

Thermo-Elasticity for Large Deformations

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Simulation Of Coupled Problems Using Finite Element Method

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Overview

- Introduction
- Motivation- Use in industries
- Why Non Linear?
- Equation Formulation
- Problem Formulation
- Results:
 - Pure Mechanical Loading
 - Pure Thermal Loading
 - Mixed Loading
 - Mesh convergence study for Maximum displacement

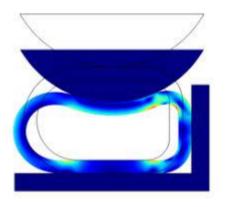
Introduction

Introduction

- Thermo-elasticity by definition addresses:
 - Mechanical deformation Elastic deformation
 - Thermal effects Temperature gradients, Heat flux
 - Their couplings Thermal expansion
- Hyper-elastic materials posses certain characteristics:
 - Fully recoverable large elastic deformations in order of 100 to 700 %
 - They show a highly nonlinear stress-strain relation
 - > Tension The material softens and then stiffens again
 - Compression A stiff response (incompressible)



Rubber



Hyper-elastic Seal

Motivation: Use in industries

Motivation: Use in industries

- Flexible adhesive joints, with very low elastic modulus and very large failure deformation
- Elastomeric pads in bridges, rail pads, car door seal, car tires
- Major stresses in Engine cylinder walls are due to temperature variation and mechanical impact
- Stresses experienced by a Pressure cooker are due to internal pressure and temperature variation
- Pressurized storage containers for liquified gases
- Industrial chimneys Self weight, pressure and temperature variation
- Turbo machinery High pressure leads to increase in temperature
- Casting and forging machinery
- Bird impact in aerospace

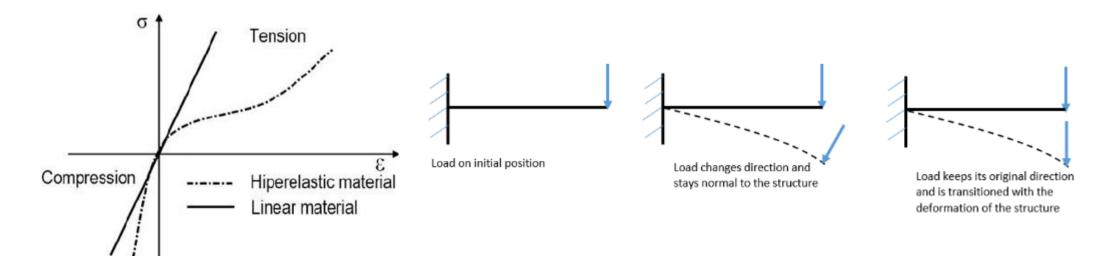




Why Non Linear?

Why Non Linear?

- Small vs large deformations
- Hyper-elastic materials (rubber and plastic) show a highly nonlinear stress-strain relation.
- Changes in geometry as the structure deforms are considered in formulating the constitutive and equilibrium equations.
- In case of contact we can have non linear boundary conditions (kinematic constraints).
 - As opposed to a linear static analysis, where the stiffness matrix remains constant, these effects result in a stiffness matrix which is not constant during the load application. As a result, a different solving strategy is required Newton Raphson Solver.



General structure of balance relations

- Local differential equations valid at every material point
 - Balance of mass

Assumption: Quasi-static : No acceleration
$$\rightarrow$$

- Balance of moment of momentum
- Balance of energy
- Entropy inequality (CDI)

$$\rho + \rho div \dot{\underline{x}} = 0$$

$$\int_{B_o} (Div \, \underline{P} + \rho_o \times (\underline{b} - \underline{\ddot{x}})) dV_o = \underline{0}$$

$$\int_{B_o} Div \, \underline{P} dV_o = \underline{0}$$

$$\underline{T} = \underline{T}^T$$

$$\int_{B} \rho \dot{\epsilon} \, dv = \int_{B} \left(\underline{T} \cdot \underline{D} + \rho r - div \underline{q} \right) dv$$

$$-\rho(\dot{\varphi} + \dot{\theta\eta}) + \underline{T} \cdot \underline{D} - \frac{1}{\theta} \underline{q} \cdot grad\theta \ge 0$$

