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## 1 Abstract

## 2 Introduction

## 3 Physical theory

### 3.1 Gravitation

In this project, a solar system will be studied. By solar system, I mean a system where only gravitational forces affect the bodies. Newtons gravitational law states that the gravitational force on a body with mass  $m$  from another body with mass  $M$  and relative position  $\vec{r}$  is given by

$$\vec{F}_G = -\frac{GmM}{\|\vec{r}\|^2}\vec{r}$$

where  $G$  is the gravitational constant,  $6.67 \cdot 10^{-11} \text{ Nm}^2/\text{s}^2$ . The direction of the force is given by the fact that gravity is an attractive force. If one of the objects is the sun, the mass is denoted by  $M_\odot$ . With the sun placed in origo,  $r$  is simply the norm of the position vector of the planet with mass  $m$ .

If there are  $n$  planets in the solar system, in addition to the sun, the sum of the forces on planet  $i$  with mass  $m_i$  is

$$\sum \vec{F}_i = \sum_{\substack{k=0 \\ k \neq i}}^n \frac{Gm_i m_k}{\|\vec{r}_k - \vec{r}_i\|^3} (\vec{r}_k - \vec{r}_i)$$

with  $m_0 = M_\odot$  and  $\vec{r}_0 = \vec{0}$ . From Newton's second law, we know that  $\sum \vec{F}_i = m_i \vec{a}_i$ , so the acceleration of planet  $i$  is given by

$$\vec{a}_i = \sum_{\substack{k=0 \\ k \neq i}}^n \frac{Gm_k}{\|\vec{r}_k - \vec{r}_i\|^3} (\vec{r}_k - \vec{r}_i) \quad (1)$$

As the sun has been fixed to origo,  $\vec{a}_0$  is set to  $\vec{0}$ .

### 3.2 Choice of units

In the solar system, seconds and meters are unpractical, as planets are millions of kilometers apart and take years to do one lap around the sun. As such, it is common to use so-called astronomical units, where 1 ua is the mean distance between the sun and the earth, and time is measured in years. To express the gravitational constant in these units, we can use that if the earth were moving in a circle around the sun, the acceleration in Newton's second law would be given by the centripetal acceleration:

$$\frac{mv^2}{r} = \frac{GmM_\odot}{r^2} \implies G = \frac{r}{M_\odot} v^2 = \frac{1 \text{ ua}}{M_\odot} \cdot \left( \frac{2\pi \cdot 1 \text{ ua}}{1 \text{ yr}} \right)^2 = \frac{4\pi^2}{M_\odot} \cdot 1 \text{ ua}^3/\text{yr}^2$$

With this change of units, equation (1) can be written as

$$\vec{a}_i = \sum_{\substack{k=0 \\ k \neq i}}^n 4\pi^2 \frac{m_k}{M_\odot} \frac{\vec{r}_k - \vec{r}_i}{\|\vec{r}_k - \vec{r}_i\|^3} \cdot 1 \text{ ua}^3/\text{yr}^2 \quad (2)$$

## 4 Mathematical theory

Equation (2), together with some initial conditions, determines the motion of the bodies in the solar system. When written out in components, the equation gives a coupled set of differential equations. This set of equations is difficult, if at all possible, to solve analytically, so numerical work is required. As per usual, the time is discretised as  $t_i = t_0 + ih$ , where  $h$  is the time step,  $h = (t_n - t_0)/n$ . Acceleration, velocity and position are discretised correspondingly.

### 4.1 Forward Euler

The Forward Euler method, also called the Explicit Euler method, and hereafter called simply the Euler method, uses a first order Taylor polynomial to approximate a solution to the

differential equation. With  $x'(t) = v(t)$  and  $v'(t) = a(t)$ , we have that

$$\begin{aligned}\vec{r}_i(t+h) &\approx \vec{r}_i(t) + h\vec{v}_i(t) \\ \vec{v}_i(t+h) &\approx \vec{v}_i(t) + h\vec{a}_i(t)\end{aligned}$$

Discretised version:

$$\begin{aligned}\vec{r}_{i,j+1} &\approx \vec{r}_{i,j} + h\vec{v}_{i,j} \\ \vec{v}_{i,j+1} &\approx \vec{v}_{i,j} + h\vec{a}_{i,j}\end{aligned}$$

To clarify the indices:  $\vec{a}_{i,j}$  is the acceleration of planet  $i$  at time step  $j$ . This is calculated from equation (2) on the preceding page.

