

Observation of  $\beta$ -Amyloid Peptide Oligomerization by Pressure-Jump NMR SpectroscopyC. Ashley Barnes,<sup>§</sup> Angus J. Robertson,<sup>§</sup> John M. Louis, Philip Anfinrud, and Ad Bax\*<sup>ID</sup>

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## Supporting Information

**ABSTRACT:** Brain tissue of Alzheimer's disease patients invariably contains deposits of insoluble, fibrillar aggregates of peptide fragments of the amyloid precursor protein (APP), typically 40 or 42 residues in length and referred to as  $A\beta^{40}$  and  $A\beta^{42}$ . However, it remains unclear whether these fibrils or oligomers constitute the toxic species. Depending on sample conditions, oligomers can form in a few seconds or less. These oligomers are invisible to solution NMR spectroscopy, but they can be rapidly (<1 s) resolubilized and converted to their NMR-visible monomeric constituents by raising the hydrostatic pressure to a few kbar. Hence, utilizing pressure-jump NMR, the oligomeric state can be studied at residue-specific resolution by monitoring its signals in the monomeric state. Oligomeric states of  $A\beta^{40}$  exhibit a high degree of order, reflected by slow longitudinal  $^{15}\text{N}$  relaxation ( $T_1 > 5$  s) for residues 18–21 and 31–34, whereas the N-terminal 10 residues relax much faster ( $T_1 \leq 1.5$  s), indicative of extensive internal motions. Transverse relaxation rates rapidly increase to ca. 1000  $\text{s}^{-1}$  after the oligomerization is initiated.

Amyloid diseases are characterized by the presence of protease-resistant aggregates of peptides or proteins. These disorders occur spontaneously, mostly later in life, and are associated with a wide range of ailments, ranging from Alzheimer's and Parkinson's disease to type 2 diabetes and amyloidosis.<sup>1</sup> It remains a matter of debate whether the mature fibrils or smaller oligomeric aggregates are the dominant toxic species *in vivo*. Alternatively, oligomers in dynamic equilibrium with mature fibrils have been implicated as the cause of neuronal dysfunction.<sup>2</sup> Atomic resolution models of many types of mature fibrils have been derived from solid state NMR data<sup>3–6</sup> and, more recently, cryo-electron microscopy.<sup>7,8</sup> These structures are characterized by the presence of canonical cross- $\beta$  structure, but exhibit a wide diversity of arrangements,<sup>3,6,9</sup> which in some cases has been correlated with disease phenotype.<sup>10,11</sup>

Despite their strong link to cellular pathology,<sup>2,12–14</sup> considerably less structural information is available about the oligomeric species. Small clusters of monomers have a high surface-to-volume ratio and therefore a high interfacial energy with water, making them inherently unstable. However, once present, nucleation theory predicts their rapid further growth, provided that the concentration of monomers is sufficiently

high.<sup>15</sup> Detailed structural characterization of such a fleeting species therefore presents a fundamental experimental challenge. Here, we introduce technology that rapidly switches  $A\beta^{40}$  between its monomeric and oligomeric states, permitting the repeated observation of oligomer formation and growth while providing structural details on this enigmatic process.

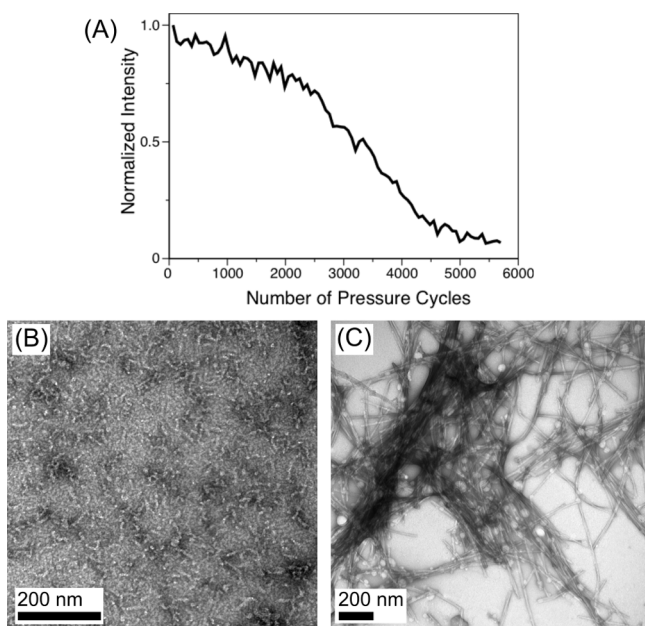
Hydrostatic pressure has long been used to denature natively folded proteins, a process attributed to the smaller molar volume of the unfolded chain relative to the native protein fold.<sup>16,17</sup> Whereas several types of amyloid fibrils similarly can be reverted to their soluble monomeric states by the application of a few kbar of pressure, this tends to be a slow process, requiring hours.<sup>18–20</sup> By contrast, we demonstrate that  $A\beta^{40}$  can not only oligomerize very rapidly, but it can also be converted back rapidly to the monomeric, intrinsically disordered state. Analogous to the study of protein folding, oligomer formation can then be tracked by recording NMR spectra<sup>21</sup> while repeatedly jumping between low pressure, which favors the aggregated state, and high pressure, where the oligomer melts and reverts to its monomeric, disordered state.

