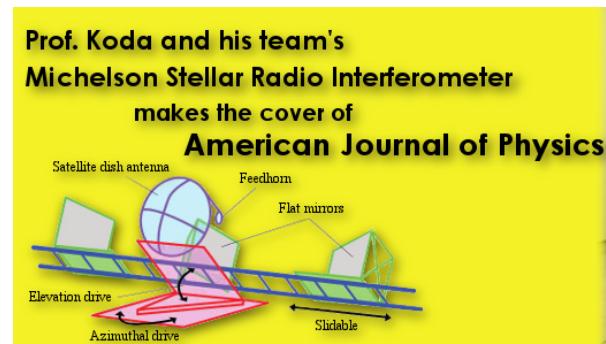


Radio Astronomy

Jin Koda

Contents

- Radio signal from the sky
- Antenna
 - Antenna response pattern
- Receiver
 - Bolometer, HEMT, SIS
 - Heterodyne receiver
- Interferometer



Radio Signal from the Sky

How much energy would we collect?

Suppose we observe a 10 Jy calibrator with the CARMA array for 1 year, 24 hours/day.

$$E = \frac{1}{2} S \eta A \Delta v t$$

- The factor of $\frac{1}{2}$ arises because we are sensitive to 1 polarization
- S = source flux density = 10 Jy = 10×10^{-26} watts $m^{-2} Hz^{-1}$
- η = aperture efficiency ~ 0.60
- A = geometrical collecting area = $6 \times 85 m^2 + 9 \times 29 m^2 = 771 m^2$
- Δv = instantaneous bandwidth = 3GHz = 3×10^9 Hz
- t = 1 year = 3×10^7 sec

Result: $E = 2 \times 10^{-6}$ joules

1 calorie = 4.2 joules heats 1 cm³ (20 drops?) of water by 1 C.

→ Must observe for 100,000 years to heat 1 drop of water by 1 C.

Antenna

Radio telescopes

IRAM 30m telescope (Spain)



Nobeyama 45m telescope (Japan)



Telescope/Antenna

Large Millimeter Telescope 50m (Mexico)



Very Large Array (New Mexico)



~1-50 GHz

Centimeter wavelength

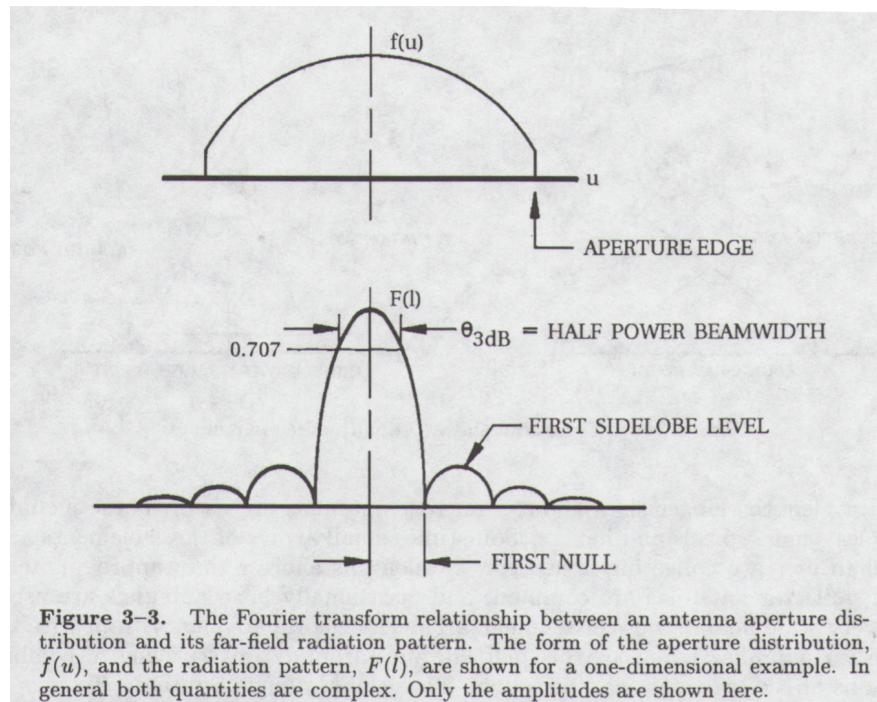
ALMA array (Chile)



~100-850 GHz

Millimeter to submillimeter wavelength

Antenna response function



- Resolution $\sim \frac{\lambda}{D}$
- Main beam & sidelobes
 - depends on surface accuracy
 - sidelobes could be very high

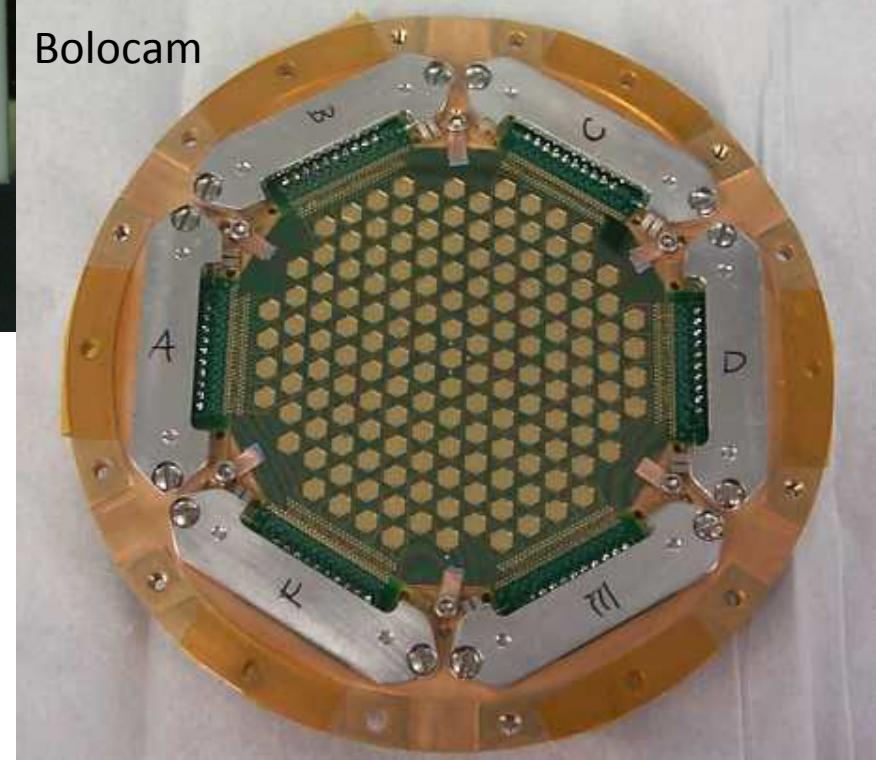
Receiver

detectors for radio astronomy

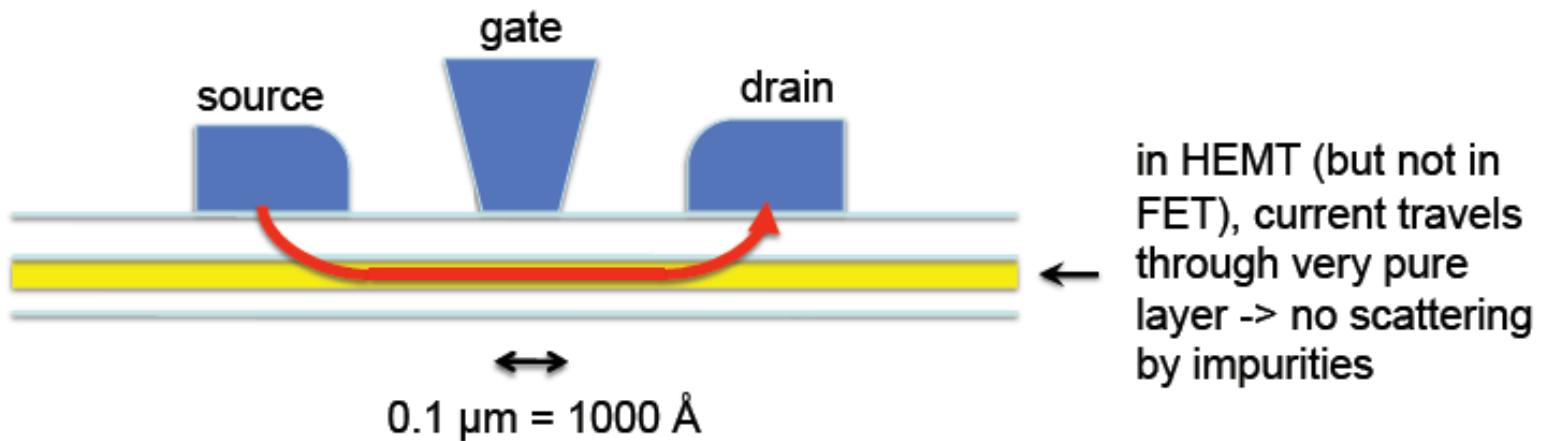
1. bolometers
 - absorbed photon increases temperature, changes resistance
 - phase of incoming signal is lost – unsuitable for aperture synthesis
 - operate at ~0.3 K
2. HEMT (High Electron Mobility Transistor) amplifiers
 - preferred below 50 GHz, good up to 115 GHz
 - operate at ~20 K
3. SIS mixers
 - mixes incoming signal with local oscillator to convert it to a lower frequency where it is amplified (by HEMT)
 - preferred for 100+ GHz
 - operate at ~4 K

Detected signal as
electromagnetic wave,
not as photon

Bolometer



High Electron Mobility Transistor (HEMT) amplifier

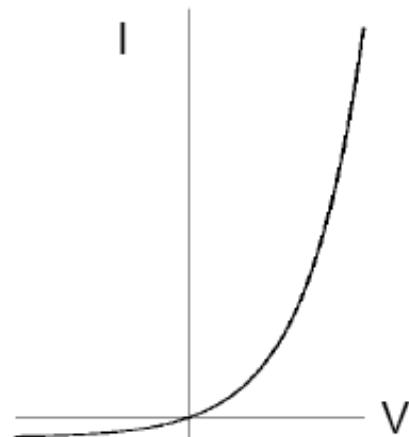


- gate voltage controls width of channel, modulates current from source to drain
- to operate at 100 GHz, charge carriers must transit under the gate in $\sim 1/10 \times 1/100 \text{ GHz} \sim 10^{-12} \text{ sec}$
- must travel $0.1 \mu\text{m}$ in $10^{-12} \text{ sec} \sim 100 \text{ km s}^{-1}$

heterodyne receiver

- converts incoming signal to a lower frequency where it can be amplified
- how? ‘mix’ the incoming signal with a strong ‘local oscillator’ in a nonlinear device to generate an ‘intermediate frequency (IF)’
- essentially the local oscillator is a clock that samples the incoming signal periodically
- example of nonlinear device: a diode

current/voltage relationship
for a diode



mixer – a nonlinear device

- linear device (superposition principle):

$$\omega_1, \omega_2 \xrightarrow{\text{linear device}} \omega_1, \omega_2$$

- nonlinear device:

$$\omega_1, \omega_2 \xrightarrow{\text{nonlinear device}} \omega_1, \omega_2, \omega_1 + \omega_2, \omega_1 - \omega_2, 2\omega_1 + \omega_2, \dots$$

- diode is an example of a nonlinear device:

$$I = I_0(e^{\alpha V} - 1) \sim I_0(\alpha V + \frac{1}{2} \alpha^2 V^2 + \dots)$$

$$V = A \cos \omega_1 t + B \cos \omega_2 t$$

$$\begin{aligned} V^2 &= A^2 \cos^2 \omega_1 t + B^2 \cos^2 \omega_2 t + 2AB \cos \omega_1 t \cos \omega_2 t + \dots \\ &= \dots + AB \cos(\omega_1 + \omega_2)t + AB \cos(\omega_1 - \omega_2)t + \dots \end{aligned}$$

- note: *amplitude* at frequency $\omega_1 - \omega_2$ is *linearly* related to amplitudes A and B

waveforms in a heterodyne receiver



local oscillator (LO)



'signal' – just
random noise for
radio astronomy



LO + signal
(voltage in diode)



current through
diode



IF after low pass
filtering

Why Large Telescope?

Why Large Telescope?

- Large collecting area
 - More photons
 - Fainter targets
- Higher resolution (at cm/mm-wavelengths)
 - Angular resolution: λ/D

Angular Resolution & Antenna response function

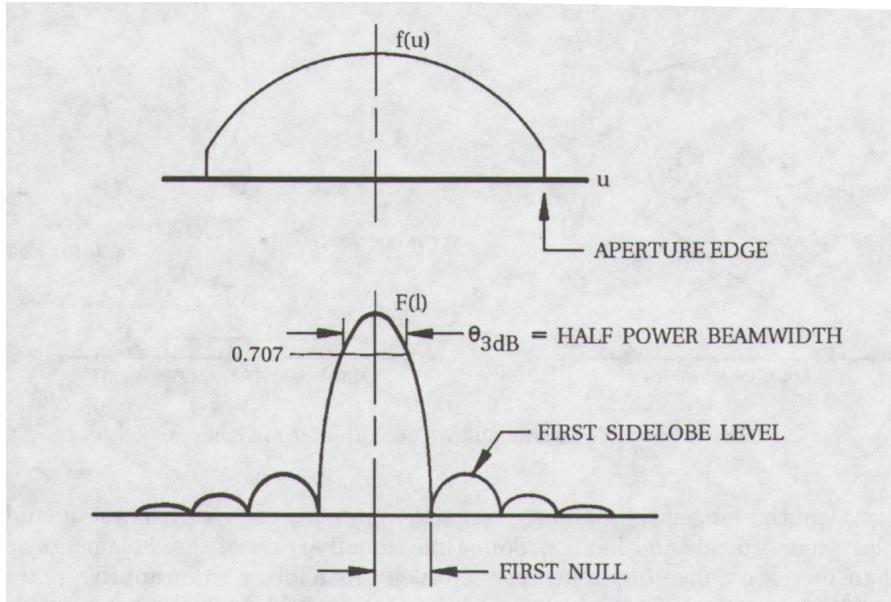
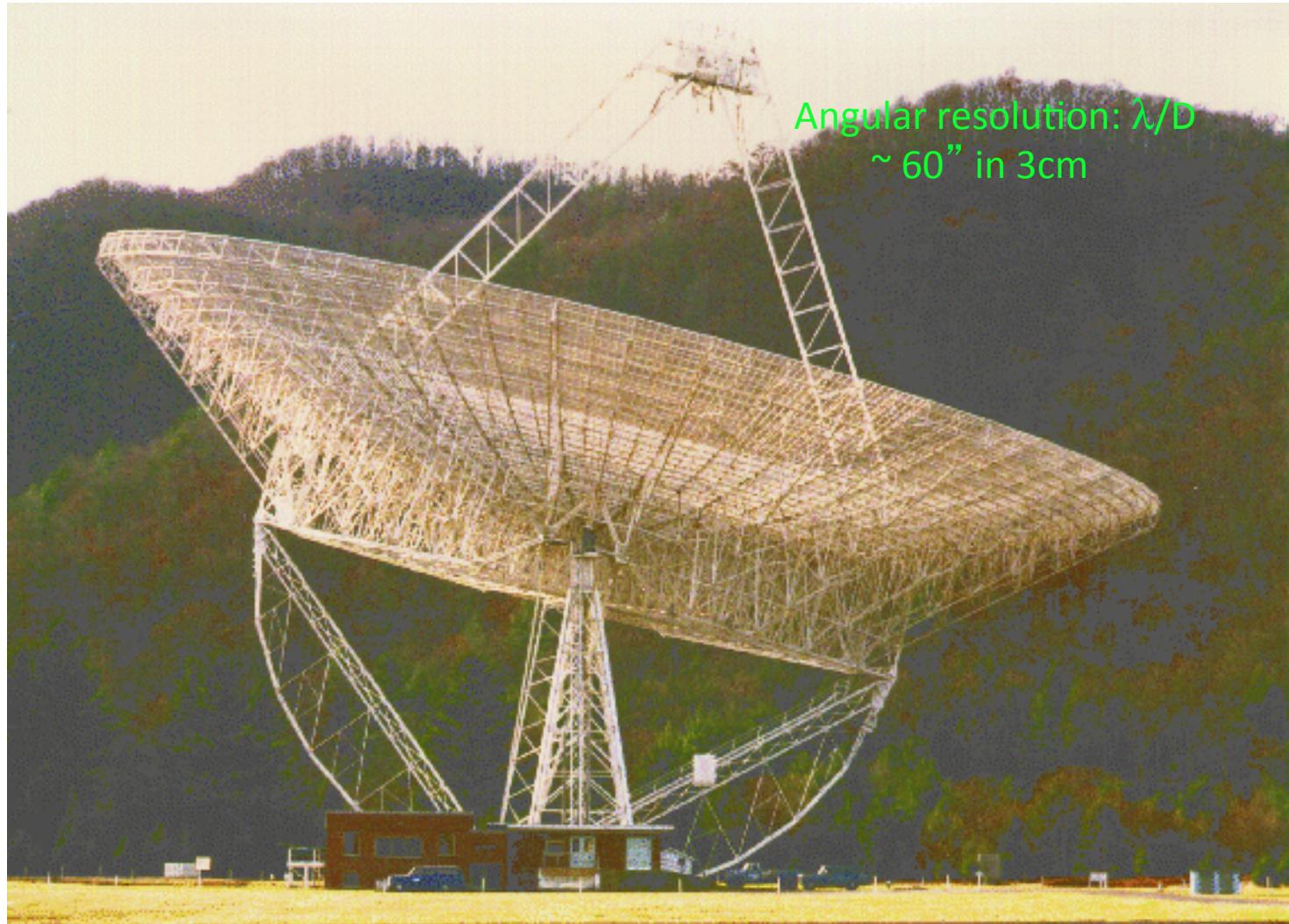


