

DRAFT-001  
June 1, 2015

## Feasibility Study for the Final State $hh \rightarrow (b\bar{b})(b\bar{b})$

This is a feasibility study for the final state  $hh \rightarrow (b\bar{b})(b\bar{b})$  using reconstruction techniques for both resolved and boosted topologies as well as multivariate methods.

# 1 Introduction

## 2 Samples

### 2.1 Background

All background samples are generated with the **SHERPA** event generator, version 2.1.1. For the explicit runcards used in the generation, see Appendix A.

The NNPDF 3.0  $n_f = 4$  LO set with strong coupling  $\alpha_S = 0.118$  is used for all samples. At the generator level the following basic cuts are applied. Each final state particle in the hard process must have  $p_T \geq 20$  GeV, and be located within  $|\eta| \leq 3.0$ . All final state particles must be separated by a minimum  $\Delta R_{\min} = 0.1$ . Factorisation and renormalisation scales are set as  $\mu_F = \mu_R = H_T/2$ . Total cross-sections and details of the samples generated are shown in Table 1.

Process	Generator	$N_{\text{evt}}$	$\sigma_{\text{tot}}$
$pp \rightarrow HH$	MG5_aMC@NLO	100K	$1.729 \times 10^{-2}$ pb
$pp \rightarrow b\bar{b}b\bar{b}$	SHERPA 2.1.1	3M	$1.121 \times 10^3$ pb
$pp \rightarrow b\bar{b}jj$	SHERPA 2.1.1	3M	$2.659 \times 10^5$ pb
$pp \rightarrow jjjj$	SHERPA 2.1.1	3M	$9.709 \times 10^6$ pb
$pp \rightarrow t\bar{t}$	SHERPA 2.1.1	3M	$2.514 \times 10^3$ pb

Table 1: Summary of generated samples to date. All **SHERPA** samples have a MC error of 0.05%.

Although suffering from a large theory uncertainty, we can compare the result of our background samples against those presented in the MG5\_aMC@NLO paper [1]. Here for comparison we require in all samples four anti- $k_T$   $R = 0.5$  jets with  $p_T \geq 80$  GeV, and the leading jet must have  $p_T \geq 100$  GeV. All jets must be within an acceptance of  $|\eta| \leq 2.5$ . In the case of the samples with  $b$  quarks in the final state, these requirements are extended to the appropriate number of  $b$ -jets. For example, in the  $2b2j$  sample there must be at least two  $b$ -jets that pass the cuts outlined above.

In Table 2 this comparison is summarised for the  $2b2j$  and  $4b$  samples. Considering the large theory errors, agreement is reasonable in both instances.

Process	$\sigma$ aMC@NLO	$\sigma$ Oxford (SHERPA)
$b\bar{b}b\bar{b}$	$5.050 \times 10^{-1}$ pb	$4.123 \times 10^{-1}$ pb
$b\bar{b}jj$	$1.852 \times 10^2$ pb	$4.239 \times 10^2$ pb
$jjjj$	-	$4.450 \times 10^4$ pb

Table 2: Comparison of LO Oxford SHERPA cross-sections with those of the aMC@NLO paper. The aMC@NLO cross-sections come with a quoted 50% theory uncertainty.

### 3 Object and Event Selection

#### 3.1 Object Definition and Selection

All jets are clustered using all final state particles after the parton shower (but before hadronisation) as input.

Small- $R$  jets are anti- $k_T$   $R = 0.4$  jets. They are required to have  $p_T > 40$  GeV and  $|\eta| < 2.5$ .

Small- $R$  subjets are anti- $k_T$   $R = 0.3$  jets. They are required to have  $p_T > 50$  GeV and  $|\eta| < 2.5$ .

Large- $R$  jets are anti- $k_T$   $R = 1.0$  jets with  $p_T > 200$  GeV and  $|\eta| < 2.0$  (the smaller pseudorapidity is an experimental requirement related to the trackjet-based calibration). They are also required to be mass-drop tagged using the FastJet default parameters:  $\mu = 0.67$  and  $y_{\text{cut}} = 0.09$ .

Small- $R$  jets are  $b$ -tagged as follows: True  $b$ -jets, which have at least one  $b$ -quark among their constituents, have a 80%  $b$ -tagging probability. For all other jets, a mistag rate of 0.01 is used.

Large- $R$  jets are  $b$ -tagged by ghost-associating anti- $k_T$   $R = 0.3$  (selected as above) to the large- $R$  jets. A large- $R$  jets is considered  $bb$ -tagged if the leading and subleading  $p_T$  associated subjets are  $b$ -tagged.

#### 3.2 Resolved Topology

The resolved selection requires the presence of at least four  $b$ -tagged, selected anti- $k_T$   $R = 0.4$  jets.

The di-Higgs system is reconstructed by considering all possibilities of forming two pairs of jets with invariant masses  $m_{j_1j_2}$  and  $m_{j_3j_4}$ , respectively, and choosing the configuration that minimises their difference  $|m_{j_1j_2} - m_{j_3j_4}|$ . Only the four leading- $p_T$  jets are considered here.

### 3.3 Boosted Topology

The boosted selection requires the presence of at least two  $bb$ -tagged, selected anti- $k_T$   $R = 1.0$  jets which are taken as the two Higgs boson candidates.

### 3.4 Intermediate Topology

The intermediate topology requires the presence of at least one  $bb$ -tagged, selected anti- $k_T$   $R = 1.0$  jet. If there is more than one, the leading one is taken as the first Higgs boson candidate. In addition, there must be at least two anti- $k_T$   $R = 0.4$  jets with a separation  $\Delta R > 1.0$  with respect to the large- $R$  jet that is selected as a Higgs boson candidate. The second Higgs boson candidate is reconstructed from the two small- $R$  jets that minimise the difference  $|m_{fj} - m_{j1j2}|$ .

### 3.5 Higgs Mass Window Cut

All Higgs candidates are required to meet the following Higgs mass window cut:  $|m_H - 125| < 40$  GeV

[THE INFORMATION BELOW HAS NOT BEEN UPDATED YET]

#### 3.5.1 Resolved $b$ -jet tagging efficiencies

In order to perform an initial validation of the  $b$ -tagging procedure used in our resolved analysis we consider first a simple flat  $b$ -jet tagging efficiency of  $\epsilon_b = 0.8$  and a light jet mis-tag rate of  $\epsilon_l = 0.01$ . In the UCL diHiggs analysis [2] the  $2b2j$  background was discarded by applying a naive selection factor of  $\epsilon_b^2 \epsilon_l^2$  to a similar  $2b2j$  cross-section as in Table 2 and finding the resulting cross-section to be considerably smaller than the equivalent selection factor applied to the  $4b$  sample.

In practice we have found with these simple efficiencies, that the actual selection efficiencies differ considerably from the naive factors of  $\epsilon_b^{n_b} \epsilon_l^{n_l}$ . In Table 3 the results of this simple  $b$ -tagging test are demonstrated. Shown in the table are the total cross-sections before and after the application of the described  $b$ -tagging procedure, the resulting percentage of the cross-section passing the  $b$ -tag, and the naive selection percentage that would result from applying powers of  $\epsilon_{b/l}$  to the total cross-section. The pass rate for samples with four  $b$  quarks in the final state is therefore considerably lower in practice than one would naively expect, and the pass rate for the  $2b2j$  sample considerably higher.