Material Modelling

- Thermodynamic consistent constitutive modelling
- Axioms of Thermodynamics
 - Principle of determinisms
 - Principle of Equipresence
 - Principle of local action
 - Material Frame indifference

 $\underline{T} = 2\rho \underline{F} \frac{\partial \varphi}{\partial \underline{C}} \underline{F}^{T}$ $\eta = -\frac{\partial \varphi}{\partial \theta}$ $\frac{\partial \varphi}{\partial grad\theta} = 0 \quad \Longrightarrow \quad \varphi = \varphi(\underline{C}, \theta)$ $-\frac{1}{\theta} q \cdot grad\theta \ge 0$

- Evaluation of entropy inequality
- Helmholtz free energy function for isotropic hyper elastic materials

$$\psi = \frac{1}{\rho_o} \left[\frac{1}{2} \lambda (\ln J)^2 - \mu \ln J + \frac{1}{2} \mu \left(\frac{I_c}{3} \right) - 3 \alpha \kappa (\ln J) (\theta - \theta_o) - \rho_o c \left(\theta \ln \frac{\theta}{\theta_o} - \theta + \theta_o \right) \right]$$

Where,

 $\lambda, \mu \rightarrow$ Material constants (Lame constants)

 $\alpha \rightarrow$ Thermal expansion coefficient, c \rightarrow Heat capacity

 $\kappa \rightarrow$ Compression Modulus , $I_c \rightarrow tr \underline{C}$

Balance of Momentum

Weak Form:

$$\int_{B_o} Div \, \underline{P}. \, \delta \underline{u} dV_o = \underline{0}$$

Final Form →

$$\int_{B_o} \underline{S} \cdot \delta \underline{E} dV_o - \int_{\partial B_o} \underline{t_0} \cdot \delta \underline{u} dA_o = 0$$

Discretized Weak Form

$$G_{u} = \sum_{i=1}^{J} \delta \underline{du}^{I} \left[\int_{B_{o}} B^{T}(\underline{u}, \underline{X}) \cdot \underline{S} dV_{o} - \int_{\partial B_{o}} N^{I} \underline{t} dA_{o} \right]$$

Linearized Weak Form

$$\Delta DG_{u} = \sum^{I} \sum^{J} \delta \underline{du}^{I} \left[\int_{B_{0}} \left(B^{T^{I}}(\underline{u}, \underline{X}) \cdot \frac{\partial \underline{S}}{\partial \underline{E}} \cdot B^{J}(\underline{u}, \underline{X}) \right) dV_{o} \Delta d\underline{u}^{J} \right] - \sum^{I} \sum^{J} \delta \underline{du}^{I} \left[\int_{B_{0}} \left(B^{T^{I}}(\underline{u}, \underline{X}) \cdot \frac{\partial \underline{S}}{\partial \theta} \cdot N^{J}(\underline{X}) \right) dV_{o} \Delta d\theta^{J} \right]$$

Balance of Momentum

•
$$\underline{S} = \mu(\underline{I} - \underline{C}^{-1}) + \lambda \ln(J) \underline{C}^{-1} - 3\alpha_s k(\theta - \theta_0) \underline{C}^{-1}$$

$$\begin{array}{l} \bullet \quad \frac{d\underline{S}}{d\underline{E}} = -\mu \frac{d\underline{C}^{-1}}{d\underline{E}} - \lambda lnJ \frac{d\underline{C}^{-1}}{d\underline{E}} - \lambda \frac{dlnJ}{d\underline{E}} \otimes \underline{C}^{-1} - 3\alpha_S \kappa (\theta - \theta_o) \frac{d\underline{C}^{-1}}{d\underline{E}} \\ \\ \text{The derivatives} \rightarrow \qquad \qquad \frac{d\underline{C}^{-1}}{d\underline{E}} = - \big[C_{lm}^{-1} C_{nk}^{-1} + C_{ln}^{-1} C_{mk}^{-1} \big] \\ \\ \frac{dlnJ}{d\underline{E}} \otimes \underline{C}^{-1} = \underline{C}^{-1} \otimes \underline{C}^{-1} \\ \\ \frac{dS_{lk}}{dE_{mn}} = \left[\mu - \lambda lnJ + 3\alpha_S \kappa (\theta - \theta_o) \right] \big[C_{lm}^{-1} C_{nk}^{-1} + C_{ln}^{-1} C_{mk}^{-1} \big] + \lambda C_{lk}^{-1} C_{mn}^{-1} \end{aligned}$$