Figure 3-3. The Fourier transform relationship between an antenna aperture distribution, $f(u)$, and its far-field radiation pattern, $F(l)$, are shown for a one-dimensional example. In general both quantities are complex. Only the amplitudes are shown here.



- Resolution \sim wavelength/antenna diameter
- Main beam & sidelobes

[Old] Green Bank Telescope (cm-wave)

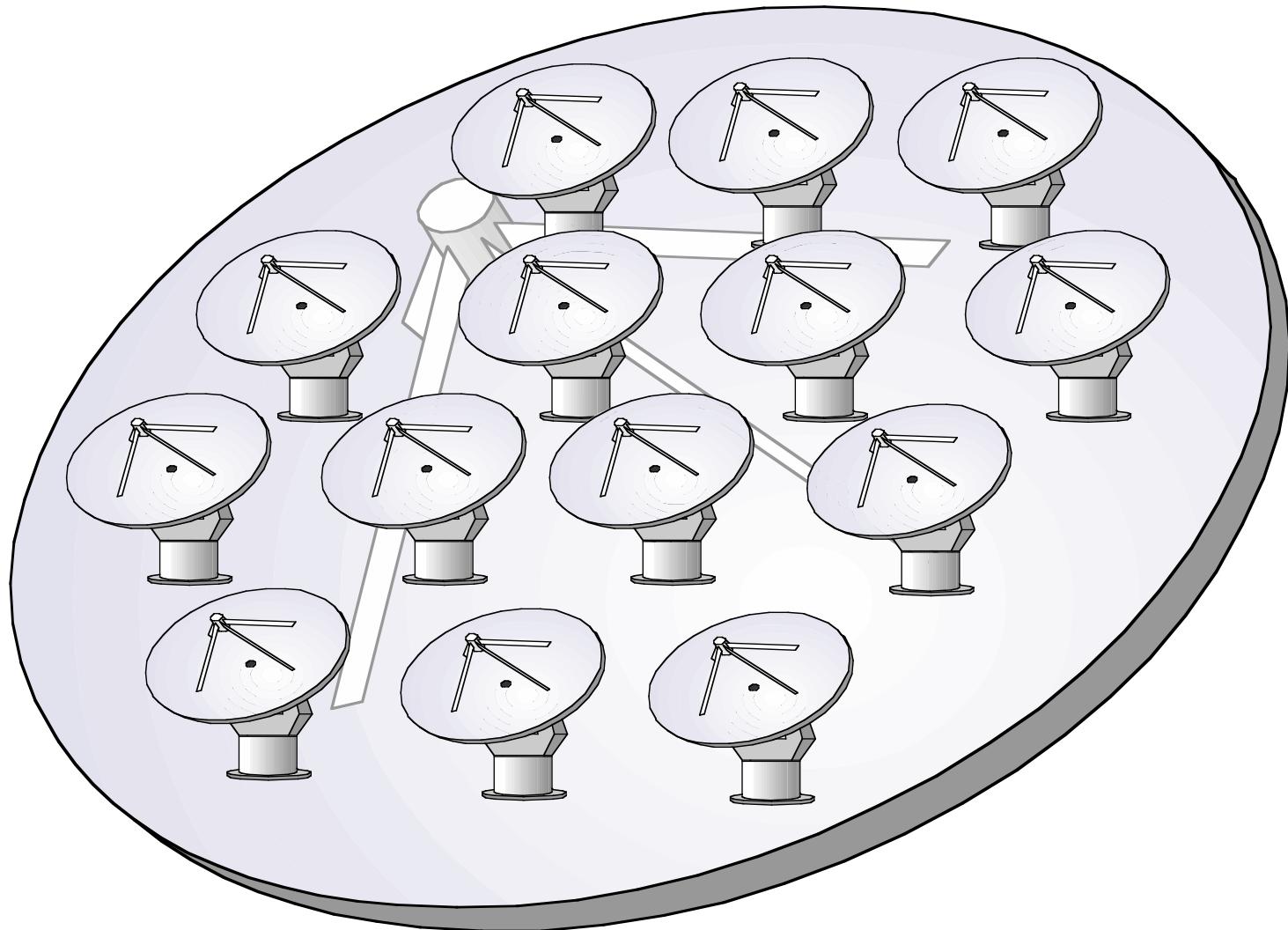


Collapsed



Interferometers

Interferometer (Many Elements)



Radio Interferometers

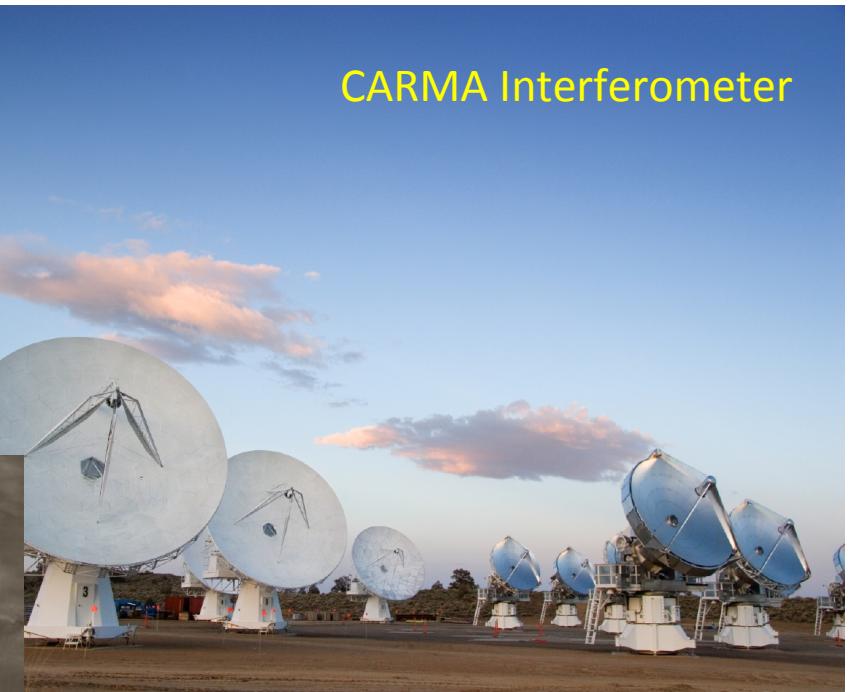
Very Large Array (VLA) at New Mexico



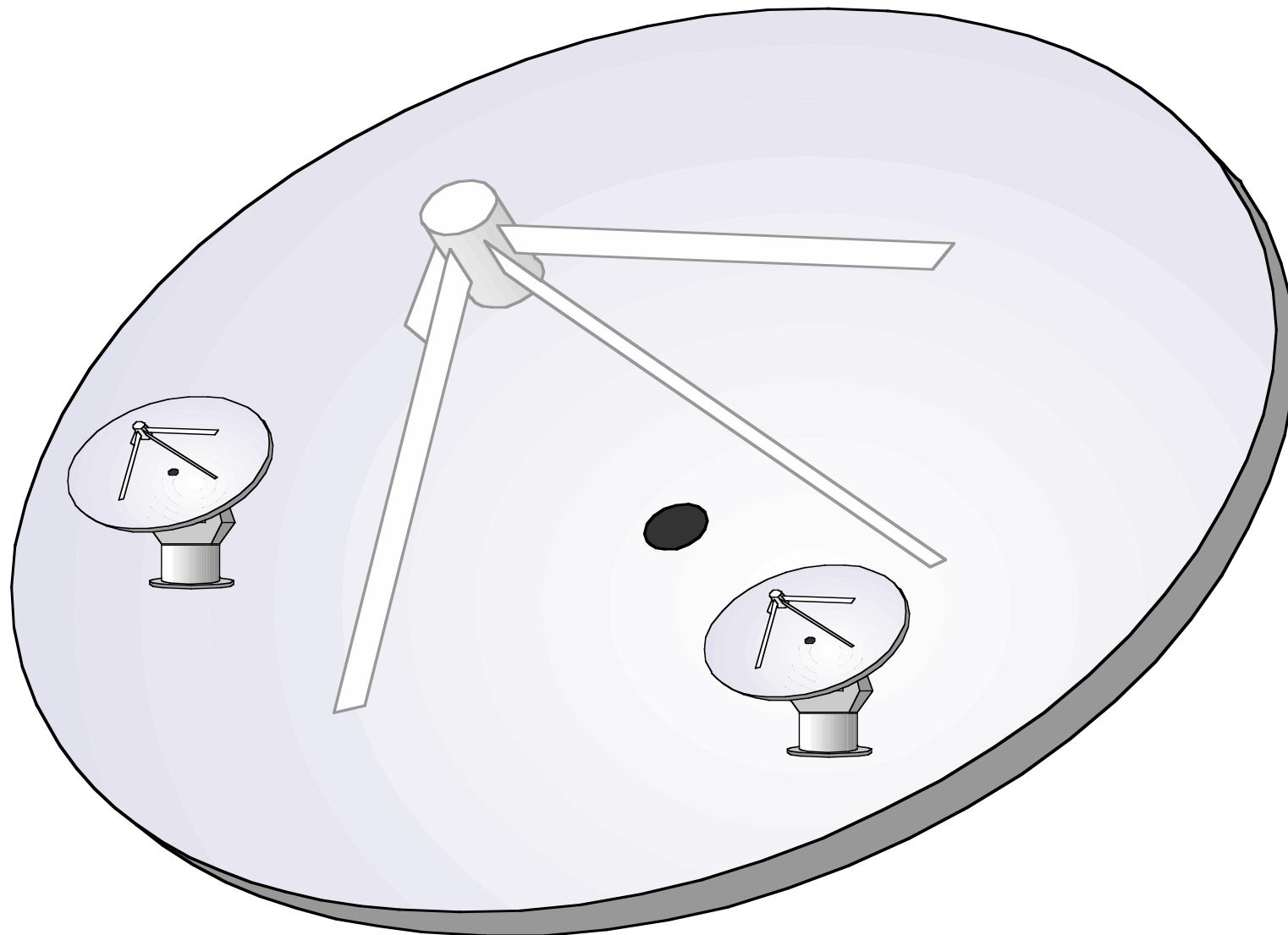
ALMA (under construction)