For natively folded proteins, the unfolding/refolding cycle can often be repeated indefinitely; for  $A\beta^{40}$  the process is repeatable only a finite number of times before a pressure-resistant, insoluble, fibrillar species forms and depletes the NMR-visible monomeric species (Figure 1). NMR measurements were carried out at high concentrations of peptide (ca. 1.3 mM), at room temperature, at relatively high ionic strength, and pH 6.0 (see SI), conditions that reduce the oligomerization time to seconds at 1 bar.

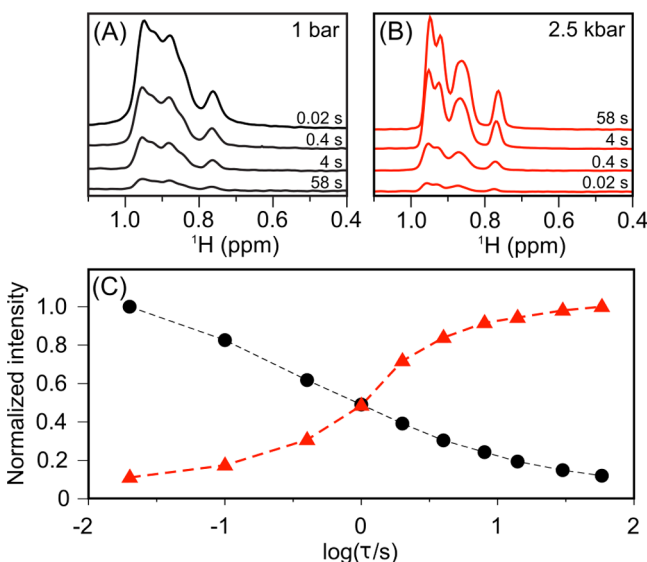
Protein folding studies by pressure-jump NMR are most informative when the times needed for folding and unfolding are much shorter than the longitudinal relaxation time of the protein's  $^{15}\text{N}$  nuclei.<sup>21</sup> Resonances of the unfolded protein recorded following a pressure jump then can report on the state of the (partially) folded protein while it develops at low pressure.<sup>22</sup> Under our conditions, 1 s at low pressure sufficed for extensive  $A\beta^{40}$  oligomer formation, with a comparable time needed for the aggregate to return to its monomeric form after jumping back to high pressure (Figure 2). The short (~600 ms)  $^{15}\text{N}$   $T_1$  relaxation times of the monomeric peptide prevent direct application of the recently developed pressure-jump protein folding experiments due to the inability to separate signals originating from peptides that oligomerized from peptides that remained free in solution at 1 bar. However,

