

Neutrino Self-Interaction and Core-Collapse Supernovae

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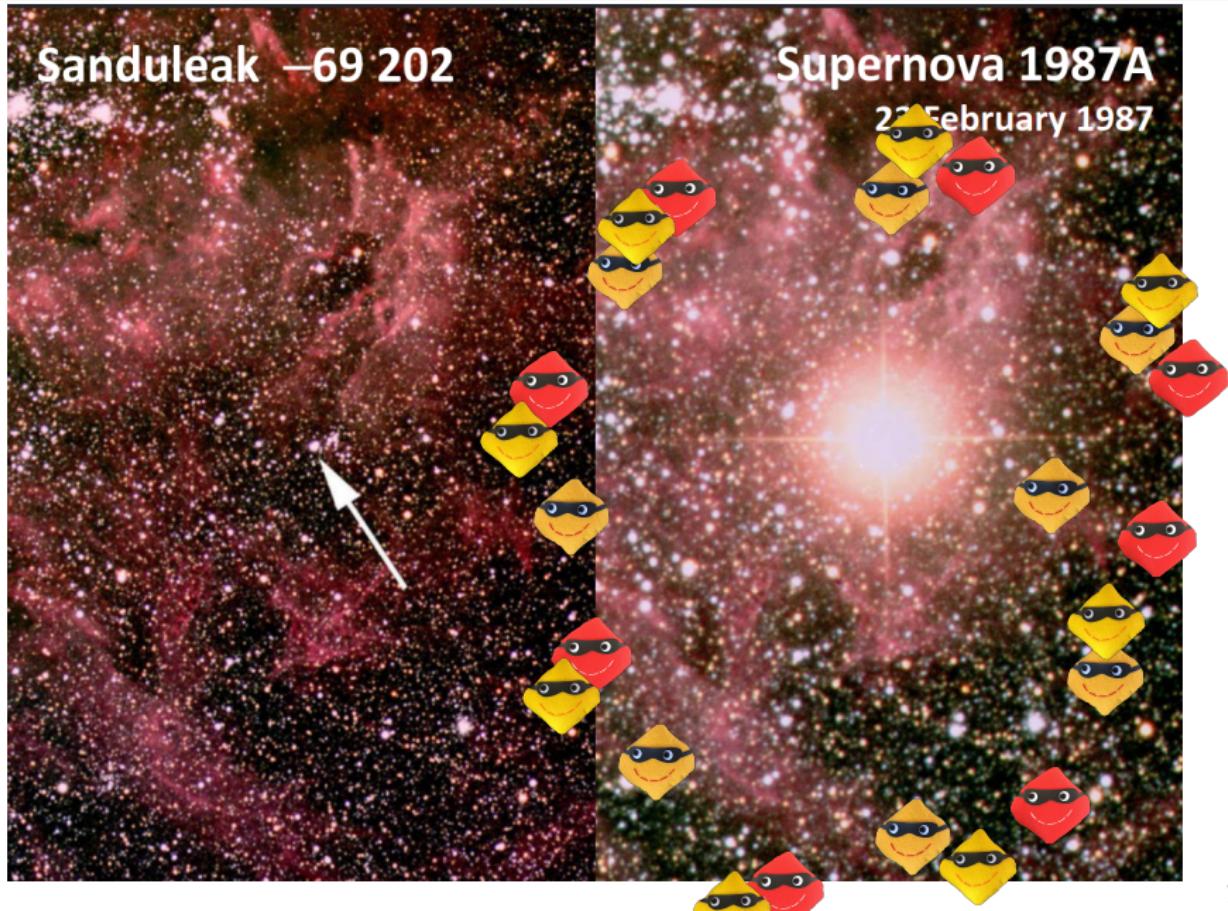
**Neutrino Theory and Experiment Workshop,
Colorado State University**
May 29, 2025



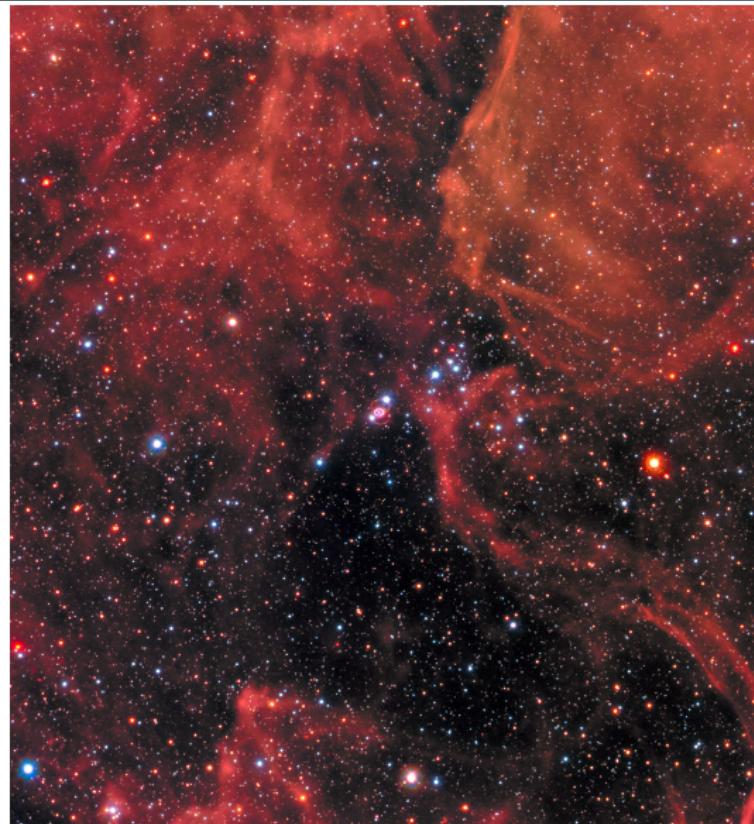
Why is studying astrophysical neutrinos crucial?



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Hubble (2017)



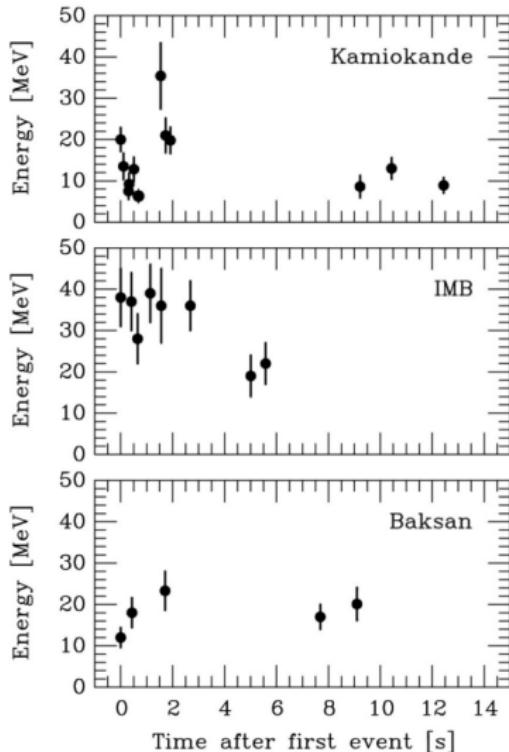
JWST (2023)

- Neutron star remnant
- Binary system

Fransson et al. (2024)

Morris & Podsiadlowski (2007), (2009)

Established track record of neutrino discoveries: SN 1987A



Courtesy of G. Raffelt



- Neutrino detection from SN 1987A:
 - confirmed the core-collapse scenario
 - 99% of the energy emitted in neutrinos
 - best limit at the time on the ν mass

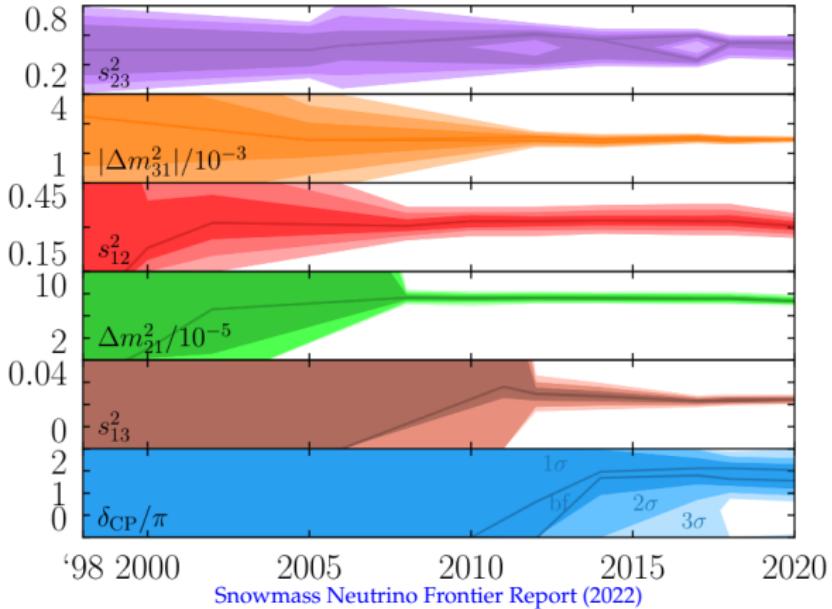
Towards Precise Neutrino Properties Measurements

We known now:

- large mixing angles
- non-zero masses

Remaining questions

- Majorana vs Dirac
- absolute masses
- degree of CP violation



Fermions

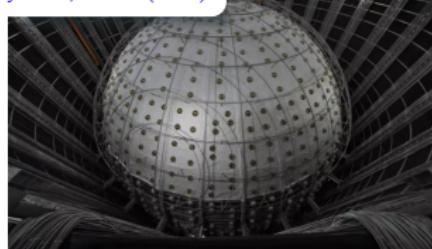
| Leptons | Quarks | | | Force carriers | | |
|---------|----------------------------|------------------------|-----------------------|----------------------|--------------------------|----------------------|
| | u _{up} | c _{charm} | t _{top} | γ _{photon} | H _{Higgs boson} | Z _{Z boson} |
| | d _{down} | s _{strange} | b _{bottom} | g _{gluon} | | |
| | ν _e electron | ν _μ muon | ν _τ tau | | | |
| | e _{electron} | μ _{muon} | τ _{tau} | W _{W boson} | | |

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = U \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}$$

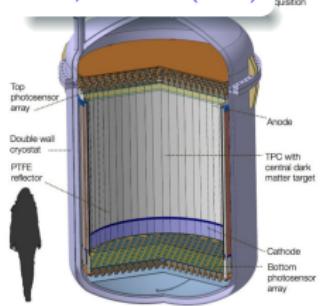
$$U = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

How to achieve full picture of neutrinos? All hands on deck!