•
$$\frac{\partial S}{\partial \theta} = -3\alpha_S k \underline{C}^{-1}$$

Balance of Energy

Weak Form:

$$\int_{B_o} \underline{q}_o \cdot Grad\delta\theta dV_o - \int_{B_o} (\rho_o c\dot{\theta} + 3\alpha_s \kappa\theta \underline{D} \cdot \underline{I} - \rho_o r) \delta\theta dV_o = \int_{\partial B_o} \underline{q}_0 \cdot \underline{N} \delta\theta dA_o$$
 Where ,
$$q_o = -J\alpha_\theta Grad\theta \underline{C}^{-1}$$

Discretized Weak Form

$$G_{\theta} = \sum^{I} \delta d\theta^{I} \left[\int_{B_{o}} B^{I} \left(\underline{\underline{u}}, \underline{X} \right) \cdot q_{0,lin} dV_{o} - \int_{B_{o}} N^{I} \rho_{0} C \theta dV_{o} - \int_{B_{o}} N^{I} 3\alpha_{s} \kappa \theta tr \underline{\dot{E}} dV_{o} - \int_{B_{o}} N^{I} \rho_{0} r dV_{o} - \int_{B_{o}} N^{I} q_{0,lin} \underline{N} dA_{o} \right]$$

Linearized Weak Form

$$\begin{split} \Delta D G_{\theta} &= \sum_{I}^{I} \sum_{J}^{J} \delta d\theta^{I} \int_{B_{o}} \left(B^{T^{I}}(\theta) . \frac{\partial q_{\{o,lin\}}}{\partial G r a d\theta} . B^{J}(\theta) \right) dV_{o} \Delta d\theta^{J} - \sum_{I}^{I} \sum_{J}^{J} \delta d\theta^{I} \int_{B_{o}} \left(N^{I} \frac{\partial Bal \; Energy}{\partial \dot{\theta}} \frac{\delta}{\beta \Delta t} . N^{J} \right) dV_{o} \Delta d\theta^{J} \\ &- \sum_{I}^{I} \sum_{J}^{J} \delta d\theta^{I} \int_{B_{o}} \left(N^{I} \frac{\partial Bal \; Energy}{\partial \theta} N^{J} \right) dV_{o} \Delta d\theta^{J} - \sum_{I}^{I} \sum_{J}^{J} \delta \theta^{I} \int_{B_{o}} \left(N^{I} \frac{\partial Bal \; Energy}{\partial \dot{\theta}} \frac{\delta}{\beta \Delta t} . B^{J}(\underline{u}, \underline{X}) \right) dV_{o} \Delta d\underline{u}^{J} \end{split}$$

Balance of Energy

•
$$\left(\frac{\partial q_{0,lin}}{\partial Grad\theta}\right) = -\alpha_{\theta}\underline{I}$$

•
$$\frac{\partial Bal\ Energy}{\partial \dot{\theta}} = \rho_0 C$$

•
$$\frac{\partial Bal\ Energy}{\partial \theta} = 3\alpha_S k \cdot tr \dot{E}$$

•
$$\frac{\partial Bal\ Energy}{\partial \underline{\dot{E}}} = 3\alpha_S k\theta \underline{I}$$

