

$y_1 = (pq)^{-\frac{1}{2}}t^{-2}AB^{-1}$  and  $y_2 = (pq)^{-\frac{1}{2}}q^{-1}t^{-2}AB^{-1}$ . For  $y_1$  the integral over  $w_1$  is pinched at  $w_1 = v^{\pm 1}$  and we proceed exactly as in appendix B to get the same result up to some overall factors

$$\begin{aligned} & \frac{1}{2} \frac{\Gamma_e(qt^{-2})\Gamma_e(pq^2t^2A^{-2}B^2)\Gamma_e(t^{-2}B^{-2})\Gamma_e(t^{-2}A^2)\Gamma_e((pq)^{-1}q^{-1}t^{-4}A^2B^{-2})}{\Gamma_e((pq)^{-1}t^{-4}A^2B^{-2})} \\ & \times \Gamma_e(t^{-2}AB^{-1}C^{\pm 1}D^{\pm 1})\Gamma_e(AB^{-1}u^{\pm 1}v^{\pm 1})\Gamma_e((qp)^{\frac{1}{2}}tBD^{\pm 1}v^{\pm 1})\Gamma_e((qp)^{\frac{1}{2}}tA^{-1}C^{\pm 1}v^{\pm 1}) \\ & \times \Gamma_e((pq)^{\frac{1}{2}}tBC^{\pm 1}u^{\pm 1})\Gamma_e((pq)^{\frac{1}{2}}tA^{-1}D^{\pm 1}u^{\pm 1}). \end{aligned} \quad (1)$$

Now we look at the pole  $y_2$ . the integral over  $w_1$  is pinched at  $w_1 = q^{\pm \frac{1}{2}}v^{\pm 1}$ . Substituting these values (??) now becomes

$$\begin{aligned} & \frac{1}{2} \frac{\Gamma_e(t^{-2})\Gamma_e(pq^3t^2A^{-2}B^2)}{\Gamma_e(pq^3t^4A^{-2}B^2)} \frac{\Gamma_e(v^2)}{\Gamma_e(qv^2)} \Gamma_e(q^{-\frac{1}{2}}AB^{-1}u^{\pm 1}(q^{\frac{1}{2}}v)^{\pm 1})\Gamma_e((qp)^{\frac{1}{2}}q^{\frac{1}{2}}tBD^{\pm 1}(q^{\frac{1}{2}}v)^{\pm 1}) \\ & \times \Gamma_e((qp)^{\frac{1}{2}}q^{\frac{1}{2}}tA^{-1}C^{\pm 1}(q^{\frac{1}{2}}v)^{\pm 1}) \oint \frac{dw_2}{4\pi iw_2} \frac{1}{\Gamma_e(w_2^{\pm 2})} \Gamma_e\left(\frac{(pq)^{\frac{1}{2}}}{t^2}q^{\frac{1}{2}}vw_2^{\pm 1}\right) \\ & \times \Gamma_e((qp)^{\frac{1}{2}}q^{\frac{1}{2}}t^2A^{-1}Bw_2^{\pm 1}u^{\pm 1})\Gamma_e(q^{-\frac{1}{2}}t^{-1}B^{-1}D^{\pm 1}w_2^{\pm 1}) \\ & \times \Gamma_e((qp)^{\frac{1}{2}}q^{\frac{1}{2}}t^2w_2^{\pm 1}v^{-1})\Gamma_e(q^{-\frac{1}{2}}t^{-1}Aw_2^{\pm 1}C^{\pm 1}) + \{v \leftrightarrow v^{-1}\}. \end{aligned}$$

The integral here can be interpreted as index of the SU(2) gauge theory with four flavors. Using the Intriligator–Pouliot duality transformation (that is  $V(\underline{s}) = \prod_{1 \leq j < k \leq 8} \Gamma_e(s_j s_k) V(\sqrt{pq}/\underline{s})$  in the notations of) the expression becomes

$$\begin{aligned} & \frac{1}{2} \frac{\Gamma_e(t^{-2})\Gamma_e(pq^3t^2A^{-2}B^2)}{\Gamma_e(pq^3t^4A^{-2}B^2)} \frac{\Gamma_e(v^2)}{\Gamma_e(qv^2)} \Gamma_e(q^{-\frac{1}{2}}AB^{-1}u^{\pm 1}(q^{\frac{1}{2}}v)^{\pm 1})\Gamma_e((qp)^{\frac{1}{2}}q^{\frac{1}{2}}tBD^{\pm 1}(q^{\frac{1}{2}}v)^{\pm 1}) \\ & \times \Gamma_e((qp)^{\frac{1}{2}}q^{\frac{1}{2}}tA^{-1}C^{\pm 1}(q^{\frac{1}{2}}v)^{\pm 1})\Gamma_e(pq^2A^{-1}Bvu^{\pm 1})\Gamma_e((pq)^{\frac{1}{2}}t^{-3}B^{-1}vD^{\pm 1}) \\ & \times \Gamma_e((pq)^{\frac{1}{2}}t^{-3}AvC^{\pm 1})\Gamma_e(pq^2t^4A^{-2}B^2)\Gamma_e((pq)^{\frac{1}{2}}tA^{-1}u^{\pm 1}D^{\pm 1})\Gamma_e((pq)^{\frac{1}{2}}tBu^{\pm 1}C^{\pm 1}) \\ & \times \Gamma_e(pq^2t^4A^{-1}Bv^{-1}u^{\pm 1})\Gamma_e(q^{-1}t^{-2}B^{-2})\Gamma_e(q^{-1}t^{-2}AB^{-1}C^{\pm 1}D^{\pm 1}) \\ & \times \Gamma_e((pq)^{\frac{1}{2}}tB^{-1}v^{-1}D^{\pm 1})\Gamma_e(q^{-1}t^{-2}A^2)\Gamma_e((pq)^{\frac{1}{2}}tAv^{-1}C^{\pm 1})\Gamma_e(pq^2) \\ & \times \oint \frac{dw_2}{4\pi iw_2} \frac{1}{\Gamma_e(w_2^{\pm 2})} \Gamma_e(q^{-\frac{1}{2}}t^2v^{-1}w_2^{\pm 1})\Gamma_e(q^{-\frac{1}{2}}t^{-2}AB^{-1}w_2^{\pm 1}u^{\pm 1}) \\ & \times \Gamma_e((pq)^{\frac{1}{2}}q^{\frac{1}{2}}tBD^{\pm 1}w_2^{\pm 1})\Gamma_e((qp)^{\frac{1}{2}}q^{\frac{1}{2}}tA^{-1}w_2^{\pm 1}C^{-1})\Gamma_e(q^{-\frac{1}{2}}t^{-2}vw_2^{\pm 1}) + \{v \leftrightarrow v^{-1}\}. \end{aligned}$$

We have zero multiplying the integral but the integral is pinched at  $w_2 = (q^{\pm \frac{1}{2}}t^{-2}v)^{\pm 1}$  due to the colliding of poles of the first and last elliptic gamma functions in the last integral. We get different contribution for each choice