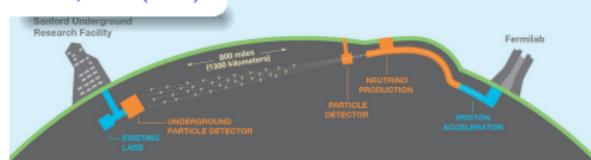
JUNO, China (2025)



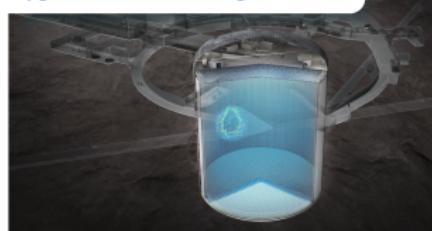
XLZD, DARWIN (20XX)



DUNE, USA (2030)



Hyper-Kamiokande, Japan (2027)



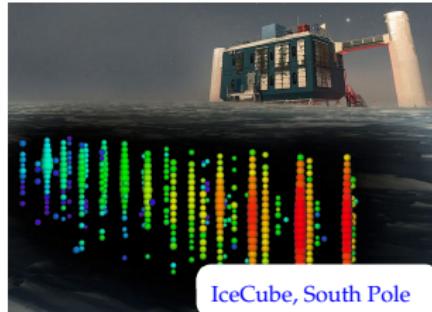
Rubin Observatory, Chile (2025)



- Complementarity with:

- reactor and accelerator searches
- electromagnetic surveys
- other astrophysical messengers

IceCube, South Pole

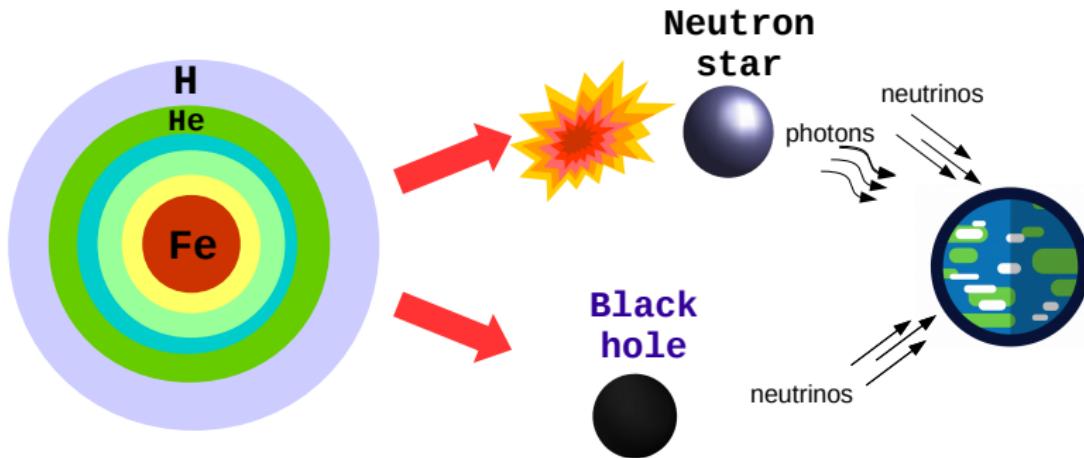


Neutrinos from Core-collapse Supernovae

Why are neutrinos important for a core-collapse supernova?

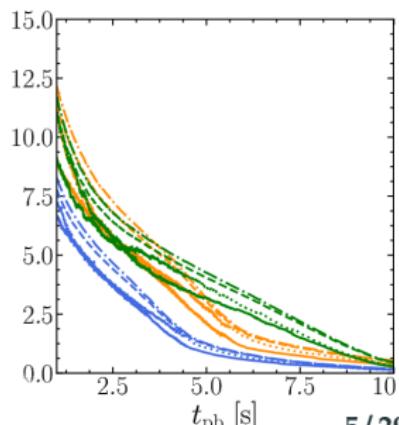
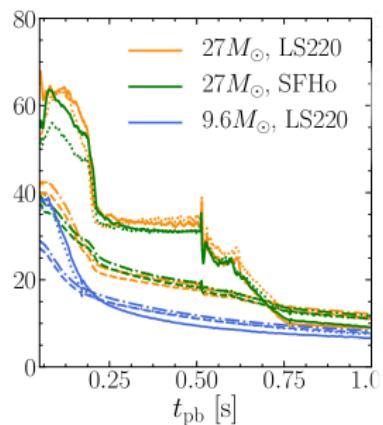
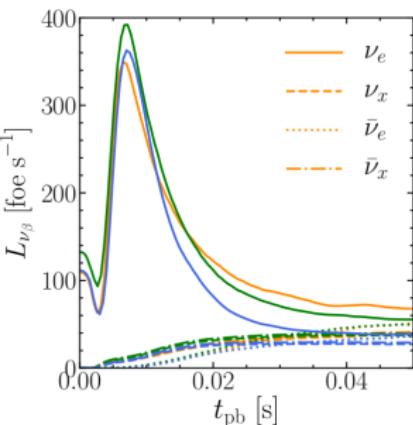
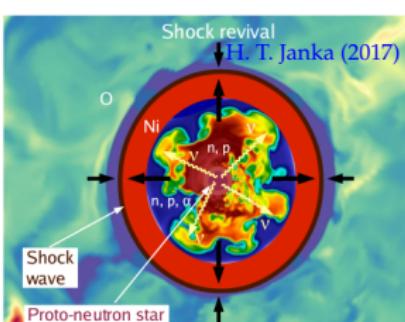
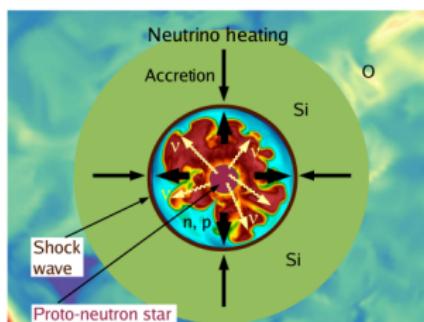
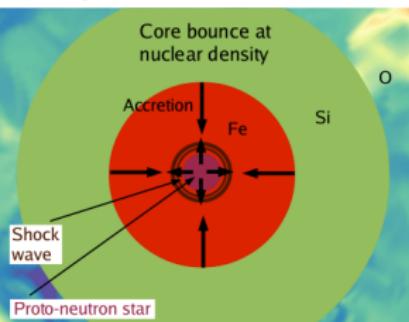
Neutrinos:

- $\sim 10^{58}$ of them emitted from a single core collapse
- only they can reveal the deep interior conditions
- only particles detectable from the collapse to a black hole



Different Phases of Supernova Explosion

- Infall phase,
 ν_e burst ~ 40 ms
- Accretion phase,
 ~ 100 ms
- Cooling phase,
 ~ 10 s



Why core-collapse supernovae are good physics probes?

Advantages

- extreme physical conditions not accessible on Earth
- within the reach of existing and upcoming detectors

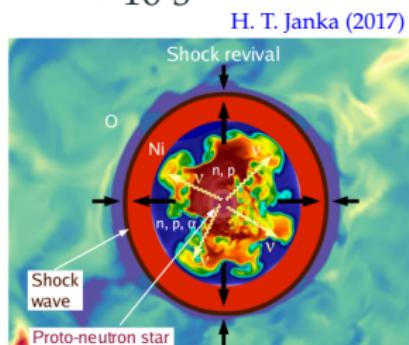
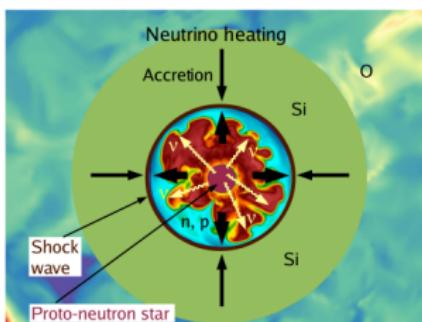
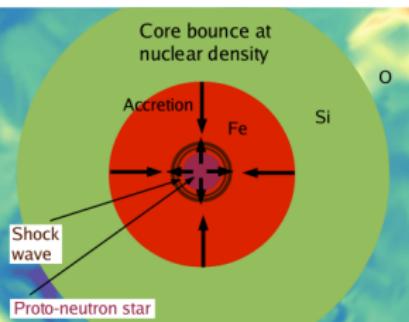
What can we learn with a variety of detectors?

- explosion mechanism Bethe & Wilson (1985),
Fischer et al. (2011)...
- nucleosynthesis Woosley et al. (1994),
Surman & McLaughlin (2003)...
- compact object formation Warren et al. (2019),
Li, Beacom et al. (2020)...
- neutrino mixing Balantekin & Fuller (2013),
Tamborra & Shalgar (2020)...
- non-standard physics McLaughlin et al. (1999),
de Gouv  a et al. (2019) ...

Neutrinos from Supernovae as Probes of New Physics

Different Phases of Supernova Explosion

- Infall phase,
 ν_e burst ~ 40 ms
- Accretion phase,
 ~ 100 ms
- Cooling phase,
 ~ 10 s



New neutrino physics affects the core-collapse supernovae:

- change diffusion time \rightarrow possible change in the star's fate
- changed diffusion time \rightarrow changed duration of the neutrino signal
- new cooling channel \rightarrow affects explosion probability

astrophysical feedback often ignored

How important is astrophysical feedback?

Do non-standard neutrino self-interactions help or inhibit supernova explosion?

In collaboration with G. Fuller, L. Graf, P. Cheong,
J. Froustey, S. Shalgar, K. Kherer, O. Scholer

PRL accepted

Do Neutrinos Have Self-Interactions?

IL NUOVO CIMENTO

VOL. XXXIII, N. 5

1º Settembre 1964

Do Neutrinos Interact between Themselves?

Z. BIALYNICKA-BIRULA

Institute of Physics, Polish Academy of Sciences - Warsaw

(ricevuto il 26 Giugno 1964)



1. – Introduction.

The neutrino is the only elementary particle, which, according to our present knowledge, does not take part in other than weak and gravitational interactions. Its role in nature is not yet fully understood and its interaction properties are only partially known.

The purpose of this note is to answer the following question: Do the present experimental data allow for the existence of interactions between neutrinos much stronger than their weak interactions? The answer to this question is positive. It turns out that such interactions even if they were 10^6 times stronger than weak interactions could not be detected with the present experimental accuracy.

Zofia Bialynicka-Birula (1964)

Lepton number violating neutrino self-interactions

Motivation - to be taken with a grain of salt:

- lepton number conservation - accidental symmetry
- potential cosmological hints

Barenboim et al. (2019), Song, Gonzalez-Garcia, Salvado (2018), ..

- strong impact on core-collapse supernova

Kolb et al. (1982), Fuller et al. (1988), Farzan et al. (2018), AMS, Tamborra (2020), ...