From Taylor's formula, the error for a first order Taylor polynomial goes as  $O(h^2)$ . This is the error made in each step — the error is accumulated, so the total error will be proportional  $h$ .

## 4.2 Velocity-Verlet

The Velocity-Verlet method, hereafter called the Verlet method, is based on a second order Taylor polynomial.

$$\begin{aligned}\vec{r}_i(t+h) &\approx \vec{r}_i(t) + h\vec{v}_i(t) + \frac{1}{2}h^2\vec{a}_i(t) \\ \vec{v}_i(t+h) &\approx \vec{v}_i(t) + h\vec{a}_i(t) + \frac{1}{2}h^2\vec{a}'_i(t)\end{aligned}$$

There is no explicit expression for  $\vec{a}'(t)$ , however it can be approximated using the good old formula

$$\vec{a}'(t) \approx \frac{\vec{a}(t+h) - \vec{a}(t)}{h}$$

Since the acceleration is independent of the velocity, the newly updated position,  $\vec{a}(t+h)$ , can be calculated using  $\vec{r}(t+h)$ . Inserting this into the expression for  $\vec{v}(t+h)$ , we get

$$\begin{aligned}\vec{v}_i(t+h) &\approx \vec{v}_i(t) + h\vec{a}_i(t) + \frac{1}{2}h(\vec{a}_i(t+h) - \vec{a}_i(t)) \\ &= \vec{v}_i(t) + \frac{1}{2}h(\vec{a}_i(t) + \vec{a}_i(t+h))\end{aligned}$$

The discretised version then becomes

$$\begin{aligned}\vec{r}_{i,j+1} &\approx \vec{r}_{i,j} + h\vec{v}_{i,j} + \frac{1}{2}h^2\vec{a}_{i,j} \\ \vec{v}_{i,j+1} &\approx \vec{v}_{i,j} + \frac{1}{2}h(\vec{a}_{i,j} + \vec{a}_{i,j+1})\end{aligned}$$

The error of a second order Taylor polynomial is given as  $O(h^3)$ . The approximation for  $\vec{a}'(t)$  has an error proportional to  $h$ , but this error is multiplied with  $h^2$  when inserted into the expression for  $\vec{v}_i(t+h)$ . As such, the error for each step is proportional to  $h^3$ . The error is again accumulated for each step, so the total error will be proportional to  $h^2$ . This is one order better than the Euler method.

## 5 Implementation

The implementation is heavily object oriented and modular by design. The Planet class ([planet.h](#) and [planet.cpp](#)) represents a body, with methods for calculating the acceleration according to equation (2) as well as updating position and velocity for both algorithms. The SolarSystem class ([solarsystem.h](#) and [solarsystem.cpp](#)) administrates the planets, and contains a method for each algorithm. [cpp\\_ui.cpp](#) sets up the solar system according to the command line arguments, and calls the solver method specified. [python\\_ui.py](#) is a user-friendly interface which runs [cpp\\_ui.cpp](#) according to the command line arguments, see [README.md](#).

### 5.1 Number of floating point operations (FLOPs)

To calculate the acceleration (common for both algorithms):

- For each of the  $n$  planets:
  - For each of the  $n - 1$  other planets:
    1. Subtract the  $x$ ,  $y$  and  $z$  components (3 FLOPs).
    2. Square the differences in the  $x$ ,  $y$  and  $z$  components (3 FLOPs).
    3. Add these differences (2 FLOPs).
    4. Take the square root (1 FLOP<sup>1</sup>).
    5. Calculate the cube of the norm (2 FLOPs).
    6. Divide  $4\pi^2$  times the mass ratio (precalculated) by the cube of the norm (1 FLOP).
    7. Multiply this factor by the differences in  $x$ ,  $y$  and  $z$  positions (3 FLOPs).
    8. Add these to the  $x$ ,  $y$  and  $z$  components of the acceleration (3 FLOP).