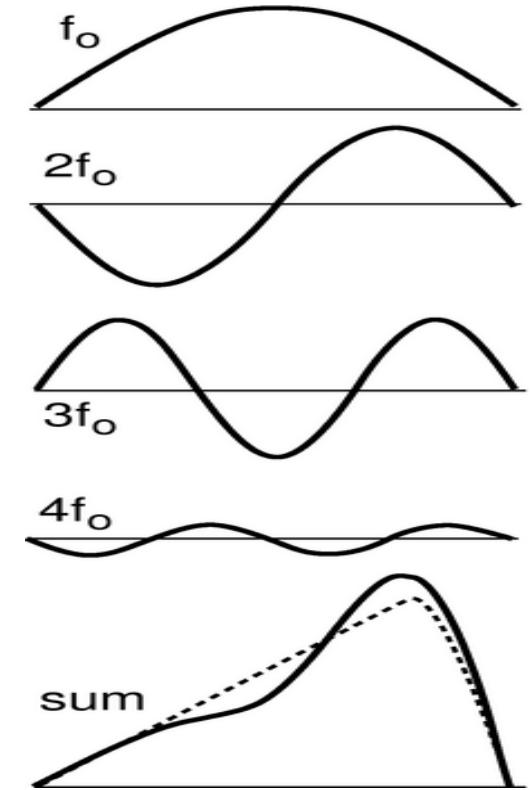
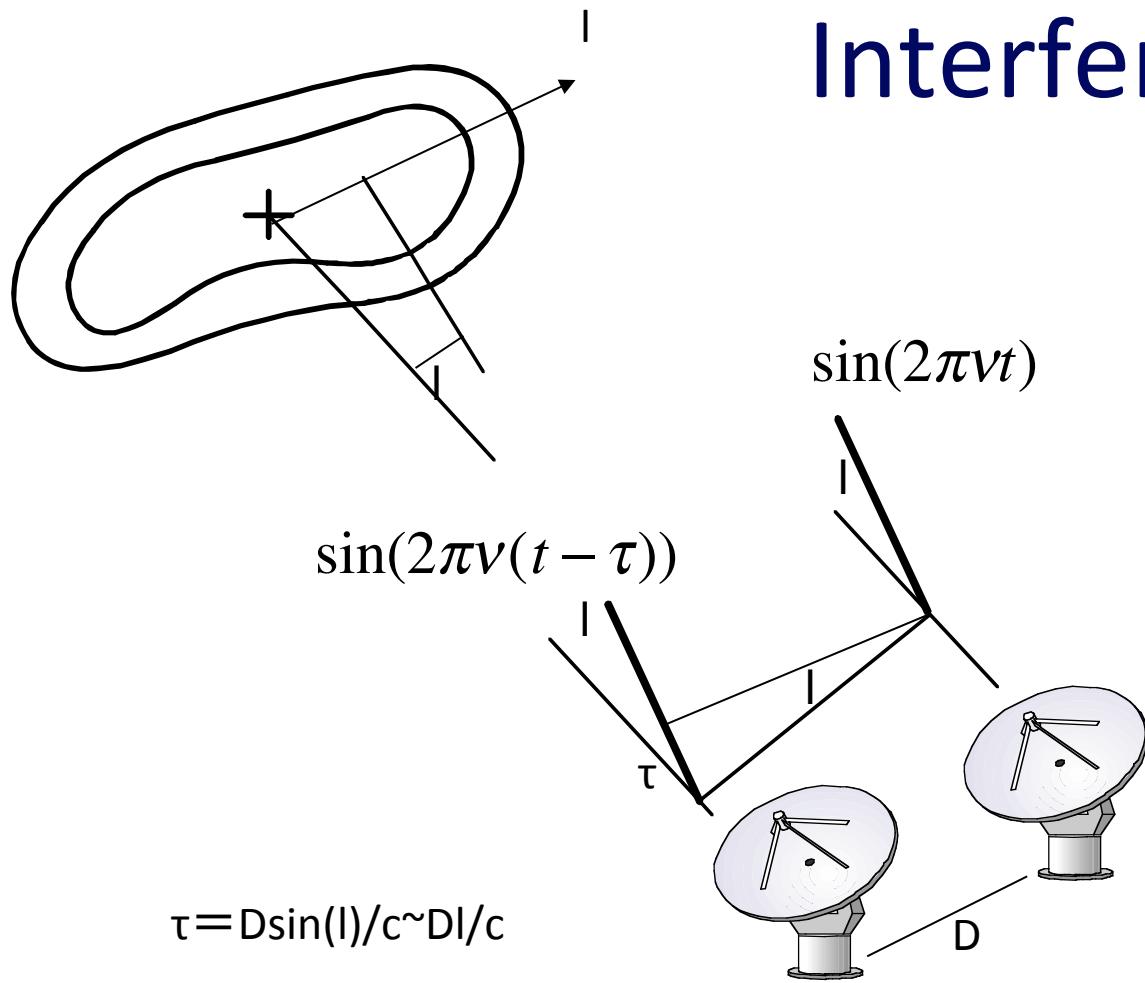
CARMA Interferometer



Interferometer (2 Elements)



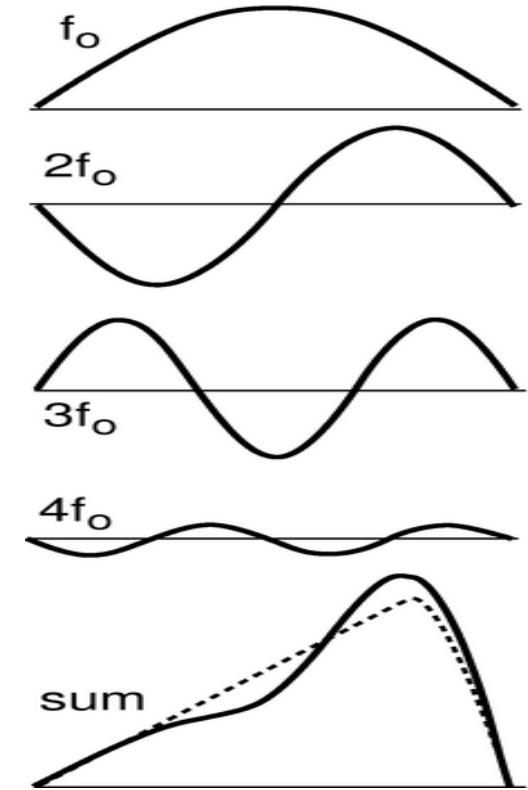
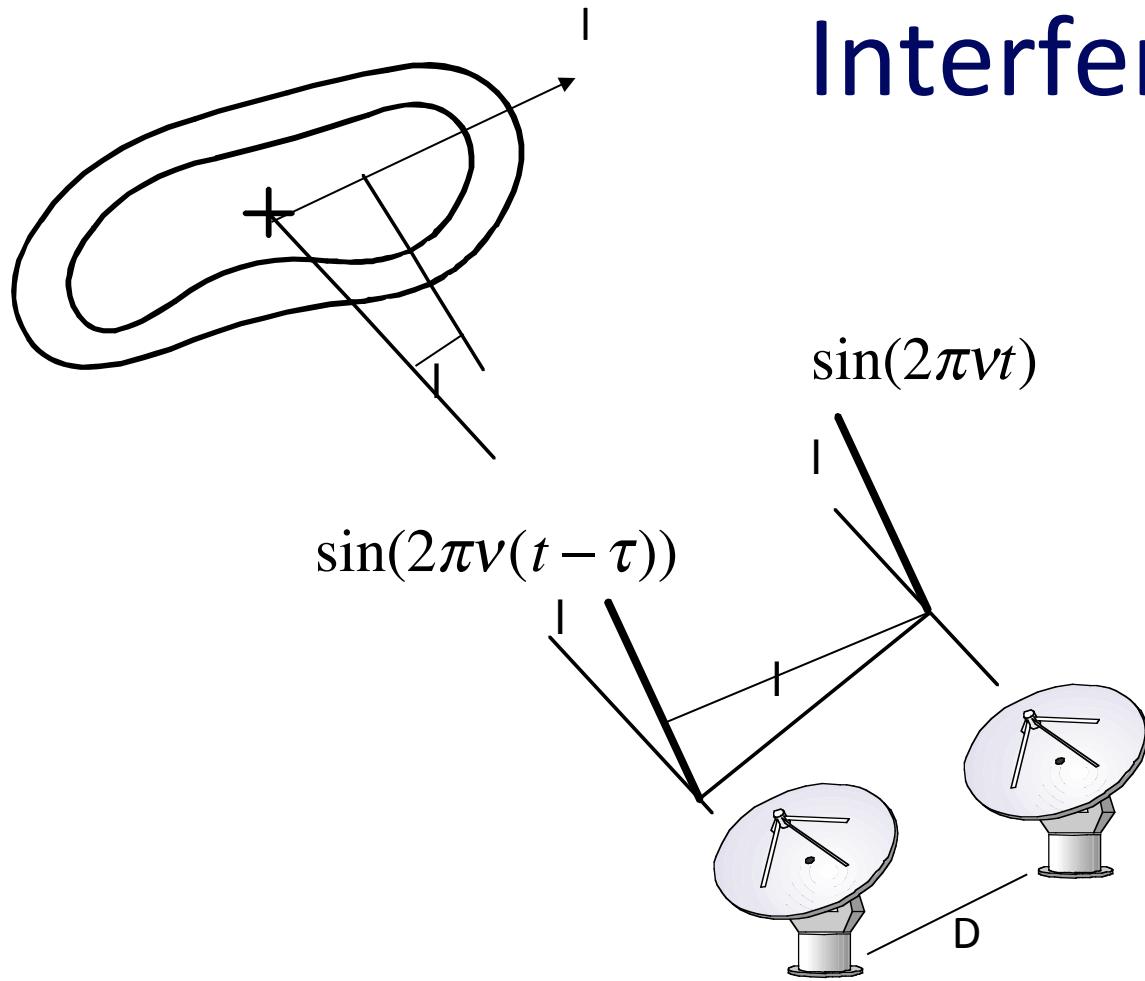
Interferometer



$$F = \sin(2\pi vt) \sin(2\pi v(t - \tau)) \\ \approx \cos(2\pi v\tau) = \cos\left(\frac{2\pi Dl}{\lambda}\right)$$

A Fourier component

Interferometer

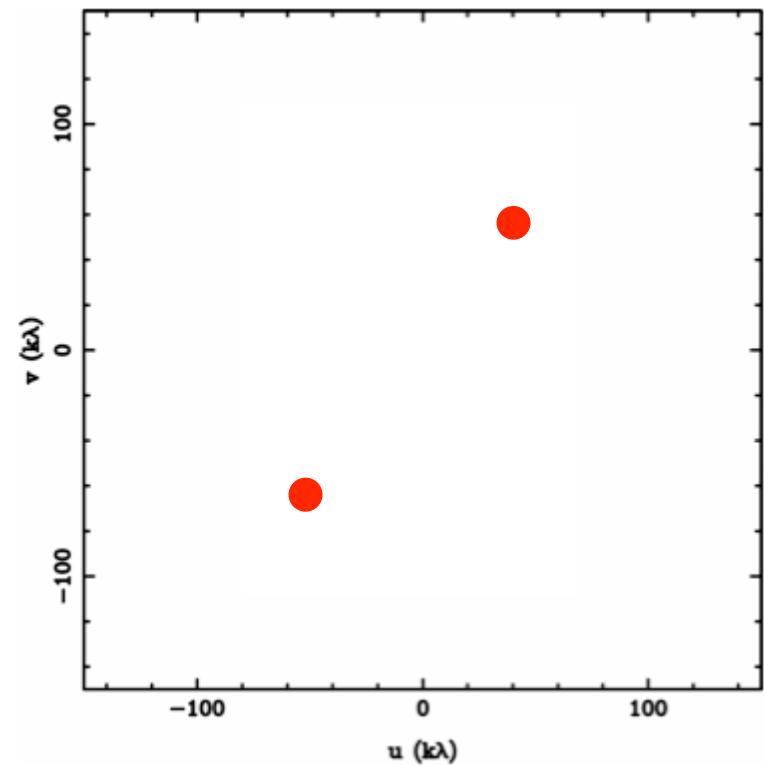
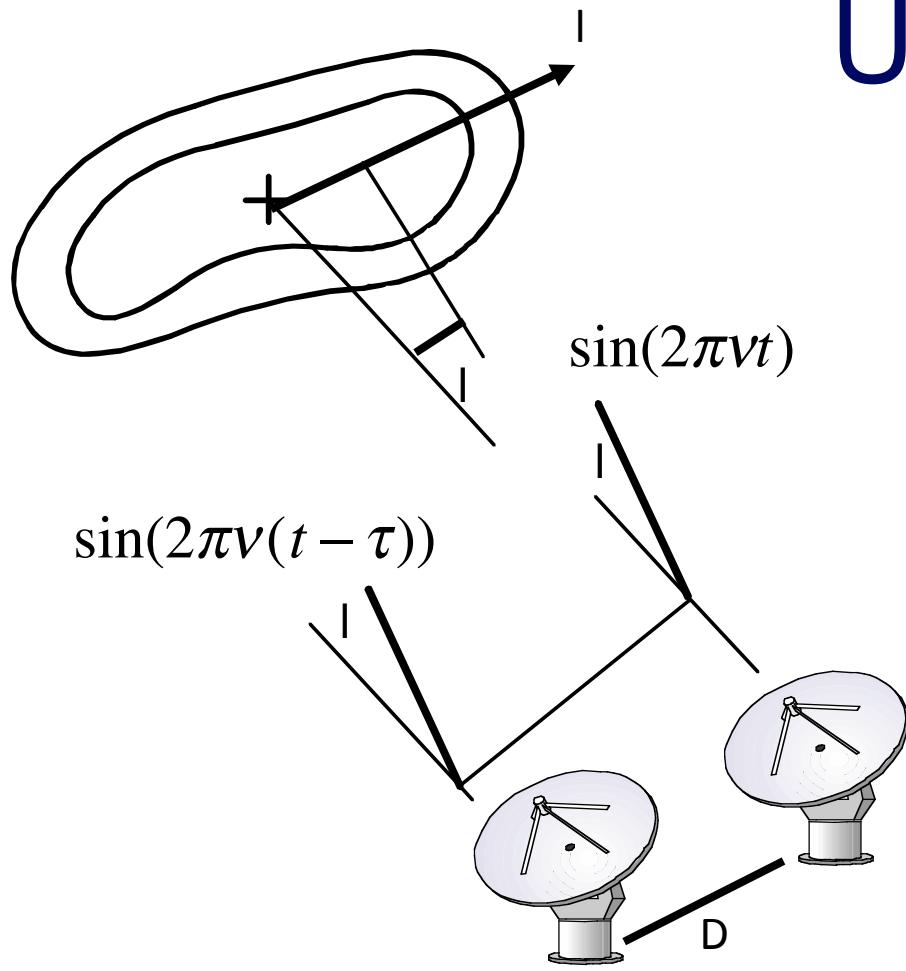


$$F = \sin(2\pi vt) \sin(2\pi v(t - \tau)) \\ \approx \cos(2\pi v\tau) = \cos\left(\frac{2\pi Dl}{\lambda}\right)$$



$$\text{Image} = \sum_D A_D \cos\left(\frac{2\pi Dl}{\lambda}\right)$$

UV-coverage

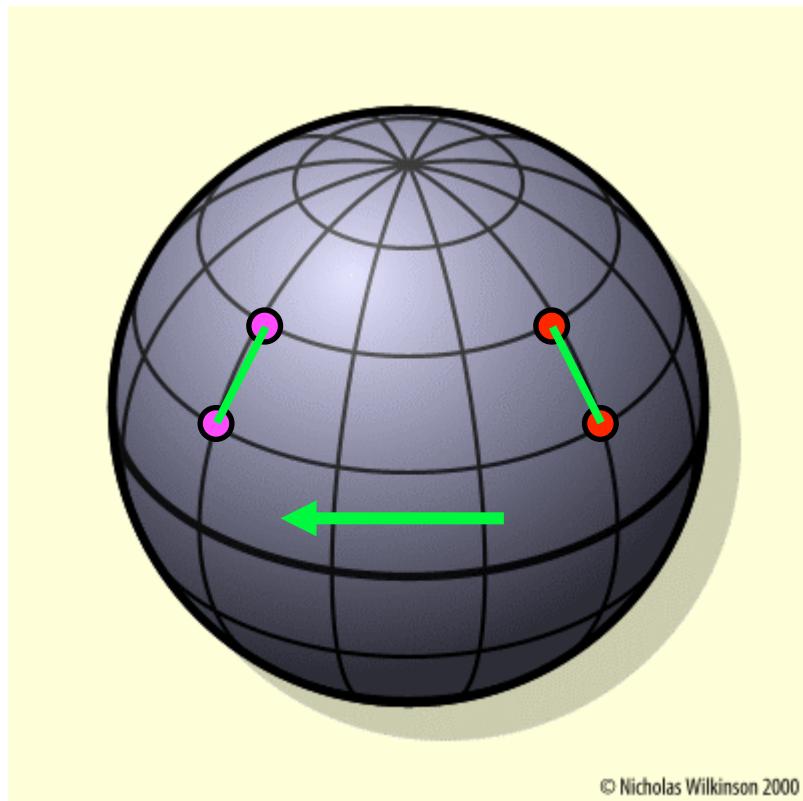


$$F = \sin(2\pi vt) \sin(2\pi v(t - \tau)) \\ \approx \cos(2\pi v\tau) = \cos\left(\frac{2\pi Dl}{\lambda}\right)$$