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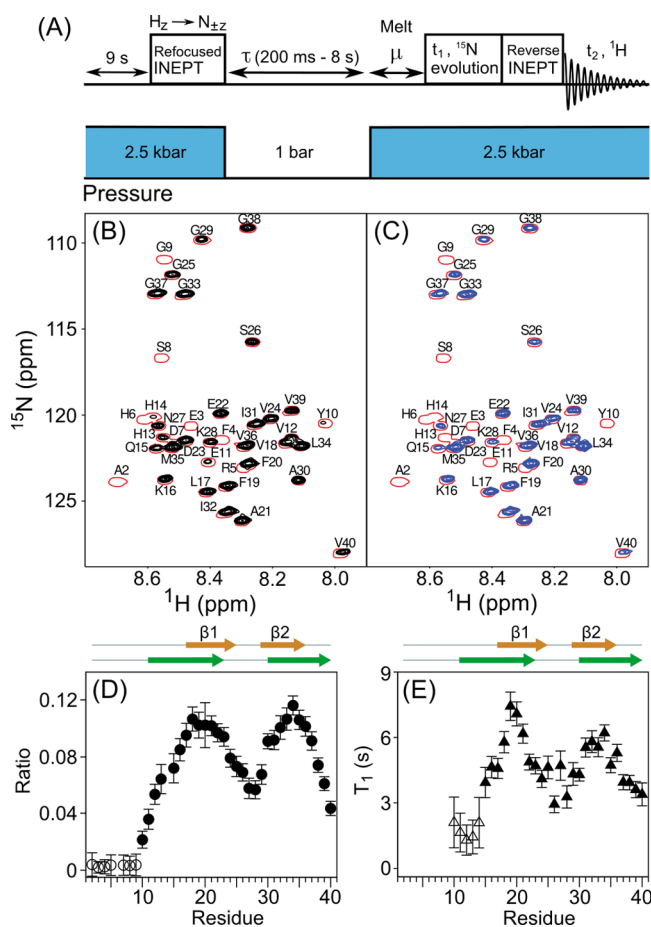
**Figure 1.** Loss of monomeric peptide signal and formation of  $A\beta^{1-40}$  ordered structures during pressure-jump NMR experiments. (A) Decay of  $A\beta^{40}$  heteronuclear single quantum coherence (HSQC) intensity, recorded 9.8 s after jumping to 2.5 kbar, as a function of the number of high-pressure/low-pressure (10 s/5 s) cycles. The last time point corresponds to 24 h after the start of the pressure-jump NMR experiment. (B, C) Transmission electron microscopy (TEM) images of negatively stained samples harvested at atmospheric pressure after a single pressure cycle (B) or after 12 000 pressure cycles (C).



**Figure 2.** Time course of  $A\beta^{40}$  aliphatic proton resonance intensity following a sudden change in pressure. (A) Spectra recorded at the indicated times after a rapid drop in pressure from 2.5 kbar to 1 bar following a 30 s sample equilibration at 2.5 kbar. (B) Spectra recorded after a jump in pressure to 2.5 kbar following 30 s equilibration at 1 bar. (C) Integrated methyl group resonance intensity (1.05–0.7 ppm) as a function of time following the pressure drop (black) or pressure jump (red).

whereas  $^{15}\text{N}$   $T_1$  values of disordered peptides and small proteins are short,  $^{15}\text{N}$  relaxation times in solids or in very large, slowly tumbling protein systems can be many seconds or even minutes.<sup>23,24</sup> Therefore, oligomeric states formed

following a pressure drop will preserve a substantial fraction of their nonequilibrium  $^{15}\text{N}$  nuclear spin polarization long after that of the monomeric  $A\beta^{40}$  peptide has vanished. At the end of the  $\sim 5$  s low pressure interval (Figure 3A), the pressure is



**Figure 3.** Observation of  $A\beta^{40}$   $^1\text{H}$ – $^{15}\text{N}$  HSQC signals during a  $T_1$ -filtered pressure-jump NMR experiment. (A) Schematic timing diagram; see SI for details. Strong, nonequilibrium  $^{15}\text{N}$   $z$  magnetization is generated by a refocused INEPT transfer<sup>28</sup> just prior to dropping the pressure. After a variable low-pressure delay (0.2–8 s) and a 350 ms “melting” interval,  $\mu$ , at high pressure,  $^{15}\text{N}$   $t_1$  evolution and  $^1\text{H}$  detection are used to generate conventional  $^{15}\text{N}$ – $^1\text{H}$  HSQC spectra. (B) Overlays of spectra recorded with a 0.2 s low-pressure duration (single red contour), and resonances that remain after a 5.5 s low-pressure interval (black contours). (C) Overlay of the spectra recorded with 0.2 s (red) and 8 s (blue) low-pressure intervals. (D) Intensity ratios observed in (B) vs residue number. Open symbols correspond to an upper limit for the nonobserved intensity. (E)  $^{15}\text{N}$  longitudinal relaxation time ( $T_1$ ) in the oligomeric state, derived from the intensity ratio of resonances observed with 5.5 s and 8 s low pressure durations.  $\beta$ -Strand positions are from Hoyer et al.<sup>25</sup> (orange) and Paravastu et al.<sup>3</sup> (green).