Process	$\sigma_{\text{tot}}$	After b-tag	% selected	% Naive selection
$H\bar{H}$	$1.729 \times 10^{-2}$ pb	$1.460 \times 10^{-3}$ pb	<b>8.444%</b>	40.96%
$b\bar{b}b\bar{b}$	$9.298 \times 10^2$ pb	78.32 pb	<b>8.423%</b>	40.96%
$b\bar{b}jj$	$2.208 \times 10^5$ pb	84.80 pb	<b>0.038%</b>	0.0064%
$jjjj$	$9.709 \times 10^6$ pb	20.85 pb	<b><math>2.147 \times 10^{-4}\%</math></b>	$1.0 \times 10^{-6}\%$
<b>S/B</b>	$1.74 \times 10^{-9}$	$7.94 \times 10^{-6}$		
<b>S/<math>\sqrt{\text{B}}</math></b>	$9.50 \times 10^{-3}$	$1.86 \times 10^{-1}$		

Table 3:  $b$ -tagging analysis performed upon the samples described in Section 2. Cross-sections are given for the samples before and after the application of the  $b$ -tagging procedure whereby the four hardest jets in the event are used, and the flat selection efficiency and mis-tag rates applied. The percentage of cross-section surviving the selection is shown, along with the naive expectation ( $\epsilon_b^{n_b} \epsilon_l^{n_l}$ ) for comparison. Finally, the signal-to-background ratio and significance are given both before and after the  $b$ -tagging, for  $\mathcal{L} = 3 \text{ ab}^{-1}$ . . Note, older samples have been used here with different generator cuts.

These differences are due to the choices made in the selection procedure. In practice the four hardest jets in the event often do not correspond directly to the four parton level final state particles. In Table 4 the number of true  $b$ -jets (that is, prior to the application of selection efficiencies and fake rates) present in the hardest four jets per event is shown. Indeed, for the samples with four  $b$ -quarks in the parton level final state, the most common scenario has only three  $b$ -jets in the hardest four jets, significantly affecting the selection efficiency.

A more accurate estimation for the passing cross-section may be obtained by multiplying the various fractions of  $b$ -jets by the corresponding efficiencies, for example the 3  $b$ -jet rate by  $\epsilon_b^3 \epsilon_l^1$ . Summing the various scenarios we then have a better estimate for the percentage of cross-section passing the selection, shown in the final column of Table 4. This agrees well with the observed rates shown in Table 3.

### 3.5.2 $b$ -jet tagging efficiencies with UCL-style selection

In the UCL analysis their jet selection does not proceed as described in the previous section, presumably as requiring the four hardest jets in the event to be  $b$ -tagged cut too aggressively into the signal process. Rather, the selection proceeds by taking all anti- $k_T$   $R = 0.5$  jets with  $p_T > 40 \text{ GeV}$  and running the  $b$ -tagging algorithm on them. The four hardest successfully  $b$ -tagged jets

Process	0 $b$ -jets	1 $b$ -jet	2 $b$ -jets	3 $b$ -jets	4 $b$ -jets	% selected
$H\bar{H}$	0.083%	2.599%	24.712%	52.651%	19.955%	<b>8.445%</b>
$bbbb$	0.858%	7.606%	27.523%	43.977%	20.036%	<b>8.434%</b>
$b\bar{b}jj$	9.115%	41.792%	47.806%	1.216%	0.071%	<b>0.038%</b>
$jjjj$	95.668%	3.789%	0.531%	0.011%	3.00E-4%	<b><math>2.17 \times 10^{-4}\%</math></b>

Table 4: Breakdown of the number of true  $b$ -jets in the hardest four clustered jets for various samples. The final column shows the predicted pass rate for the  $b$ -tagging given the distribution of events in the preceding columns.

are then selected for further analysis. With such a method implemented for our samples and with our previous selection efficiencies and fake rates, Table 5 summarises the results.

Process	$\sigma_{\text{tot}}$ (4 Jets $p_T > 40$ GeV)	After $b$ -tag	% selected
$H\bar{H}$	$1.081 \times 10^{-2}$ pb	$8.1147 \times 10^{-4}$ pb	7.467%
$bbbb$	$7.239 \times 10^1$ pb	3.961 pb	5.472%
$b\bar{b}jj$	$2.137 \times 10^4$ pb	4.779 pb	0.0224%
<b>S/B</b>	$5.04 \times 10^{-7}$	$9.28 \times 10^{-5}$	
<b>S/<math>\sqrt{\text{B}}</math></b>	$1.27 \times 10^{-1}$	$4.75 \times 10^{-1}$	

Table 5: Summary of the  $b$ -tagging procedure of the UCL analysis. Presented as in Table 3. In this case the total cross-sections are presented after a cut requiring four anti- $k_T$  R=0.5 jets with  $p_T > 40$  GeV in the final state. The percentage selected column is given for the  $b$ -tagging stage *only* for comparison. Note, older samples have been used here with different generator cuts.

Once again, it is clear that more of the  $b\bar{b}jj$  sample is making it past the  $b$ -tagging selection than originally expected. It seems unlikely that this background can be neglected in the resolved analysis based upon a simple factor of  $\epsilon_b^2 \epsilon_l^2$  as used by the UCL group.

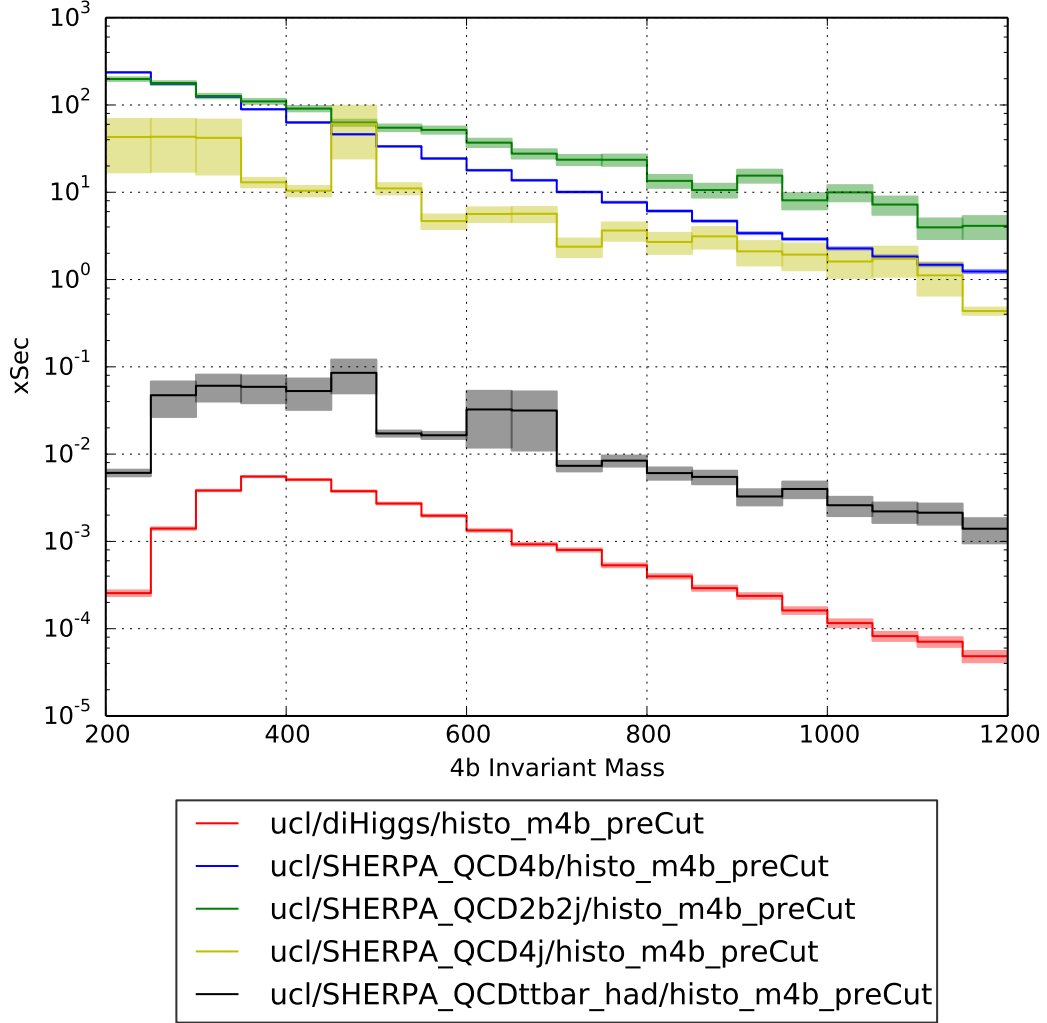


Figure 1: Contributions from the various samples to the observed four  $b$  invariant mass system after applying the  $b$ -tagging procedure outlined in Section 3.5.1.