Time Discretization: Newmark Beta

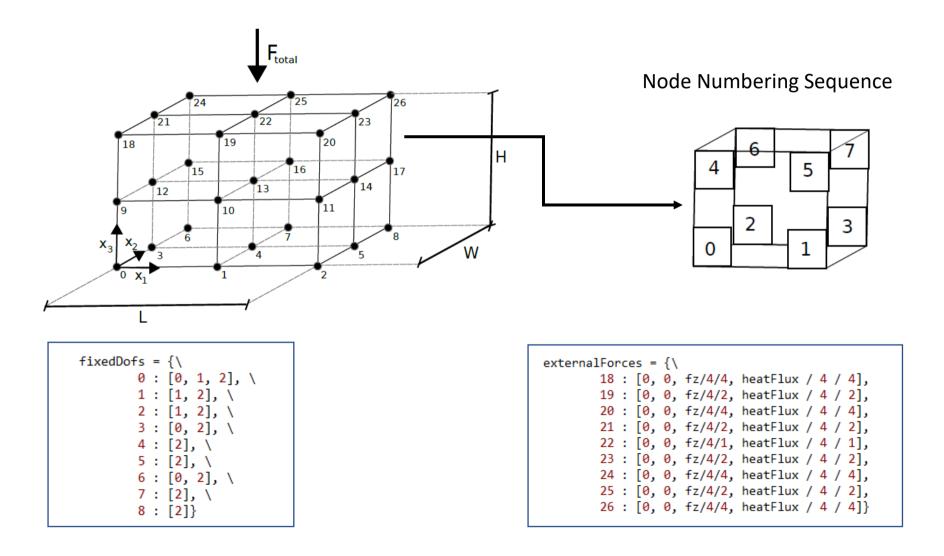
$$\underline{\dot{X}}_{t+\Delta t} = \underline{\dot{X}} + \left[(1 - \delta) \underline{\ddot{X}}_t + \delta \left(\underline{\ddot{X}}_{t+\Delta t} \right) \right] \Delta t$$

$$\Delta \underline{\dot{x}} = \frac{\delta}{\beta \Delta t} \Delta \underline{x}$$

$$\underline{X}_{t+\Delta t} = \underline{X} + \underline{\dot{X}}_t \Delta t + \left[\left(\frac{1}{2} - \beta \right) \underline{\ddot{X}}_t + \beta \left(\underline{\ddot{X}}_{t+\Delta t} \right) \right] \Delta t^2$$

$$\Delta \underline{\ddot{x}} = \frac{1}{\beta \Delta t^2} \Delta \underline{x}$$

Finite Element Mesh, Boundary Conditions and Loads



Equation System:

$$\Delta DG_{u} = \sum_{i}^{I} \sum_{\beta}^{J} \delta \underline{u}^{i} \left[\int_{B_{o}} \left(B^{T^{I}}(\underline{u}, \underline{X}) \cdot \frac{\partial \underline{S}}{\partial \underline{E}} \cdot B^{J}(\underline{u}, \underline{X}) \right) dV_{o} \Delta d\underline{u}^{J} \right] - \sum_{i}^{I} \sum_{\beta}^{J} \delta \underline{u}^{i} \left[\int_{B_{o}} \left(B^{T^{I}}(\underline{u}, \underline{X}) \cdot \frac{\partial \underline{S}}{\partial \underline{\theta}} \cdot N^{J}(\underline{X}) \right) dV_{o} \Delta d\theta^{J} \right]$$

$$\left[\delta d\underline{u} \quad \delta d\theta \right] \left\{ \left[\underbrace{\underline{K}uu}_{\underline{K}\theta u} \quad \underbrace{\underline{K}u\theta}_{K\theta \theta} \right] \left[\Delta d\underline{u}_{\underline{\theta}} \right] - \left[\underbrace{\underline{R}u}_{R\theta} \right] \right\} = \underline{\mathbf{0}}$$

$$\Delta DG_{\theta} = \sum_{i}^{I} \sum_{\beta}^{J} \delta d\theta^{I} \int_{B_{o}} \left(B^{T^{I}}(\theta) \cdot \frac{\partial q_{\{o,lin\}}}{\partial Grad\theta} \cdot B^{J}(\theta) \right) dV_{o} \Delta d\theta^{J}$$

$$- \sum_{i}^{I} \sum_{\beta}^{J} \delta d\theta^{I} \int_{B_{o}} \left(N^{I} \rho_{o} C \cdot \frac{\delta}{\beta \Delta t} \cdot N^{J} \right) dV_{o} \Delta d\theta^{J}$$

$$- \sum_{i}^{I} \sum_{\beta}^{J} \delta d\theta^{I} \int_{B_{o}} \left(N^{I} \frac{\delta}{\beta \Delta t} 3\alpha_{S} \kappa \theta \underline{I} \cdot B^{J}(\underline{u}, \underline{X}) \right) dV_{o} \Delta d\underline{u}^{J}$$

