Series expanding the bands of La₃ Ni₂ O₇

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Loading the library

In[*]:= AppendTo[\$Path, StringDelete[NotebookDirectory[], "/La3Ni207"] <> "Package"];
<< BandUtilitiesExtras`</pre>

Loading and analysing the tight binding model

We first load the tight binding model (with symmetrisation) on to a matrix and use it to define the Hamiltonian

To compare with the unsymmetrised model, we load it as well

The effect of the symmetrisation on the weak splitting

Eigenvalues of the model at the $(\pi, 0)$

The band structures

For the sake of plotting the band structure and series expanding, we redefine the Hamiltonian

```
ln[806]:= Hamltn1New[k] := Hamltn1[k[1], k[2]]
```

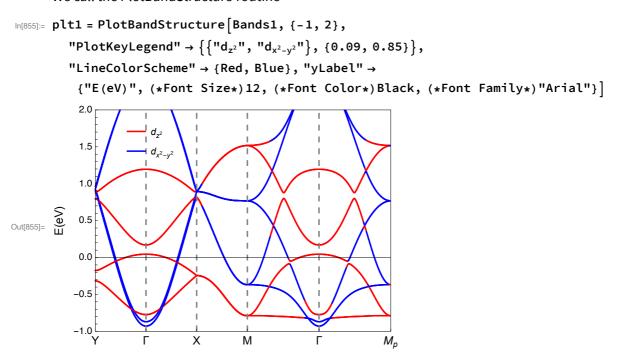
We can now plot the band structure. Firstly let us define the path

and generate the band structure data

```
In[808]= Bands1 = GenerateBandStructure[Hamltn1New, HSymmPts, "OrbitalProjection" → True,
          "OrbitalGrouping" \rightarrow \{\{1, 3, 5, 7\}, \{2, 4, 6, 8\}\}, "TotalPoints" \rightarrow 400];
```

We have grouped all the d_{z^2} orbitals together, and likewise the $d_{x^2-\nu^2}$ orbitals as well, for the sake of orbital projections.

We call the PlotBandStructure routine



Locating the (possible) critical point located between the M and Γ points

For the LocateCriticalPoint routine, we first identify an approximate range for "critical" points in bands 3 and 4.

```
ln[813]:= kp1 = HSymmPts[4, 1] + 0.5 (HSymmPts[5, 1] - HSymmPts[4, 1]);
      kp2 = HSymmPts[4, 1] + 0.75 (HSymmPts[5, 1] - HSymmPts[4, 1]);
In[816]= kcrit3 = LocateCriticalPoint[Hamltn1New, kp1, kp2, 3, "MaxIterations" → 20] [[1]]
      The value of MaxIterations has been reset to 20.
      Point found after 14 iteration(s).
Out[816]= \{1.31036, 1.31036\}
In[817]:= kcrit4 = LocateCriticalPoint[Hamltn1New,
          kp1, kp2, 4, "MaxIterations" \rightarrow 30, "Resolution" \rightarrow 10<sup>-6</sup> [[1]
      The value of MaxIterations has been reset to 30.
      Point found after 20 iteration(s).
Out[817]= \{1.24668, 1.24668\}
```

Below we will see that these are not genuine critical points. The directional derivative along the MF

direction vanishes but not the gradient itself.

Series expanding the band structure

We will attempt to series expand the band structure at some points of interest near the Fermi level.

The Y point

We notice a VHS at the Y point on band 2. We confirm this first with the symmetrised model and then the unsymmetrised model

```
In [818]:= Taylor Expand Band [Hamltn1New, \{0, \pi\} // N, 2, 4(*, "Resolution" \rightarrow 10<sup>-8</sup>,
          "Dimension"\rightarrow 2*), "Messages" \rightarrow False /. \{p1 \rightarrow k_x, p2 \rightarrow k_y\}
-0.00183652\ k_{x}^{4}-0.00180483\ k_{x}^{3}\ k_{y}+0.172471\ k_{x}^{2}\ k_{y}^{2}-0.0130803\ k_{x}\ k_{y}^{3}-0.388893\ k_{y}^{4}\big\}\big\}
```