$18n(n - 1)$  FLOPs in total.

#### For the Euler method:

- For each of the  $n$  planets:
  1. Multiply the  $x$ ,  $y$  and  $z$  components of the acceleration with the time step (3 FLOPs).
  2. Add these to the  $x$ ,  $y$  and  $z$  components of the velocity (3 FLOPs).
  3. Multiply the  $x$ ,  $y$  and  $z$  components of the velocity with the timestep (3 FLOPs).
  4. Add these to the  $x$ ,  $y$  and  $z$  components of the position (3 FLOPs).
- Calculate the acceleration ( $18n(n - 1)$  FLOPs).

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<sup>1</sup>Or possibly more.

This gives a total of

$$12n + 18n(n - 1) = 18n^2 - 6n \quad (3)$$

FLOPs per time step.

**For the Verlet method:**

- For each of the  $n$  planets:
  1. Multiply the time step with the  $x$ ,  $y$  and  $z$  components of the velocity (3 FLOPs).
  2. Multiply half the square of the time step (precalculated) with the  $x$ ,  $y$  and  $z$  components of the acceleration (3 FLOPs).
  3. Add these to the  $x$ ,  $y$  and  $z$  components of the position (6 FLOPs).
- Calculate the new acceleration ( $18n(n - 1)$  FLOPs).
- For each of the  $n$  planets:
  1. Add the  $x$ ,  $y$  and  $z$  components of the new and old accelerations (3 FLOPs).
  2. Multiply these by half the time step (precalculated) (3 FLOPs).
  3. Add the results to the  $x$ ,  $y$  and  $z$  components of the velocity (3 FLOPs).
  4. Set the new acceleration to be the old acceleration, and set the new acceleration to  $\vec{0}$  (0 FLOPs).

This gives a total of

$$12n + 18n(n - 1) + 9n = 18n^2 + 3n \quad (4)$$

FLOPs per time step.

## 6 Results

### 6.1 Earth-sun-system

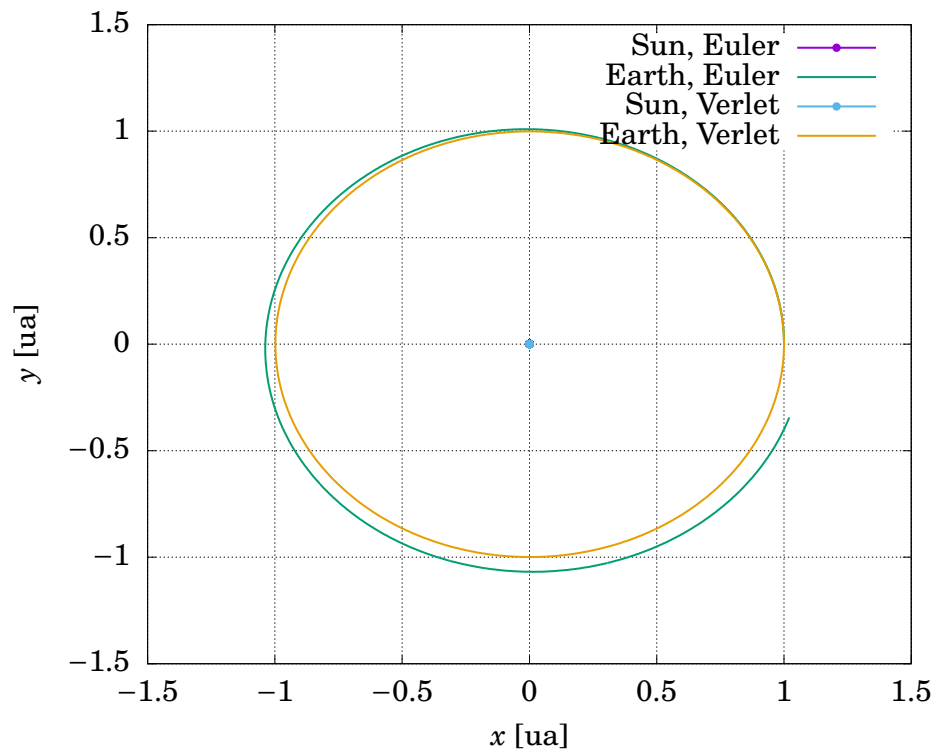


Figure 1: Simulation of the earth-sun-system with both the Forward Euler and the Velocity-Verlet algorithms. The simulation was run for 1 year with 1000 integration points, which clearly is insufficient for the Forward Euler method. The simulation was run by [earthsun.sh](https://github.com/andersjohansson/earthsun.sh).

