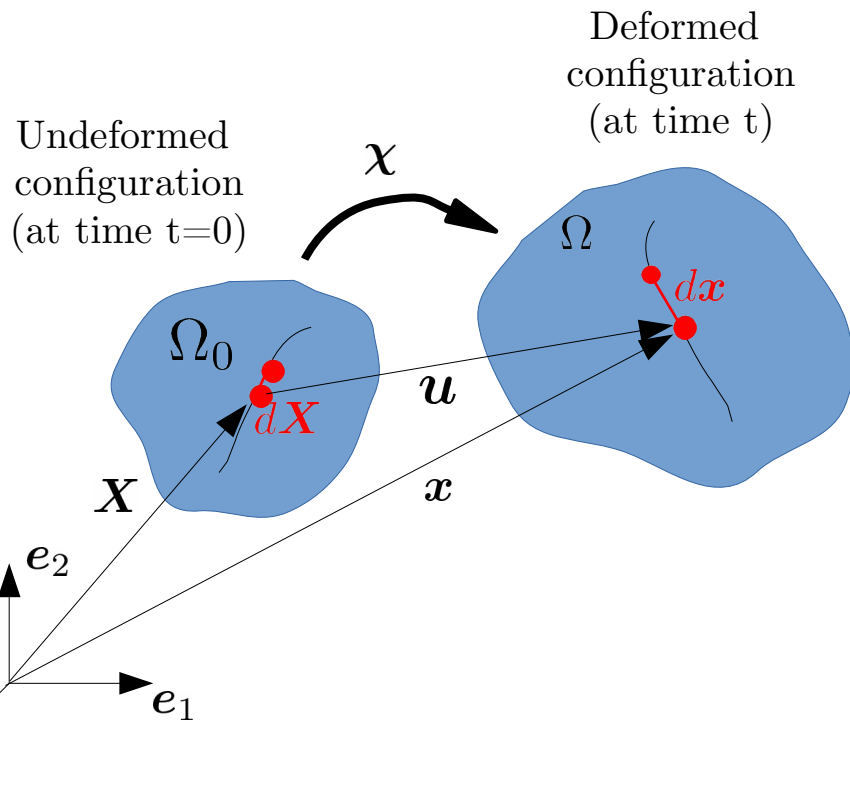


ME531: Advanced Mechanics of Solids

Motion, Strain and Stress

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Deformation gradient



Consider a line element $d\mathbf{X}$ in Ω_0 which deforms to a line element $d\mathbf{x}$. *Deformation gradient* tensor maps the line element $d\mathbf{X}$ to $d\mathbf{x}$ as,

$$d\mathbf{x} = \mathbf{F}d\mathbf{X} \text{ or } dx_i = F_{iJ}dX_J,$$

where, $\mathbf{F} = \frac{\partial \mathbf{x}}{\partial \mathbf{X}}$ or $F_{iJ} = \frac{\partial x_i}{\partial X_J}$,
or $\mathbf{F} = \text{Grad } \mathbf{x} = \nabla_{\mathbf{X}} \mathbf{x}$

is an invertible tensor. \mathbf{F}^{-1} is the *inverse deformation gradient* tensor, defined as

$$\mathbf{F}^{-1} = \frac{\partial \mathbf{X}}{\partial \mathbf{x}} \text{ or } F_{Ij}^{-1} = \frac{\partial X_I}{\partial x_j}.$$

or $\mathbf{F}^{-1} = \text{grad } \mathbf{X} = \nabla_{\mathbf{x}} \mathbf{X}$

It relates line elements as,

$$d\mathbf{X} = \mathbf{F}^{-1}d\mathbf{x} \text{ or } dX_I = F_{Ij}^{-1}dx_j.$$

Example

Consider a two dimensional motion given by two equations as

$$x_1 = 4 - 2X_1 - X_2$$

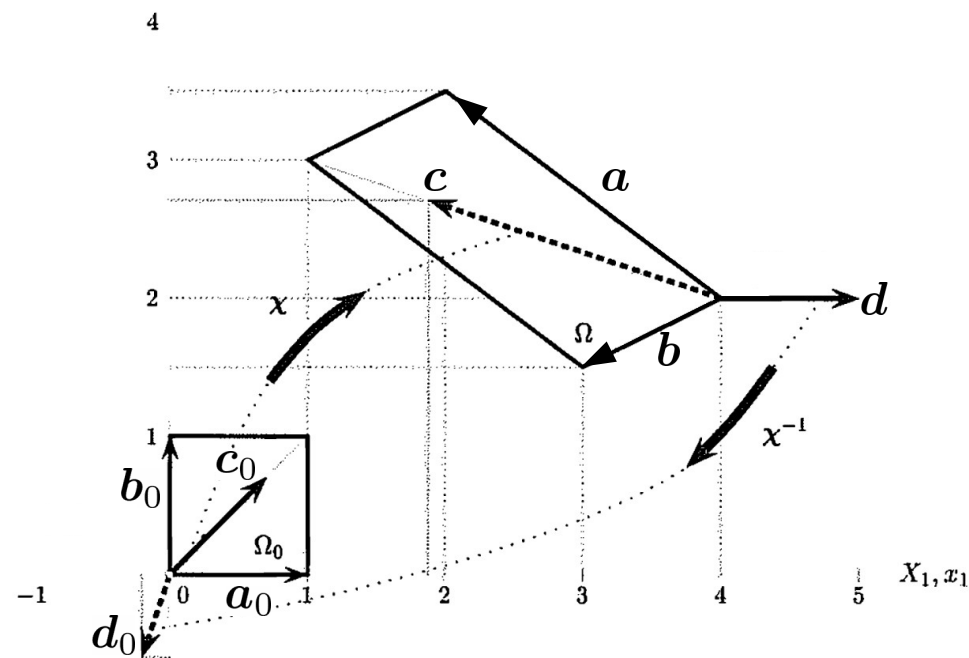
$$x_2 = 2 + 1.5X_1 - 0.5X_2$$

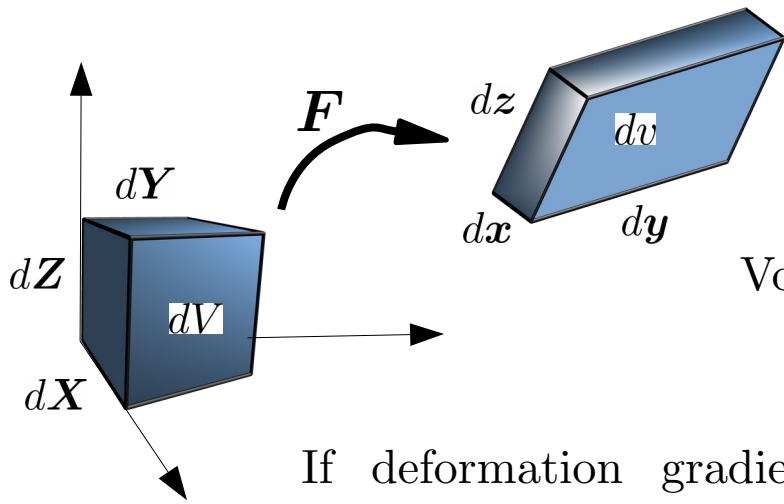
Deformation gradient for the given motion is calculated as,

$$[\mathbf{F}] = \begin{bmatrix} -2 & -1 \\ 1.5 & -0.5 \end{bmatrix}, \text{ and } [\mathbf{F}]^{-1} = \frac{1}{5} \begin{bmatrix} -1 & 2 \\ -3 & -4 \end{bmatrix}$$

Now consider following vectors $\mathbf{a}_0 = [1,0]$, $\mathbf{b}_0 = [0,1]$ and $\mathbf{c}_0 = [0.707, 0.707]$ in the reference configuration. Current or deformed vectors can be obtained as $\mathbf{a} = \mathbf{F}\mathbf{a}_0$, $\mathbf{b} = \mathbf{F}\mathbf{b}_0$ and $\mathbf{c} = \mathbf{F}\mathbf{c}_0$.

Consider another vector $\mathbf{d} = [1,0]$ in the current configuration. Reference configuration \mathbf{d}_0 for vector \mathbf{d} can be obtained as $\mathbf{d}_0 = \mathbf{F}^{-1}\mathbf{d}$.





Consider an infinitesimally small volume element with reference volume dV going under deformation and the current volume is dv .

Volume of the element in current configuration can be given as,

$$dv = d\mathbf{z} \cdot (d\mathbf{x} \times d\mathbf{y}) \text{ or } dv = e_{ijk} dz_i dx_j dy_k$$

If deformation gradient for the motion is \mathbf{F} , then edges in the reference configuration can be mapped to the edges in undeformed configuration as,

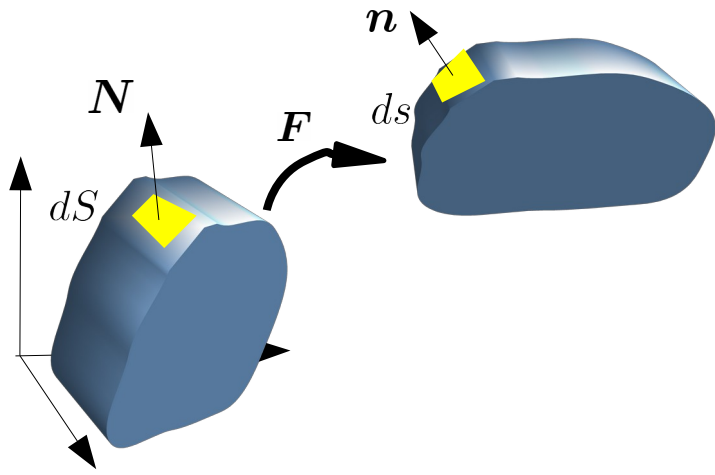
$$dx_i = F_{ij} dX_J, dy_i = F_{ij} dY_J, \text{ and } dz_i = F_{ij} dZ_J$$

Substituting above in the expression for volume, we get

$$dv = e_{ijk} (F_{ip} dZ_p) (F_{jq} dX_q) (F_{kr} dY_r)$$

It can be shown that, $e_{ijk} A_{ip} A_{jq} A_{kr} = e_{pqr} \det \mathbf{A}$.

Thus, $dv = e_{pqr} dZ_p dX_q dY_r \det \mathbf{F} \Rightarrow dv = J dV$, where $J = \det \mathbf{F}$ is called the *Jacobian* of deformation gradient tensor. The Jacobian is the ratio of elemental volume deformed configuration to its volume in undeformed configuration.



Deformation gradient \mathbf{F} is used to map a vector in undeformed configuration to the deformed configuration. However, the same does not apply to the unit normal vector \mathbf{N} to an infinitesimal surface element dS . To find the relation between the normal vector \mathbf{N} and \mathbf{n} (deformed configuration of \mathbf{N}) we use the following relation.

$$dv = JdV$$

We now present the volume dV as $d\mathbf{S} \cdot d\mathbf{X}$, where $d\mathbf{X}$ is a line element in undeformed configuration, which deform to the line element $d\mathbf{x}$, and hence $dv = d\mathbf{s} \cdot d\mathbf{x}$.

Now, we can write,

$$d\mathbf{s} \cdot d\mathbf{x} = Jd\mathbf{S} \cdot d\mathbf{X}$$

Line element $d\mathbf{x}$ and $d\mathbf{X}$ are related as, $d\mathbf{x} = \mathbf{F}d\mathbf{X}$. Using this relation we can write,

$$d\mathbf{s} \cdot \mathbf{F}d\mathbf{X} = Jd\mathbf{S} \cdot d\mathbf{X} \Rightarrow (\mathbf{F}^T d\mathbf{s} - Jd\mathbf{S}) \cdot d\mathbf{X} = 0$$

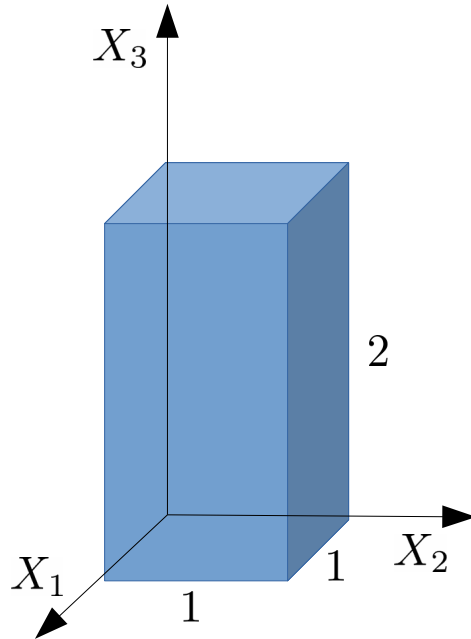
This relation holds for any arbitrary $d\mathbf{X}$, which implies $\mathbf{F}^T d\mathbf{s} - Jd\mathbf{S} = 0$.

Above gives us the **Nanson's formula** relating an area element in undeformed configuration to the area element in deformed configuration as

$$d\mathbf{s} = J\mathbf{F}^{-T}d\mathbf{S}.$$

Exercise

Plot the deformed shape of the cube which goes under the following motion.



$$x_1 = X_1 \cos(\pi X_3/2) - X_2 \sin(\pi X_3/2)$$

$$x_2 = X_1 \sin(\pi X_3/2) + X_2 \cos(\pi X_3/2)$$

$$x_3 = X_3$$

Displacement gradient

Using displacements we can write, $\mathbf{x}(\mathbf{X}, t) = \mathbf{X} + \mathbf{U}(\mathbf{X}, t)$ or $x_i = X_i + U_i$.

From the definition of deformation gradient,

$$F_{iJ} = \frac{\partial x_i}{\partial X_J} = \frac{\partial}{\partial X_J}(X_i + U_i) = \frac{\partial X_i}{\partial X_J} + \frac{\partial U_i}{\partial X_J}$$

$$F_{iJ} = \delta_{ij} + \frac{\partial U_i}{\partial X_J}, \text{ or } \mathbf{F} = \mathbf{I} + \nabla_{\mathbf{X}} \mathbf{U}$$

where, $\nabla_{\mathbf{X}} \mathbf{U} = \frac{\partial U_i}{\partial X_J}$ is called *displacement gradient tensor*.

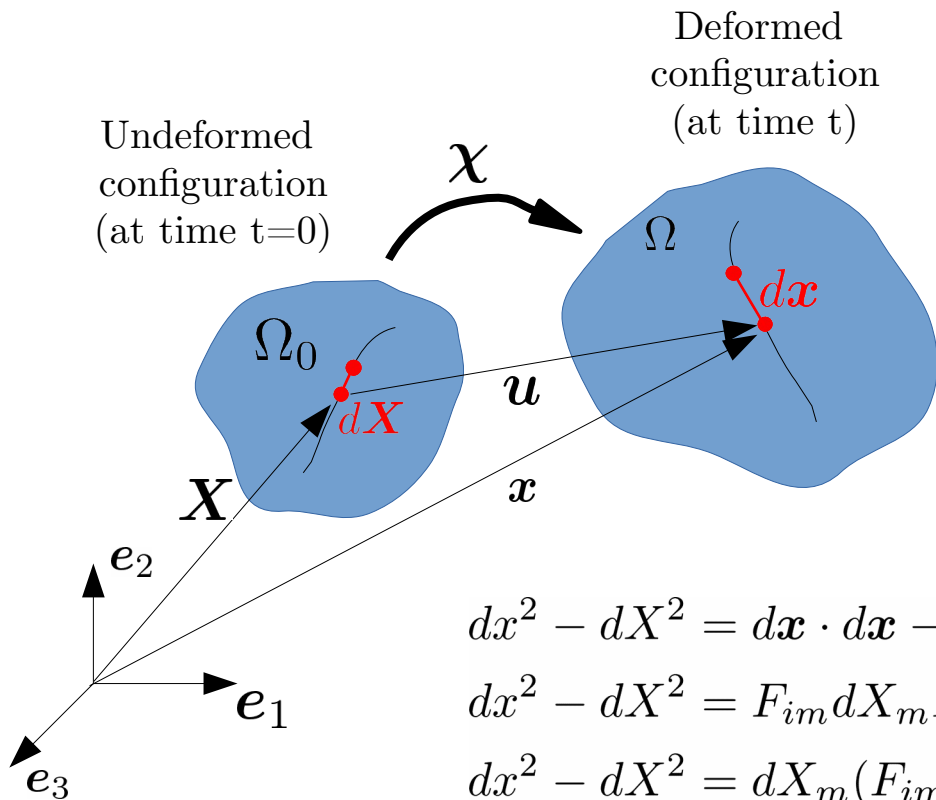
In material description, displacement gradient tensor is

$$\nabla_{\mathbf{X}} \mathbf{U}(\mathbf{X}, t) = \mathbf{F}(\mathbf{X}, t) - \mathbf{I} \text{ or } \frac{\partial U_i}{\partial X_J} = F_{iJ} - \delta_{ij}.$$

In spatial description, displacement gradient tensor is

$$\nabla \mathbf{u}(\mathbf{x}, t) = \mathbf{I} - \mathbf{F}^{-1}(\mathbf{x}, t) \text{ or } \frac{\partial u_i}{\partial x_J} = \delta_{ij} - F_{iJ}^{-1}.$$

Strain



Squared length of the line element $d\mathbf{X}$ is $dX^2 = d\mathbf{X} \cdot d\mathbf{X}$.

Similarly, squared length of the line element $d\mathbf{x}$ is $dx^2 = d\mathbf{x} \cdot d\mathbf{x}$.

Let's examine the difference between the squared length of $d\mathbf{x}$ and $d\mathbf{X}$ as,

$$dx^2 - dX^2 = d\mathbf{x} \cdot d\mathbf{x} - d\mathbf{X} \cdot d\mathbf{X} = dx_i dx_i - dX_j dX_j,$$

$$dx^2 - dX^2 = F_{im} dX_m F_{in} dX_n - dX_m \delta_{mn} dX_n,$$

$$dx^2 - dX^2 = dX_m (F_{im} F_{in} - \delta_{mn}) dX_n = d\mathbf{X} \cdot (\mathbf{F}^T \mathbf{F} - \mathbf{I}) d\mathbf{X},$$

$$dx^2 - dX^2 = dX_m (F_{im} F_{in} - \delta_{mn}) dX_n = d\mathbf{X} \cdot 2\mathbf{E} d\mathbf{X}.$$

where, $\mathbf{E} = \frac{1}{2} (\mathbf{F}^T \mathbf{F} - \mathbf{I})$ or $E_{ij} = \frac{1}{2} (F_{ki} F_{kj} - \delta_{ij})$ is called *Green-Lagrange* strain tensor.

Using the relation, $\mathbf{F} = \mathbf{I} + \nabla_X \mathbf{U}$ or $F_{ij} = \delta_{ij} + \frac{\partial U_i}{\partial X_j}$, we write,

$$E_{ij} = \frac{1}{2} \left[\left(\delta_{ki} + \frac{\partial U_k}{\partial X_i} \right) \left(\delta_{kj} + \frac{\partial U_k}{\partial X_j} \right) - \delta_{ij} \right]$$

$$E_{ij} = \frac{1}{2} \left[\frac{\partial U_i}{\partial X_j} + \frac{\partial U_j}{\partial X_i} + \frac{\partial U_k}{\partial X_i} \frac{\partial U_k}{\partial X_j} \right] \quad \text{or} \quad \mathbf{E} = \frac{1}{2} \left[\nabla_X \mathbf{U} + (\nabla_X \mathbf{U})^T + (\nabla_X \mathbf{U})^T \nabla_X \mathbf{U} \right]$$

Eulerian description of the Green-Lagrange strain tensor is called *Euler-Almansi strain tensor* and defined as,

$$\mathbf{e} = \frac{1}{2} (\mathbf{I} - \mathbf{F}^{-T} \mathbf{F}^{-1}) \quad \text{or} \quad e_{ij} = \frac{1}{2} (\delta_{ij} - F_{ki}^{-1} F_{kj})$$

$$\mathbf{e} = \frac{1}{2} \left[\nabla \mathbf{u} + (\nabla \mathbf{u})^T - (\nabla \mathbf{u})^T \nabla \mathbf{u} \right] \quad \text{or} \quad e_{ij} = \frac{1}{2} \left[\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} - \frac{\partial u_k}{\partial x_i} \frac{\partial u_k}{\partial x_j} \right]$$