

# Mapping Density Functional Theorem onto Heisenberg model

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## MOTIVATION

Employing Quantum Mechanics in modelling thermal properties of new materials has a significant disadvantage - far from the thermodynamic limit we do not have temperature defined. Hence to get transitional temperatures e.g. Curie (Néel)  $T_C$  ( $T_N$ ) for a proposed **magnetic material**, we are limited to the model calculations, e.g. the established atomistic spin-dynamics (ASD) simulations [1,2]. This in turn requires us to find the explicit form of the magnetic **Heisenberg Hamiltonian** (cf. *Some Math* for more details)

$$\mathcal{H} \equiv -\frac{1}{2} \sum_{i \neq j} J_{ij} \mathbf{s}_i \cdot \mathbf{s}_j,$$

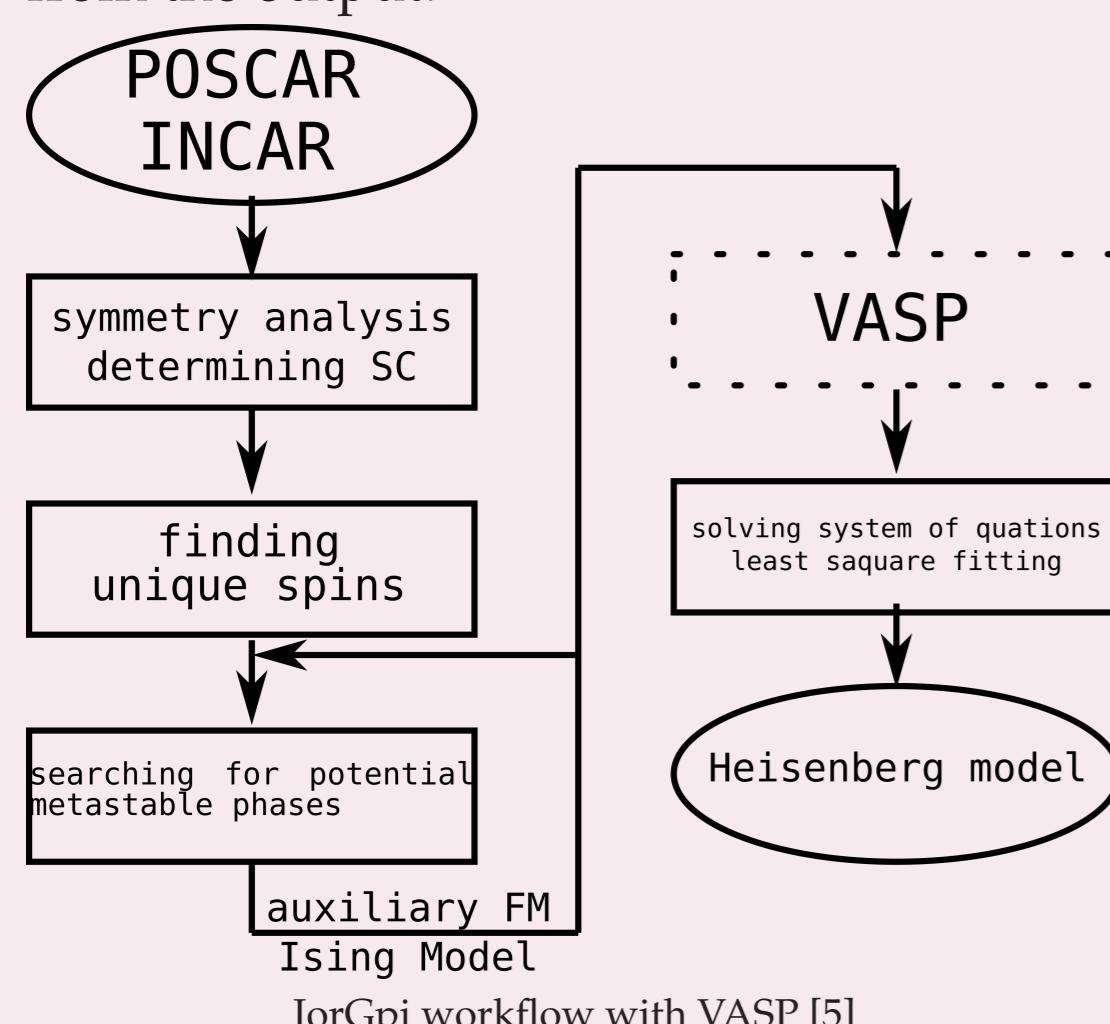
where  $J_{ij}$  are exchange interaction magnitudes of unitary vectors  $\{\mathbf{s}_i\}$  on a given lattice.

We are aiming at creating a **generic approach** removing the burden of designing the Hamiltonian from user while **minimising the computational cost** (in contrast to the Korringa-Kohn-Rostoker (KKR) Green function formalism [3] and the frozen-magnon approach [4]).

$\hookrightarrow$

## MAPPING SOFTWARE: JORG $\pi$

The aim is to process the **ab-initio** model of magnetic material, generate a number of possibly stable states with rotated spin, and calculate the **effective Heisenberg model** from the output.



Jorgpi requires Python 3.6 with  
 numpy 1.16.0, scipy 1.0.0, spglib 1.12.0, matplotlib 3.0.0, setuptools 40.8.0, defusedxml -  
 The bottleneck section - **Ising model solver** (cf. *Metastable states*) utilizes C++17 and GNU Scientific Library.

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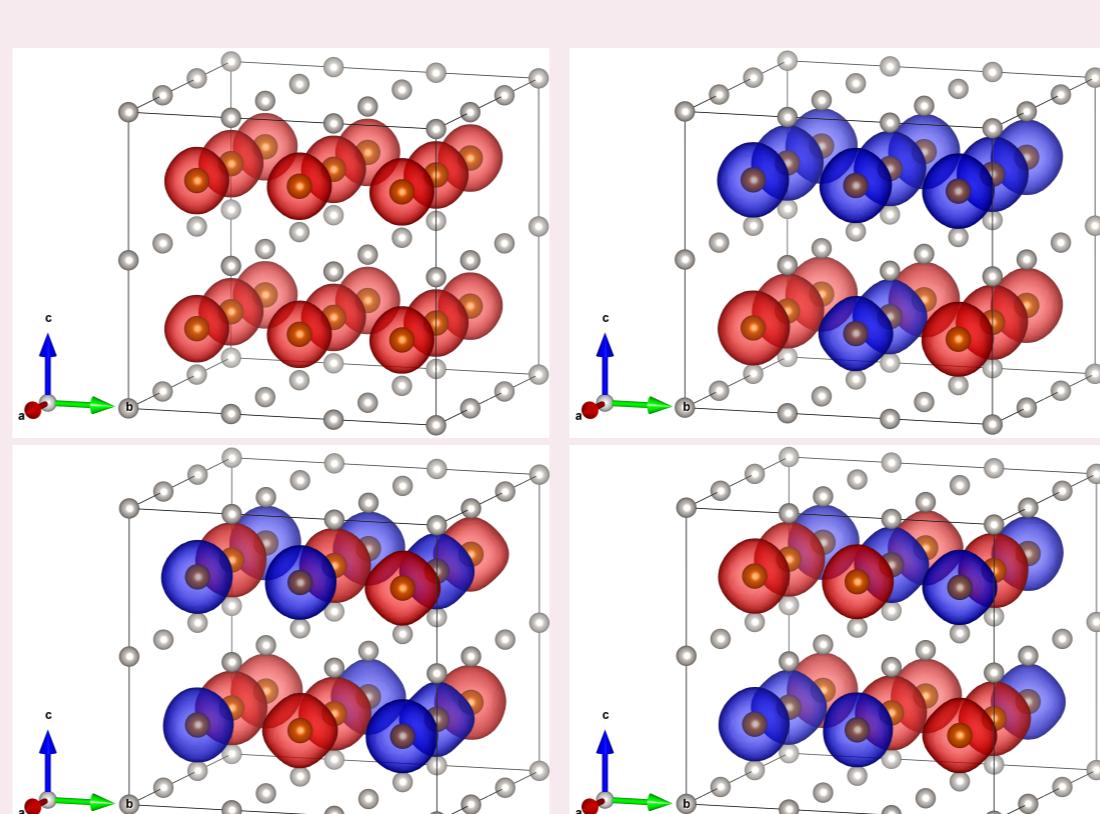
usage: __main__.py [-h] [-I INPUT] [-i INCAR] [-o OUTPUT] [-c cutoff] [-n neighbor] [-w Wyckoffs] [-r reference] [-e elements] [-g group] [-p period] [-b block] [-f P,D,F] [-x extra-dimensions]
                   [-s symmetry] [-r redundant] [-soc spin-orbit] [-refined refined] [-x extra-dimensions EXTRA-DIMENSIONS]