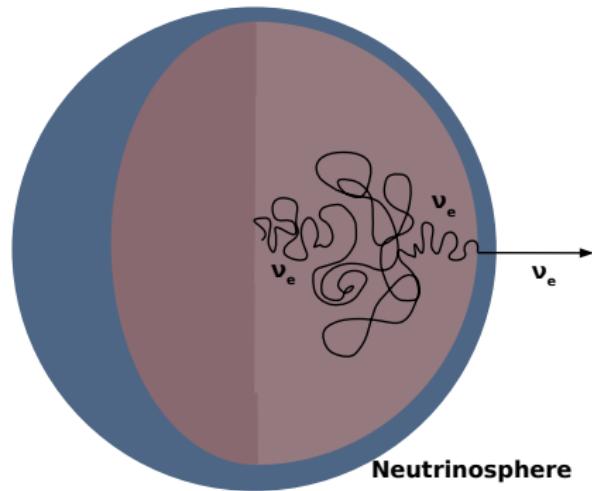
New Interaction Lagrangian

$$\mathcal{L}^\phi = g_{\phi,\alpha\beta} \phi \overline{\nu_{L,\alpha}} \nu_{L,\beta}^c$$

Probability of the New Interaction

$$\sigma_{\nu\text{SI}} \approx \frac{G_{\nu\text{SI}}^2}{8\pi} E_\nu^1 E_\nu^2 (1 - \cos \theta)$$

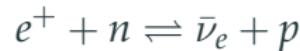
Neutrino Trapping and β -equilibrium



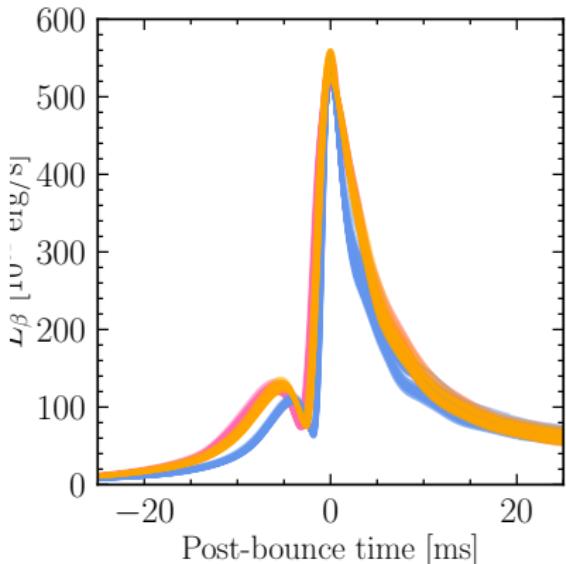
Neutrino trapping



β -equilibrium



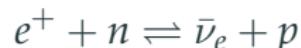
Neutrino Trapping and β -equilibrium



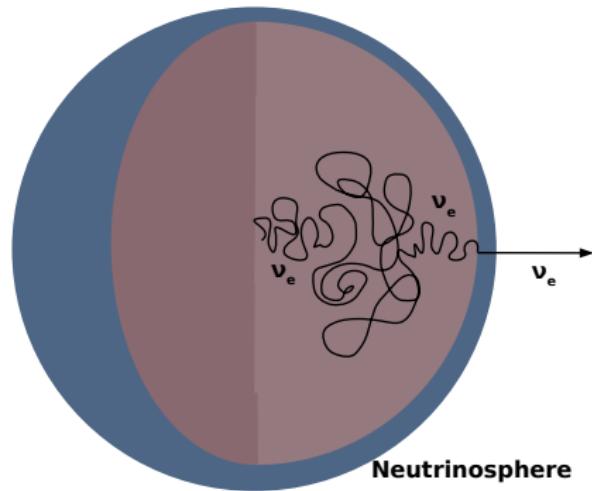
Neutrino trapping



β -equilibrium



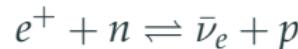
Neutrino Trapping and β -equilibrium



Neutrino trapping



β -equilibrium



LNV ν SI Implementation:

Thermalize the population of ν and $\bar{\nu}$ once $\rho \sim 10^{11} - 10^{12} \text{ g cm}^{-3}$

$$\nu_e \rightleftharpoons \nu_e, \bar{\nu}_e, \nu_\mu, \bar{\nu}_\mu, \nu_\tau, \bar{\nu}_\tau, \quad \nu_e \rightleftharpoons \nu_e, \bar{\nu}_e, \nu_x, \bar{\nu}_x, \quad \nu_e \rightleftharpoons \nu_e, \bar{\nu}_e$$

LNV ν SI in supernovae illustrated by cats

Without LNV ν SI



Pics credit: Tony Zhouw who's courageously cat sitting Aurora and Beetle

LN ν SI in supernovae illustrated by cats

With LN ν SI



Pics credit: Tony Zhouw who's courageously cat sitting Aurora and Beetle

Static, Homogenous and Isotropic Boltzmann Equation

Boltzmann Equation

$$\frac{df_\nu}{dt} = (1 - f_\nu) j_\nu - f_\nu \chi_\nu ,$$

Electron fraction evolution - weak rates



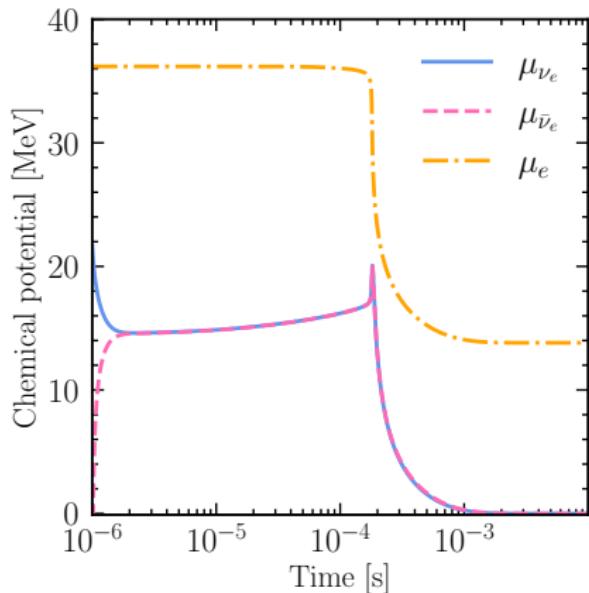
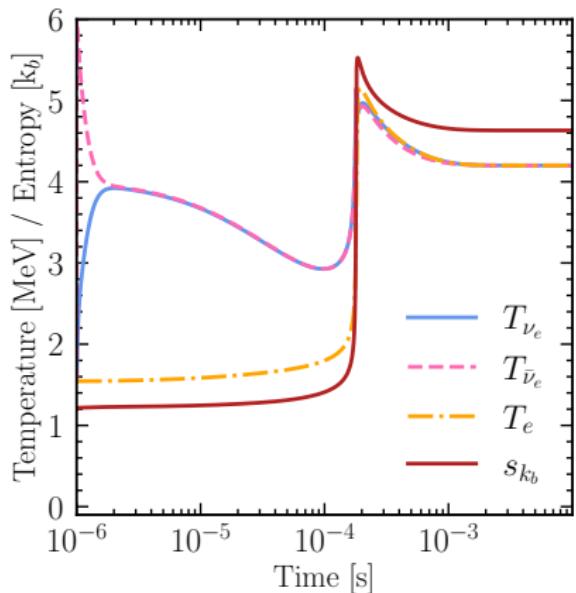
$$\frac{dY_e}{dt} = R_{\nu_e} - R_{\bar{\nu}_e} - R_{e^-} + R_{e^+} , \quad e^+ + n \rightleftharpoons \bar{\nu}_e + p$$

Temperature and chemical potential evolution for leptons

$$\frac{dT_i}{dt} = \left(\frac{\partial \rho_i}{\partial \mu_i} \frac{dn_i}{dt} - \frac{\partial n_i}{\partial \mu_i} \frac{d\rho_i}{dt} \right) / \left(\frac{\partial n_i}{\partial T_i} \frac{\partial \rho_i}{\partial \mu_i} - \frac{\partial n_i}{\partial \mu_i} \frac{\partial \rho_i}{\partial T_i} \right) ,$$

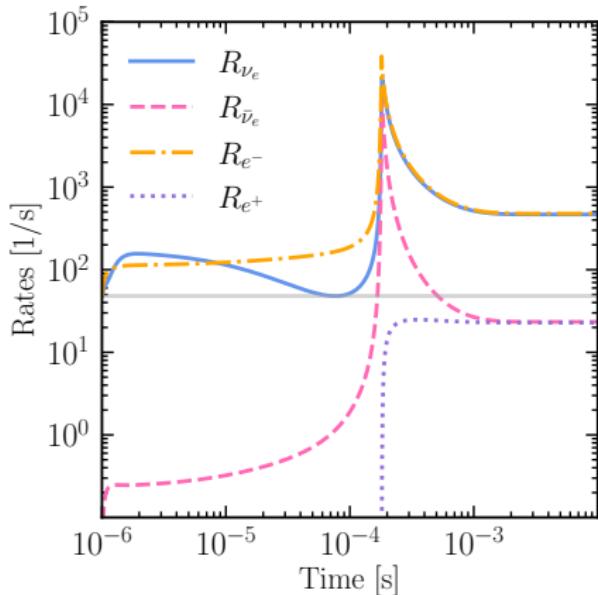
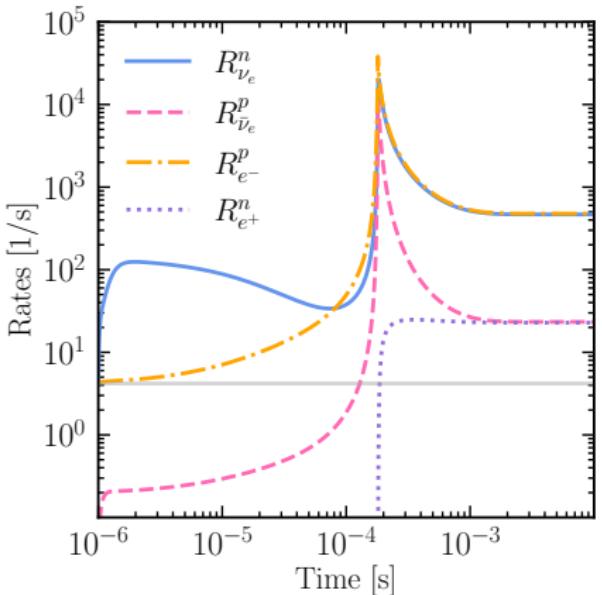
$$\frac{d\mu_i}{dt} = \left(\frac{\partial \rho_i}{\partial T_i} \frac{dn_i}{dt} - \frac{\partial n_i}{\partial T_i} \frac{d\rho_i}{dt} \right) / \left(\frac{\partial n_i}{\partial \mu_i} \frac{\partial \rho_i}{\partial T_i} - \frac{\partial n_i}{\partial T_i} \frac{\partial \rho_i}{\partial \mu_i} \right) .$$

Evolution of Thermodynamical Quantities



- new interactions quickly equilibrate ν_e and $\bar{\nu}_e$ seas
- enhanced ν_e and e^- captures heat up the matter
- similar results for all flavors equilibration

Weak reaction rates



- initial increase in $\nu_e + n$, $\nu_e + A$ and $e^- + A$
- enhanced ν_e and e^- captures heat up the matter
- similar results for all flavors equilibration