Image $= \sum_D A_D \cos\left(\frac{2\pi Dl}{\lambda}\right)$

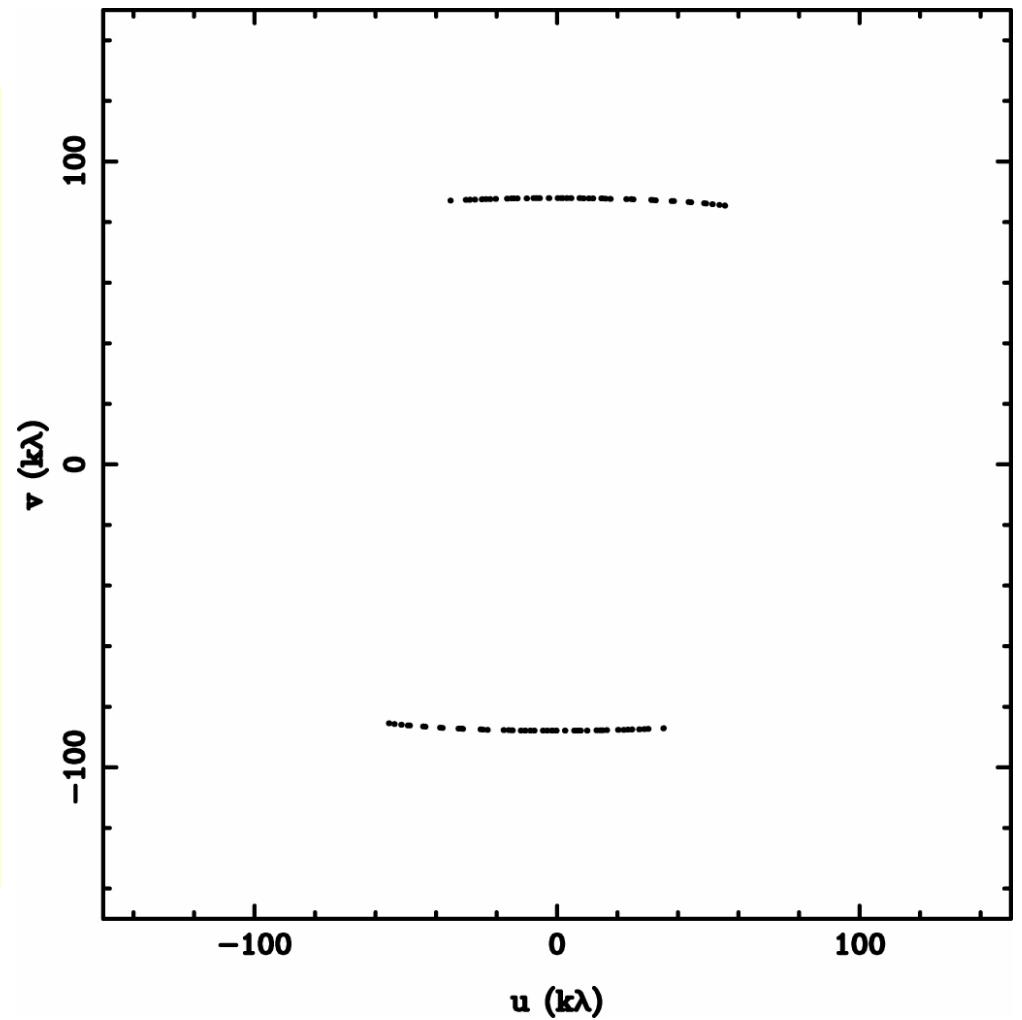
Synthesis imaging (2 antennas)

Earth (rotating)



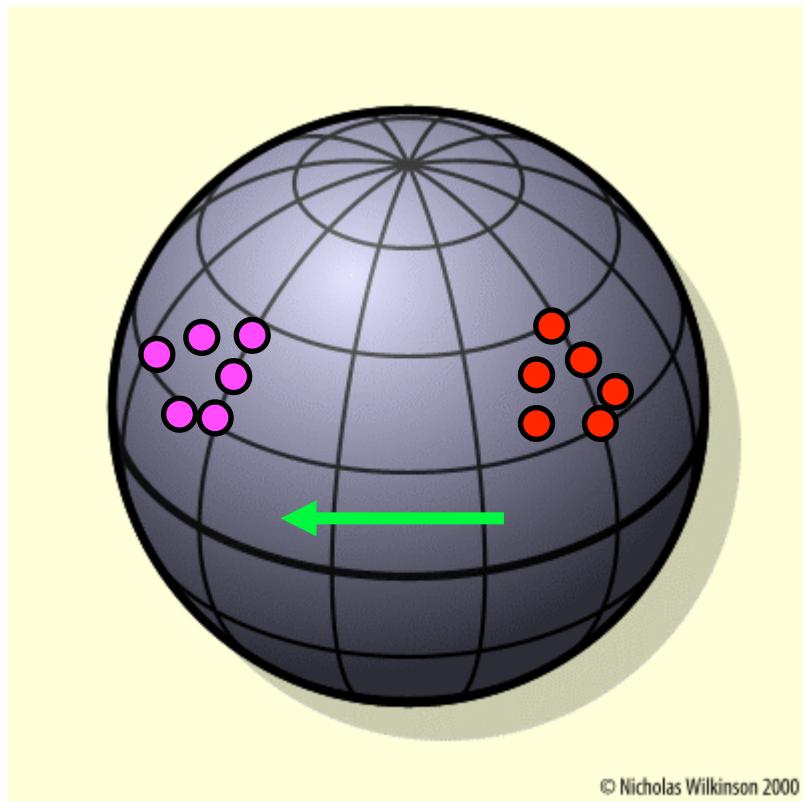
Look down the earth
from your favorite target

Data coverage in Fourier space



Synthesis imaging (6 antennas)

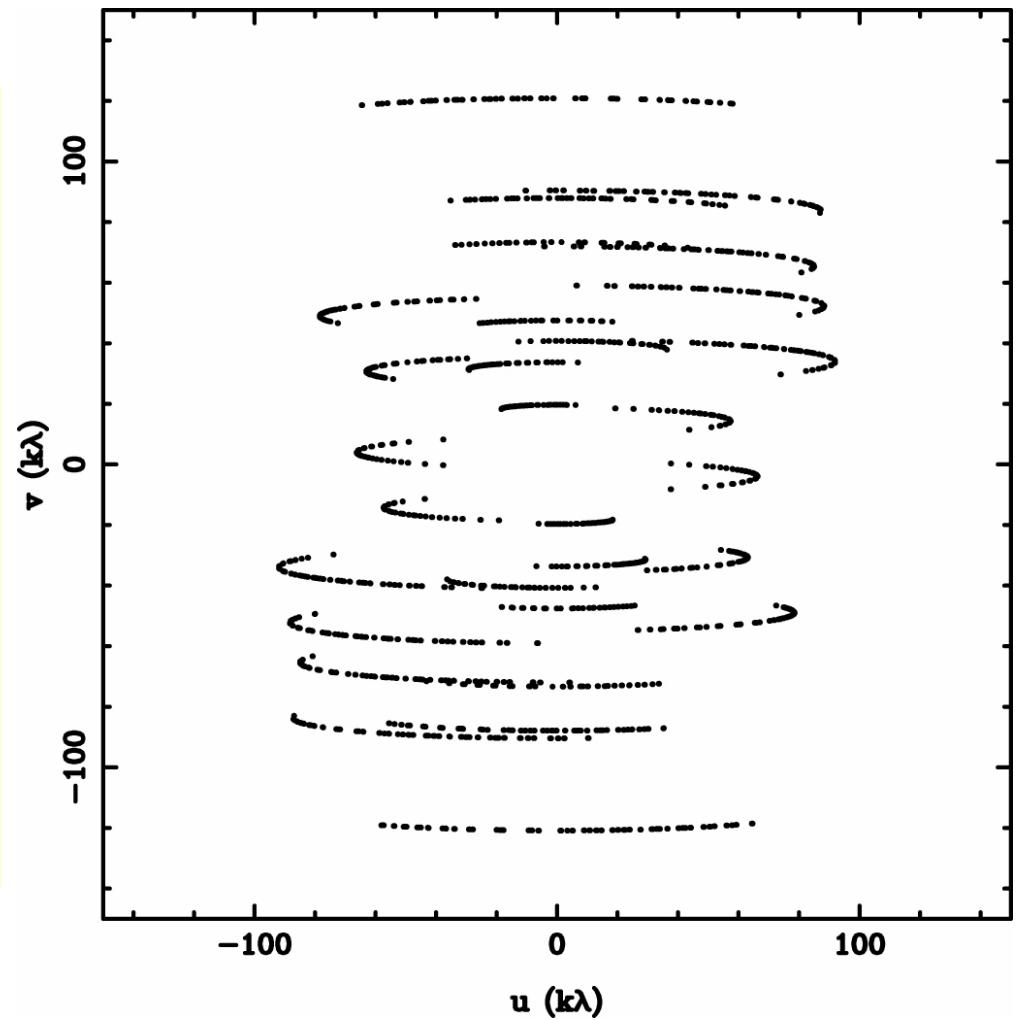
Earth (rotating)



$$\# \text{ of ant pairs} = N(N+1)/2$$

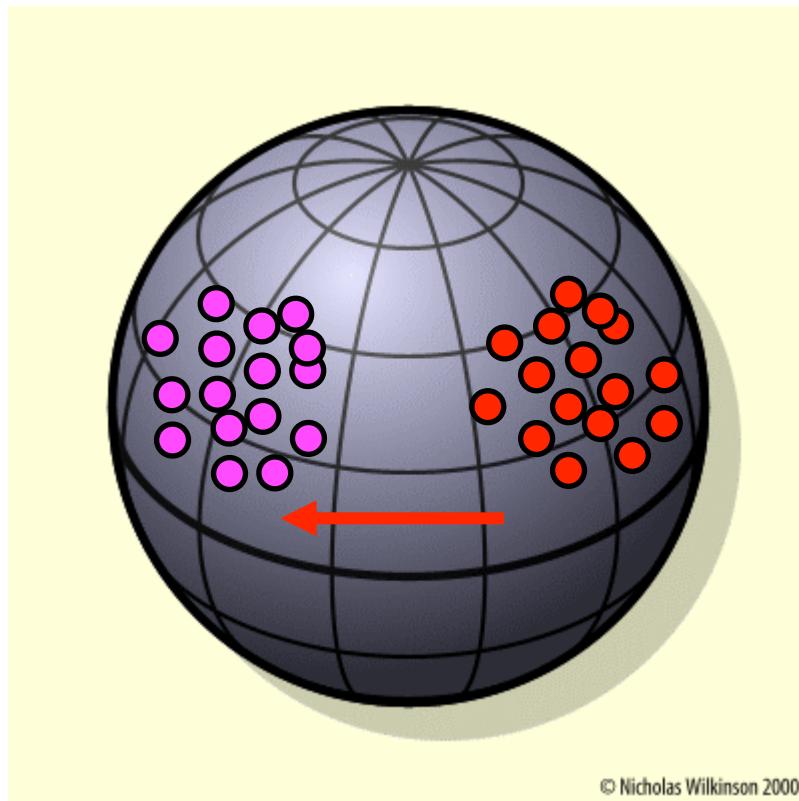
N=6 → 15 pairs

Data coverage in Fourier space



Synthesis imaging (15 antennas)

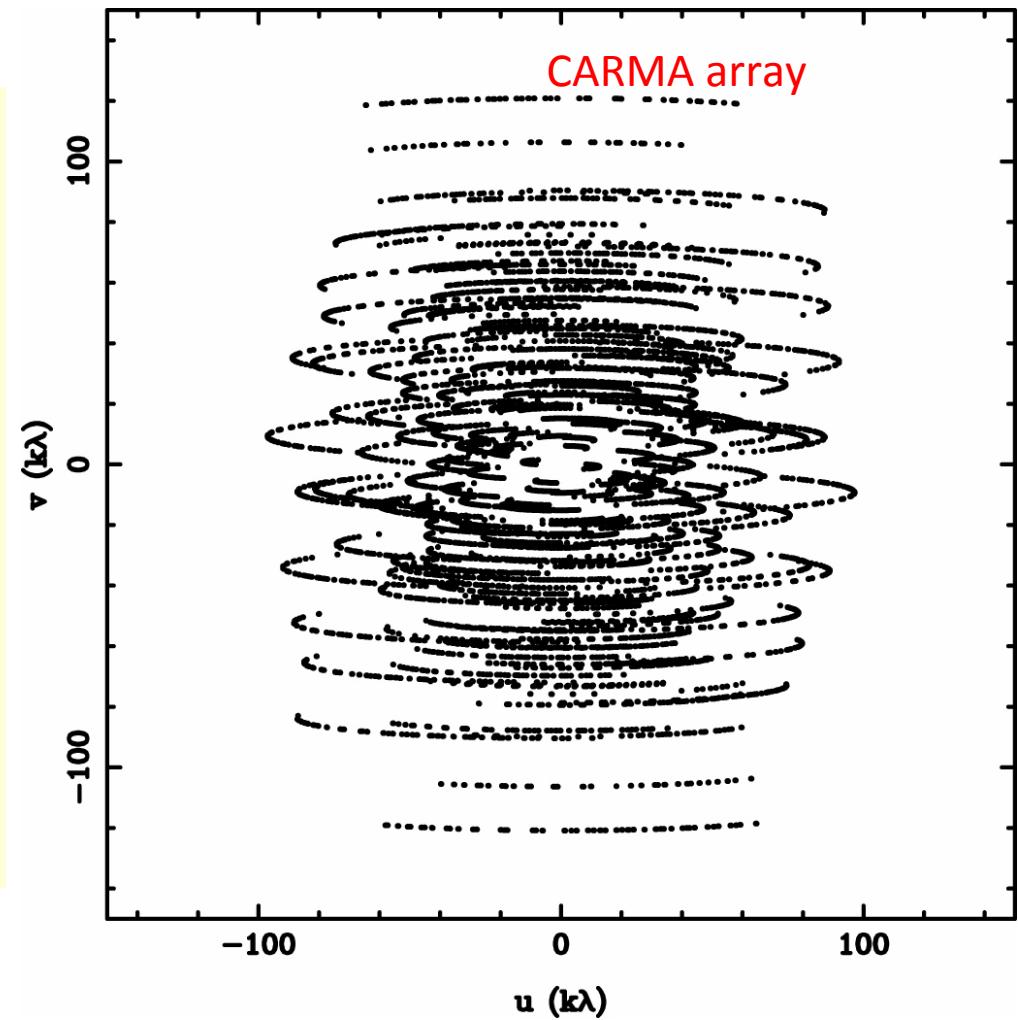
Earth (rotating)



$$\# \text{ of ant pairs} = N(N+1)/2$$

N=15 → 105 pairs!

