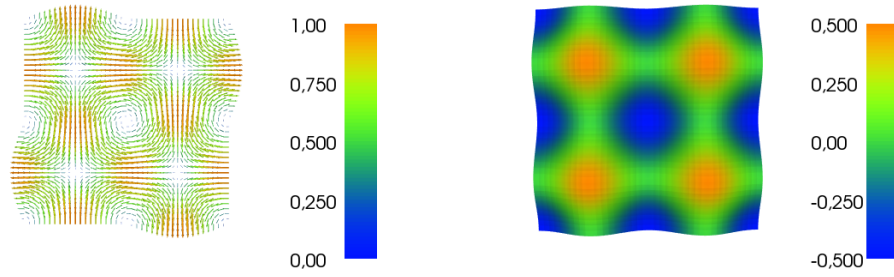
## Setting up 2D problem

For the 2D case, the computational domain is set as  $\Omega \in [-1, 1]^2$ . Further on we will set the flow conditions as

Physical Quantity	Value
Reynolds Number, Re	1000
Characteristic length, L	2
Characteristic velocity, $V_0$	1
Space discretization, N	[8, 16, 32, 64]
Time step $\Delta t$	[0.1, 0.01, 0.001]
End time, T	1.0

Using our analytical solution for time  $t = 0$ , we set up the initial condition for the domain  $\Omega$

Initial velocity field



## Results

Experiments run for Time  $T = 1$  for different choices of  $dt$  and N

*IPCS*

N	$L_2$ dt = 0.1	Runtime	$L_2$ dt = 0.01	Runtime	$L_2$ dt = 0.001	Runtime
8	0.721623	0.410195	0.180146	0.652783	0.178604	3.73357
16	5.34035	0.171748	0.0184505	0.871824	0.0184569	6.47402
32	0.321217	0.361791	0.00122458	1.90135	0.00122848	17.9395
64	0.0046837	1.11411	0.000913817	6.26709	0.000103177	62.2844

*Chorin*

N	$L_2$ dt = 0.1	Runtime	$L_2$ dt = 0.01	Runtime	$L_2$ dt = 0.001	Runtime
8	0.840517	1.382	0.210405	0.382298	0.201774	3.26483
16	4.81879	0.173606	0.0314729	0.698007	0.0292556	6.25138
32	0.0543076	0.342292	0.00307992	1.78643	0.00226862	16.8787
64	0.00329629	1.16073	0.000444121	6.02408	0.000178841	71.846

## Setting up 3D problem

The computational domain is defined as a cube with sides of length  $2\pi L$ ,  $-\pi L \leq x, y, z \leq \pi L$ . We set our physical quantities as follows.

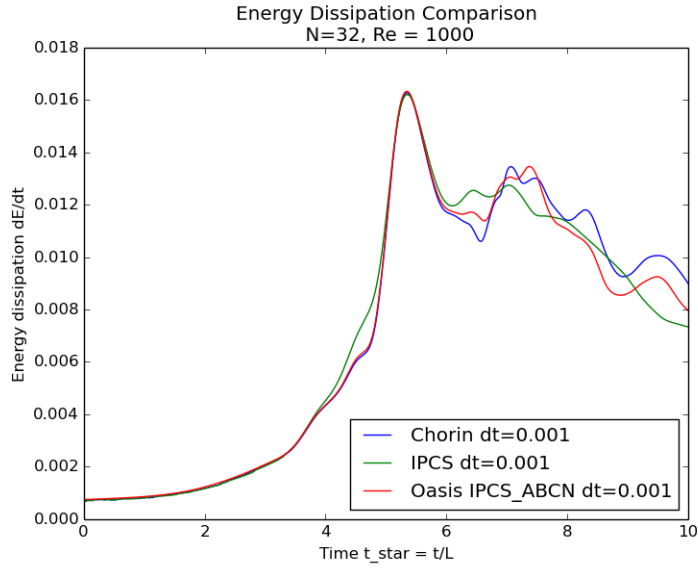
Physical Quantity	Value
Reynolds Number, $Re$	1000
Characteristic length, $L$	1
Characteristic velocity, $V_0$	1
Time step $\Delta t$	0.001
End time, $T$	10.0

Further exploration in the coarseness of the discretiza-

tion in space and time will be investigated aswell.