Fast LNV ν SI - Approximate Evolution

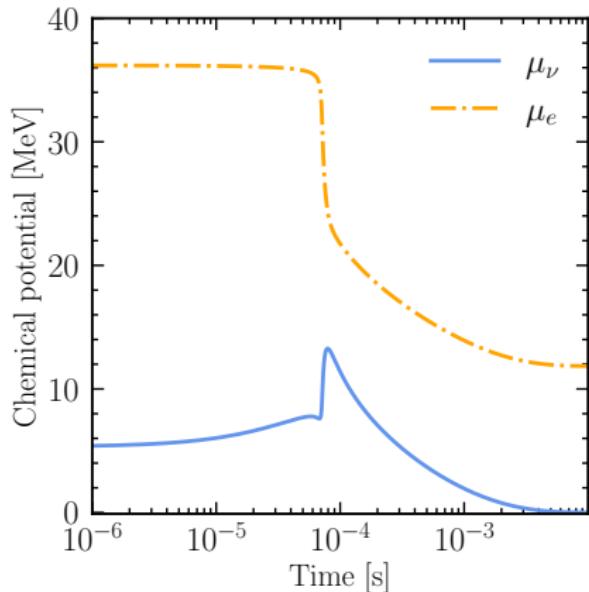
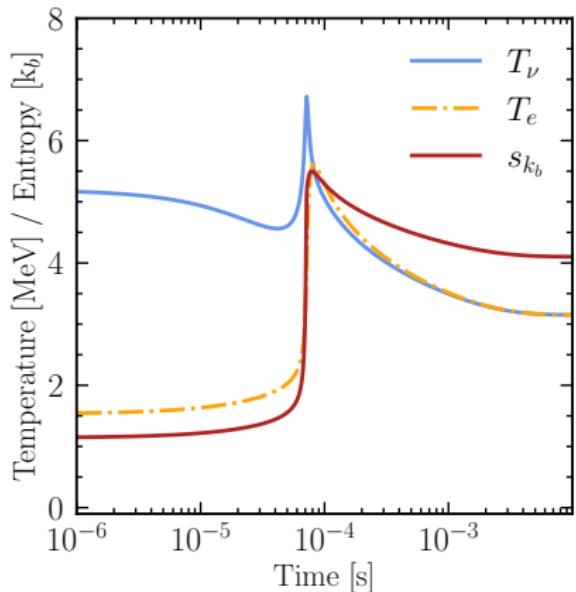
LNV ν SI timescale much faster than weak timescale →
a single ν species evolution

$$\sum_{\alpha} \left(\frac{dn_{\nu\alpha}}{dt} + \frac{dn_{\bar{\nu}\alpha}}{dt} \right) = \frac{\delta n_{\nu}}{\delta t} \quad \text{sum over charged-current}$$
$$\sum_{\alpha} \left(\frac{d\rho_{\nu\alpha}}{dt} + \frac{d\rho_{\bar{\nu}\alpha}}{dt} \right) = \frac{\delta \rho_{\nu}}{\delta t} \quad \text{weak interactions}$$

$$\frac{dT_{\nu}}{dt} = \frac{\frac{\partial \rho_{\nu}}{\partial \mu_{\nu}} \frac{\delta n_{\nu}}{\delta t} - \frac{\partial n_{\nu}}{\partial \mu_{\nu}} \frac{\delta \rho_{\nu}}{\delta t}}{2N_F \left(\frac{\partial n_{\nu}}{\partial T_{\nu}} \frac{\partial \rho_{\nu}}{\partial \mu_{\nu}} - \frac{\partial n_{\nu}}{\partial \mu_{\nu}} \frac{\partial \rho_{\nu}}{\partial T_{\nu}} \right)}$$

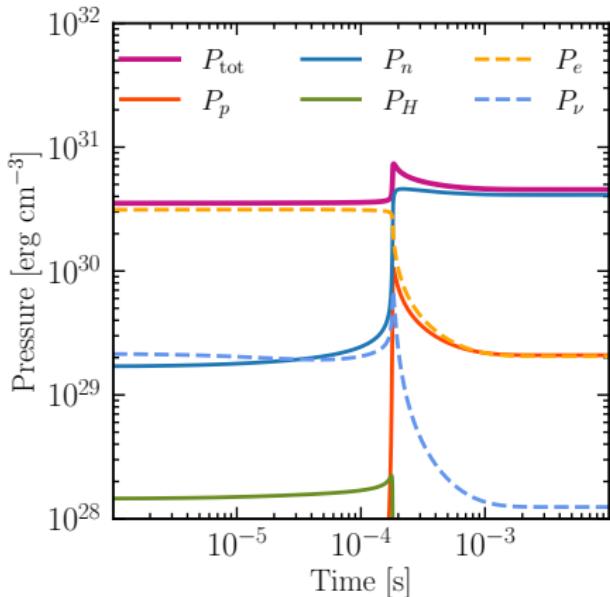
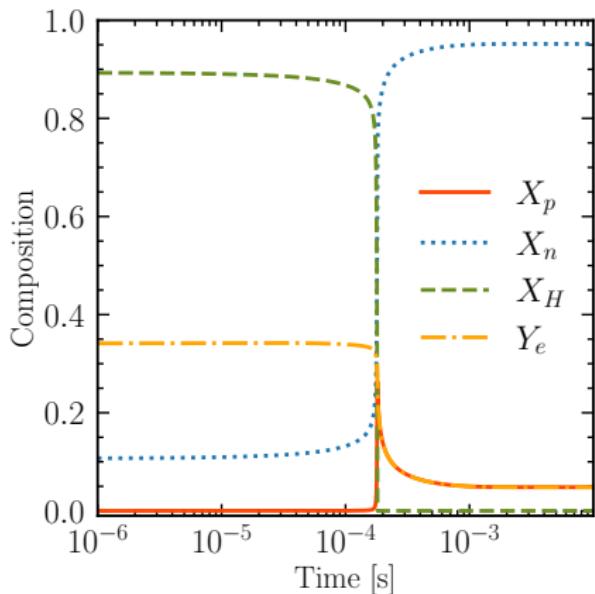
$$\frac{d\mu_{\nu}}{dt} = \frac{\frac{\partial \rho_{\nu}}{\partial T_{\nu}} \frac{\delta n_{\nu}}{\delta t} - \frac{\partial n_{\nu}}{\partial T_{\nu}} \frac{\delta \rho_{\nu}}{\delta t}}{2N_F \left(\frac{\partial n_{\nu}}{\partial \mu_{\nu}} \frac{\partial \rho_{\nu}}{\partial T_{\nu}} - \frac{\partial n_{\nu}}{\partial T_{\nu}} \frac{\partial \rho_{\nu}}{\partial \mu_{\nu}} \right)}$$

Evolution of Thermodynamical Quantities



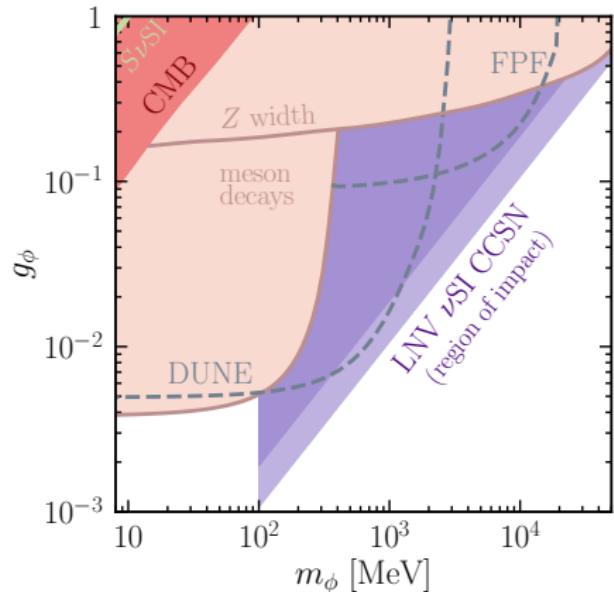
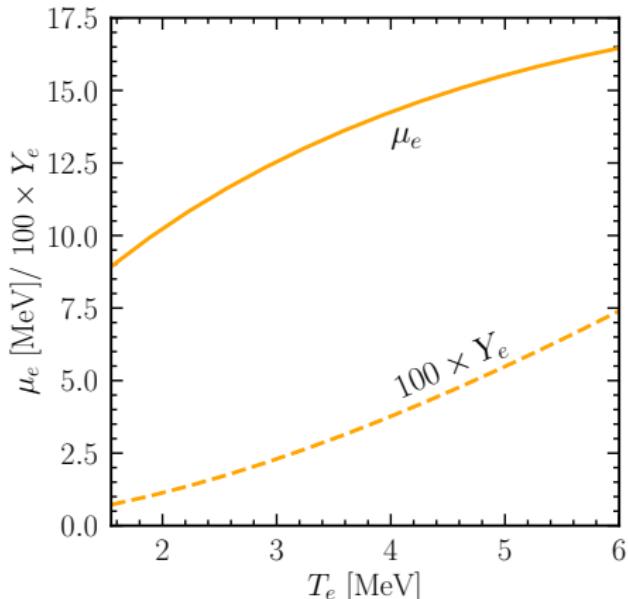
- the same qualitative results for all six flavor equilibration

Composition and Pressure Support of the Core



- s_{k_b} - entropy generation shifts composition towards no heavy nuclei
$$X_H \propto s_{k_B}^{1-\langle A \rangle} n_p^Z n_n^N \exp(E_b/T_e)$$
- enhanced deleptonization changes the pressure support of the core

New β -equilibrium with LNV ν SI



- regardless of the final T_e the new equilibrium has a very low Y_e

$$\mu_e = \delta m_{np} - T_e \ln \left(\frac{Y_e}{1-Y_e} \right)$$
, with $Y_e = \frac{1}{\pi^2 \rho} \int_0^\infty dp_e p_e^2 f_e(E_e, T_e, \mu_e)$
- complementarity with future accelerator-based experiments

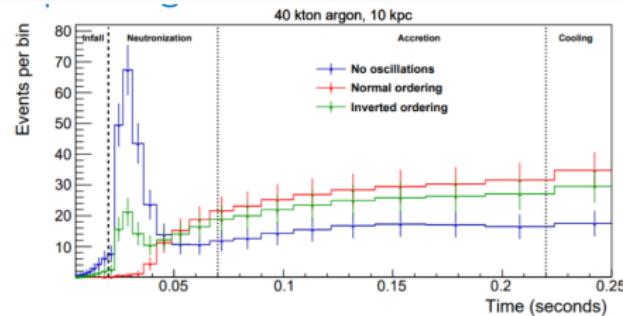
Why focus only on a single rare event?

Images: Kurzgesagt



Single galactic SN event

- rare event
- precise information about one star



D. Pershey slides

Ready for rare supernovae too?