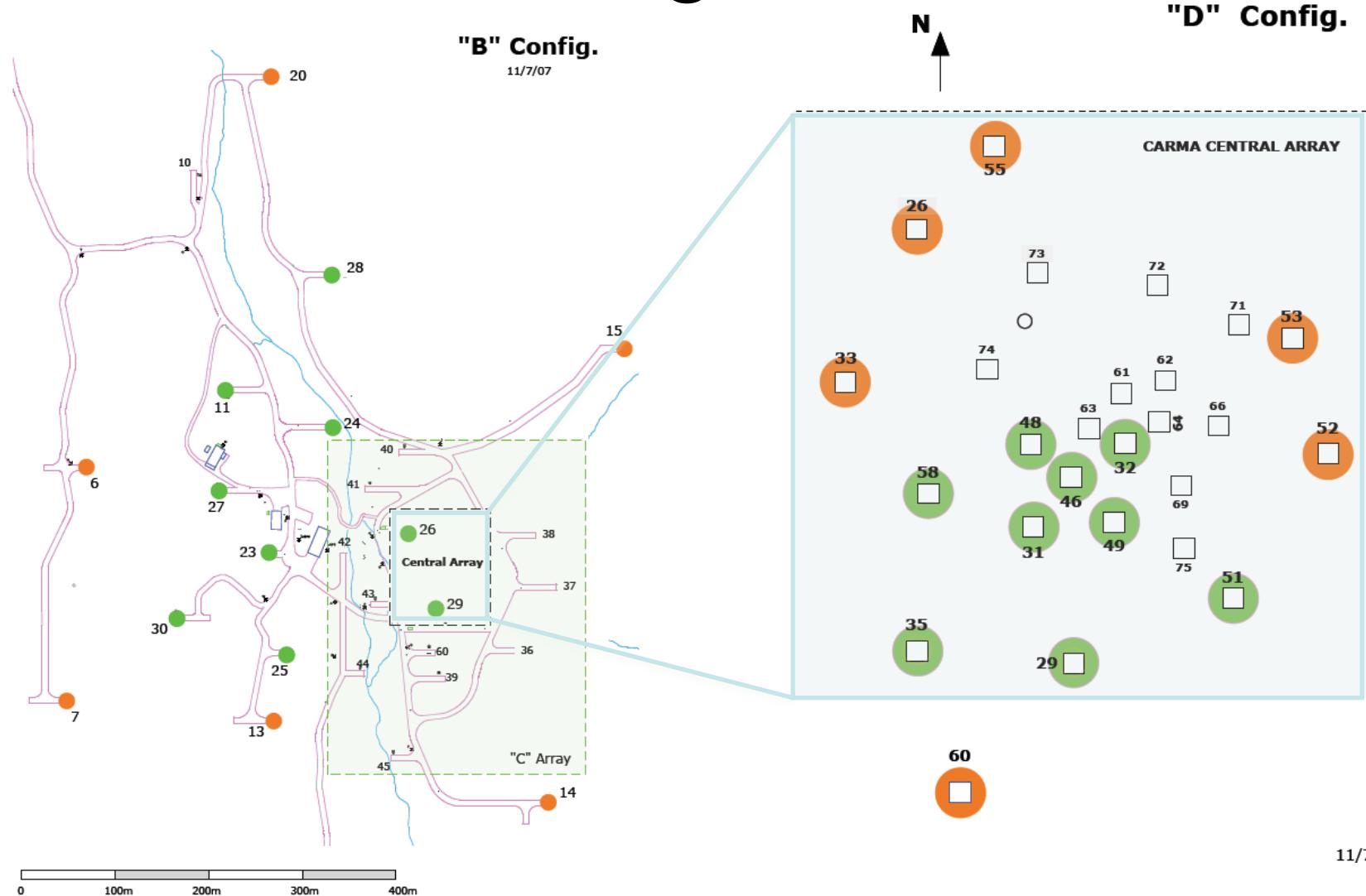
Data coverage in Fourier space



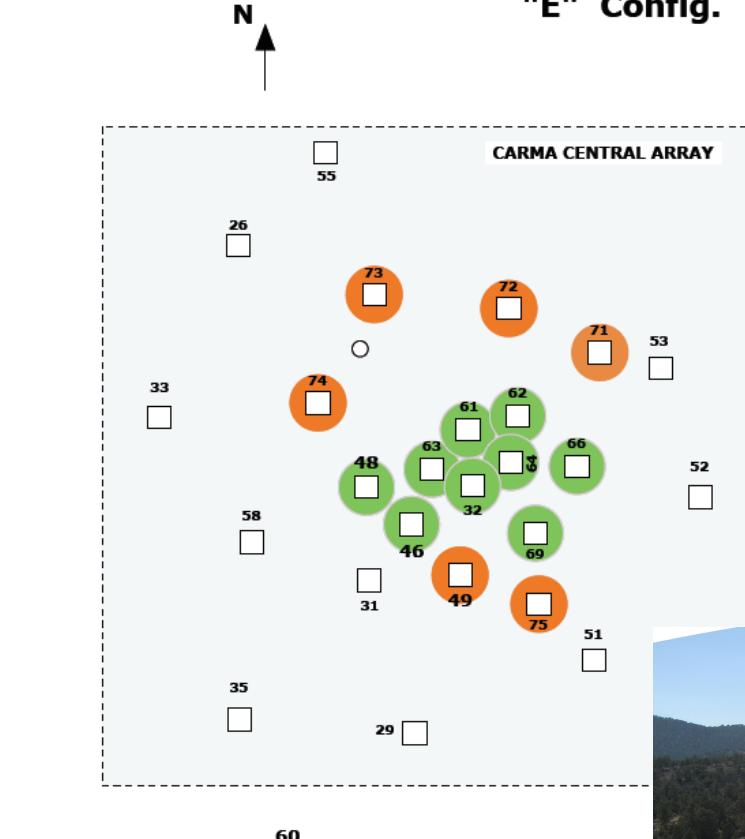
Configuration Change



B-configuration



"E" Config.

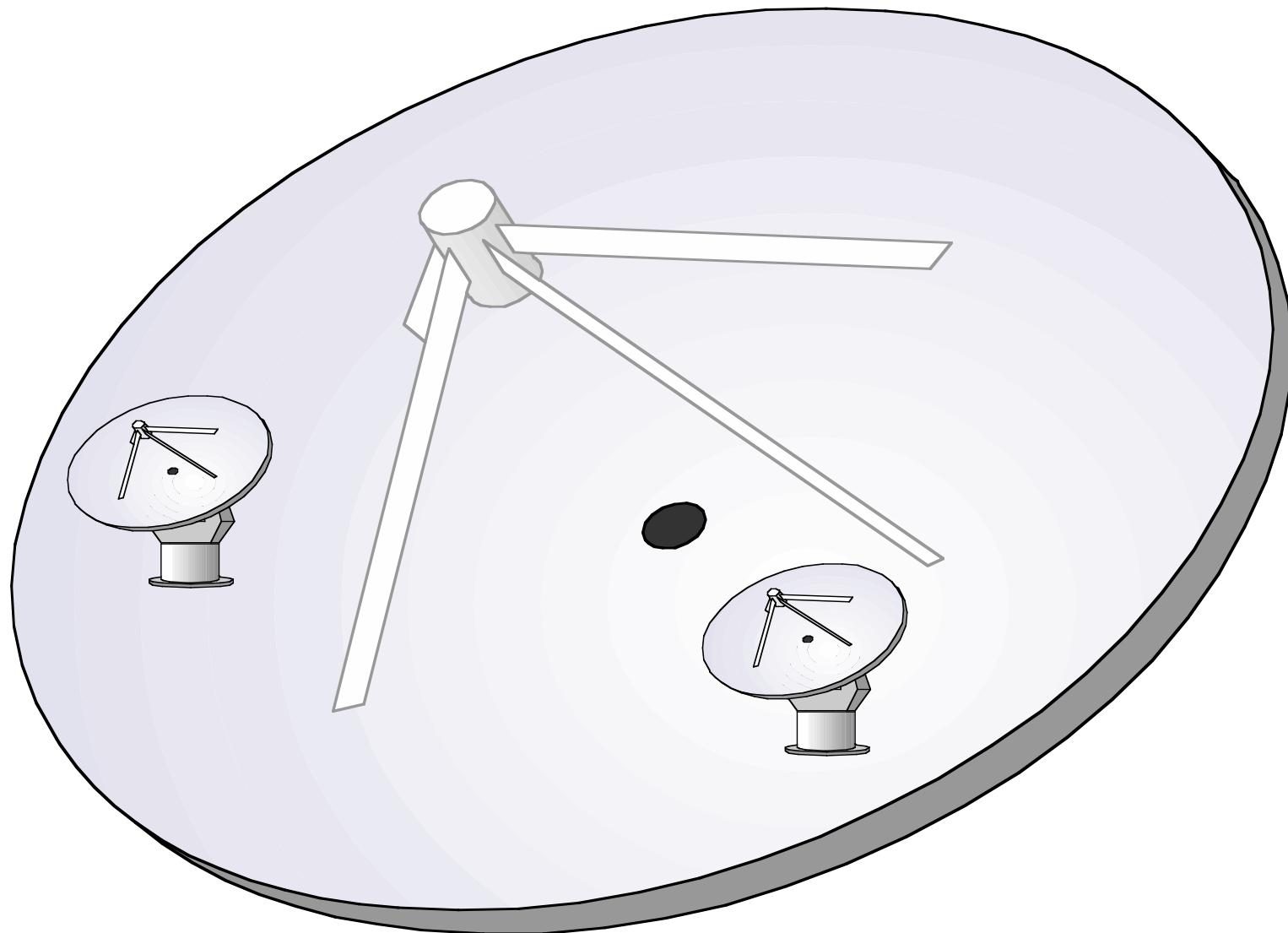


E-configuration

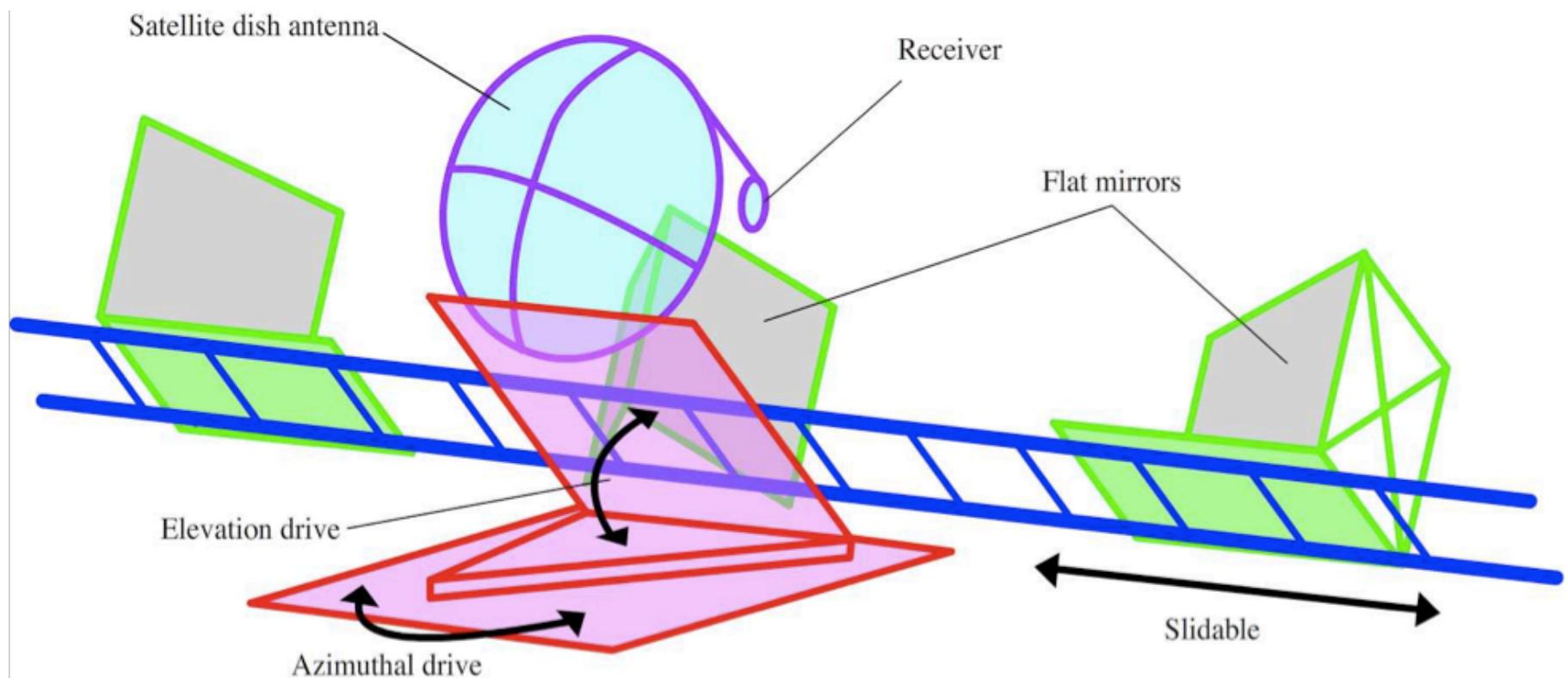


Stony Brook Radio Interferometer

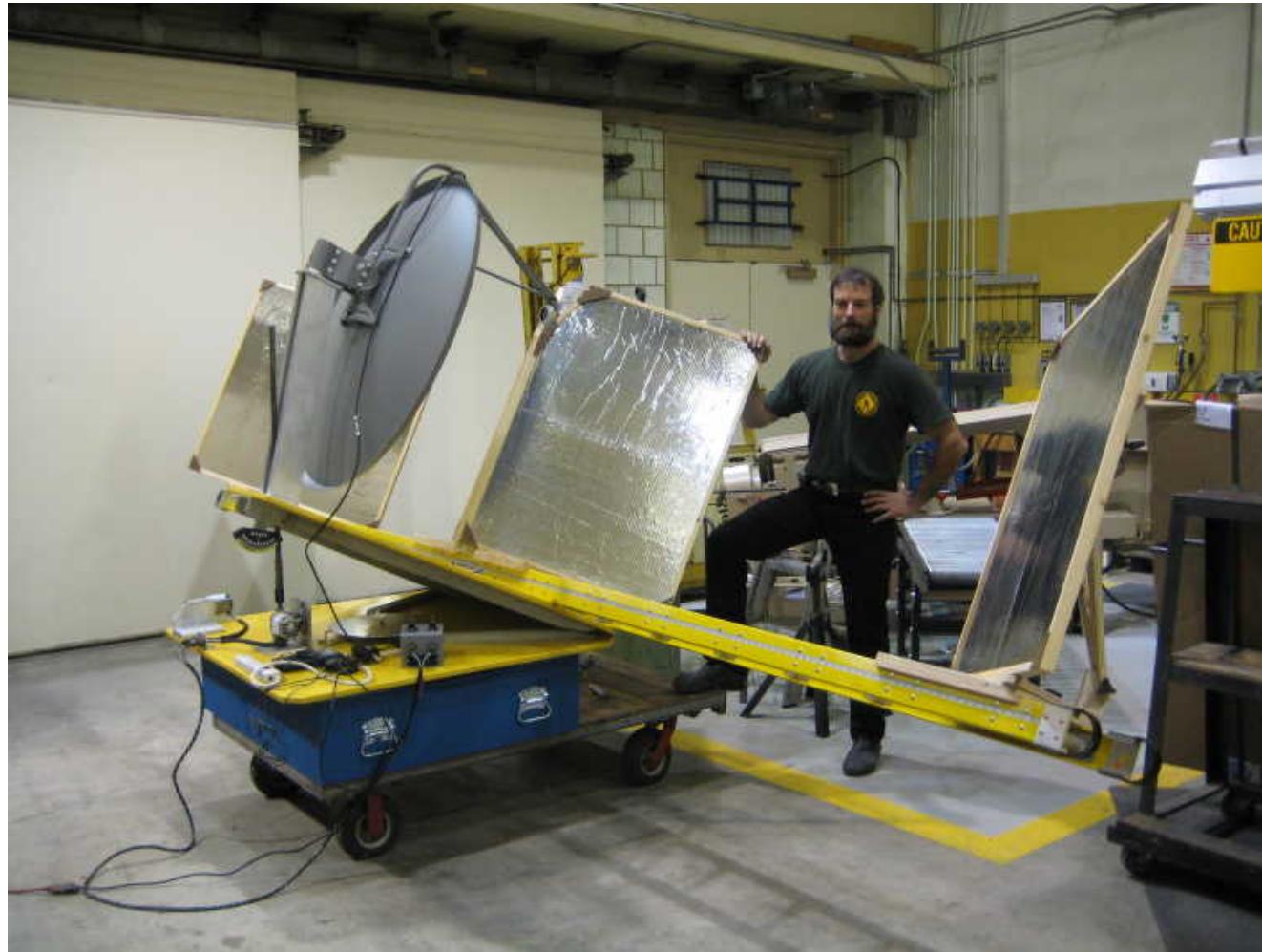
Interferometer (2 Elements)