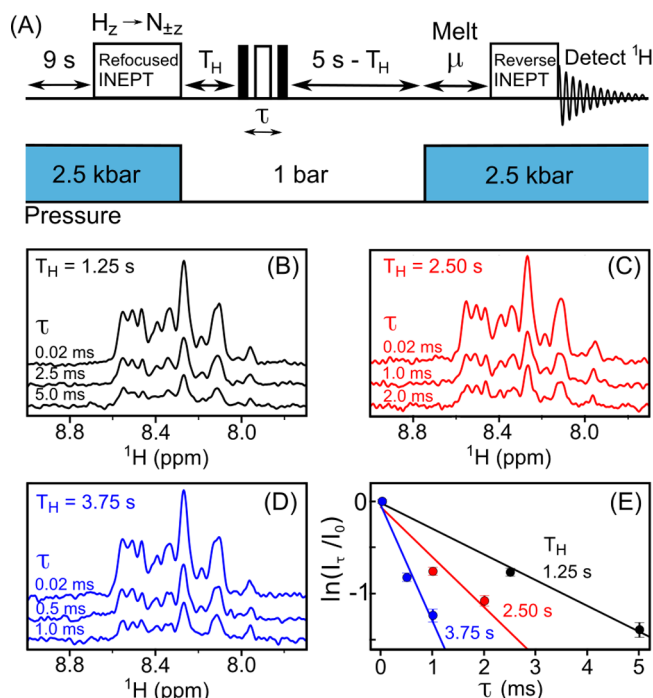
jumped back to high, causing the oligomers to redissolve. While some relaxation occurs during this 350 ms “melting period”, the fraction that is retained during the subsequent  $^{15}\text{N}$   $t_1$  evolution period will be visible in the two-dimensional  $^{15}\text{N}$ – $^1\text{H}$  correlation map, and can therefore unveil structural properties of the oligomeric state. Effectively, the ca. 5 s low pressure duration between initial  $^{15}\text{N}$  polarization and HSQC detection serves as a  $T_1$  relaxation filter: residues that become immobilized due to oligomerization during the low-pressure

interval retain a substantial fraction of their signal and remain visible to NMR; peptides that remain monomeric are attenuated more than 1000-fold and are therefore invisible.

Comparison of  $A\beta^{40}$   $^{15}\text{N}$ – $^1\text{H}$  correlation spectra recorded with either a 200 ms or a 5.5 s low-pressure duration shows the selective loss of signals from residues A2–G9 that did not rapidly gain a substantial degree of order during the low-pressure period (Figure 3B,D), contrasting with high intensities, in particular for L17–D23 and I30–G37. Longitudinal  $^{15}\text{N}$  relaxation in the very slowly tumbling, oligomeric species is dominated by internal motions. These  $T_1$  values can be derived from spectral intensities obtained with 5.5 s and 8 s low pressure intervals (2-point  $T_1$  measurement, Figure 3C). As expected, on average these  $T_1$  values correlate with the fractional signal intensity recovered after the  $T_1$  filter (Figure 3D,E). However, the rates at which different residues lose mobility during oligomerization varies. For example, the  $T_1$  values of C-terminal residues G37–V40 are all in the 3.5–4 s range, whereas the fraction of recovered magnetization is 2-fold higher for G37 than for V40. This result indicates that V40 becomes ordered later than G37, i.e., only after the oligomer has increased in size. The highest fraction of recovered magnetization, corresponding to residues that become most ordered at an early stage of the oligomerization process, coincides with the two antiparallel strands in the X-ray structure of an affibody-trapped  $A\beta^{40}$ ,<sup>25</sup> which has been linked to structural features seen in oligomers and protofibrils.<sup>26</sup>

Growth of the oligomeric species can be monitored by  $^{15}\text{N}$   $R_2$  transverse relaxation measurements. In the macromolecular limit, these rates are dominated by  $J(0)$  spectral density terms, which means that  $R_2$  values are proportional to  $S^2\tau_c$ , where  $S^2$  is the generalized Lipari–Szabo order parameter for internal motion,<sup>27</sup> ranging between 0 (total disorder) and 1 (fully rigid), and  $\tau_c$  is the rotational correlation time. Only residues with long  $T_1$   $^{15}\text{N}$  relaxation times are observed in our measurements, in practice requiring  $S^2 \geq \sim 0.7$ . Therefore,  $R_2$  provides a direct measure for  $\tau_c$  of the oligomers during their growth phase.