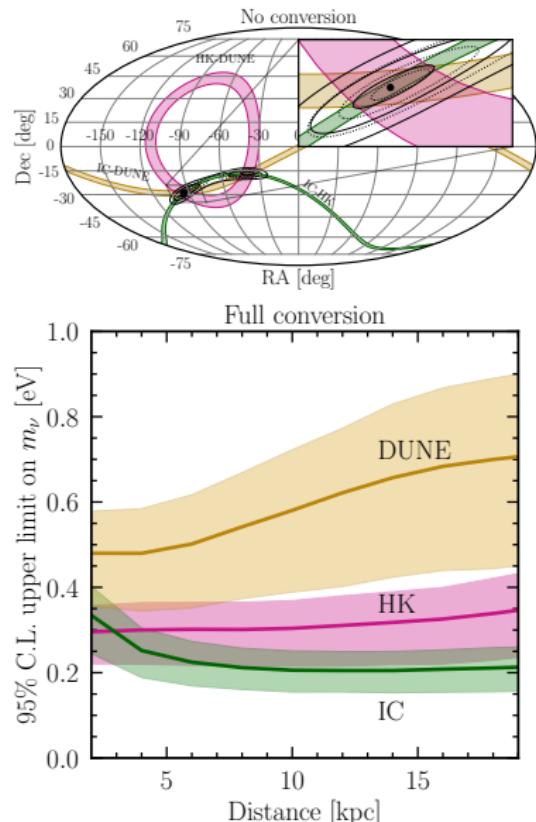
Why focus only on a single rare event?

Images: Kurzgesagt



Single galactic SN event

- rare event
- precise information about one star



Pitik, Heimsoth, AMS, Balantekin (2022)

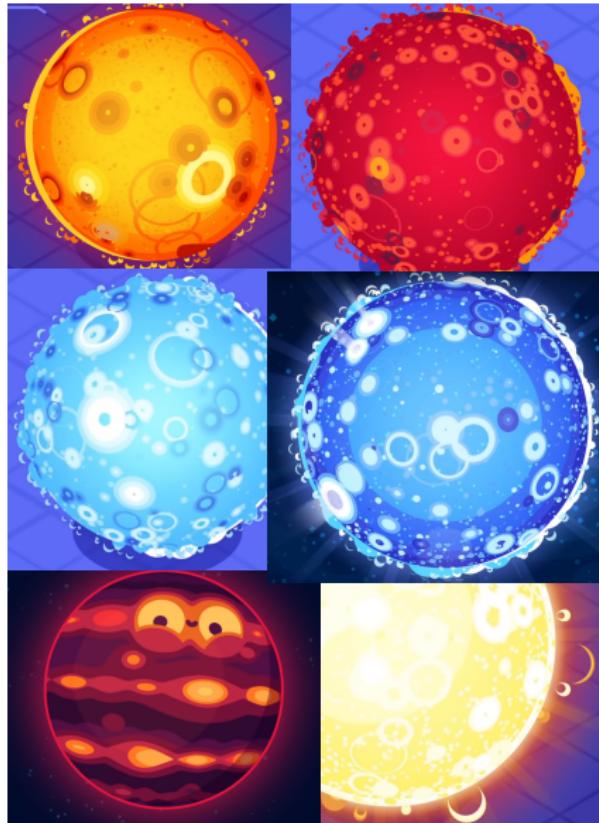
Why focus only on a single rare event?

Images: Kurzgesagt



Single galactic SN event

- rare event
- precise information about one star



Multiple SN events (larger distances)

- accumulation of events
- will detect in coming years

Diffuse supernova neutrino background

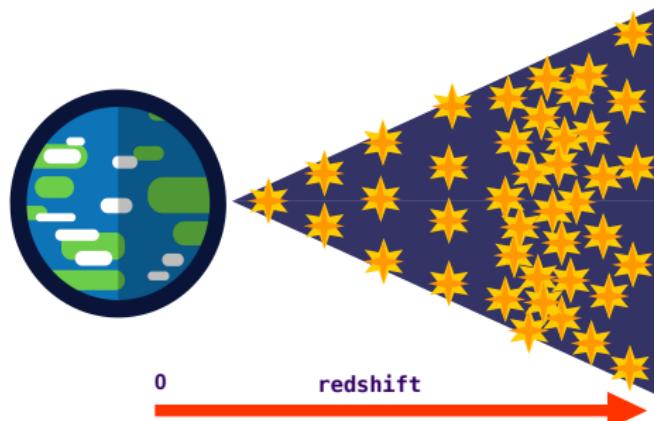
$$\Phi_{\nu_\beta}(E) = \frac{c}{H_0} \int dM \int dz \frac{R_{\text{SN}}(z, M)}{\sqrt{\Omega_M(1+z)^3 + \Omega_\Lambda}} [f_{\text{CC-SN}} F_{\nu_\beta, \text{CC-SN}}(E', M) + f_{\text{BH-SN}} F_{\nu_\beta, \text{BH-SN}}(E', M)]$$

Diagram illustrating the components of the diffuse supernova neutrino background flux:

- cosmological supernovae rate**: Represented by a pink arrow pointing to the term $\frac{R_{\text{SN}}(z, M)}{\sqrt{\Omega_M(1+z)^3 + \Omega_\Lambda}}$.
- fraction of neutron-star-forming progenitors**: Represented by a red arrow pointing to the term $f_{\text{CC-SN}}$.
- neutrino flux from a single star**: Represented by a magenta arrow pointing to the term $F_{\nu_\beta, \text{CC-SN}}(E', M)$.
- fraction of black-hole-forming progenitors**: Represented by a blue arrow pointing to the term $f_{\text{BH-SN}}$.

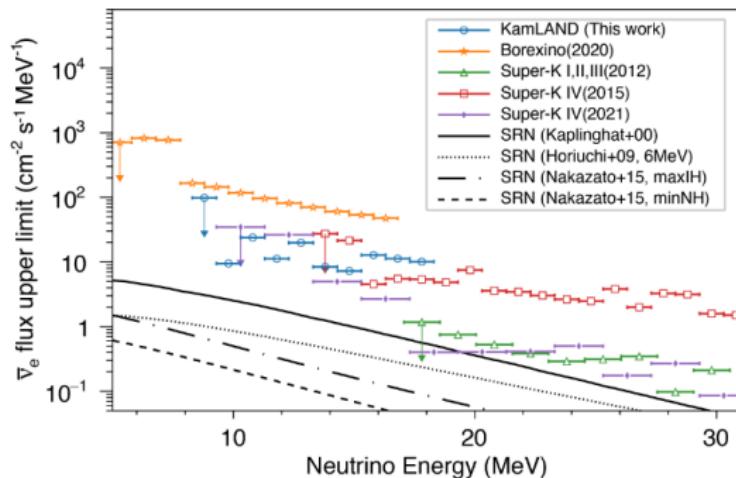
The DSNB is sensitive to:

- $R_{\text{SN}}, f_{\text{BH-SN}}$
- neutrino flavor evolution
- equation of state
- mass accretion rate in BH-SN
- non-standard physics



Diffuse supernova neutrino background: current limits

SK collab. (2021)

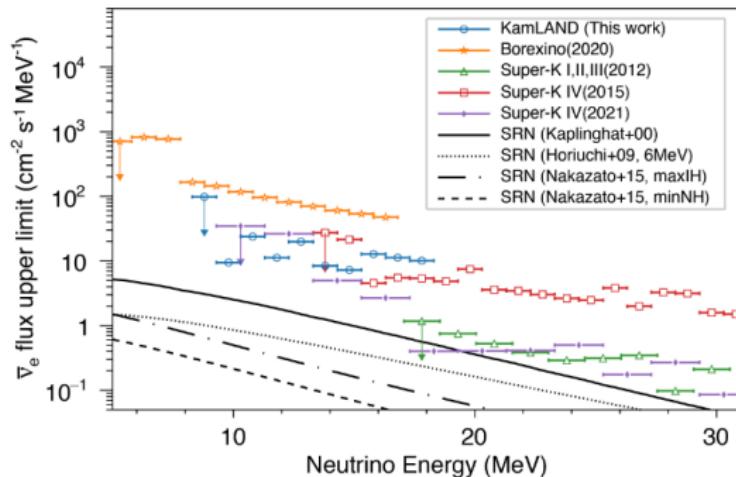


DSNB limits:

- $\bar{\nu}_e \approx 2.7 \text{ cm}^{-2} \text{s}^{-1}$ for $E_\nu > 17.3 \text{ MeV}$ SK collab. (2021), SK collab. (2023)
soon detected by SK (Gd) Beacom, Vagins (2004) and JUNO JUNO collab. (2021)
- $\nu_e \approx 19 \text{ cm}^{-2} \text{s}^{-1}$ for $E_\nu \in [22.9, 36.9 \text{ MeV}]$ SNO collab. (2020)
possibly detectable by DUNE Møller, AMS, et al. (2018), Zhu et al. (2019)
- $\nu_x \approx 750 \text{ cm}^{-2} \text{s}^{-1}$ for $E_\nu > 19.3 \text{ MeV}$ Lunardini, Peres (2008)

Diffuse supernova neutrino background: current limits

SK collab. (2021)



DSNB limits:

- $\bar{\nu}_e \approx 2.7 \text{ cm}^{-2} \text{ s}^{-1}$ for $E_\nu > 17.3 \text{ MeV}$ SK collab. (2021), SK collab. (2023)
soon detected by SK (Gd) Beacom, Vagins (2004) and JUNO JUNO collab. (2021)
- $\nu_e \approx 19 \text{ cm}^{-2} \text{ s}^{-1}$ for $E_\nu \in [22.9, 36.9 \text{ MeV}]$ SNO collab. (2020)
possibly detectable by DUNE Møller, AMS, et al. (2018), Zhu et al. (2019)
- $\nu_x \lesssim 100 \text{ cm}^{-2} \text{ s}^{-1}$ for $E_\nu > 19.3 \text{ MeV}$ AMS, Beacom, Tamborra (2021)

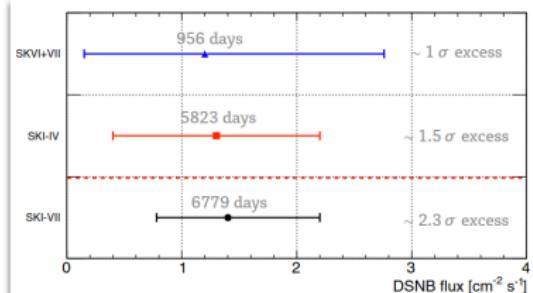
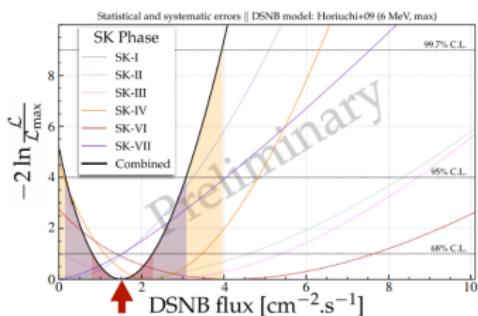
Tension from zero assumption

Spectral-fitting analysis



Spectrum fitting analysis to extract significance

- Total 6779 days of SK (5823 d pure-water and 956 d Gd-water) combined
- Analysis threshold: $E_\nu > 17.3$ MeV
- Suppress uncertainty of background prediction by fitting both $N_n=1$, $N_n \neq 1$



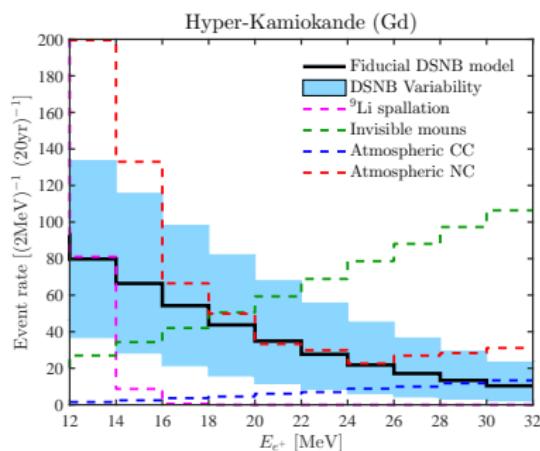
Highlight:

- Sensitivity of SK-Gd ~1000 days exposure is already comparable level it with ~6000 days of pure-water SK
 - Best fit of whole SK observation is $1.4^{+0.8}_{-0.6} \text{ cm}^{-2} \text{s}^{-1}$ for $E_\nu > 17.3$ MeV
- exhibit $\sim 2.3 \sigma$ excess!!