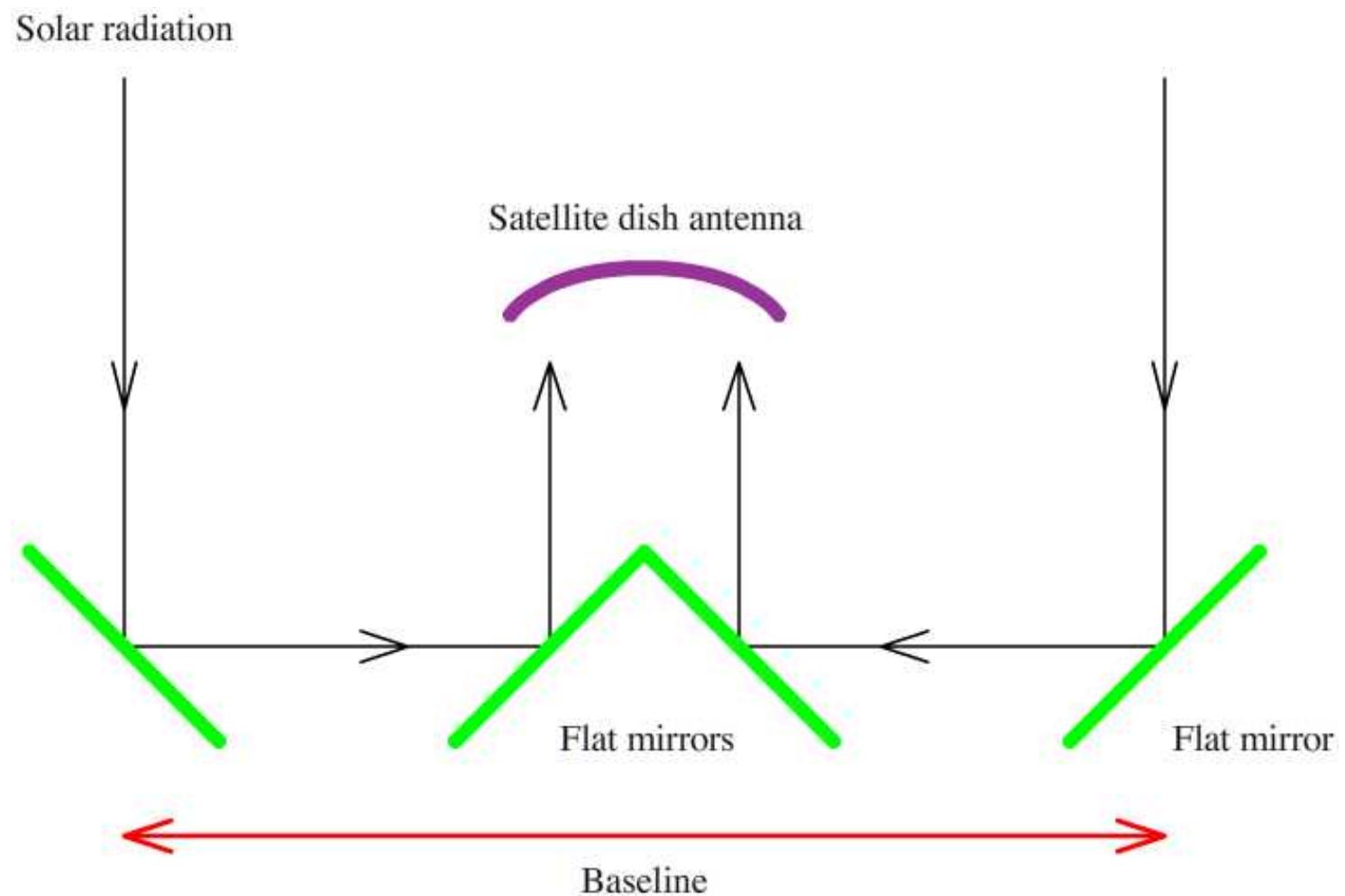


Stony Brook Radio Interferometer



Stony Brook Radio Interferometer





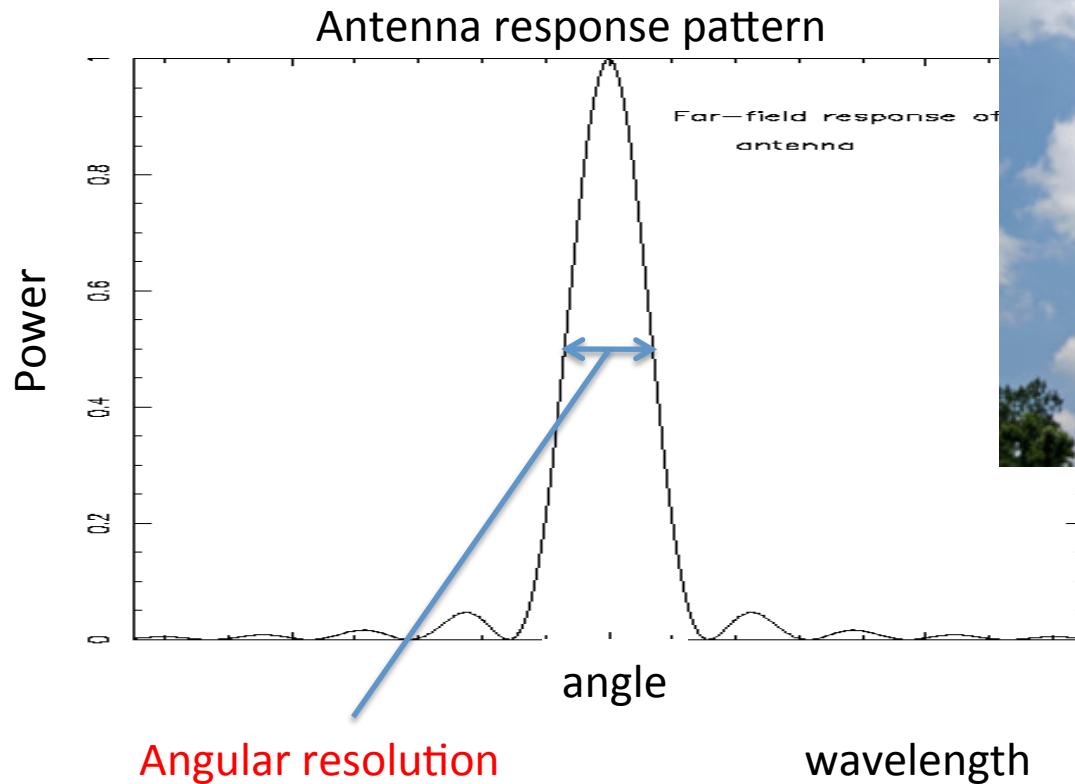


Components

Telescope Structure



Antenna



wavelength

$$\lambda = 2.6\text{cm}$$

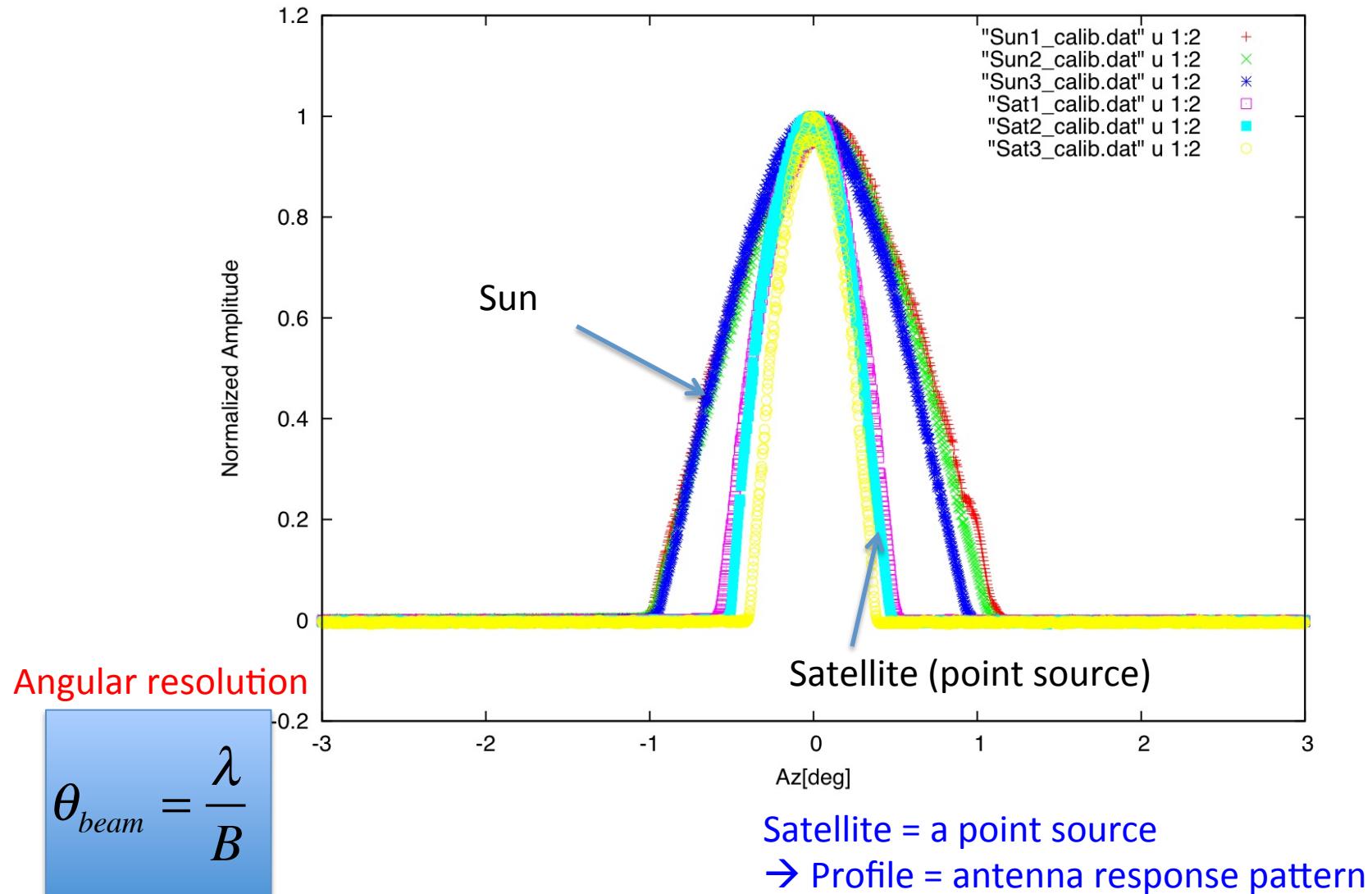
Antenna diameter

$$B = 1\text{m}$$

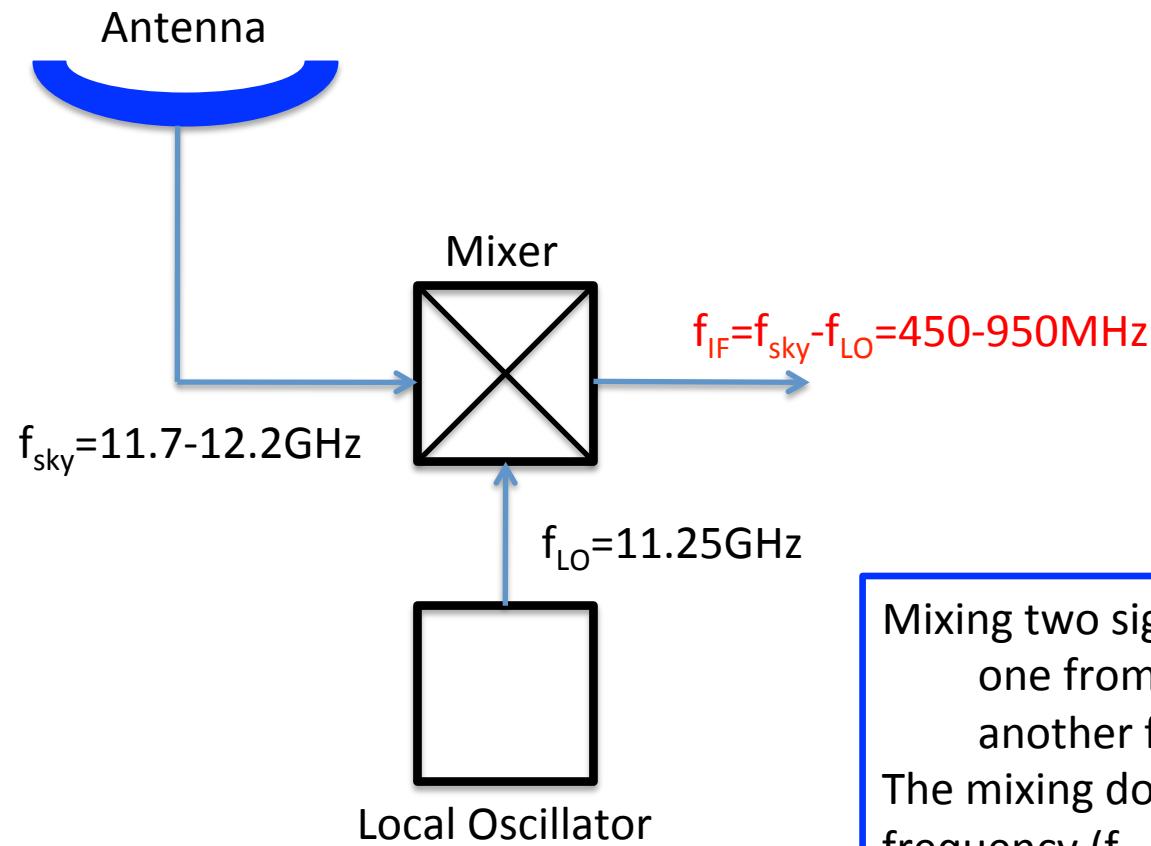
$$\theta = 1.45 \text{ deg}$$

$$\theta_{beam} = \frac{\lambda}{B}$$

Antenna Response Pattern



Heterodyne Receiver I



Mixing two signals:
one from the sky
another from the local oscillator.
The mixing down-converts the high sky
frequency (f_{sky}) to a lower frequency ($f_{\text{IF}}=f_{\text{sky}}-f_{\text{LO}}$). The low frequency is easier to handle.