Implementation of Finite Element Code

• Shape functions: Tri-Linear Functions

Green Lagrange strain tensor

$$E_{ij} = \frac{1}{2} (u_{i,j} + u_{j,i} + u_{k,i} u_{k,j})$$

$$\boldsymbol{E}(\boldsymbol{u}) = \begin{bmatrix} E_{11} \\ E_{22} \\ E_{33} \\ 2E_{12} \\ 2E_{23} \\ 2E_{13} \end{bmatrix} = \begin{bmatrix} u_{1,1} \\ u_{2,2} \\ u_{3,3} \\ u_{1,2} + u_{2,1} \\ u_{2,3} + u_{3,2} \\ u_{1,3} + u_{3,1} \end{bmatrix} + \begin{bmatrix} 1/2 \left(u_{1,1} \ u_{1,1} + u_{2,1} \ u_{2,1} + u_{3,1} \ u_{3,1} \right) \\ 1/2 \left(u_{1,2} \ u_{1,2} + u_{2,2} \ u_{2,2} + u_{3,2} \ u_{3,2} \right) \\ 1/2 \left(u_{1,3} \ u_{1,3} + u_{2,3} \ u_{2,3} + u_{3,3} \ u_{3,3} \right) \\ u_{1,1} \ u_{1,2} + u_{2,1} \ u_{2,2} + u_{3,1} \ u_{3,2} \\ u_{1,2} \ u_{1,3} + u_{2,2} \ u_{2,3} + u_{3,2} \ u_{3,3} \\ u_{1,1} \ u_{1,3} + u_{2,1} \ u_{2,3} + u_{3,1} \ u_{3,3} \end{bmatrix}$$

• Differential operator (L) which maps the displacement vector variation to the Green Lagrange strain by addition of constant part and deformation-dependent part

$$\mathbf{D}_{\varepsilon} = \begin{bmatrix} \frac{\partial}{\partial X_{1}} & 0 & 0 \\ 0 & \frac{\partial}{\partial X_{2}} & 0 \\ 0 & 0 & \frac{\partial}{\partial X_{3}} \\ \frac{\partial}{\partial X_{2}} & \frac{\partial}{\partial X_{1}} & 0 \\ 0 & \frac{\partial}{\partial X_{3}} & \frac{\partial}{\partial X_{2}} \\ \frac{\partial}{\partial X_{3}} & 0 & \frac{\partial}{\partial X_{1}} \end{bmatrix} + \begin{bmatrix} u_{1,1} \frac{\partial}{\partial X_{1}} & u_{2,1} \frac{\partial}{\partial X_{1}} & u_{3,1} \frac{\partial}{\partial X_{1}} \\ u_{1,2} \frac{\partial}{\partial X_{2}} & u_{2,2} \frac{\partial}{\partial X_{2}} & u_{3,2} \frac{\partial}{\partial X_{3}} \\ u_{1,3} \frac{\partial}{\partial X_{3}} & u_{2,3} \frac{\partial}{\partial X_{3}} & u_{2,3} \frac{\partial}{\partial X_{1}} + u_{2,1} \frac{\partial}{\partial X_{2}} & u_{3,2} \frac{\partial}{\partial X_{1}} + u_{3,1} \frac{\partial}{\partial X_{2}} \\ u_{1,3} \frac{\partial}{\partial X_{1}} + u_{1,1} \frac{\partial}{\partial X_{2}} & u_{2,2} \frac{\partial}{\partial X_{1}} + u_{2,1} \frac{\partial}{\partial X_{2}} & u_{3,2} \frac{\partial}{\partial X_{1}} + u_{3,1} \frac{\partial}{\partial X_{2}} \\ u_{1,3} \frac{\partial}{\partial X_{1}} + u_{1,1} \frac{\partial}{\partial X_{3}} & u_{2,3} \frac{\partial}{\partial X_{2}} + u_{2,2} \frac{\partial}{\partial X_{3}} & u_{3,3} \frac{\partial}{\partial X_{2}} + u_{3,2} \frac{\partial}{\partial X_{3}} \\ u_{1,3} \frac{\partial}{\partial X_{1}} + u_{1,1} \frac{\partial}{\partial X_{3}} & u_{2,3} \frac{\partial}{\partial X_{1}} + u_{2,1} \frac{\partial}{\partial X_{3}} & u_{3,3} \frac{\partial}{\partial X_{1}} + u_{3,1} \frac{\partial}{\partial X_{3}} \end{bmatrix}$$