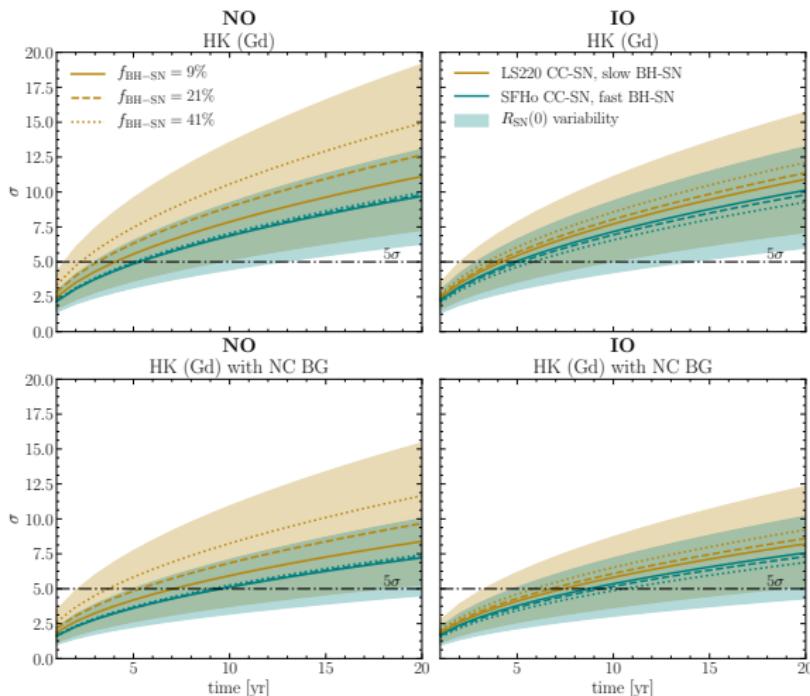
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Slide credit: Masayuki Harada talk at Neutrino 2024

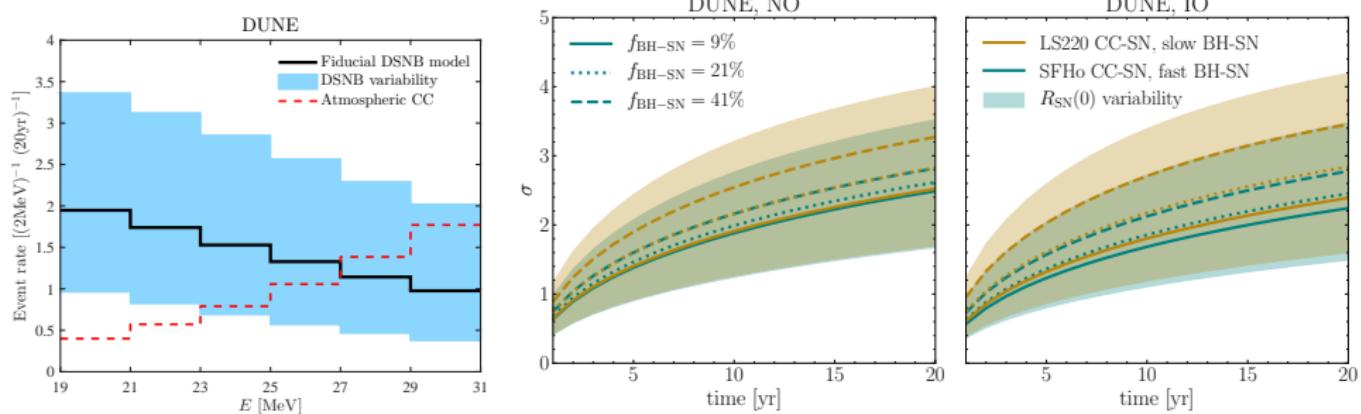
HK DSNB detection perspectives



- fiducial volume: 374 kton
- efficiency: 67%

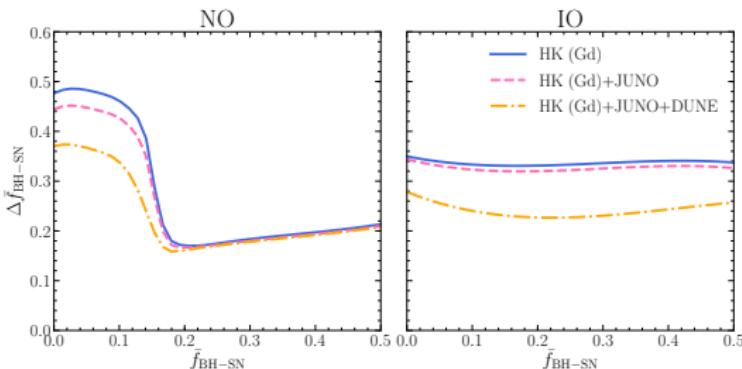


DUNE DSNB detection perspectives



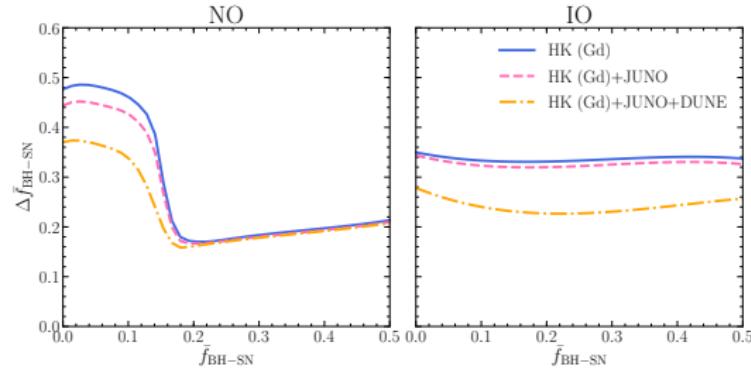
- fiducial volume: 40 kton
- efficiency: 86%

Expected 1σ uncertainty: fraction of BH forming progenitors



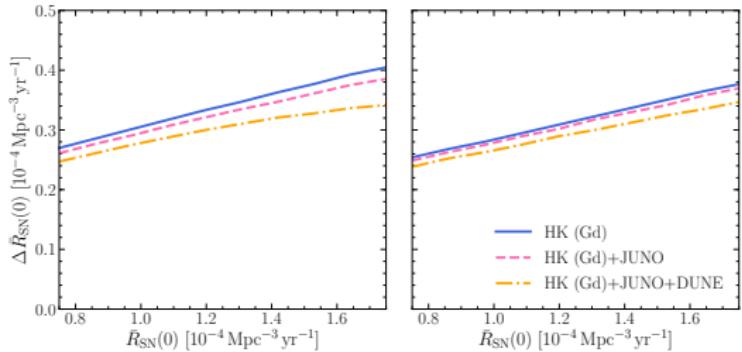
- The high uncertainty comes from $f_{\text{BH-SN}}$ –mass accretion rate degeneracy
- DUNE is sensitive to neutrinos → helps to reduce the uncertainty

Expected 1σ uncertainty: local supernova rate



- The high uncertainty comes from $f_{\text{BH-SN}}$ -mass accretion rate degeneracy
- DUNE is sensitive to neutrinos → helps to reduce the uncertainty

- Relative error of 20%-33% independent of the mass ordering.



Astrophysical uncertainties affecting the DSNB

- Neutrino Flux from an "Average Supernova"
Lunardini (2009), Horiuchi et al. (2018), Kresse et al. (2018), ...
- Cosmological Supernovae Rate
Beacom (2010), Horiuchi et al. (2011), Ando et al. (2023), Ekanger et al. (2024)...
- Initial Mass Function
Ziegler, Edwards, **AMS**, Tamborra, Horiuchi, Ando, Freese (2022)
- Fraction of Black-Hole-Forming Progenitors
Lunardini (2009), Lien et al. (2010), Keehn & Lunardini (2012), Priya & Lunardini (2017),
Møller, **AMS**, Tamborra, Denton (2018), Horiuchi et al. (2018), Kresse et al. (2018), ...
- Binary Interactions
Horiuchi, Kinugawa, Takiwaki, Takahashi (2021)
Sanduleak and Betelgeuse in binary systems? Morris & Podsiadlowski (2007), (2009), Goldberg et al (2024), MacLeod et al (2024)
- **Low-Energy Atmospheric Neutrinos**
Detection of sub-100 MeV in JUNO **AMS** & Beacom (2023)

Non exhaustive list of references

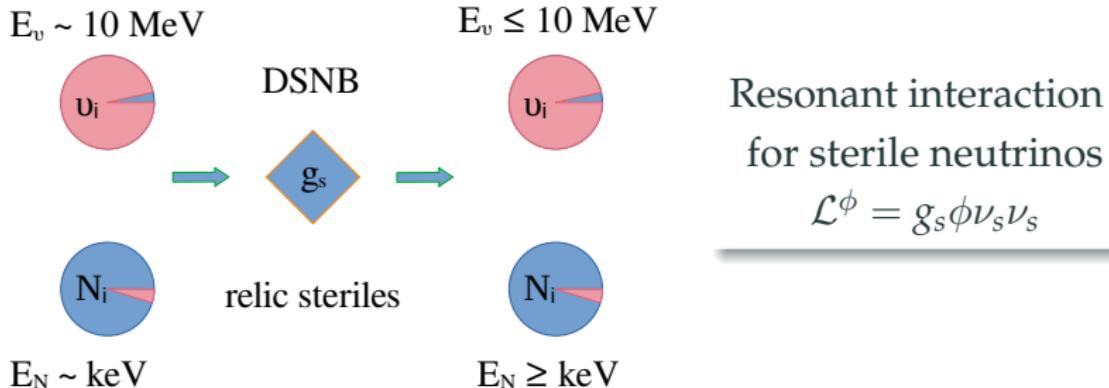
**How to probe new physics with
these uncertainties?**

Probing self-interacting sterile neutrino dark matter with the DSNB

In collaboration with B. Balantekin, G. Fuller, and A. Ray

Phys.Rev.D in 108 (2023) 12, 123011

Do KeV-mass Sterile Neutrinos Have Self-Interactions?



$$\sigma(E_\nu) = \frac{g_s^4}{4\pi} \frac{s}{(s - m_\phi^2)^2 + m_\phi^4 \Gamma_\phi^2} \approx \frac{\pi g_s^2}{m_\phi^2} E_\nu \delta(E_R - E_\nu), \text{ where } E_R = m_\phi^2 / 2m_s$$