### 3.5.3 $bb$ -jet rejection for the four hardest jets

### 3.5.4 $bb$ -jet rejection for the UCL-style selection

## 3.6 Boosted Topology

In the boosted topology, the decay products of each Higgs boson are merged into a single large- $R$  jet with a two-prong substructure. These Higgs jets are

Process	$\sigma_{\text{tot}}$ (4 Jets $p_T > 40$ GeV)	After b-tag	After $bb$ rejection	% selected
$H\bar{H}$	17.2911	1.4655	1.38429	%
$bbbb$	1.12149e+06	101152	95723.1	%
$b\bar{b}jj$	2.65895e+08	98116.7	88956.2	%
$jjjj$	1.22738e+10	626267	604.205	%
<b>S/B</b>				
<b>S/<math>\sqrt{B}</math></b>				

Table 6: Summary of the  $b$ -tagging and  $bb$ -rejection cuts on the four leading jets.

Process	$\sigma_{\text{tot}}$ (4 Jets $p_T > 40$ GeV)	After b-tag	After $bb$ rejection	% selected
$H\bar{H}$	10.7756	0.796946	0.744208	7.467%
$bbbb$	101152	4789.57	4438.72	%
$b\bar{b}jj$	1.72539e+07	9962.27	8564.06	%
$jjjj$	6.10999e+08	1.07288e-06	1.07288e-06	%
<b>S/B</b>				
<b>S/<math>\sqrt{B}</math></b>				

Table 7: Summary of the  $b$ -tagging and  $bb$ -rejection cuts on the four leading jets.

reconstructed as anti- $k_T$   $R = 1.0$  jets for which two substructure variables, the first  $k_T$  splitting scale  $d_{12}$  and the 2-subjettiness ratio  $\tau_{21}$ , are calculated. *Pre-cut* histograms are written out after requiring the presence of at least two large- $R$  in the event but without applying any further kinematic cuts. *Post-cut* histograms are filled after the following additional cuts: Both jets are required to have  $p_T > 100$  GeV and  $|\eta| > 2.5$ . Moreover, each jet is required to have at least two b-tagged anti- $k_T$   $R = 0.3$  jets matched to it via *ghost association*: To this end, the constituents of a given large- $R$  jet is reclustered using the jet algorithm and radius parameter of the original jet and so-called *ghost jets*, each corresponding to a small- $R$  jet in the event, are added to the input for the cluster sequence. The 4-vector of a ghost jet is obtained from a given small- $R$  jet by setting its transverse momentum and mass to negligibly small values but retaining its direction in  $\eta$  and  $\phi$ . A small- $R$  jet is considered matched to the large- $R$  jet if its ghost is found among the constituents of the reclustered jet. Ghost association allows to unambiguously match small- $R$  to large- $R$  jets, even in dense environments where a simple matching based on the distance

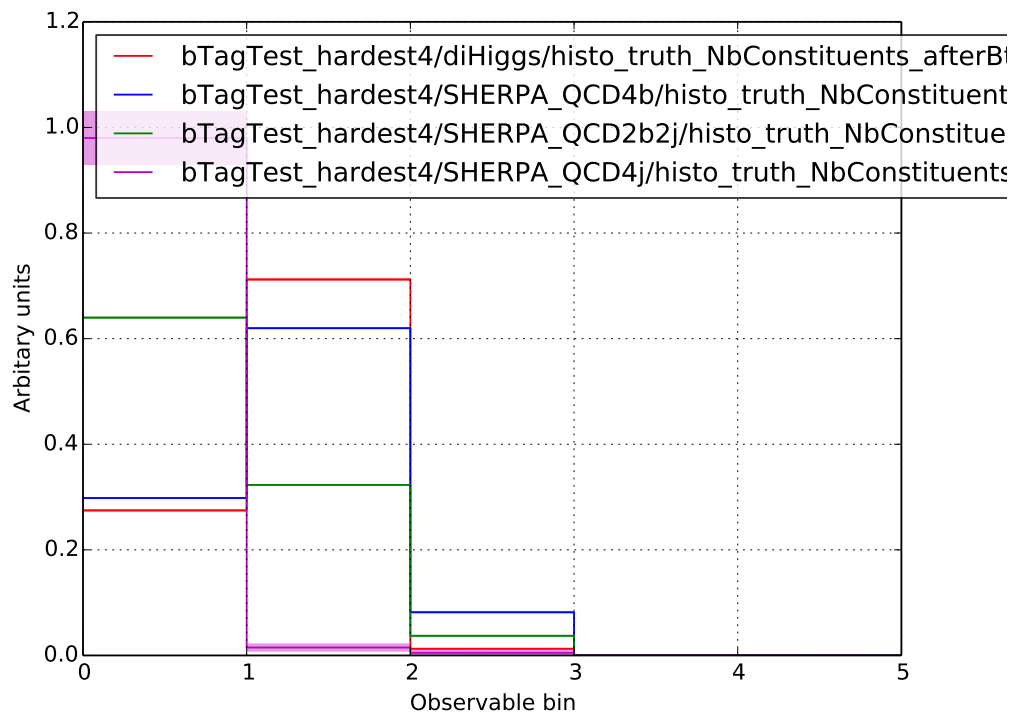
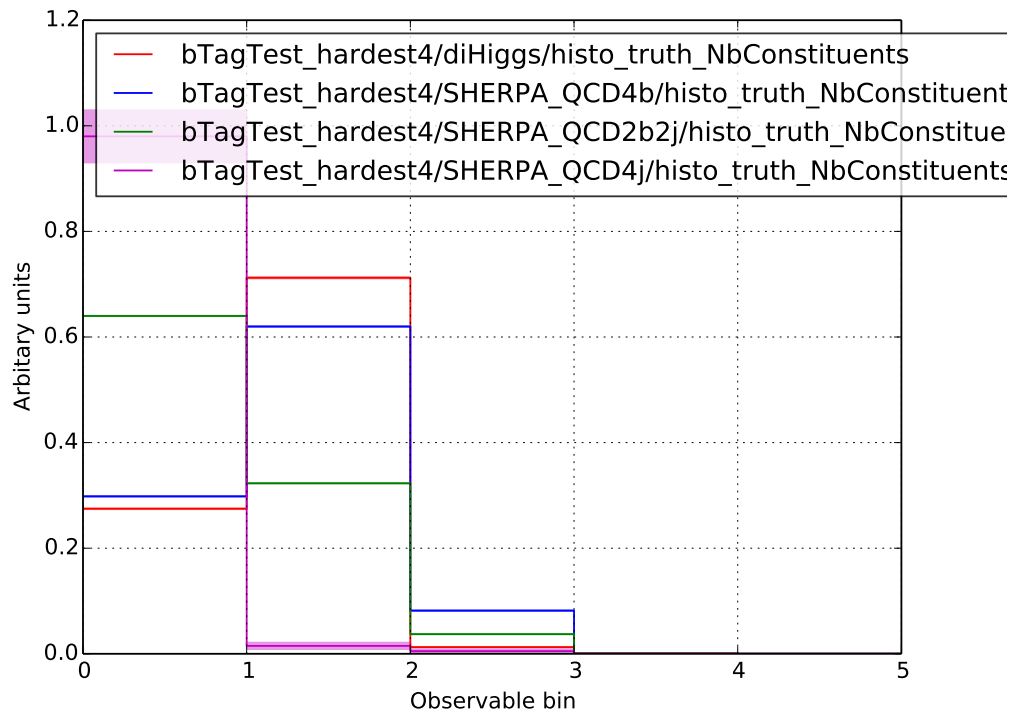