So this is an ordinary saddle with log diverging DOS

The Γ point

The Γ point seems to host a maximum on band 4, near the Fermi level

```
In[819]= TaylorExpandBand [Hamltn1New, \{0, 0\} // N, 4, 4(*, "Resolution" \rightarrow 10<sup>-8</sup>,
                                                                                "Dimension"\rightarrow 2*), "Messages" \rightarrow False] /. \{p1 \rightarrow k_x, p2 \rightarrow k_y\}
\text{Out[819]= } \left\{ \left\{ \text{0.0456275, 0, -0.0291695} \; k_x^2 + \text{0.00225983} \; k_x \; k_y - \text{0.0287074} \; k_y^2, \; \text{0, -0.000610605} \; k_x^4 - \text{0.00225983} \; k_x \; k_y + \text{0.00287074} \; k_y^2, \; \text{0, -0.000610605} \; k_x^4 + \text{0.00225983} \; k_x \; k_y + \text{0.00287074} \; k_y^2, \; \text{0, -0.000610605} \; k_x^4 + \text{0.00225983} \; k_x \; k_y + \text{0.00287074} \; k_y^2, \; \text{0, -0.000610605} \; k_x^4 + \text{0.00225983} \; k_x \; k_y + \text{0.00287074} \; k_y^2, \; \text{0, -0.000610605} \; k_x^4 + \text{0.00225983} \; k_x \; k_y + \text{0.00287074} \; k_y^2, \; \text{0, -0.000610605} \; k_x^4 + \text{0.00225983} \; k_x \; k_y + \text{0.00287074} \; k_y^2, \; \text{0, -0.000610605} \; k_x^4 + \text{0.00225983} \; k_x \; k_y + \text{0.00287074} \; k_y^2, \; \text{0, -0.0006106005} \; k_x^4 + \text{0.00225983} \; k_x \; k_y + \text{0.00287074} \; k_y^2, \; \text{0, -0.0006106005} \; k_x^4 + \text{0.00225983} \; k_x \; k_y + \text{0.00287074} \; k_y^2, \; \text{0, -0.0006106005} \; k_x^4 + \text{0.00225983} \; k_x \; k_y + \text{0.00287074} \; k_y^2, \; \text{0, -0.0006106005} \; k_x^4 + \text{0.00225983} \; k_x \; k_y + \text{0.00287074} \; k_y^2, \; \text{0, -0.0006106005} \; k_x^4 + \text{0.00225983} \; k_x \; k_y + \text{0.00225983} \; k_x + \text{0.00225983} \; k_x + \text{0.00225983} \; k_y + \text{0.00225983} \; 
                                                                                          0.00464882 \, k_x^3 \, k_y - 0.00605766 \, k_x^2 \, k_y^2 - 0.00310384 \, k_x \, k_y^3 + 0.000612562 \, k_y^4 \} \,
```

Indeed, this is a maximum that once again has log diverging DOS below the maximum energy and zero DOS above.

The X point

The X point is the one of the most interesting of the lot since it appears to have a pair of Van Hove singularities in both the models. Let us confirm that explicitly

```
In [822]:= Taylor ExpandBand [Hamltn1New, \{\pi, 0\} // N, 4, 2, "Resolution" \rightarrow 10^{-3}
            (*,"Dimension"\rightarrow 2*), "Messages" \rightarrow False /. \{p1 \rightarrow k_x, p2 \rightarrow k_y\}
Out[822]= \{\{0.811291, -0.217837 \, k_x, -2.78856 \, k_x^2 + 0.00737567 \, k_x \, k_y - 2.13615 \, k_y^2\},
          \{0.811702, 0.217837 \, k_x, -2.80123 \, k_x^2 + 0.00653128 \, k_x \, k_y - 2.14467 \, k_y^2 \} \}
```

The series expansion indicates an interesting anomaly: neither of the bands have a VHS or for that matter a critical point thanks to the linear in k_x term. This might seem to contradict the π rotation symmetry at the X point that would normally force an isolated band to have critical point. But this is not necessary when we have multiple bands intersecting. In fact we see that the pair of bands together respect the π rotation symmetry at the X point!