- sterile component in the DSNB ν_i interacts with the mostly sterile relic background of N_i

bigger parameter space for keV serile neutrino dark matter with self-interactions:

Maria D. Astros and S. Vogl (2023), T. Bringmann et al. (2022)

Modeling secret neutrino interactions in DSNB

Modified DSNB flux

$$\phi_\alpha(E_\nu) \simeq \sum_{i=1}^3 |U_{\alpha i}|^2 \int_0^{z_{\max}} dz \frac{P_i(E_\nu, z)}{H(z)} \times R_{\text{SN}}(z) F_{\text{SN}}^i(E_\nu(1+z))$$

Probability of interaction

$$P_i(E_\nu, z) = e^{-\tau_i(E_\nu, z)}$$

$$\tau_i(E_\nu, z) \simeq \tau_R \Theta(z - z_R) = \frac{\Gamma_R(z_R)}{(1 + z_R) H(z_R)} \Theta(z - z_R)$$

where $z_R = E_R/E_\nu - 1$,

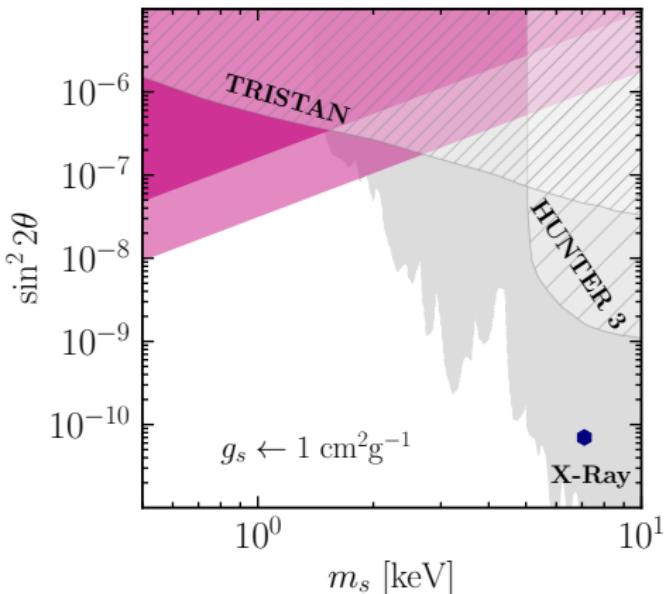
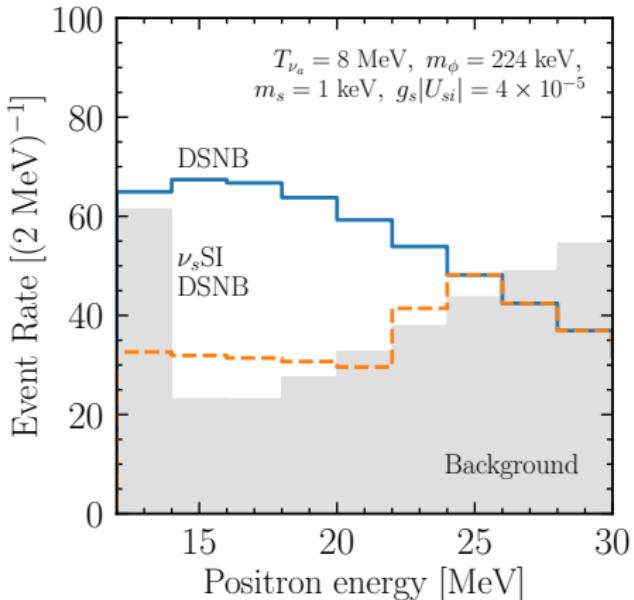
interaction rate $\Gamma_R(z_R) \simeq |U_{si}|^2 n_{\nu_s}(z_R) \sigma_R$,

and sterile neutrino number density $n_{\nu_s}(z_R) = n_{\nu_s}(1 + z_R)^3$

similar studies for active neutrino self-interactions and eV-mass sterile neutrinos:

Goldberg et al. (2005), Baker et al. (2007), Farzan, Palomares-Ruiz (2014), Reno et al. (2018), Creque-Sarbinowski et al. (2021) 27 / 29

Secret neutrino interactions: DSNB



- Sterile neutrino self-interactions may result in features in DSNB
- Overlap with the TRISTAN experiment parameter space (DUNE?)
- Reduction of the astrophysical uncertainties helps but not by a lot

Conclusions

Core-collapse supernovae

- serve as testing grounds in constraining standard and new physics
- reliable limits, only when the sources are accurately modeled

Detection of astrophysical neutrino fluxes

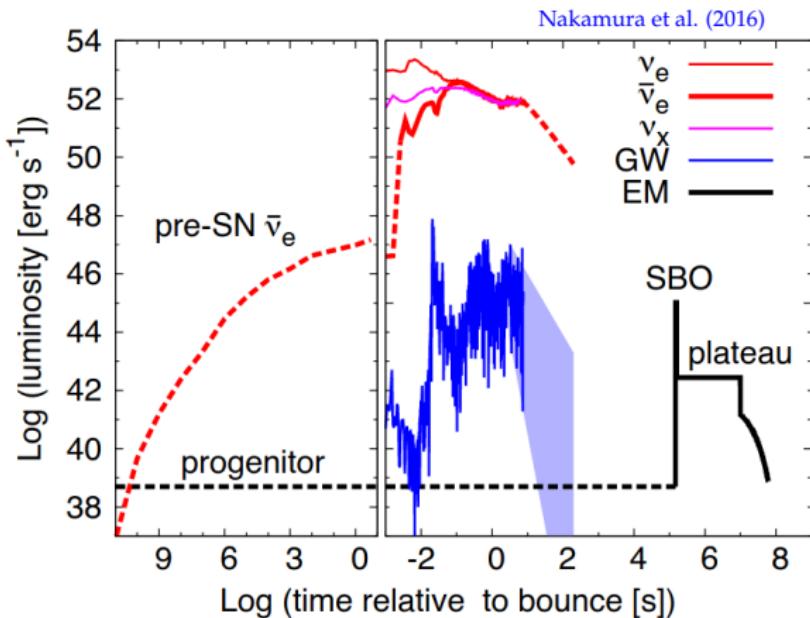
- brings us closer to fully understanding the physics inside the sources
- help us to probe potential new physics scenarios

Exciting times ahead!

Thank you for the attention!

Backup

Core-Collapse Supernova Light Curve



Partial Derivatives for the Fermi-Dirac distributions

The partial derivatives for the Fermi-Dirac distributions are given by [Escudero \(2020\)](#)

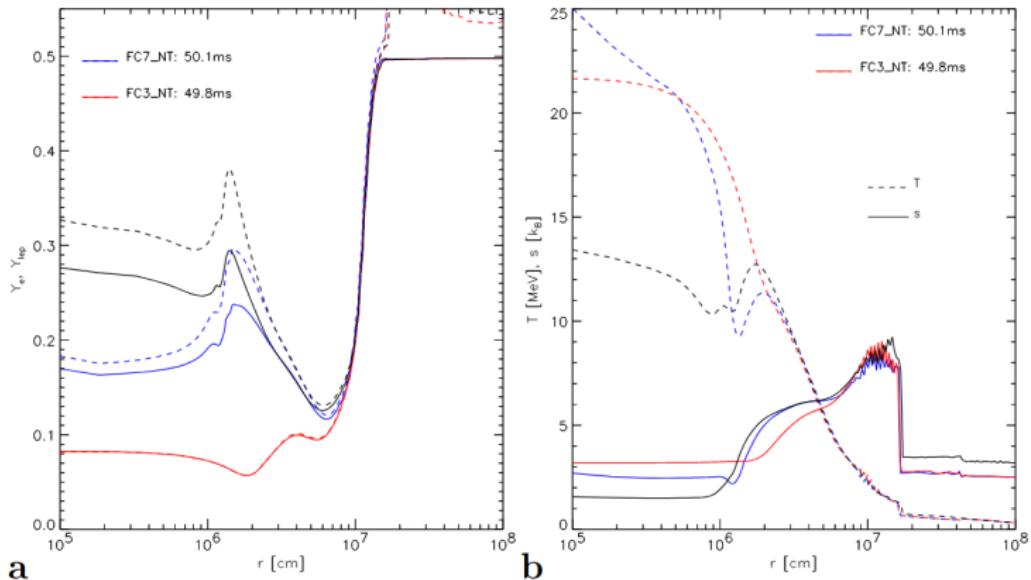
$$\frac{\partial n}{\partial T} = \frac{g}{2\pi^2} \int_m^\infty dE E \sqrt{E^2 - m^2} \frac{(E - \mu)}{4T^2} \cosh^{-2} \left(\frac{E - \mu}{2T} \right), \quad (1a)$$

$$\frac{\partial \rho}{\partial T} = \frac{g}{2\pi^2} \int_m^\infty dE E^2 \sqrt{E^2 - m^2} \frac{(E - \mu)}{4T^2} \cosh^{-2} \left(\frac{E - \mu}{2T} \right), \quad (1b)$$

$$\frac{\partial n}{\partial \mu} = \frac{g}{2\pi^2} \int_m^\infty dE E \sqrt{E^2 - m^2} \left[2T \cosh \left(\frac{E - \mu}{T} \right) + 2T \right]^{-1}, \quad (1c)$$

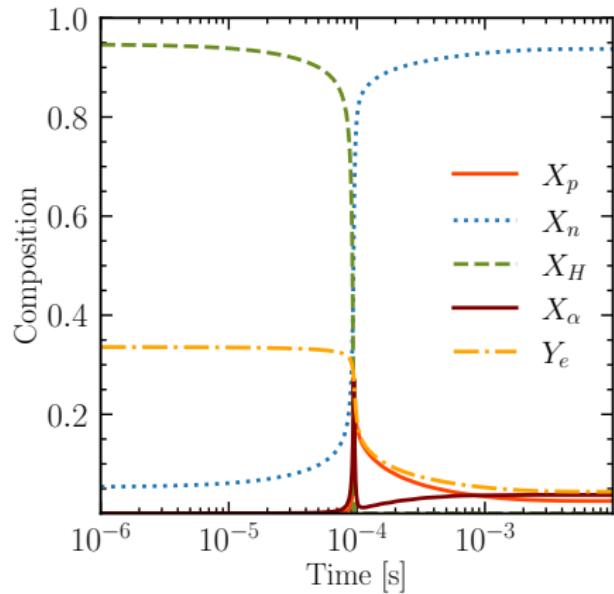
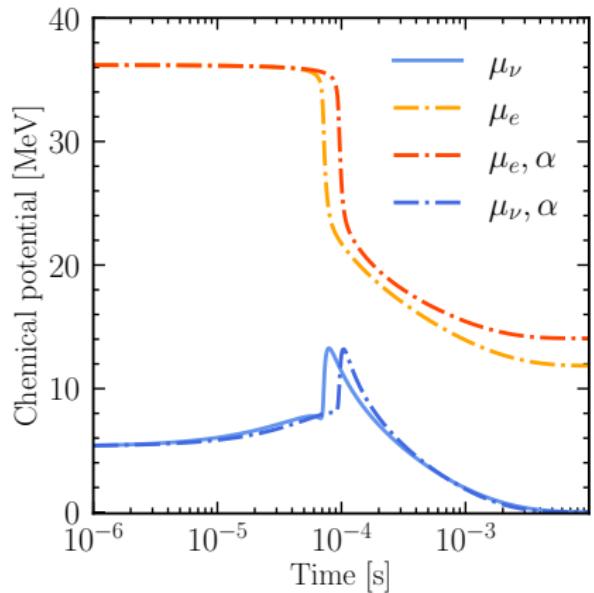
$$\frac{\partial \rho}{\partial \mu} = \frac{g}{2\pi^2} \int_m^\infty dE E^2 \sqrt{E^2 - m^2} \left[2T \cosh \left(\frac{E - \mu}{T} \right) + 2T \right]^{-1} \quad (1d)$$