Implementation of Finite Element Code – Important snippets

```
# Loop for implementation of linear entries

self.epsilon[gp][0] = self.gradU[gp][0][0]

self.epsilon[gp][1] = self.gradU[gp][1][1]

self.epsilon[gp][2] = self.gradU[gp][2][2]

self.epsilon[gp][3] = self.gradU[gp][0][1] + self.gradU[gp][1][0]

self.epsilon[gp][4] = self.gradU[gp][1][2] + self.gradU[gp][2][1]

self.epsilon[gp][5] = self.gradU[gp][2][0] + self.gradU[gp][0][2]

# Loop to add all non linear entrieds into E

for j in range(0,3):

self.epsilon[gp][0] += 0.5 * self.gradU[gp][j][0] * self.gradU[gp][j][0]

self.epsilon[gp][1] += 0.5 * self.gradU[gp][j][1] * self.gradU[gp][j][1]

self.epsilon[gp][2] += 0.5 * self.gradU[gp][j][2] * self.gradU[gp][j][2]

self.epsilon[gp][3] += self.gradU[gp][j][1] * self.gradU[gp][j][2]

self.epsilon[gp][4] += self.gradU[gp][j][1] * self.gradU[gp][j][2]

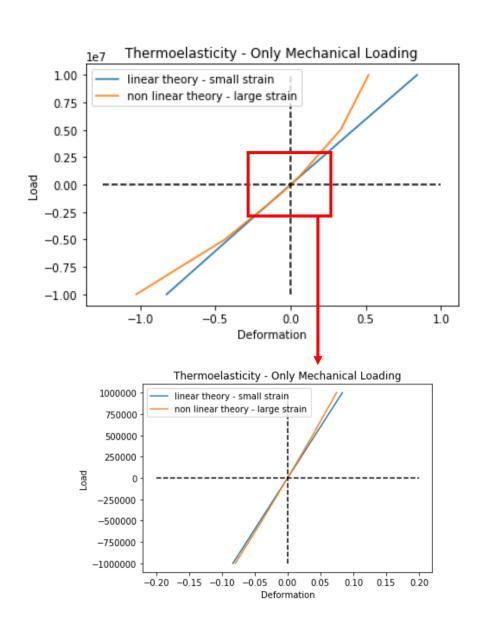
self.epsilon[gp][5] += self.gradU[gp][j][0] * self.gradU[gp][j][2]
```

```
# Linear theory - Heat flux is deformation independent
self.qlin[gp][0] = -self.alphaT * self.gradTheta[gp][0]
self.qlin[gp][1] = -self.alphaT * self.gradTheta[gp][1]
self.qlin[gp][2] = -self.alphaT * self.gradTheta[gp][2]
```