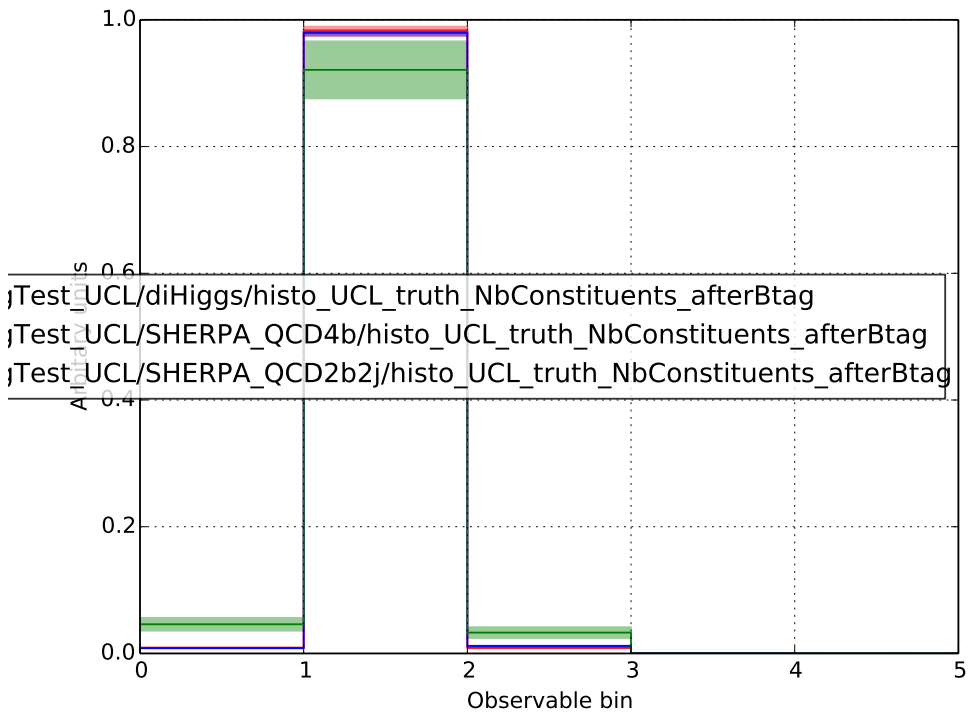
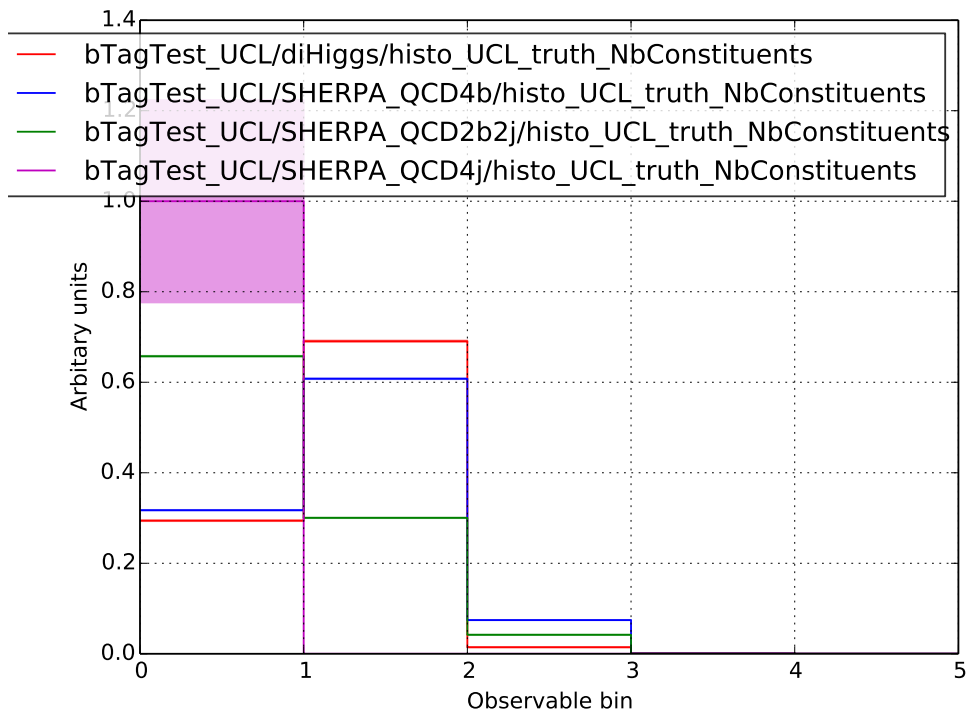
$\Delta R$  between jets may lead to multiple matchings.

### **3.7 Boosted Topology with Variable- $R$ Jets**

The same selection as in Subsection 3.6 is applied but this time using Variable- $R$  jets with the following parameters:  $\rho = 500$  GeV,  $R_{max} = 1.0$ ,  $R_{min} = 0.2$ .

## **4 Multivariate Tools**





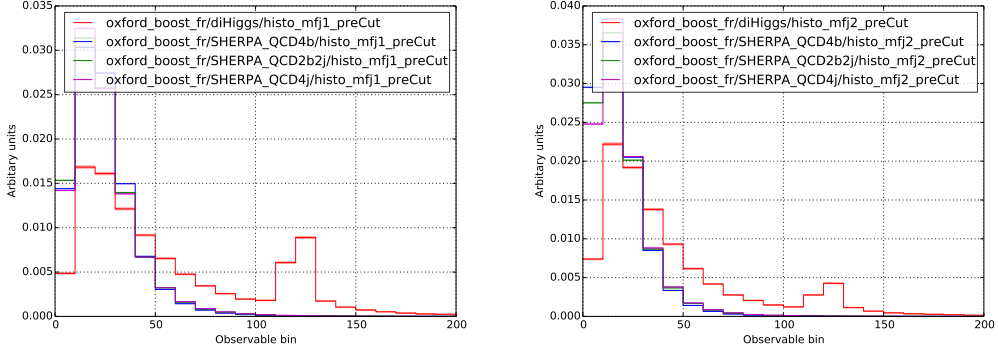


Figure 4: Mass distributions for the leading and subleading anti- $k_T$   $R = 1.0$  jets before the applying the boosted selection.

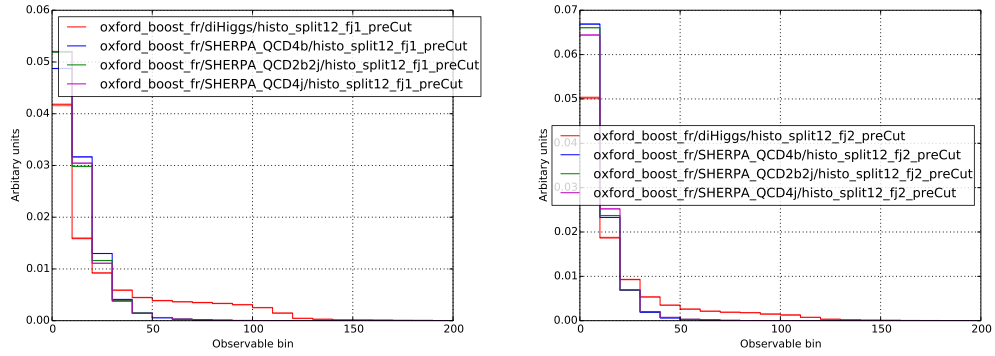


Figure 5:  $d_{12}$  distributions for the leading and subleading anti- $k_T$   $R = 1.0$  jets before the applying the boosted selection.