$$\sin(2\pi f_1 t) \sin(2\pi f_2 t) = \frac{1}{2} \cos[2\pi(f_1 - f_2)t] - \frac{1}{2} \cos[2\pi(f_1 + f_2)t]$$

Down-converted frequency

Heterodyne Receiver I

Non-linear response of an amplifier (inside the receiver)

$$F(v) = \underline{\alpha_1 v + \alpha_2 v^2} + \alpha_3 v^3 + \dots$$

Let's think only about the first two terms

$$v_{out} = F(\underline{A_1 \sin \omega_1 t + A_2 \sin \omega_2 t})$$

Signals from the sky and local oscillator

$$v_{out} = \alpha_1(A_1 \sin \omega_1 t + A_2 \sin \omega_2 t) + \alpha_2(A_1 \sin \omega_1 t + A_2 \sin \omega_2 t)^2 + \dots$$

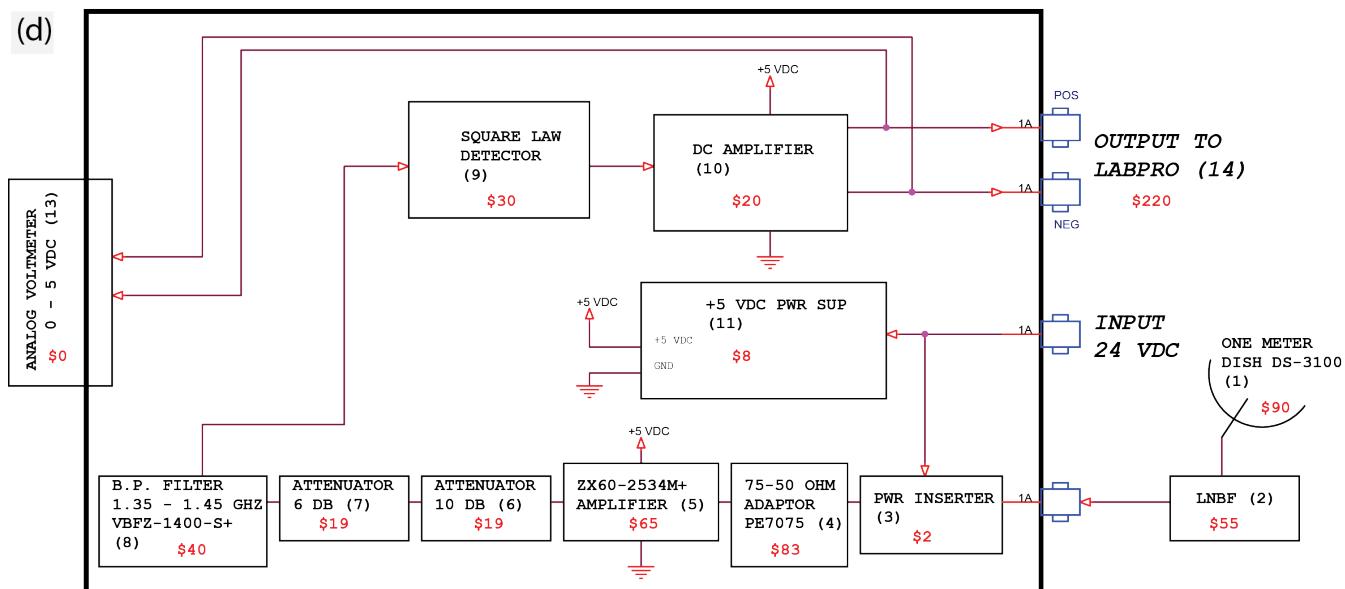
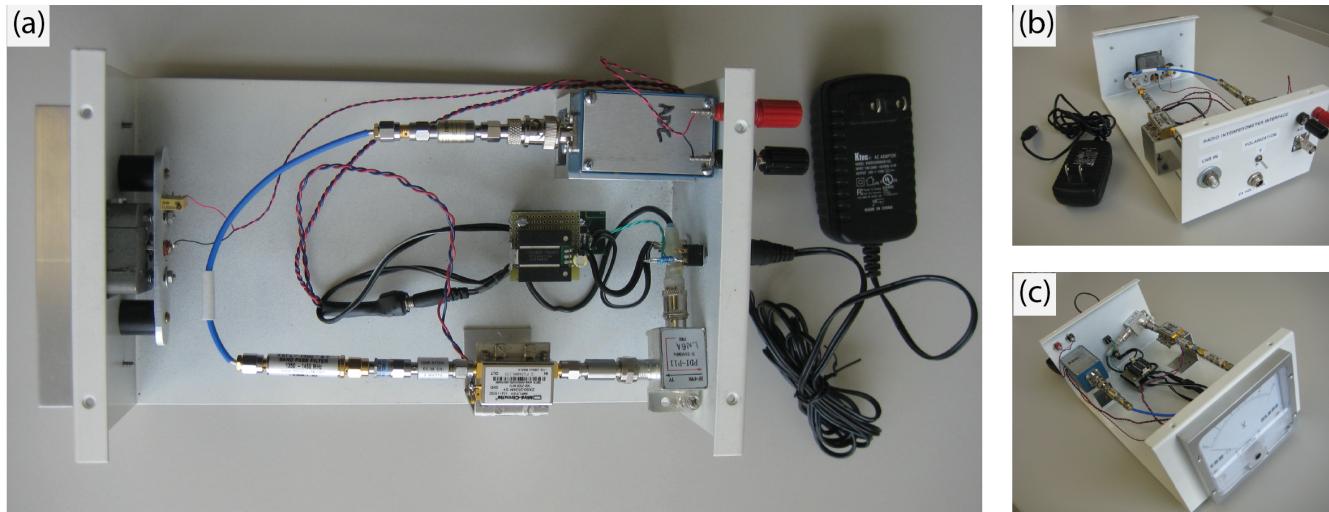
Trigonometric identities

$$\begin{aligned} v_{out} &= \alpha_1(A_1 \sin \omega_1 t + A_2 \sin \omega_2 t) \\ &\quad + \alpha_2\left(\frac{A_1^2}{2}[1 - \sin 2\omega_1 t] + A_1 A_2[\cos(\omega_1 t - \omega_2 t) - \cos(\omega_1 t + \omega_2 t)]\right. \\ &\quad \left. + \frac{A_2^2}{2}[1 - \sin 2\omega_2 t]\right) + \dots \end{aligned}$$

Filter out unnecessary frequencies

$$v_{out} = \underline{\alpha_2 A_1 A_2 \cos(\omega_1 - \omega_2)t - \alpha_2 A_1 A_2 \cos(\omega_1 + \omega_2)t} + \dots$$

Detector



A/D Converter

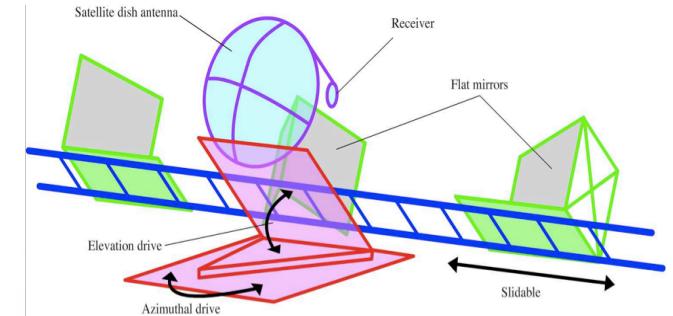


Analog signal → Digital signal → Computer

Mirrors and Skin Depth

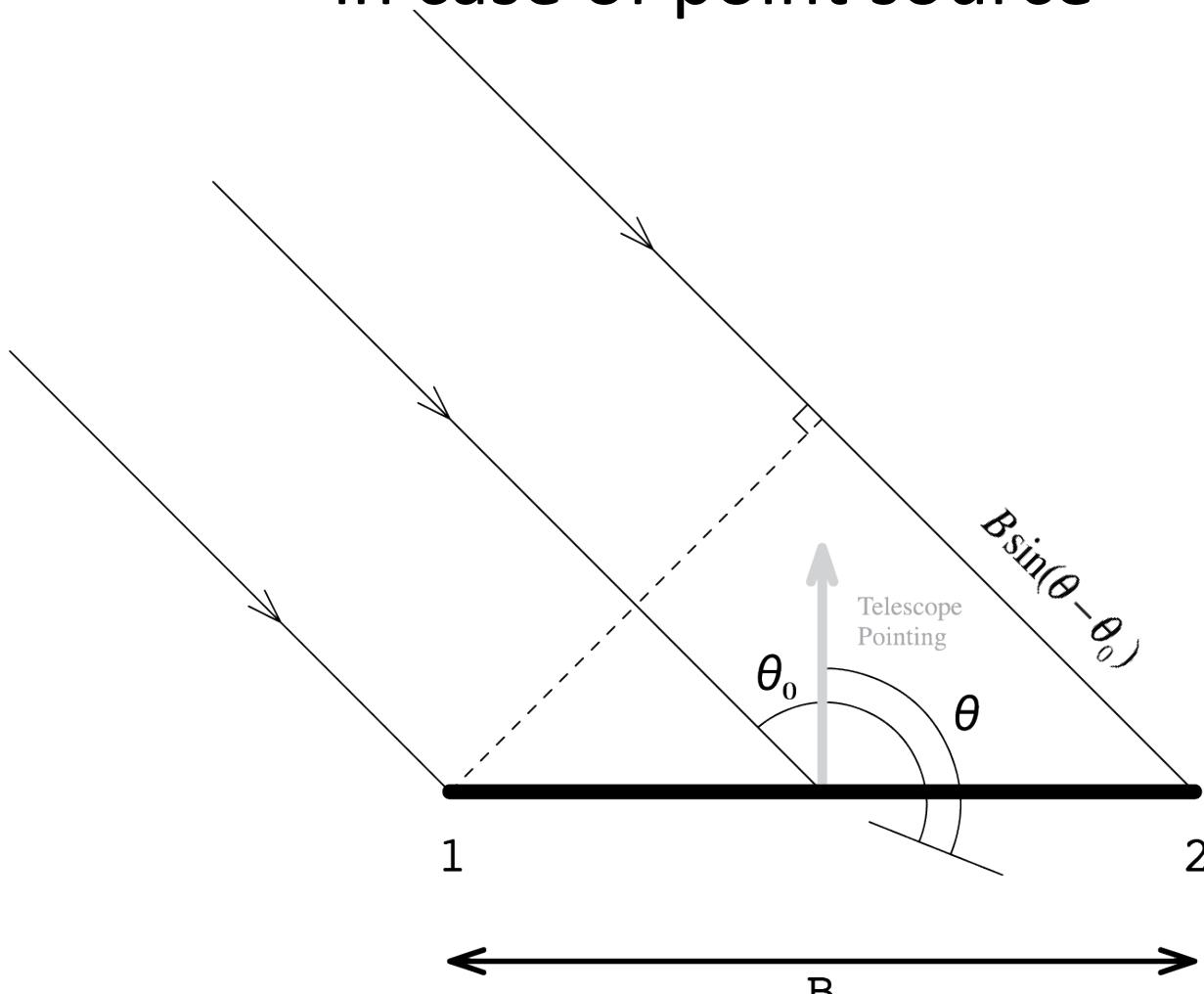
- Material that reflects 2.6cm radio signals
 - Aluminum: skin depth $\sim 0.8\text{micron}$
- Necessary surface accuracy
 - $\lambda/10 \sim 3\text{mm}$

Kitchen aluminum
foil is good enough.



Equations

Mixing Signals from Two Mirrors in case of point source



$$E_1(\theta) = E(\theta_0) \cos[2\pi v t]$$

Signal from Mirror 1

$$E_2(\theta) = E(\theta_0) \cos[2\pi v(t - \tau)]$$

Signal from Mirror 2

Total Power in case of point source

Total power – we observe this!

$$P(\theta) = \langle E_{tot}^2(\theta) \rangle = \langle (E_1(\theta) + E_2(\theta))^2 \rangle$$

$$= \left\langle E^2(\theta_0) \left[\cos^2(2\pi v t) + \cos^2(2\pi v(t - \tau)) + 2 \cos(2\pi v t) \cos(2\pi v(t - \tau)) \right] \right\rangle$$

$$= \left\langle E^2(\theta_0) \left[\frac{1 + \cos(4\pi v t)}{2} + \frac{1 + \cos(4\pi v(t - \tau))}{2} + \cos(2\pi v(2t - \tau)) + \cos(2\pi v\tau) \right] \right\rangle$$

$$= \left\langle E^2(\theta_0) [1 + \cos(2\pi v\tau)] \right\rangle$$

$$\downarrow \quad \tau = \frac{B \sin(\theta - \theta_0)}{c} \quad B_\lambda = \frac{B}{\lambda} \quad \sin(\theta - \theta_0) \approx \theta - \theta_0$$

Time average $\rightarrow 0$

$$P(\theta) \approx \left\langle E^2(\theta_0) [1 + \cos[2\pi B_\lambda(\theta - \theta_0)]] \right\rangle$$

Total Power

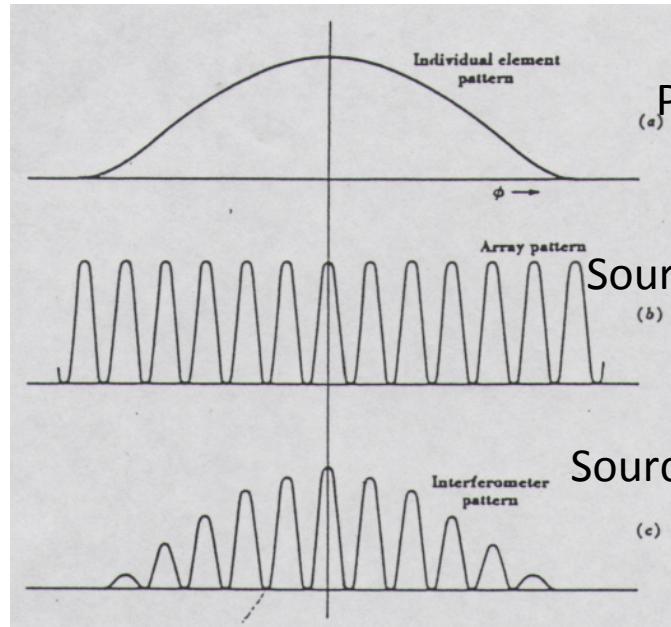
In case of general astronomical objects

$$P(\theta) \approx \int \underline{\varepsilon(\theta_0)} [1 + \cos[2\pi B_\lambda(\theta - \theta_0)]] d\theta_0$$