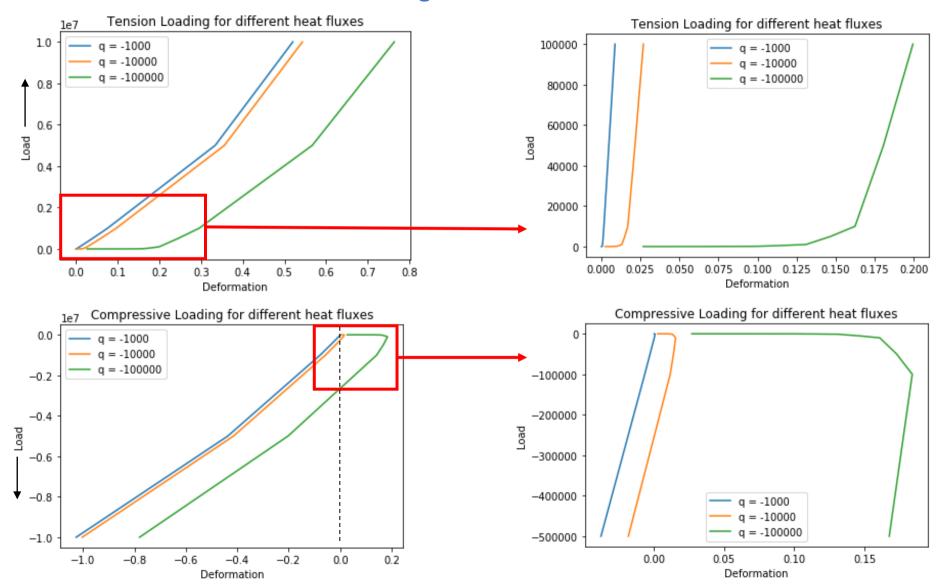
```
# mC = ( lmu - llambda * ln(J) + 3.0 * alphaS * kappa * (theta - thetaO)) * (C^-1 x C^-1)^T + llambda * (C^-1 x C^-1)
invCinvCT = zeros((3, 3, 3, 3), dtype=float64)
invCinvC = zeros((3, 3, 3, 3), dtype=float64)
mCconst1 = self.lmu - self.llambda * log(J) + 3.0 * self.alphaS * self.kappa * (self.theta[gp] - self.theta0[gp])
mCconst2 = self.11ambda
# Formulation of (C^-1 \times C^-1)^T and (C^-1 \times C^-1)
index = [[0,0], [1,1], [2,2], [0,1], [1,2], [0,2]]
for 1 in range (0.3):
    for k in range (0,3):
                                     \# k = a
        for m in range (0,3):
            for n in range (0.3): \# n = 5
                invCinvCT[1][k][m][n] = invCauchy[1][m] * invCauchy[n][k] + invCauchy[1][n] * invCauchy[m][k]
                invCinvC[1][k][m][n] = invCauchy[1][k] * invCauchy[m][n]
 # Formulation of mC which is changing at every gauss point of the element
for i in range (0.6):
    for j in range (0,6):
        self.mC[gp][i][i] = mCconst1 * invCinvCT[index[i][0]][index[i][1]][index[i][0]][index[i][1]] \
                            + mCconst2 * invCinvC[index[i][0]][index[i][1]][index[j][0]][index[j][1]]
```

Load vs. Deformation – Pure Mechanical Loading

- Small strain theory → linearised E and q
 → linear results throughout
- Large strain theory → non linear E formulation
 → additional geometric stiffness
- Tensile Load → deformation is reducing in comparison with the linear theory
- Compressive Load → non linearity of curve is not pronounced



Load vs. Deformation – Mixed Loading



Load vs. Deformation – Mixed Loading

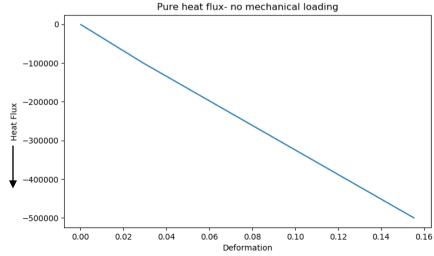
• Thermo-elasticity with small strains considers linearised E and linearised q.

| | LINEAR | | NON-LINEAR | |
|---------|---------------------------|-----------------------|---------------------------|-----------------------|
| α | Z - displacement | | | |
| | 1 st Load Step | Last Load Step | 1 st Load Step | Last Load Step |
| 0.0001 | -0.005 <u>17115</u> | 0.12 <u>4248263</u> | -0.005 <u>2514249</u> | 0.12 <u>56020314</u> |
| 0.00001 | -0.007 <u>902588</u> | 0.017 <u>9542813</u> | -0.007 <u>7980955</u> | 0.017 <u>8963899</u> |
| 0 | -0.008332539 <u>7</u> | -0.008332539 <u>0</u> | -0.008332156 <u>7</u> | -0.008332156 <u>8</u> |

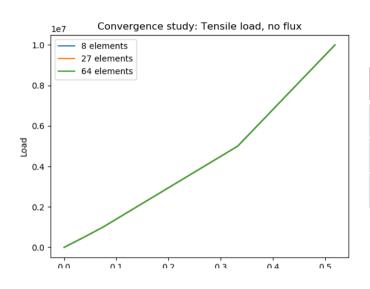
• Fz = -100000 & q = -500000

Heat Flux vs. Deformation – Pure Thermal Loading

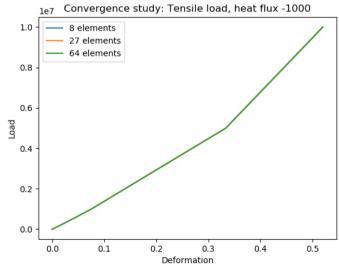
We get a linear graph → linear heat flux theory.