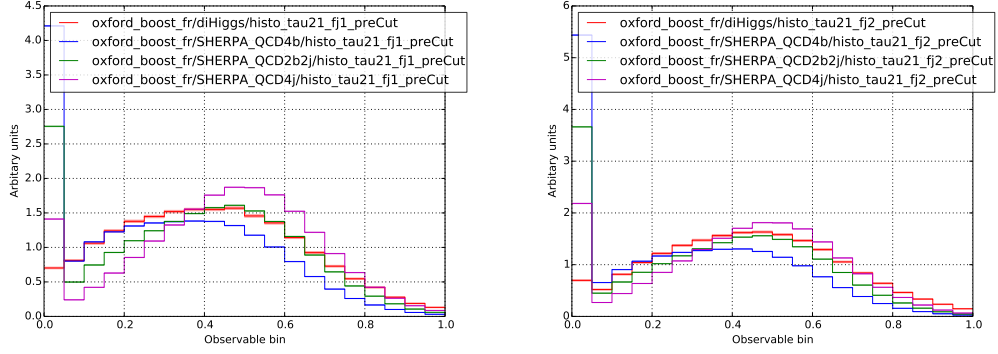


Figure 6:  $\tau_{21}$  distributions for the leading and subleading anti- $k_T$   $R = 1.0$  jets before the applying the boosted selection.

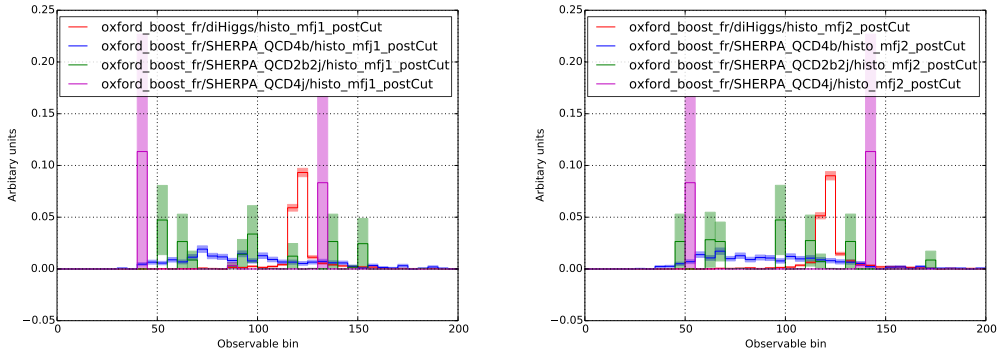


Figure 7: Mass distributions for the leading and subleading anti- $k_T$   $R = 1.0$  jets after the applying the boosted selection.

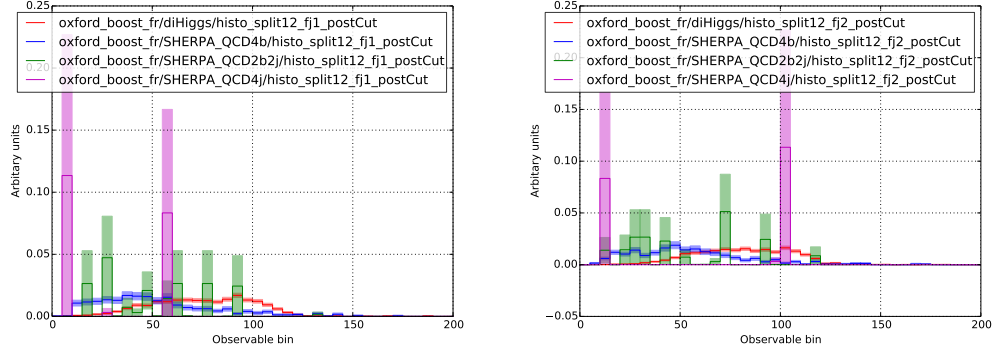


Figure 8:  $d_{12}$  distributions for the leading and subleading anti- $k_T$   $R = 1.0$  jets after the applying the boosted selection.

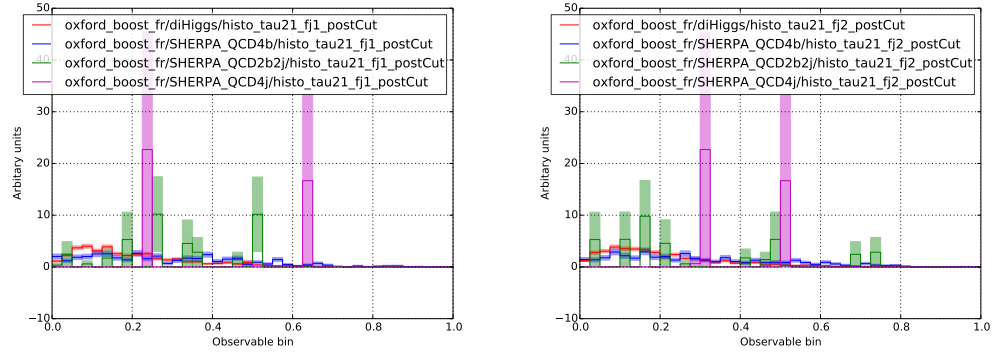


Figure 9:  $\tau_{21}$  distributions for the leading and subleading anti- $k_T$   $R = 1.0$  jets after the applying the boosted selection.

## 5 Results

### 5.1 Example resolved analysis

In this section we present the results of a resolved style analysis upon the full set of available samples for the purposes of illustration. The cuts used in this example analysis are as so. *b-tag* cuts those events which fail to tag four anti- $k_T$   $R = 0.5$   $b$ -jets with  $b$  quark  $p_T > 15$  GeV. Flat  $b$ -jet efficiencies and mistag rates of 80% and 1% respectively. A  $c$  mistag rate of 17% is also applied. *jet*  $p_{T,\eta}$  requires that the four jets each have  $p_T > 40$  GeV and  $|\eta| < 2.5$ . *di-jet*  $p_T$  requires that each dijet pair representing a higgs candidate has  $p_T > 150$  GeV. *di-jet*  $\Delta\eta$  requires that the two dijets are separated by no more than 1.5 units in  $\eta$ . Finally  *$m_H$  window* requires that the invariant mass of each dijet is within 15 GeV of  $m_H$ . The remaining cross-section after each of these cuts, along with  $S/B$  and  $S/\sqrt{B}$  are given in Table 8.

Cut	$HH$ [fb]	$4b$ [fb]	$2b2j$ [fb]	$4j$ [fb]	$t\bar{t}$ [fb]	$S/B$	$S/\sqrt{B}$
b-tag	1.516	102271	135255	42934	202.9	5.40E-6	0.157
jet $p_{T,\eta}$	0.723	3847	10676	5045	122.7	3.67E-5	0.282
di-jet $p_T$	0.371	553.7	1838	926.0	22.45	1.11E-4	0.352
di-jet $\Delta\eta$	0.336	324.8	991.4	312.6	16.39	2.04E-4	0.454
$m_H$ window	0.190	7.632	13.60	0.502	2.675	7.78E-3	2.106

Table 8: Table demonstrating the cutflow of an example resolved analysis applied to the full sample set.  $S/\sqrt{B}$  is given for HL-LHC luminosities.