#### 6.1.1 Visual error analysis

When given a tangential velocity of  $2\pi$  ua/yr, the earth should return to its initial position after one year.

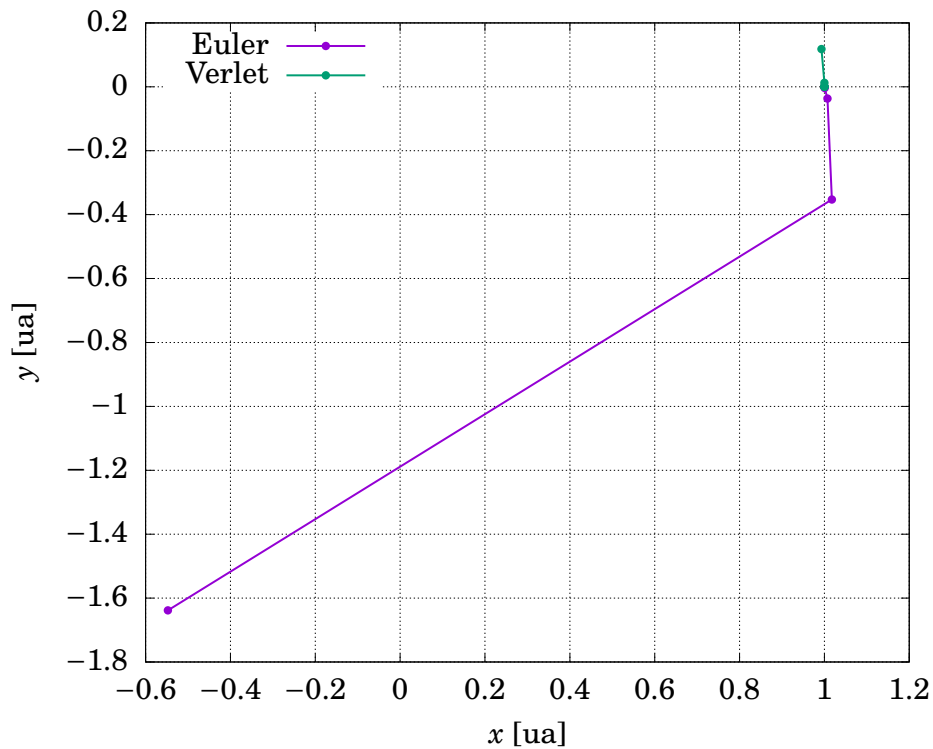


Figure 2: The end point of the earth after simulating for one year, with  $10^n$  timesteps for  $n = 2, 3, \dots$ . It is clear that the Verlet method converges much more quickly than Euler's method. The data is generated by [errortearthsun.sh](https://github.com/AndersJohansson/errortearthsun.sh).

## 6.2 Algorithm timing

Table 1: Benchmarks for both algorithms for the earth-sun-system. Generated by [timetable.py](https://github.com/AndersJohansson/timetable.py).

| $n$       | Euler time            | Verlet time           | Ratio |
|-----------|-----------------------|-----------------------|-------|
| 100       | $2.1 \cdot 10^{-5}$ s | $2.6 \cdot 10^{-5}$ s | 1.24  |
| 1000      | $7.3 \cdot 10^{-5}$ s | $9.7 \cdot 10^{-5}$ s | 1.33  |
| 10000     | 0.00053 s             | 0.00067 s             | 1.26  |
| 100000    | 0.00524 s             | 0.00659 s             | 1.26  |
| 1000000   | 0.0517 s              | 0.0655 s              | 1.26  |
| 10000000  | 0.598 s               | 0.769 s               | 1.29  |
| 100000000 | 5.77 s                | 6.68 s                | 1.16  |