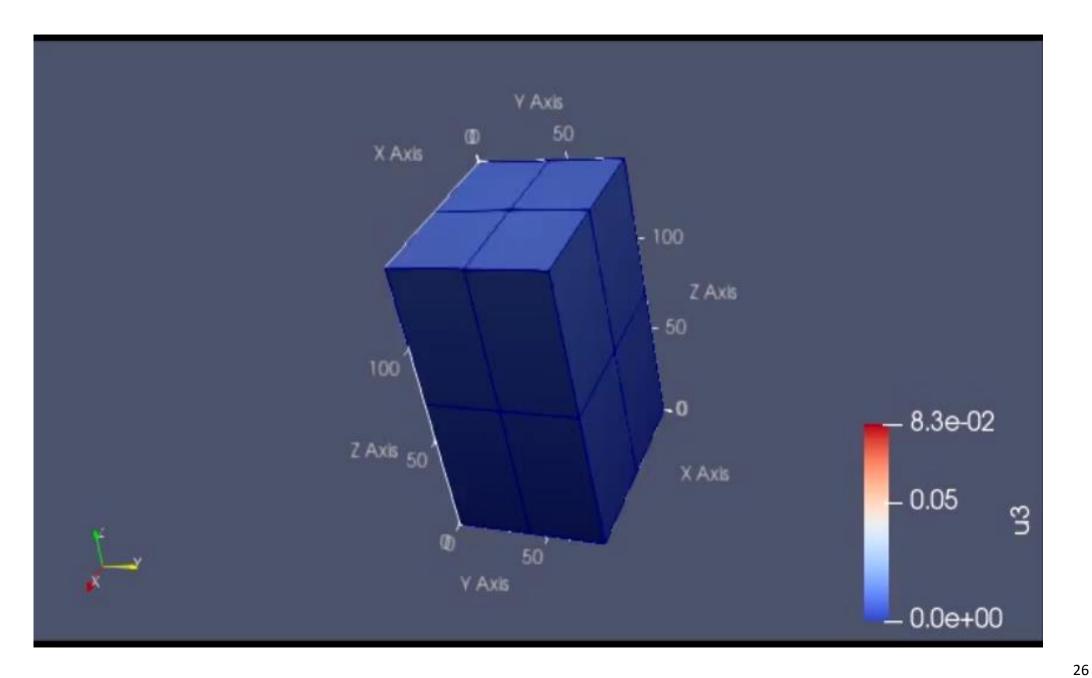
Mesh Convergence Study



| Load | 8 Elements | 27 Elements | 64 Elements |
|------|----------------------|----------------------|----------------------|
| 1e6 | 0.0744 <u>642054</u> | 0.0744 <u>597857</u> | 0.0744 <u>597857</u> |
| 5e6 | 0.332 <u>3761857</u> | 0.332 <u>364331</u> | 0.332 <u>364331</u> |
| 1e7 | 0.5184983087 | 0.5185321238 | 0.5185690033 |

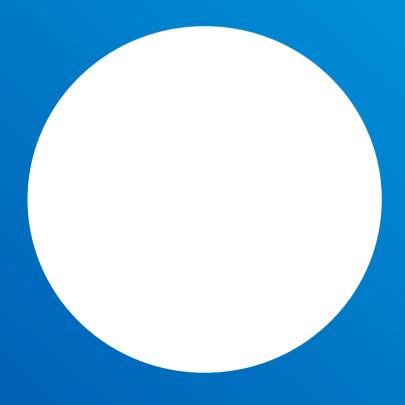


| Load | 8 Elements | 27 Elements | 64 Elements |
|------|--------------|--------------|--------------|
| 1e6 | 0.0754666549 | 0.0756117875 | 0.0760184981 |
| 5e6 | 0.3335196949 | 0.3336648635 | 0.3340613036 |
| 1e7 | 0.5197699119 | 0.5199717548 | 0.5203819076 |





Vielen Dank!



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