### 5.2 Example boosted analysis

Here we shall discuss an example boosted analysis. In this case *Fatjet cuts* refers to the requirement of two  $R = 1$  anti- $k_T$  jets with  $p_T > 250$  GeV and  $|\eta| < 2.5$ . *2 Subjects* requires two anti- $k_T$   $R = 0.3$  subjects in each fat jet. *b-Tagging* cuts those events in which the two hardest subjects are not successfully  $b$ -tagged. Here efficiency and mistag factors are applied for  $b$  and  $c$  jets as a function of subject  $p_T$ . *Mass-drop* requires that the two fat jets pass a BDRS mass-drop tagger with parameters  $\mu = 0.67$  and  $y_{\text{cut}} = 0.09$ . Finally  *$m_H$  window* only passes events where the two fat jet masses are within 15 GeV of  $m_H$ . In Table 9 we see the results of these cuts.

Cut	$HH$ [fb]	$4b$ [fb]	$2b2j$ [fb]	$4j$ [fb]	$t\bar{t}$ [fb]	$S/B$	$S/\sqrt{B}$
Fatjet cuts	2.802	3045	1.13E6	4.57E7	9808	5.98E-8	0.022
2 Subjets	2.740	2694	1.01E6	4.21E7	9787	6.35E-8	0.023
b-Tagging	0.070	39.76	296.9	352.9	7.586	1.00E-5	0.145
Mass-drop	0.068	36.84	256.0	305.1	7.514	1.12E-4	0.151
$m_H$ window	0.045	2.134	10.83	0.630	0.438	3.20E-3	0.658

Table 9: Table demonstrating the cutflow of an example boosted analysis applied to the full sample set.  $S/\sqrt{B}$  is given for HL-LHC luminosities.

### 5.3 MVA

We now apply a basic MVA to the final output of the resolved and boosted analyses. The MVA used is a feedforward artificial neural network with  $N_{\text{par}} \times 5 \times 3 \times 1$  architecture, where  $N_{\text{par}}$  is the number of input parameters for the network. The MVA was trained for 50,000 GA generations with a cross-entropy error function. In Figure 10 the neural network classification of events is demonstrated for both the resolved and boosted cases. In Figure 11 the ROC curves for the two analyses are shown, demonstrating that the neural network is able to perform well in both cases. Figure 12 shows the distribution of NN weights in both cases.

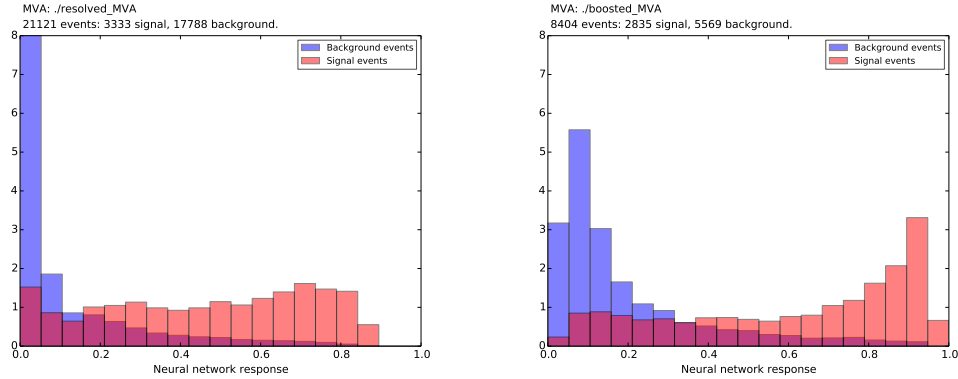


Figure 10: Neural network response demonstrated for the case of the resolved (left) and boosted (right) analyses. The frequency of background and signal events is plotted as a function of neural network output.



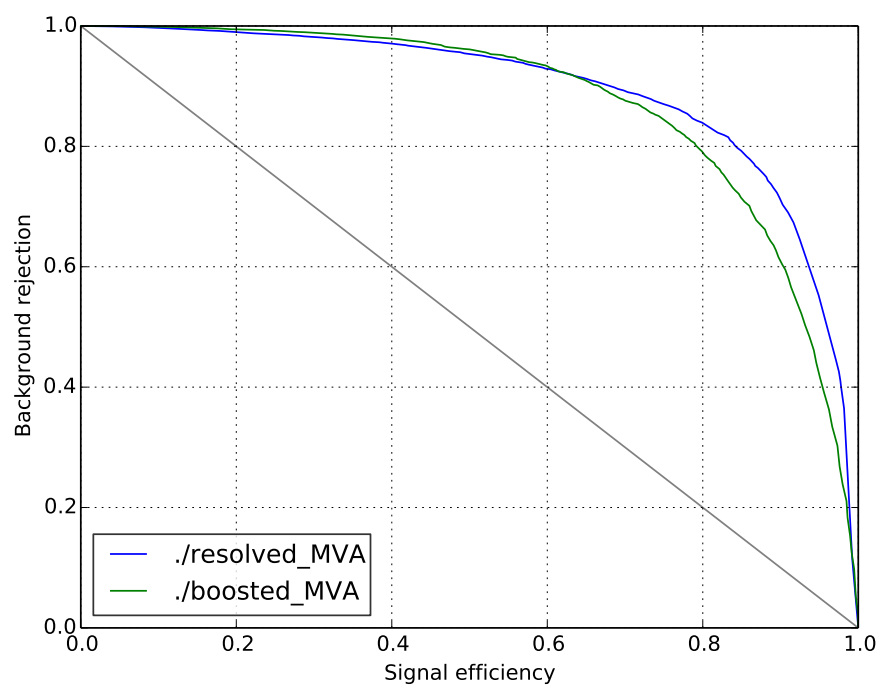


Figure 11: ROC curve for the neural network discriminant in the boosted (green) and resolved (blue) cases.