$^{15}\text{N}$   $R_2$  values were probed at three time points after the initiation of oligomerization by the pressure drop (1.25, 2.5, and 3.75 s) by inserting a Hahn-echo  $R_2$  block within the 5 s low-pressure  $T_1$  filter (Figure 4A). Even at the early, 1.25 s time point, transverse relaxation occurs at a rapid rate of ca.  $275\text{ s}^{-1}$ , about 50-fold faster than for the monomeric peptide (Figure 4). Although signal-to-noise was insufficient to measure the  $R_2$  values in a residue-specific manner (which would require 2D spectra), the decay of the various resonances in the one-dimensional spectra appears quite homogeneous. This makes it unlikely that conformational exchange, which typically shows large residue-by-residue variations, has a strong impact on these  $R_2$  rates. Assuming an average  $S^2 \approx 0.85$  value and an assumption of approximately isotropic tumbling,  $R_2 = 275\text{ s}^{-1}$  yields an estimate for the tumbling time of  $\tau_c \approx 200\text{ ns}$ , or a particle size of ca. 350 kDa, which corresponds to a mass of roughly 80 peptides. When doubling the time interval at which the Hahn-echo block is inserted to 2.5 s after the pressure drop,  $R_2$  has increased to ca.  $600\text{ s}^{-1}$ . A further increase to  $R_2 \approx 1000\text{ s}^{-1}$  is seen at  $T = 3.75\text{ s}$ , a value pointing to a size of ca. 1.3 MDa, or about 300 peptides. The transverse decay appears somewhat nonexponential, with the decay slightly faster at short echo times,  $\tau$ , and slower at longer  $\tau$  values. This behavior is seen at all three  $T_H$  values and indicates that aggregate sizes are heterogeneous, with reported



**Figure 4.** Measurement of transverse  $^{15}\text{N}$  relaxation rates of  $A\beta^{1-40}$  oligomers during a pressure-jump NMR experiment. (A) Schematic timing diagram. The scheme differs from Figure 3A by insertion of a Hahn-echo block of duration  $\tau$  at time  $T_H$  into the low-pressure interval, and the absence of a  $t_1$  evolution period. See SI for details. (B–D) Amide regions of the  $^1\text{H}$  spectra detected for different Hahn-echo delay durations,  $\tau$ , inserted at (B)  $T_H = 1.25\text{ s}$ , (C)  $T_H = 2.5\text{ s}$ , and (D)  $T_H = 3.75\text{ s}$  after the drop to low pressure. Each spectrum results from 256 transients, recorded in an interleaved manner, such that, to a very good approximation, sample aging affects all spectra equally. (E) Plots of the intensity decay observed in panels B–D as a function of  $\tau$ .

values reflecting rough averages under the conditions of our measurement.

Our pressure-jump NMR experiments provide a residue-specific recording of the initiation and growth of  $A\beta^{40}$  oligomers.  $A\beta$  aggregates can adopt different structures, which may relate to disease phenotype,<sup>10,11</sup> and as demonstrated here, even pressure-resistant amyloids can develop. Morphologically, the  $A\beta^{40}$  oligomers harvested after a single pressure drop on a fresh  $250\text{ }\mu\text{M}$  sample (Figure S3), have an appearance that does not show regular amyloid fibril structure and is more reminiscent of the features seen in TEM images of oligomers taken after rapidly neutralizing a sample initially prepared at pH 12.<sup>29</sup> Furthermore, these oligomers formed after a single pressure-jump do not dissolve after a 5-fold dilution unless the pressure is raised to 2.5 kbar (Figure S4). Given their morphological similarity to the pH-quenched oligomers<sup>29</sup> (as assessed via TEM images), the pressure-jump oligomers likely contain antiparallel  $\beta$ -strand arrangements, thereby creating a high energy barrier toward forming the parallel  $\beta$ -sheet arrangements seen in mature fibrils. The ability of pressure-jump NMR to provide site-specific information on the very earliest processes of aggregation opens fundamental new opportunities to study this critical process. Pressure-jump NMR in conjunction with nuclear Overhauser effect measurement, paramagnetic labeling, and stroboscopic chemical shift



measurements of  $A\beta^{40}$  are currently in progress to gain further structural information.

## ■ ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/jacs.9b06970.

Experimental procedures and supporting figures with additional data (PDF)

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### Notes

The authors declare no competing financial interest.

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