Find minimal number of unique spin flips
optional arguments:
  -h, --help            show this help message and exit
  -I INPUT, --input INPUT  input POSCAR file
  -i INCAR, --INCAR INCAR, -I INCAR           input INCAR file
  -o OUTPUT, --output OUTPUT   output directory
  -c cutoff, --R-CUTOFF   off distance (in Å) for calculations
  -n neighbor, --N-NEIGHBOR   a range of the last Neighbor taken into account
  -w Wyckoffs, --Wyckoffs   narrows down the atomic selection to the atoms in Wyckoff positions defined by string (eg. 'abc')
  -r reference, --reference REFERENCE   number of reference atom in inputfile
  -e elements ELEMENTS, --elements ELEMENTS   string of all elements taken into account (eg. 'CuO')
  -g group GROUP, --group GROUP   group number (eg. 1 => 'H1UNaRbCsF')
  -p period PERIOD, --period PERIOD   period of the supercell (eg. 2,3,4,5,6,7,8,9,10,11,12,13,14,15,16,17,18)
  -b block BLOCK, --block BLOCK   block name (eg. P ==> 'SRCSNNSDFSAIS1SPSSC1GosGeSa'
  -s symmetry, --S symmetry   creates a redundant system of equations for final calculation of the Heisenberg exchange interaction
  -r redundant, --redundant   (work-in-progress) is spin-orbit coupling enabled
  -soc spin-orbit, --SOC spin-orbit   should use refined supercell (default False)
  -refined refined, --refined refined   string 'X' or 'Z' of extra cell copies in each directions
  -x extra-dimensions EXTRA-DIMENSIONS, --extra-dimensions EXTRA-DIMENSIONS   string 'X' or 'Z' of extra cell copies in each directions
  
```

## EXEMPLARY RESULTS - FEPT

Space group: P4/mmm (123):  
 $a = b = 2.66084 \text{ \AA}; c = 3.68635 \text{ \AA}; \alpha = \beta = \gamma = 90^\circ$

	$\mu_{Fe} (\mu_B)$	$\mu_{Pt} (\mu_B)$		
FLEUR [6]	2.866	0.384		
VASP [5]	2.823	0.318		
Jorgpi	<b>2.849</b>	<b>0.315</b>		
Distance (Å)	FLEUR [6]	SPR-KKR [7]	Jorgpi	Ref. [8]
$J_{[100]}$ (meV)	2.661	25.73	31.52	<b>35.77</b>
$J_{[001]}$ (meV)	3.686	8.82	-3.72	<b>13.28</b>
$J_{[110]}$ (meV)	3.763	19.89	17.56	<b>19.24</b>
$J_{[101]}$ (meV)	4.546	8.20	8.58	<b>15.67</b>
$J_{[111]}$ (meV)	5.268	-15.32	-10.58	<b>-21.07</b>
				-19.75



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## SOME MATH & PHYSICS

We assume that there exist a block-diagonal effective Hamiltonian  $\hat{\mathcal{H}}_{\text{eff}}$  corresponding to the resultant state of our DFT calculations

$$\hat{\mathcal{H}}_{\text{eff}} = \begin{pmatrix} \hat{\mathcal{H}}_{\text{magnetic}} & \\ & \hat{\mathcal{H}}_{\text{remainedder}} \end{pmatrix}.$$

where  $\hat{\mathcal{H}}_{\text{magnetic}} = \hat{\mathcal{H}}(\{\hat{\mathbf{M}}_i\})$  depends on the on-site magnetization operators  $\{\hat{\mathbf{M}}_i\}$ . For the case of colinear-spin system this simplifies to

$$\Delta E_{n \in N} \stackrel{\Delta M \ll M}{\approx} \sum_{\langle i,j \rangle} -\mathbf{J}_{ij} M_i^0 M_j^n,$$

where we have a set of  $N$  metastable magnetic states with magnetization norm  $M_i^n \equiv M_i^0 + \Delta M_i^n$  ( $i, j$  are sites with exactly one spin flipped).

## METASTABLE STATES

range	# sites	# possible states
1 NN	2	$2^1 - 1 = 1$
2 NN	8	$2^7 - 1 = 127$
3 NN	16	$2^{15} - 1 = 32767$
5 NN	54	$2^{53} - 1 \sim 10^{16}$
8 NN	128	$2^{127} - 1 \sim 10^{38}$
13 NN	250	$2^{249} - 1 \sim 10^{75}$
21 NN	432	$2^{431} - 1 \sim 10^{130}$

Number of possible states (assuming spin-reversal symmetry) of a Bravais bcc lattice versus range of interactions.

In the idealized magnetic ground state the ordering of magnetic moment directions  $\mathbf{s}_i^0$  is **unambiguously defined by sign** of exchange interaction magnitudes  $J_{ij}$ .

I.e., to find metastable states we use a **ferromagnetic (!)** 3D Ising model with a solver using Simulated Annealing.

$$H = - \sum_{i \neq j} K_{ij} \sigma_i \sigma_j, \quad \& K_{ij} \equiv J_{ij} \mathbf{s}_i^0 \cdot \mathbf{s}_j^0 > 0.$$

## ACKNOWLEDGMENTS

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