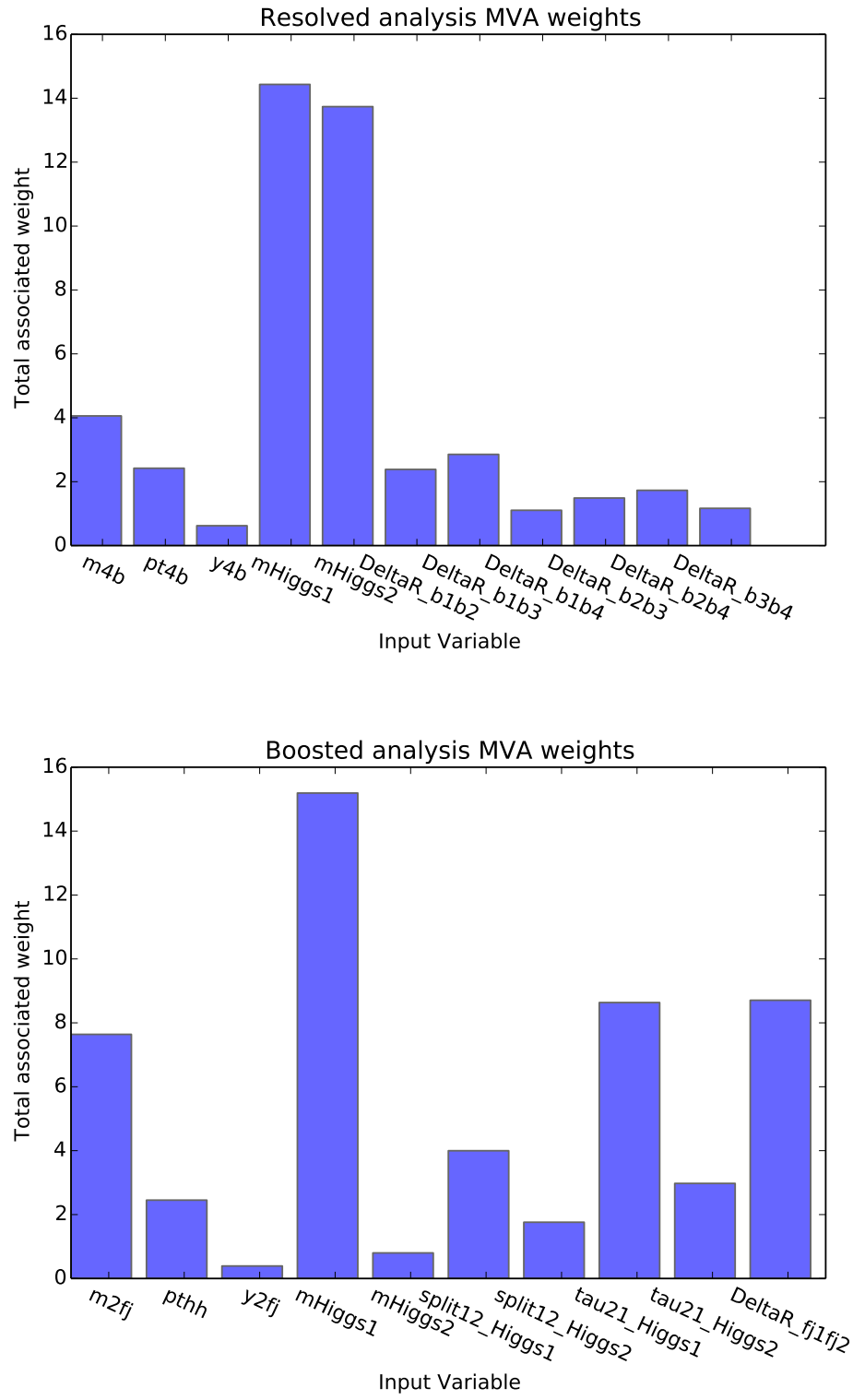


Figure 12: Distribution of weights per input variable in the final ANN analysis. The above figure shows the weight distribution in the case of the resolved analysis, while the lower plot shows the boosted analysis.

Figure 12 demonstrates that the dijet invariant mass is a crucial variable for the resolved analysis MVA. As no detector effects were modelled in this analysis, it is reasonable to assume that the MVA has greater discriminating power than would be feasible in a real analysis. To investigate this, we now consider a variant of the analysis used in Table 8 where the  $m_H$  window criterion has been relaxed to  $80 < m_H < 170$ . Firstly the same analysis is repeated with the larger  $m_H$  window. This is then compared to the same analysis but with a Gaussian smearing with  $\sigma = 10$  GeV applied to the dijet invariant mass used in the MVA. In the left panel of Figure 13 the smearing is demonstrated. Despite the smearing, the MVA applied to this analysis still vastly prefers the dijet masses as discriminators, as is demonstrated in the right panel of Figure 13.

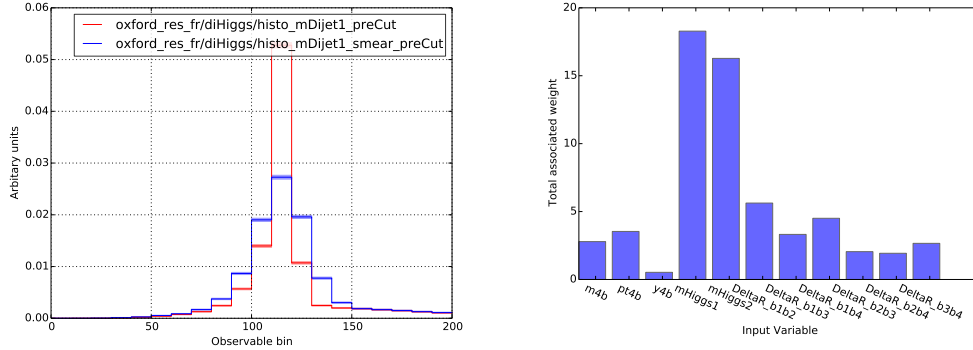


Figure 13: The left figure demonstrates the gaussian smearing applied to the dijet invariant mass distribution. The right figure demonstrates the ANN weight distribution in the MVA.

While the MVA still derives most of its discriminating power from the dijet masses, the smearing reduces its effectiveness considerably. In Figure 14 the ROC curves of the MVA applied to the smeared and unsmeared distributions are shown, where it is clear that the overall performance of the MVA is significantly hampered by the smeared invariant masses. This is echoed in the resulting significance of the analysis as demonstrated in Figure 15

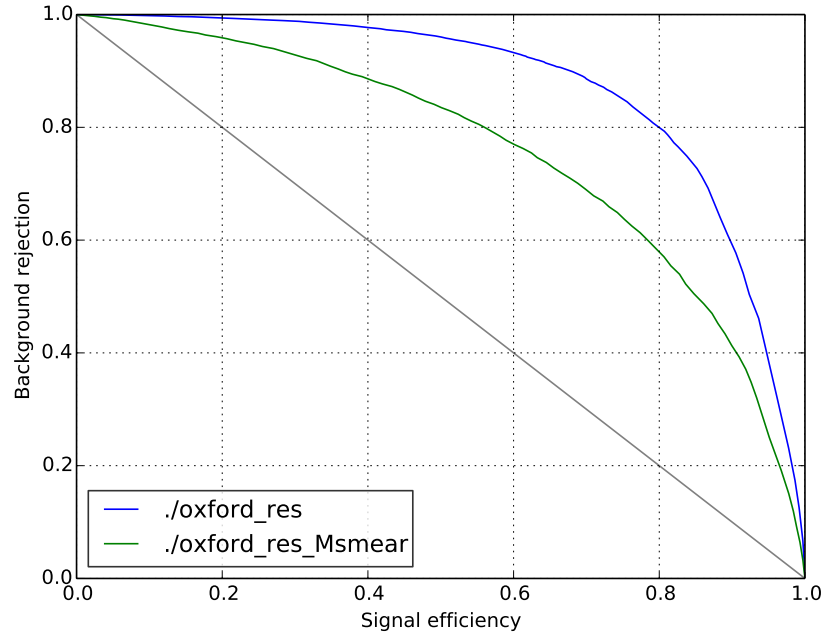


Figure 14: ROC curves of the MVA applied with the unsmeared (blue) and smeared (green) dijet invariant mass distributions.

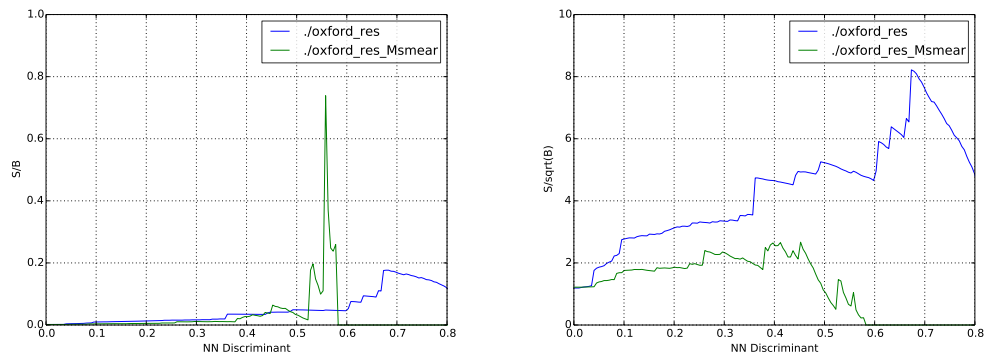


Figure 15:  $S/B$  and  $S\sqrt{B}$  (HLLHC) for the MVA applied to the unsmeared (blue) and smeared (green) dijet invariant mass distributions.

## 6 Conclusion

## References

- [1] J. Alwall, R. Frederix, S. Frixione, V. Hirschi, F. Maltoni, O. Mattelaer, H.-S. Shao and T. Stelzer *et al.*, JHEP **1407** (2014) 079 [arXiv:1405.0301 [hep-ph]].
- [2] D. Wardrope, E. Jansen, N. Konstantinidis, B. Cooper, R. Falla and N. Norjoharuddeen, “Non-resonant Higgs pair production in the  $b\bar{b}b\bar{b}$  final state at the LHC,” arXiv:1410.2794 [hep-ph].

