

The expectation values of the operator  $\mathcal{O}_{\phi^3}$  and the trace of the energy-momentum tensor are independent of  $\Omega$ . As a check, we note that the four one-point functions satisfy the following operator relation, which is associated to the violation of conformal invariance by non-zero classical beta functions,

$$\langle T^i_i \rangle = - \sum_{\mathcal{O}} (d - \Delta_{\mathcal{O}}) \phi_{\mathcal{O}} \langle \mathcal{O} \rangle \quad (1)$$

### 0.0.1 Derivative of the free energy

Following , we may now compute the derivative of  $F$  with respect to  $\alpha$  as follows. First we note that

$$\frac{dF}{d\alpha} = \frac{dS_{\text{ren}}}{d\alpha} = \lim_{\epsilon \rightarrow 0} \int d^5x \sum_{\text{fields } \Phi} \frac{\delta(\sqrt{\gamma} \mathcal{L}_{\text{ren}})}{\delta \Phi} \frac{d\Phi}{d\alpha} \Big|_{z=\epsilon} \quad (2)$$

In our case, the terms appearing in the sum over fields are

$$\begin{aligned} \frac{\delta(\sqrt{\gamma} \mathcal{L}_{\text{ren}})}{\delta \sigma} &= \sqrt{\gamma} \langle O_{\sigma} \rangle \epsilon^3 + \dots & \frac{\delta(\sqrt{\gamma} \mathcal{L}_{\text{ren}})}{\delta \phi^0} &= \sqrt{\gamma} \langle O_{\phi}^0 \rangle \epsilon^4 + \dots \\ \frac{\delta(\sqrt{\gamma} \mathcal{L}_{\text{ren}})}{\delta \phi^3} &= \sqrt{\gamma} \langle O_{\phi}^3 \rangle \epsilon^3 + \dots & \frac{\delta(\sqrt{\gamma} \mathcal{L}_{\text{ren}})}{\delta \gamma^{ij}} &= \frac{1}{2} \sqrt{\gamma} \langle T_{ij} \rangle \epsilon^5 + \dots \end{aligned} \quad (3)$$

The dots represent terms of strictly lower order in  $\epsilon$ . Furthermore, from the form of the UV asymptotic expansions , we have

$$\begin{aligned} \frac{d\sigma}{d\alpha} &= \frac{3}{4} \alpha \epsilon^2 + O(\epsilon^3) & \frac{d\phi^0}{d\alpha} &= \epsilon + O(\epsilon^3) \\ \frac{d\phi^3}{d\alpha} &= \left(1 - \alpha \frac{df_k}{d\alpha}\right) e^{-f_k} \epsilon^2 + O(\epsilon^3) & \frac{d\gamma^{ij}}{d\alpha} &= -2 \frac{df_k}{d\alpha} e^{-2f_k} \epsilon^2 + O(\epsilon^2) \end{aligned} \quad (4)$$

Combining the pieces , with the results for the one-point functions in , we find that the contribution of the metric in is suppressed by  $\epsilon^2$  compared to other terms. The derivative of the free energy is then

$$\begin{aligned} \frac{dF}{d\alpha} &= \lim_{\epsilon \rightarrow 0} \int d^5x \sqrt{\gamma} \epsilon^5 \left[ \frac{3}{2} \beta e^{-f_k} + \frac{1}{2} \beta e^{-f_k} \left(1 - \alpha \frac{df_k}{d\alpha}\right) + O(\epsilon) \right] \\ &= \text{vol}_0(S^5) \frac{1}{2} \beta e^{4f_k} \left(4 - \alpha \frac{df_k}{d\alpha}\right) \end{aligned} \quad (5)$$

where  $\text{vol}_0(S^5) = \pi^3$  is the volume of a unit  $S^5$ . The  $\Omega$  dependence in the one-point functions cancels out, consistent with the fact that  $F$  itself is independent of  $\Omega$ . We thus obtain the final result

$$\frac{dF}{d\alpha} = \frac{\pi^2}{8 G_6} \beta e^{4f_k} \left(4 - \alpha \frac{df_k}{d\alpha}\right) \quad (6)$$