Energy density

Point source (broadcast satellite)

$$\varepsilon(\theta_0) = \varepsilon_0 \delta(\theta_0 - \theta_c)$$



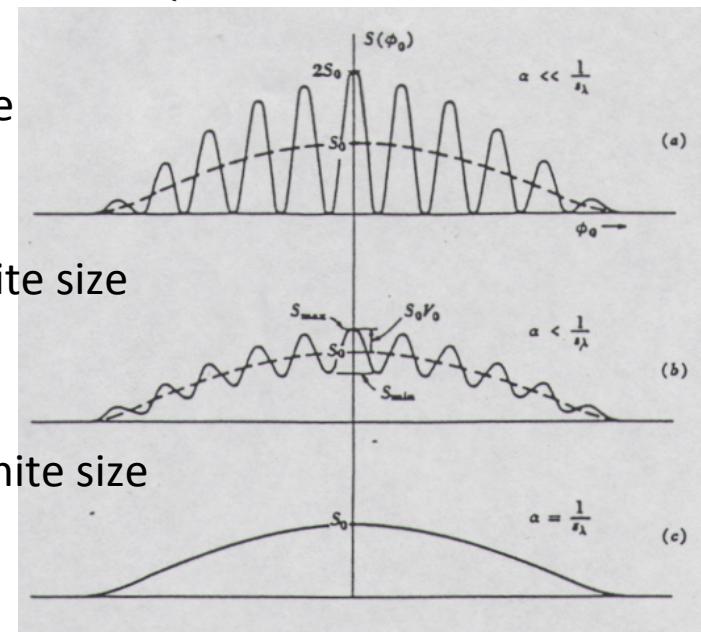
Disk (Sun)

$$\varepsilon(\theta^0) = \begin{cases} 0 & (\text{rest}) \\ \varepsilon^0 & (-\alpha \setminus \zeta < \theta < \alpha \setminus \zeta) \end{cases}$$

Point source

Source with finite size

Source with infinite size



Visibility

Get info from $P(\theta)$.

$$P(\theta) \approx \int \varepsilon(\theta_0) [1 + \cos[2\pi B_\lambda(\theta - \theta_0)]] d\theta_0$$

$$\equiv S_0 [1 + V(\theta, B_\lambda)]$$

$$S_0 \equiv \int \varepsilon(\theta_0) d\theta_0$$
$$\varepsilon'(\theta_0) \equiv \varepsilon(\theta_0) / S_0$$

$$V(\theta, B_\lambda) \equiv \int \varepsilon'(\theta_0) \cos[2\pi B_\lambda(\theta - \theta_0)] d\theta_0$$

$$= \cos[2\pi B_\lambda \theta] \int \varepsilon'(\theta_0) \cos[2\pi B_\lambda \theta_0] d\theta_0$$
$$+ \sin[2\pi B_\lambda \theta] \int \varepsilon'(\theta_0) \sin[2\pi B_\lambda \theta_0] d\theta_0$$

$$\equiv \underline{V_0(B_\lambda)} \cos[2\pi B_\lambda(\theta - \Delta\theta)]$$

“visibility”.

Visibility = F. T. of Energy Distribution

What is visibility?

By definition:

$$V_0(B_\lambda) \cos[2\pi B_\lambda \Delta\theta] = \int \varepsilon'(\theta_0) \cos(2\pi B_\lambda \theta_0) d\theta_0 \quad (1)$$

$$V_0(B_\lambda) \sin[2\pi B_\lambda \Delta\theta] = \int \varepsilon'(\theta_0) \sin(2\pi B_\lambda \theta_0) d\theta_0 \quad (2)$$

(1)+i(2)

$$V_0(B_\lambda) = \underbrace{\int \varepsilon'(\theta_0) \exp(i2\pi B_\lambda \theta_0) d\theta_0}_{\text{Fourier transformation of the energy distribution}} \times \underbrace{\exp[-i2\pi B_\lambda \Delta\theta]}_{\text{Phase shift}}$$

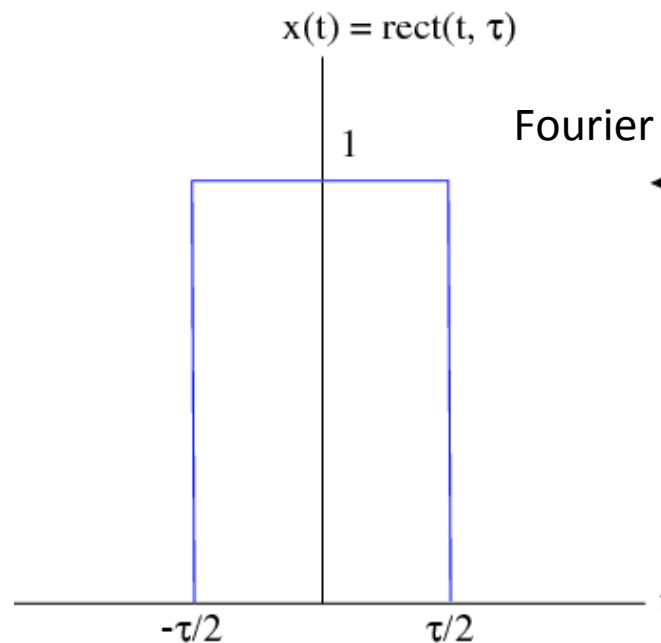
Reminder: we observe the power $P(\theta)$,

$$P(\theta) \equiv S_0 \left[1 + \underbrace{V_0(B_\lambda) \cos[2\pi B_\lambda (\theta - \Delta\theta)]}_{\text{amplitude of sinusoidal power change.}} \right]$$

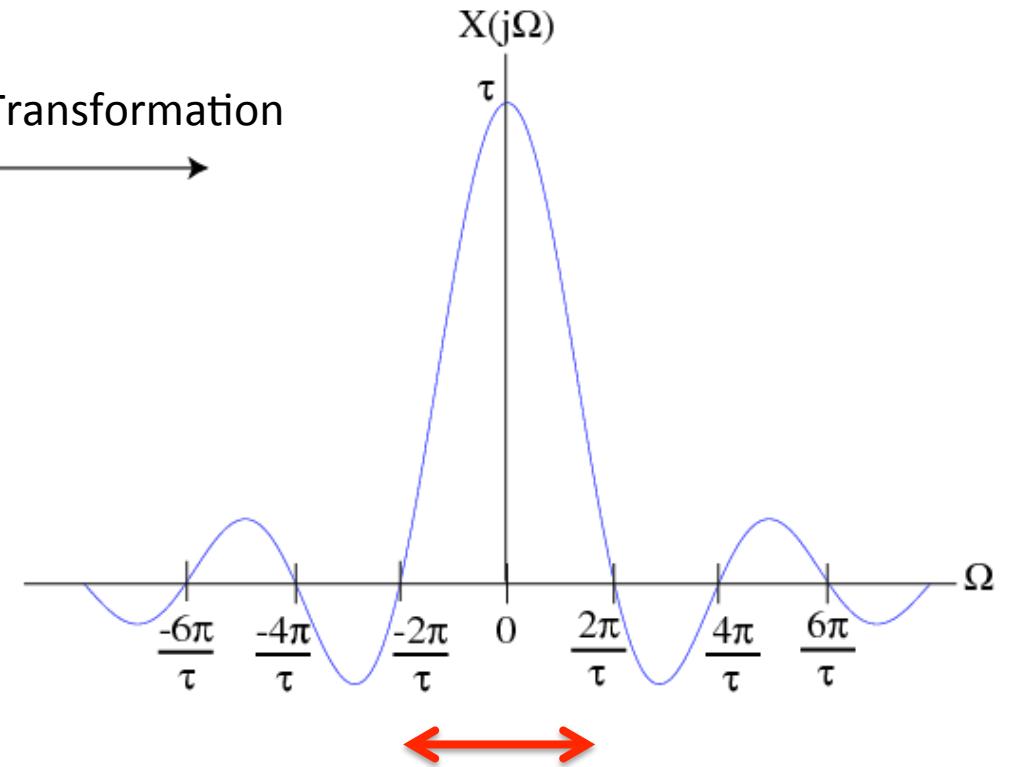
Visibility = F.T. of energy distribution across the target = amplitude of sinusoidal power change.

Sun ~ box shape in 2-d slice

The Sun in the sky



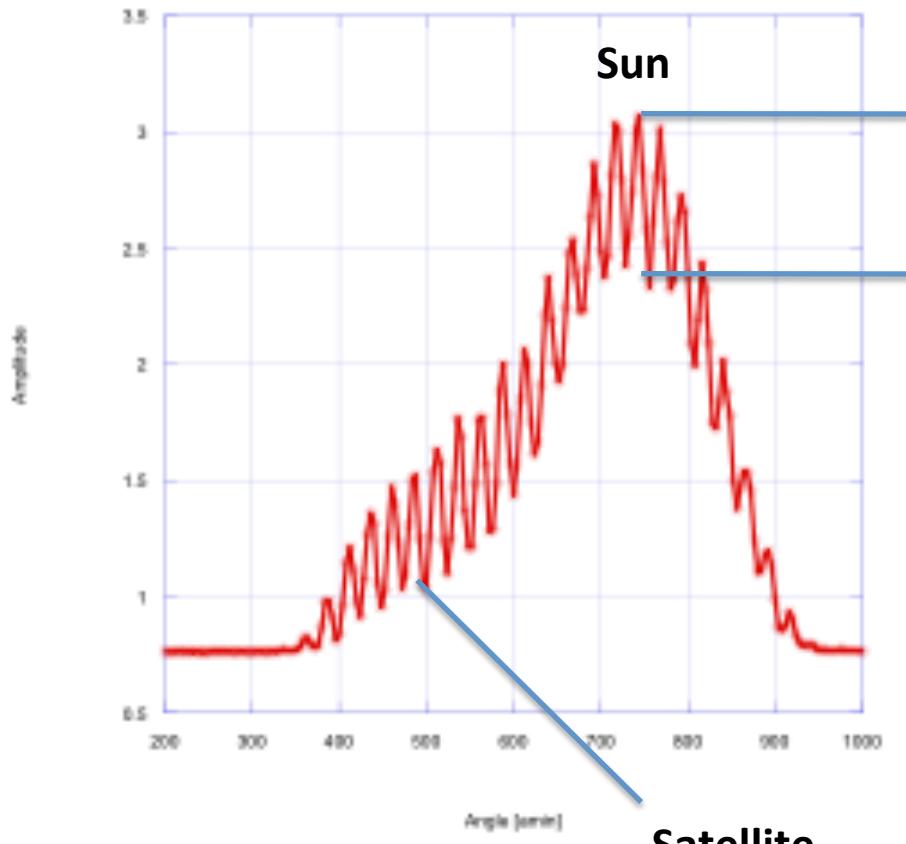
Fourier Transformation



This width gives us the diameter of the Sun
-- do calculations yourself --

How to Measure Visibilities

$$P(\theta) = S_0 [1 + V_0(B_\lambda) \cos[2\pi B_\lambda (\theta - \Delta\theta)]]$$



which happened to be next to the Sun
at the time of measurement.

$$P_{\max} = S_0 [1 + V_0(B_\lambda)]$$

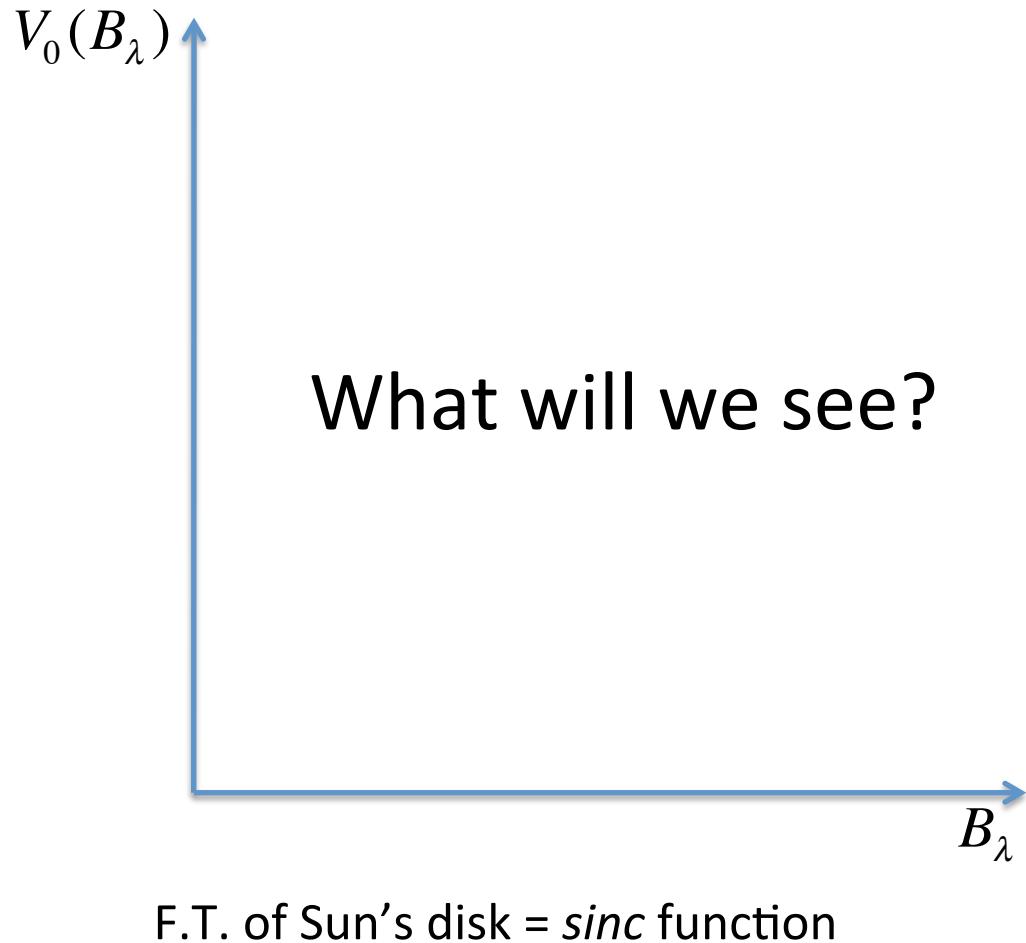
$$P_{\min} = S_0 [1 - V_0(B_\lambda)]$$

$$V_0(B_\lambda) = \frac{P_{\max} - P_{\min}}{P_{\max} + P_{\min}}$$

One Fourier component

Changing B (baseline length)
→ Many Fourier components

What Do We Expect?



Work on
equations and
figure this out.
Try the lab
measurements
and see if you
see what you
expect.