## A Event sample runcards

### A.1 Background: QCD 4b

```
(run){  
  EVENTS 3M;  
  EVENT_GENERATION_MODE U;  
  ME_SIGNAL_GENERATOR Comix;  
  
  EVENT_OUTPUT HepMC_Short[SHERPA_QCD_4b];  
  
  BEAM_1 2212; BEAM_ENERGY_1 7000;  
  BEAM_2 2212; BEAM_ENERGY_2 7000;  
  
  FRAGMENTATION=Off # disable hadronisation  
  MI_HANDLER=None # disable multiple parton interactions  
  
  SCF:=1; ### default scale factor  
  SCALES VAR{SCF*H_T2/4};  
  
  PDF_LIBRARY LHAPDFSherpa;  
  PDF_SET NNPDF30_lo_as_0118_nf_4.LHgrid;  
  PDF_SET_VERSION 0;  
  
  MASSIVE[5] 1;  
  MASS[5] 4.75;  
  
}(run);  
(processes){  
  Process 93 93 -> 5 -5 5 -5;  
  Order_EW 0;  
  End process;  
  
}(processes);  
(selector){  
  PT 5 20 7000  
  PT -5 20 7000  
  PseudoRapidity 5 -3.0 3.0  
  PseudoRapidity -5 -3.0 3.0
```

```

DeltaR 5 5 0.1 10000
DeltaR 5 -5 0.1 10000
DeltaR -5 -5 0.1 10000

}(selector);



## A.2 Background: QCD 2b2j



(run){
  EVENTS 3M;
  EVENT_GENERATION_MODE U;
  ME_SIGNAL_GENERATOR Comix;

  EVENT_OUTPUT HepMC_Short[SHERPA_QCD_2b2j];

  BEAM_1 2212; BEAM_ENERGY_1 7000;
  BEAM_2 2212; BEAM_ENERGY_2 7000;

  FRAGMENTATION=Off # disable hadronisation
  MI_HANDLER=None # disable multiple parton interactions

  SCF:=1; ### default scale factor
  SCALES VAR{SCF*H_T2/4};

  PDF_LIBRARY LHAPDFSherpa;
  PDF_SET NNPDF30_lo_as_0118_nf_4.LHgrid;
  PDF_SET_VERSION 0;

  MASSIVE[5] 1;
  MASS[5] 4.75;

}(run);

(processes){
  Process 93 93 -> 93 93 5 -5;
  Order_EW 0;
  End process;

}(processes);

(selector){
  PT 5 20 7000

```

```

PT -5 20 7000
PseudoRapidity 5 -3 3
PseudoRapidity -5 -3 3

PT 93 20 7000
PseudoRapidity 93 -3 3

DeltaR 93 93 0.1 10000
DeltaR 5 93 0.1 10000
DeltaR -5 93 0.1 10000
DeltaR 5 5 0.1 10000
DeltaR -5 5 0.1 10000
DeltaR -5 -5 0.1 10000
}(selector);

```

### A.3 Background: QCD 4j

```

(run){
  EVENTS 3M;
  EVENT_GENERATION_MODE U;
  ME_SIGNAL_GENERATOR Comix;

  EVENT_OUTPUT HepMC_Short[SHERPA_QCD_4j];

  BEAM_1 2212; BEAM_ENERGY_1 7000;
  BEAM_2 2212; BEAM_ENERGY_2 7000;

  FRAGMENTATION=Off # disable hadronisation
  MI_HANDLER=None # disable multiple parton interactions

  SCF:=1; ### default scale factor
  SCALES VAR{SCF*H_T2/4};

  PDF_LIBRARY LHAPDFSherpa;
  PDF_SET NNPDF30_lo_as_0118_nf_4.LHgrid;
  PDF_SET_VERSION 0;

  MASSIVE[5] 1;
  MASS[5] 4.75;

```



```

}(run);

(processes){
  Process 93 93 -> 93 93 93 93;
  Order_EW 0;
  End process;
}(processes);

(selector){
  PT 93 20 7000
  PseudoRapidity 93 -3 3
  DeltaR 93 93 0.1 10000
}(selector);

```

## A.4 Background: QCD ttbar

```

(run){
  EVENTS 3M;
  EVENT_GENERATION_MODE U;
  ME_SIGNAL_GENERATOR Comix;

  EVENT_OUTPUT HepMC_Short[SHERPA_QCD_ttbar];

  BEAM_1 2212; BEAM_ENERGY_1 7000;
  BEAM_2 2212; BEAM_ENERGY_2 7000;

  FRAGMENTATION=Off # disable hadronisation
  MI_HANDLER=None # disable multiple parton interactions

  SCF:=1; ### default scale factor
  SCALES VAR{SCF*H_T2/4};

  PDF_LIBRARY LHAPDFSherpa;
  PDF_SET NNP30_lo_as_0118_nf_4.LHgrid;
  PDF_SET_VERSION 0;

  MASSIVE[5] 1;
  MASS[5] 4.75;

  # Stable mode 0 specified that both particle and antiparticle are unstable

```

```

STABLE[6] = 0;
STABLE[24] = 0;

# Enable only fully hadronic decay modes
HARD_DECAYS=1;
HDH_NO_DECAY={24,12,-11}|{24,14,-13}|{24,16,-15}|{-24,-12,11}|{-24,-14,13}|{-24,

}(run);

(processes){
  Process 93 93 -> 6 -6;
  Order_EW 0;
  End process;
}(processes);

(selector){
  PT 6 20 7000
  PT -6 20 7000
  PseudoRapidity 6 -3 3
  PseudoRapidity -6 -3 3
  DeltaR 6 -6 0.1 10000
}(selector);

```

## B Decay tables for $t\bar{t}$

Decay table for :  $W^+$ .

Total width: 2.10561 GeV

Flavour width: 2.06 GeV

{24,2,-1}	$W^+ \rightarrow u \bar{d}$	0.702203	GeV, BR= 33.3492 %
{24,4,-3}	$W^+ \rightarrow c \bar{s}$	0.701546	GeV, BR= 33.318 %
{24,12,-11}	$W^+ \rightarrow \nu_e e^+$	0.234068	GeV, BR= 11.1164 %
{24,14,-13}	$W^+ \rightarrow \nu_\mu \mu^+$	0.234066	GeV, BR= 11.1163 %
{24,16,-15}	$W^+ \rightarrow \nu_\tau \tau^+$	0.233725	GeV, BR= 11.1001 %

Decay table for :  $W^-$ .

Total width: 2.10561 GeV

Flavour width: 2.06 GeV

{-24,-2,1}	$W^- \rightarrow \bar{u} b$	0.702203	GeV, BR= 33.3492 %
{-24,-4,3}	$W^- \rightarrow \bar{c} s$	0.701546	GeV, BR= 33.318 %
{-24,-12,11}	$W^- \rightarrow \bar{\nu}_e e^-$	0.234068	GeV, BR= 11.1164 %
{-24,-14,13}	$W^- \rightarrow \bar{\nu}_\mu \mu^-$	0.234066	GeV, BR= 11.1163 %
{-24,-16,15}	$W^- \rightarrow \bar{\nu}_\tau \tau^-$	0.233725	GeV, BR= 11.1001 %

Decay table for :  $t$ .

Total width: 1.59689 GeV

Flavour width: 1.5 GeV

{6,24,5}	$t \rightarrow W^+ b$	1.59689	GeV, BR= 100 %
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Decay table for :  $t\bar{b}$ .

Total width: 1.59689 GeV

Flavour width: 1.5 GeV

{-6,-24,-5}	$t\bar{b} \rightarrow W^- \bar{b} b$	1.59689	GeV, BR= 100 %
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