

Notatki do ćwiczeń

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1 Brakety

$$|\psi\rangle = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \end{pmatrix} = \sum_i \psi_i |i\rangle, \quad \text{baza } \{|i\rangle\}: \text{ortonormalna: } \langle i|j\rangle = \delta_{ij}, \text{zupełna: } \sum_i |i\rangle\langle i| = \mathbb{1}; \quad (1)$$

$$\langle\psi| = \begin{pmatrix} \psi_1^* & \psi_2^* & \psi_3^* & \dots \end{pmatrix} = \sum_i \psi_i^* \langle i|; \quad (2)$$

$$\langle\psi|\phi\rangle = \begin{pmatrix} \psi_1^* & \psi_2^* & \psi_3^* & \dots \end{pmatrix} \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \\ \vdots \end{pmatrix} = \left(\sum_i \psi_i^* \langle i| \right) \left(\sum_j \phi_j |j\rangle \right) = \sum_{ij} \psi_i^* \phi_j \delta_{ij} = \sum_i \psi_i^* \phi_i \quad (3)$$

$$|\psi\rangle\langle\phi| = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \vdots \end{pmatrix} \begin{pmatrix} \phi_1^* & \phi_2^* & \phi_3^* & \dots \end{pmatrix} = \begin{pmatrix} \psi_1 \phi_1^* & \psi_1 \phi_2^* & \psi_1 \phi_3^* & \dots \\ \psi_2 \phi_1^* & \psi_2 \phi_2^* & \psi_2 \phi_3^* & \dots \\ \psi_3 \phi_1^* & \psi_3 \phi_2^* & \psi_3 \phi_3^* & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix} \quad (4)$$

2 Operatory

Element macierzowy operatora:

$$A_{nm} = \langle n|\hat{A}|m\rangle \quad (5)$$

Operator w bazie $\{|n\rangle\}$:

$$\hat{A} = \sum_{nm} \langle n|\hat{A}|m\rangle |n\rangle\langle m| = \sum_{nm} A_{nm} |n\rangle\langle m| = \begin{pmatrix} \langle 1|\hat{A}|1\rangle & \langle 1|\hat{A}|2\rangle & \langle 1|\hat{A}|3\rangle & \dots \\ \langle 2|\hat{A}|1\rangle & \langle 2|\hat{A}|2\rangle & \langle 2|\hat{A}|3\rangle & \dots \\ \langle 3|\hat{A}|1\rangle & \langle 3|\hat{A}|2\rangle & \langle 3|\hat{A}|3\rangle & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix} \quad (6)$$

Działanie operatora na ket:

$$\hat{A}|\psi\rangle = \left(\sum_{nm} A_{nm} |n\rangle\langle m| \right) \left(\sum_k \psi_k |k\rangle \right) = \sum_n \underbrace{\sum_m A_{nm} \psi_m}_{(\hat{A}|\psi\rangle)_n} |n\rangle = |\psi'\rangle \quad (7)$$

Wartość oczekiwana w stanie $|\psi\rangle$:

$$\langle A \rangle_\psi = \langle A \rangle = \langle \psi | A | \psi \rangle = \sum_{nm} A_{nm} \langle \psi | n \rangle \langle m | \psi \rangle = \sum_{nm} \sum_{ij} A_{nm} \psi_i^* \delta_{in} \psi_j \delta_{mj} = \sum_{nm} \psi_n^* A_{nm} \psi_m \quad (8)$$

Zagadnienie własne:

$$\hat{A}|a\rangle = a|a\rangle \quad (9)$$

Operator w bazie stanów własnych:

$$\hat{A} = \sum_{aa'} \langle a | \hat{A} | a' \rangle |a\rangle \langle a'| = \sum_{aa'} a \delta_{aa'} |a\rangle \langle a'| = \sum_a a |a\rangle \langle a| \quad (10)$$

3 (Anty-)commutator

Własności komutatora $[A, B] = AB - BA$ and anty-komutatora $\{A, B\} = AB + BA$:

1. $[\alpha A + \beta B, C] = \alpha[A, C] + \beta[B, C]$, gdzie α, β to stałe;
2. $[A, B] = -[B, A]$, $\{A, B\} = \{B, A\}$;
3. $[AB, C] = A[B, C] + [A, C]B$;
4. $[AB, C] = A\{B, C\} - \{A, C\}B$.

4 Kwantowanie kanoniczne

$$[\hat{x}_i, \hat{p}_j] = i\hbar \delta_{ij}, \quad [\hat{x}_i, \hat{x}_j] = [\hat{p}_i, \hat{p}_j] = 0; \quad (11)$$

$$\text{operator położenia: } \hat{x}_i = x_i; \quad (12)$$

$$\text{operator pędu: } \hat{p}_i = -i\hbar \frac{\partial}{\partial x_i}; \quad (13)$$

5 Zasada nieoznaczoności

Dla hermitowskich operatorów: $A = A^\dagger, B = B^\dagger$ ($\Delta A = \sqrt{\langle A - \langle A \rangle \rangle^2}$):

$$\Delta A \Delta B \geq |\langle \frac{1}{2i} [A, B] \rangle| \quad (14)$$

6 Zmiana bazy—transformacje

Dane są dwie zupełne, ortonormalne bazy $\{|i\rangle\}_{i=1,2,3,\dots}$ oraz $\{|\tilde{i}\rangle\}$ oraz transformacja unitarna U :

$$U|i\rangle = |\tilde{i}\rangle, \quad U = \sum_j |\tilde{j}\rangle \langle j| \quad (15)$$

Elementy macierzowe U :

$$U_{ij} = \langle i | \left(\sum_{j'} |\tilde{j}'\rangle \langle j'| \right) | j \rangle = \langle i | \tilde{j} \rangle \quad (16)$$

Postać U w bazie $|i\rangle$

$$U = \sum_{ij} U_{ij} |i\rangle \langle j| = \sum_{ij} \langle i | \tilde{j} \rangle |i\rangle \langle j| = \begin{pmatrix} \langle 1 | \tilde{1} \rangle & \langle 1 | \tilde{2} \rangle & \langle 1 | \tilde{3} \rangle & \dots \\ \langle 2 | \tilde{1} \rangle & \langle 2 | \tilde{2} \rangle & \langle 2 | \tilde{3} \rangle & \dots \\ \langle 3 | \tilde{1} \rangle & \langle 3 | \tilde{2} \rangle & \langle 3 | \tilde{3} \rangle & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix} \quad (17)$$

Transformacja stanu $|\psi\rangle = \sum_i \psi_i |i\rangle$:

$$|\psi\rangle = \sum_i \left(\underbrace{\sum_j |\tilde{j}\rangle \langle \tilde{j}|}_{\mathbb{1}} \right) \psi_i |i\rangle = \sum_j \sum_i \underbrace{\langle \tilde{j}|i\rangle}_{U_{ji}^\dagger} \psi_i |\tilde{j}\rangle = \sum_j (U^\dagger |\psi\rangle)_j |\tilde{j}\rangle \quad (18)$$

Transformacja U jest unitarna i zachowuje normę $|\psi\rangle$:

$$UU^\dagger = \left(\sum_i |\tilde{i}\rangle \langle i| \right) \left(\sum_j |j\rangle \langle \tilde{j}| \right) = \sum_{ij} |\tilde{i}\rangle \delta_{ij} \langle \tilde{j}| = \mathbb{1}; \quad \langle \tilde{\psi} | \tilde{\psi} \rangle = \langle \psi | U^\dagger U | \psi \rangle = \langle \psi | \psi \rangle \quad (19)$$

Transformacja operatora: