

ME531: Advanced Mechanics of Solids

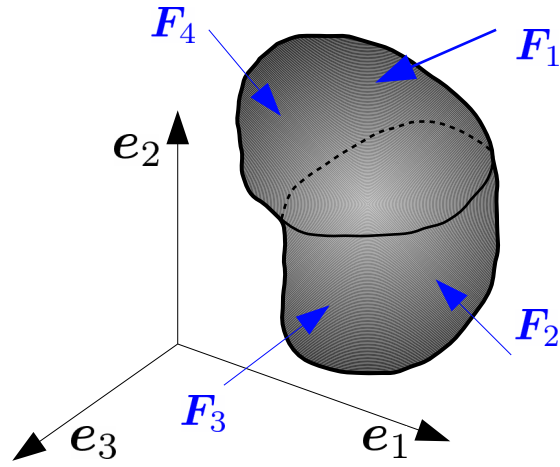
Motion, Strain and Stress

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Stress

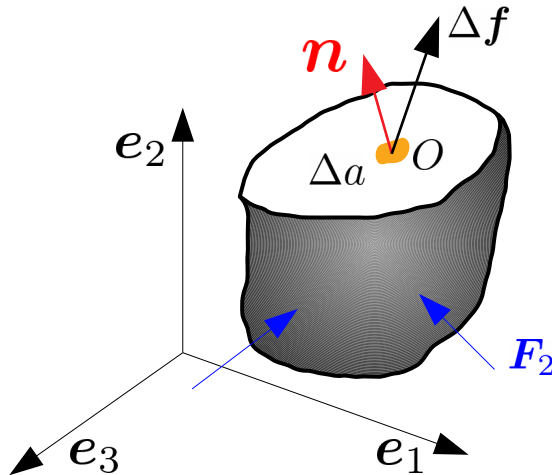


Consider a body in the deformed configuration acted upon by several external forces. Let the body be cut by an imaginary plane passing through point O .

Now, consider an elemental area ΔA in the neighborhood of point O . The reaction force on the area is $\Delta \mathbf{F}$, and normal to the area is \mathbf{n} .

We define a traction vector \mathbf{t} corresponding to normal \mathbf{n} and point O is defined as,

$$\mathbf{t}(\mathbf{x}, \mathbf{n}) = \lim_{\Delta a \rightarrow 0} \frac{\Delta \mathbf{f}}{\Delta a} = \frac{d\mathbf{f}}{da}$$



where $\mathbf{t}(\mathbf{x}, \mathbf{n})$ is called *Cauchy-traction vector* (force per unit surface area in the current configuration), exerted on the area da having normal \mathbf{n} . Traction vectors are also referred to as *surface traction*, *contact forces*, *stress vectors* or *loads*.

Cauchy's stress theorem

There exist a unique second-order tensor fields so that,

$$\mathbf{t}(\mathbf{x}, \mathbf{n}) = \mathbf{n} \cdot \boldsymbol{\sigma}(\mathbf{x}), \quad \text{or} \quad \mathbf{t}(\mathbf{x}, \mathbf{n}) = \boldsymbol{\sigma}^T(\mathbf{x}) \cdot \mathbf{n}, \quad \text{or} \quad t_i = \sigma_{ji} n_j.$$

where $\boldsymbol{\sigma}$ is known as *Cauchy stress tensor*. This relation between the stress tensor and the traction vector is known as *Cauchy's stress theorem*. It states that if traction vector \mathbf{t} depends upon the outward unit normal \mathbf{n} then it must be linear in \mathbf{n} . It immediately follows that

$$\mathbf{t}(\mathbf{x}, \mathbf{n}) = -\mathbf{t}(\mathbf{x}, -\mathbf{n}),$$

for all unit normal vectors \mathbf{n} . This is known as *Newton's (third) law of action and reaction*.

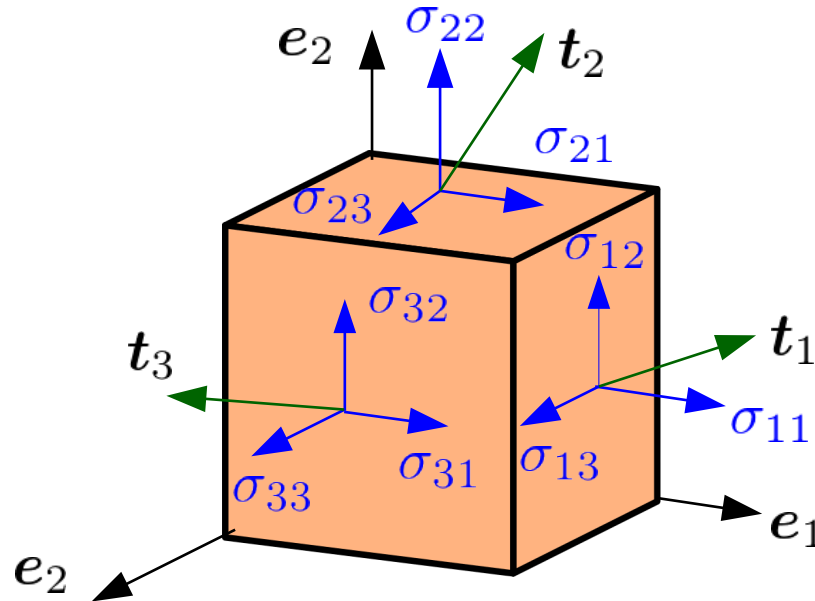
Now, we find the traction vectors at planes having basis vectors $\{e_i\}$ as normal,

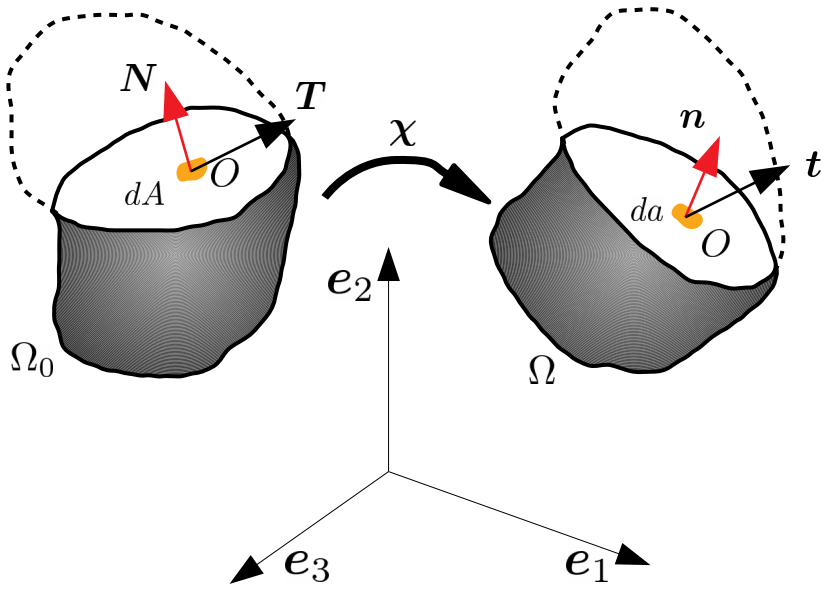
$$t(x, e_1) = \sigma(x)e_1 = \sigma_{11}e_1 + \sigma_{12}e_2 + \sigma_{13}e_3,$$

$$t(x, e_2) = \sigma(x)e_2 = \sigma_{21}e_1 + \sigma_{22}e_2 + \sigma_{23}e_3,$$

$$t(x, e_3) = \sigma(x)e_3 = \sigma_{31}e_1 + \sigma_{32}e_2 + \sigma_{33}e_3.$$

State of stress at a point can be represented by considering an infinitesimal cubic element surrounding it as,





Now, consider a body under stress in its reference and current configuration. Traction vector in the current configuration and the corresponding normal is \mathbf{t} and \mathbf{n} , respectively. Corresponding normal and (pseudo) traction vector in the reference configurations are \mathbf{T} and \mathbf{N} respectively.