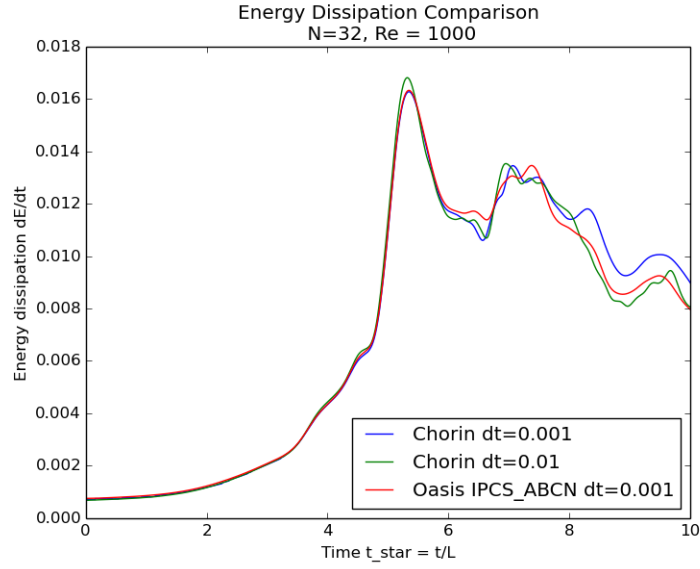
## Results and Comments

Solver	Chorin	IPCS	Oasis
Computer	Abel	Abel	Enkidu
Simulaton hours	11.14	9.30	INFO
Cores	3	3	1
Elements	P2-P1	P2-P1	P2-P1
DOF V	786432	786432	INFO
DOF P	32768	32768	INFO

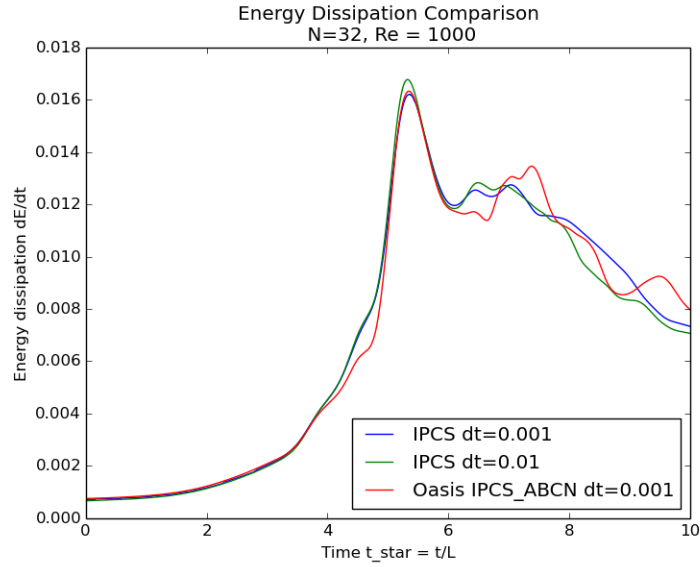


Our numerical computations yields some interesting results. Tracking the energy dissipation towards the distinctive top at  $t \approx 5.8s$ , we observe a small deviation in the IPCS solver in the time domain  $t = [4, 5]$ . Apart from this this time domain, the solvers match the Oasis computations rather good. Continuing downhill after the distinctive top, the solvers deviate from another at  $t \approx 6s$ . We observe that the initial vortexes at this points gets highly turbulent as the energy within the fluid is driven into smaller vortexes. According to the Oasis solver, it seems that my implementations are not able to "catch"all of the fluctuations within the fluid.

Chorin comparison of  $dt = [0.01, 0.001]$



IPCS comparison of  $dt = [0.01, 0.001]$



Moving on to comparing coarseness of the time discretization, we compared  $dt = 0.01, 0.1$  for both solvers. The case of  $dt = 0.1$  diverged after fairly few steps for both solvers. Due to the high Reynolds number chosen for this case, the solvers can't intercept all the fluctuations within the fluid for such a large choice of  $dt$ . For the  $dt = 0.01$  experiment, we observe that the data cling on the finer discretization in time  $dt = 0.001$  fairly well upon reaching the distinctive top in the energy dissipation plot. Here both solvers seems to overestimate the dissipation peak. before falling in with the finer discretization towards the separation point from the Oasis solver at  $t \approx 6$ , which we mentioned in the main experiment.

For the coarseness of the space discretization, comparisons between  $N = 16, 32, 64$  were tried out. For the  $N = 16$  case, the numerical calculations also diverged.....Venter på data