Furthermore, from the band structure the starkly linear behaviour of the dispersion at the X point in

the X-Γ direction is quite apparent (see below)

The M point

The bands of interest are 3 and 4 which are located below the Fermi level

```
In [857]:= Taylor ExpandBand [Hamltn1New, \{\pi, \pi\} // N, 3, 2, "Resolution" \rightarrow 10<sup>-3</sup>
              (*,"Dimension"\rightarrow 2*), "Messages" \rightarrow False /. \{p1 \rightarrow k_x, p2 \rightarrow k_y\}
Out[857]= \left\{ \left\{ -0.367799, -0.0124961 \, k_x, \, 0.201839 \, k_x^2 + 0.191851 \, k_y^2 \right\} \right\}
           \left\{-0.367794, 0.0124961 \, k_x, 0.201839 \, k_x^2 + 0.191851 \, k_y^2\right\}\right\}
```

Similar to the case of the X point, we once again have a small linear dispersion for each band which nevertheless satisfy the π rotation symmetry jointly.

The "critical" point located on the MΓ line

Having previously found the location of these points for bands 3 and 4, we compute the

```
ln[825]:= Tayl1 = TaylorExpandBand[Hamltn1New, kcrit3, 3, 2, "Resolution" \rightarrow 10<sup>-4</sup>
                (*, "Dimension"→2*), "Messages" → False /. \{p1 \rightarrow k_x, p2 \rightarrow k_y\}
Out[825]= \left\{ \left\{ -0.0838476, 0.00609075 \, k_x - 0.00601557 \, k_y, \right\} \right\}
              -0.820285 k_x^2 - 1.66029 k_x k_y - 0.797172 k_y^2 \}
In[826]:= Tayl2 = TaylorExpandBand [Hamltn1New, kcrit4, 4, 2, "Resolution" → 10<sup>-4</sup>
                (*,"Dimension"\rightarrow 2*), "Messages" \rightarrow False /. \{p1 \rightarrow k_x, p2 \rightarrow k_y\}
\text{Out} \texttt{[826]=} \quad \left\{ \left. \left\{ -0.0493854 \text{,} \\ -0.0140342 \text{ k}_{x} + 0.0140353 \text{ k}_{y} \text{,} \\ 0.592806 \text{ k}_{x}^{2} + 1.36729 \text{ k}_{x} \text{ k}_{y} + 0.788481 \text{ k}_{y}^{2} \right\} \right\}
```

We notice once again that the bands in question do not have a critical point owing to the presence of a non-zero linear part for the dispersion. It is not surprising however that the band structure seemed to indicate a pair of critical points since the directional derivatives along the MF direction (i.e the (1,1) direction) vanish as can be seen explicitly:

```
ln[829]:= Chop[Tayl1[1, 2]] /. \{k_x \to 1, k_y \to 1\}, 10^{-4}]
Out[829]= 0
In[830]:= Chop[Tayl2[1, 2]] /. \{k_x \to 1, k_y \to 1\}, 10^{-4}]
Out[830]= 0
```

DOS behaviour near the high symmetry points, close to the Fermi level

We can now mark the DOS behaviour based on the inference drawn from the series expansions derived above. We first note the locations of the critical points on the band structure plot we previously generated:

 M_p

```
In[833]:= Bands1[[2]]
Out[833]= \left\{ \{0, Y\}, \{\pi, \Gamma\}, \{2\pi, X\}, \{3\pi, M\}, \left\{3\pi + \sqrt{2}\pi, \Gamma\right\}, \left\{3\pi + 2\sqrt{2}\pi, M_p\right\} \right\}
In[856]:= Show[{plt1, Graphics[
              Text[Style["log\nDOS", FontFamily \rightarrow "Times", FontSize \rightarrow 14], \{\pi, -0.1\}]],
            Graphics[Text[Style["log\nDOS", FontFamily → "Times", FontSize → 14],
                \{\pi / 6.5, 0.02\}]], Graphics[Text[
                Style["finite\nDOS", FontFamily \rightarrow "Times", FontSize \rightarrow 14], {2.175 \pi, -0.13}]],
            Graphics \lceil \text{Text} \lceil \text{Style} \rceil "finite\nDOS", FontFamily \rightarrow "Times", FontSize \rightarrow 14],
                \left\{\left(3+\sqrt{2} \middle/ 2.3\right)\pi, -0.04\right\}\right]\right\}, ImageSize \rightarrow Full
               1.5
               1.0
Out[856]= (Se)
               0.5
                   log
              0.0 DOS
                                                                                       finite
                                    log
                                                          finite
                                                                                       DOS
                                   DOS
                                                          DOS
             -0.5
```

-1.0 L Y