Proxy "Internal Deleptonization"



M. Rampp et al. (2002)

LS 220 Equation of State: impact of α particles



Combined likelihood analyses

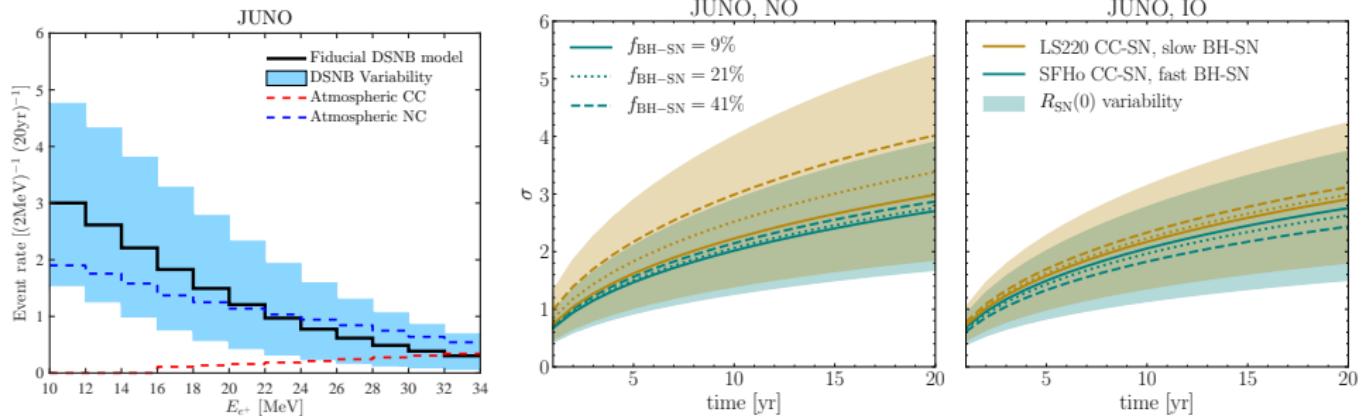
Significance test

$$\chi^2 = \min_A \left(\sum_j \chi_{A,j}^2 + \chi_{\text{HK}}^2 + \chi_{\text{JUNO}}^2 + \chi_{\text{DUNE}}^2 \right)$$

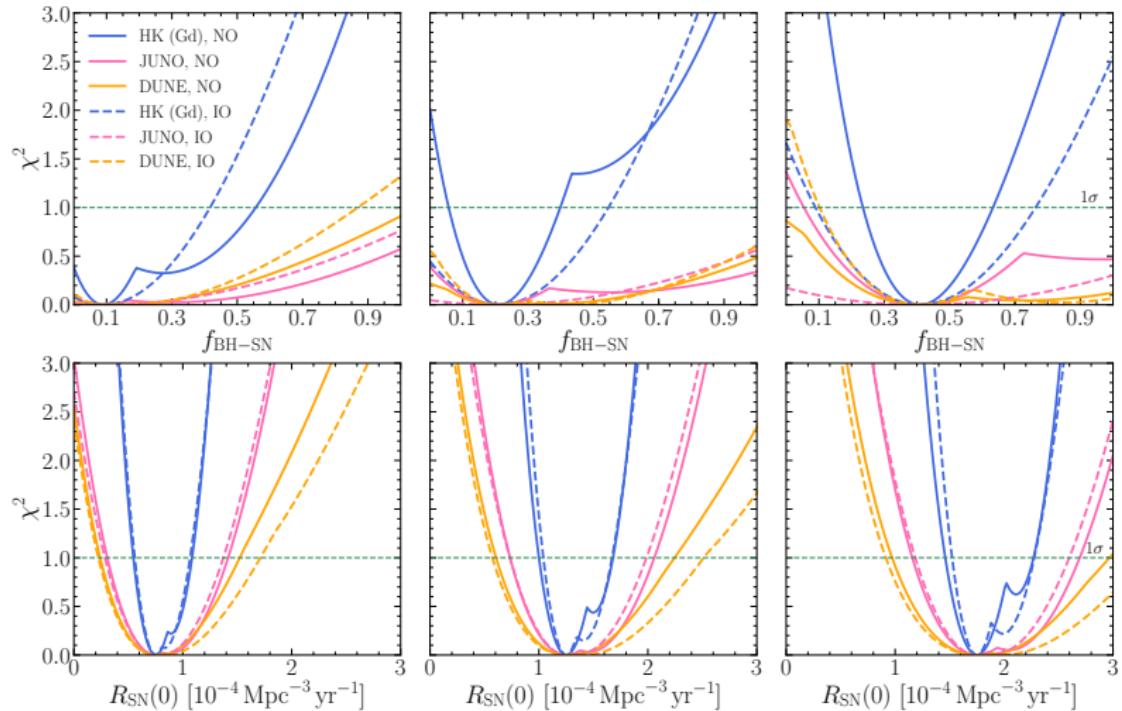
The set of parameters to be marginalized over:

- $f_{\text{BH-SN}}, \Delta f_{\text{BH-SN}} = 0.2$
- $R_{\text{SN}}(0), \Delta R_{\text{SN}(0)} = 0.25 \times 10^{-4} \text{ Mpc}^{-3} \text{ yr}^{-1}$
- background normalization uncertainty, $\Delta_{\text{BG}} = 20\%$
- liquid argon cross section uncertainty, $\Delta \sigma_{\text{LAr}} = 15\%$
- mass accretion rate - equation of state uncertainty

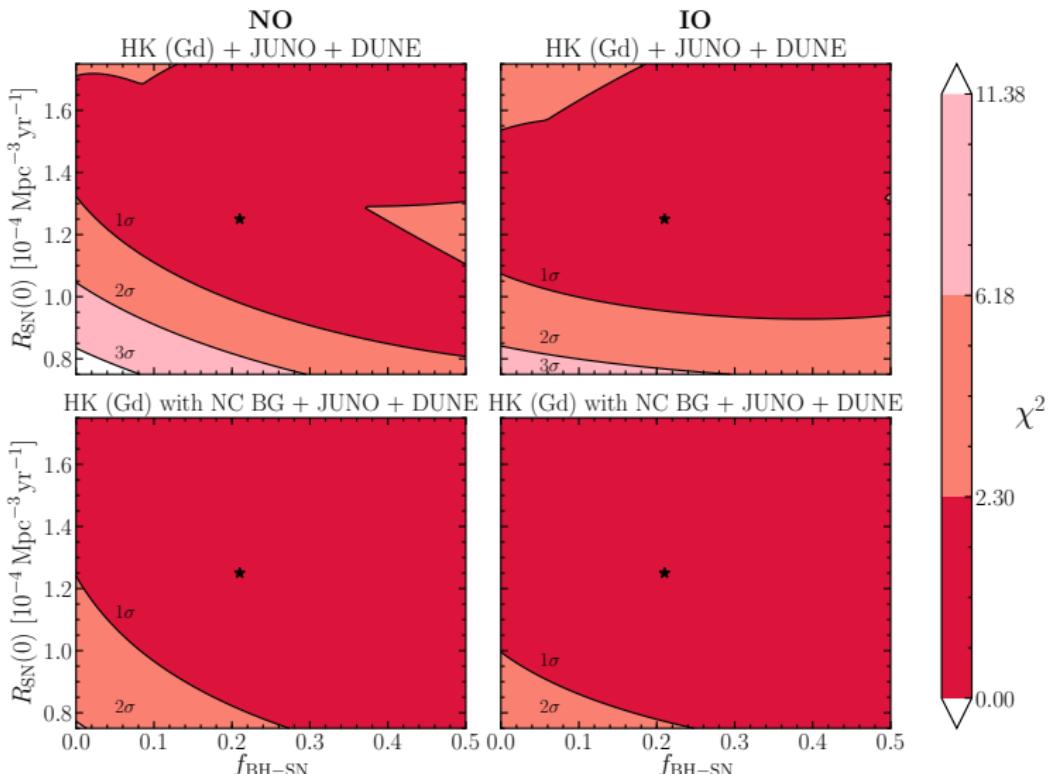
JUNO DSNB detection perspectives



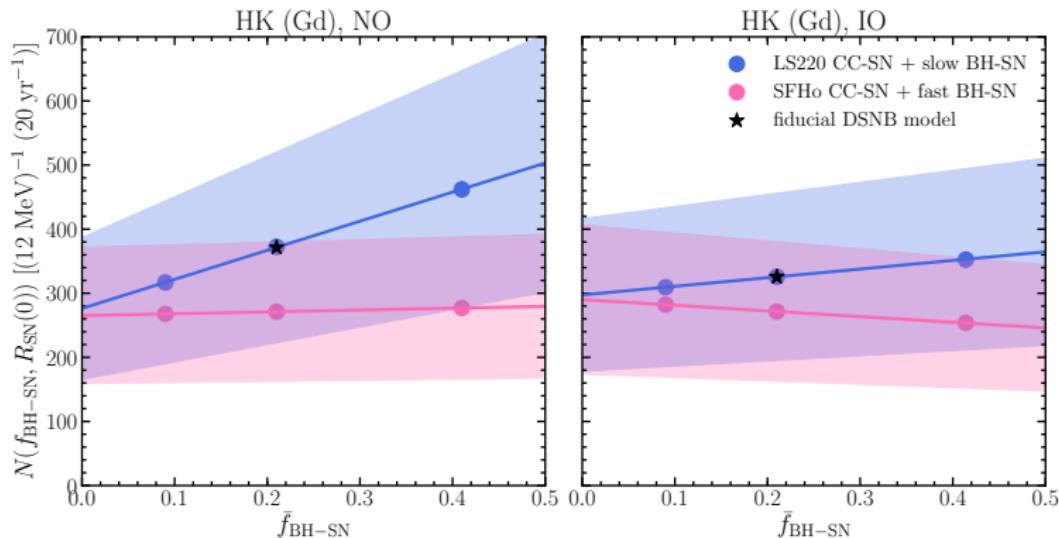
Expected 1D χ^2 as a function of $f_{\text{BH-SN}}$ and $R_{\text{SN}}(0)$



χ^2 for the fraction of BH forming progenitors - local supernova rate plane



Number of events in HK (Gd) energy window



HK detection perspectives HK report