\mathbf{T} is known as *first Piola-Kirchhoff traction vector* and works in the same directions as \mathbf{t} . It should be noted that \mathbf{T} does not describe the actual intensity as it works on the reference volume Ω_0 and is a function of reference position \mathbf{X} . Cauchy's stress theorem is also applicable for traction vector $\mathbf{T}(\mathbf{X}, \mathbf{N})$ and normal vector \mathbf{N} , which is as follows,

$$\mathbf{T}(\mathbf{X}, \mathbf{N}) = \mathbf{N} \cdot \mathbf{P}(\mathbf{X}), \quad \text{or} \quad T_i = P_{ji} N_j,$$

Here, \mathbf{P} is a second order tensor and known as *first-Piola Kirchhoff stress tensor*.³⁷

It should be noted that the force on the element area can be expressed as,

$$d\mathbf{f} = \mathbf{T}dA = \mathbf{t} da.$$

Above relation can be used to establish a relation between the Cauchy stress tensor, and the first Piola-Kirchhoff stress tensor. We use the Cauchy's stress theorem and write the previous equation as,

$$\mathbf{N} \cdot \mathbf{P}(\mathbf{X})dA = \mathbf{n} \cdot \boldsymbol{\sigma}(\mathbf{x})da.$$

Using Nanson's relations, we write,

$$\mathbf{P}(\mathbf{X}) = J\mathbf{F}^{-1}\boldsymbol{\sigma}(\mathbf{x}), \quad P_{ij} = JF_{ik}^{-1}\sigma_{kj}.$$

We can also write,

$$\boldsymbol{\sigma} = J^{-1}\mathbf{F}\mathbf{P}, \quad \sigma_{ij} = J^{-1}F_{ik}P_{kj}.$$

Later, we will prove that the Cauchy stress tensor is symmetric under the assumption of zero resultant couples. i.e.

$$\boldsymbol{\sigma} = J^{-1}\mathbf{F}\mathbf{P} = \boldsymbol{\sigma}^T, \quad \sigma_{ij} = J^{-1}P_{ij} = F_{ik}\sigma_{kj},$$

which implies that $\mathbf{F}\mathbf{P} = \mathbf{P}^T\mathbf{F}^T$, which also suggests that the tensor \mathbf{P} in general is not symmetric.

Example

Deformation of a body is described by,

$$x_1 = -6X_2, \quad x_2 = 0.5X_1, \quad x_3 = 1/3 X_3.$$

The Cauchy stress tensor for certain part of the body is given by the matrix representation as,

$$[\boldsymbol{\sigma}] = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 50 & 0 \\ 0 & 0 & 0 \end{bmatrix} \text{ kg/cm}^2$$

Determine the Cauchy traction vector \mathbf{t} and the first Piola-Kirchhoff traction vector \mathbf{T} acting on a plane, which is characterized by the outward unit normal $\mathbf{n} = \mathbf{e}_2$ in the current configuration.

For the given deformation,

$$[\mathbf{F}] = \begin{bmatrix} 0 & -6 & 0 \\ 1/2 & 0 & 0 \\ 0 & 0 & 1/3 \end{bmatrix}, \quad [\mathbf{F}]^{-1} = \begin{bmatrix} 0 & 2 & 0 \\ -1/6 & 0 & 0 \\ 0 & 0 & 3 \end{bmatrix} \quad \text{and } \det \mathbf{F} = 1.$$

The components for first Piola-Kirchoff stress tensor will be

$$[\mathbf{P}] = J[\mathbf{F}]^{-1}[\boldsymbol{\sigma}] = \begin{bmatrix} 0 & 100 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \text{ kN/cm}^2.$$

Outward unit normal \mathbf{N} in the reference configuration will be related to \mathbf{n} by Nanson's formula as,

$$\text{Thus, for } \mathbf{n}=\mathbf{e}_2, \quad \mathbf{N}dS = J^{-1}\mathbf{F}^T\mathbf{n}ds \Rightarrow \mathbf{N}dS = \frac{\mathbf{e}_1}{2}ds.$$

Hence, $\mathbf{N} = \mathbf{e}_1$ and $dS = ds/2$.

Finally, using Cauchy's stress theorem,

$$\{\mathbf{t}\} = [\boldsymbol{\sigma}]^T [\mathbf{n}] = \begin{bmatrix} 0 \\ 50 \\ 0 \end{bmatrix} \text{ kN/cm}^2, \quad \{\mathbf{T}\} = [\mathbf{P}]^T [\mathbf{N}] = \begin{bmatrix} 0 \\ 100 \\ 0 \end{bmatrix} \text{ kN/cm}^2.$$

i.e. $\mathbf{t} = 50\mathbf{e}_2$ and $\mathbf{T} = 100\mathbf{e}_2$.

It can be observed that both \mathbf{t} and \mathbf{T} have same direction, but the magnitudes of \mathbf{T} is twice that of \mathbf{t} , as deformed area is half the undeformed area.