

ME531: Advanced Mechanics of Solids

Motion, Strain and Stress

Anshul Faye

afaye@iitbhilai.ac.in

Room No. # 106

Procedure for polar decomposition of \mathbf{F}

1. First calculate \mathbf{C} as, $\mathbf{C} = \mathbf{F}^T \mathbf{F}$.
2. Find the eigenvalues and eigenvectors of \mathbf{C} as λ_i and \mathbf{N}_i .
3. Eigenvalues of \mathbf{U} is then $\sqrt{\lambda_i}$ and eigenvector remain same as \mathbf{N}_i .
4. Tensor \mathbf{U} can be obtained as,

$$\mathbf{U} = \sum_{i=1}^3 \sqrt{\lambda_i} \mathbf{N}_i \mathbf{N}_i, \quad \text{where } \mathbf{N}_i = N_i^1 \mathbf{e}_1 + N_i^2 \mathbf{e}_2 + N_i^3 \mathbf{e}_3.$$

5. Now, \mathbf{R} can be obtained as, $\mathbf{R} = \mathbf{F} \mathbf{U}^{-1}$. Here \mathbf{U} is in $\{\mathbf{e}_i\}$ basis vector, which in the matrix form can be obtained as $[\mathbf{V}] [\mathbf{U}_d] [\mathbf{V}]^T$. $[\mathbf{U}_d]$ is the diagonal form of stretch tensor with eigenvectors as basis and $[\mathbf{V}]$ is a matrix with i^{th} column corresponding to the i^{th} eigenvector.
6. Now, \mathbf{v} can be determined as, $\mathbf{v} = \mathbf{F} \mathbf{R}^T$.

Exercise

For the given motion, find \mathbf{F} , \mathbf{R} , \mathbf{U} , and \mathbf{v} tensors.

$$x_1 = 4 - 2X_1 - X_2$$

$$x_2 = 2 + 1.5x_1 - 0.5X_2$$

Also plot the configuration of a unit square having two opposite corners at (0,0) and (1,1) after the application of (i) \mathbf{U} and then \mathbf{R} , (ii) \mathbf{R} and then \mathbf{v} .

Material and Spatial derivatives

Material field: Field which is a function of material coordinates \mathbf{X} and time t .

Spatial field: Field which is a function of spatial coordinates \mathbf{x} and t .

Material time derivative of a material field:

Consider a smooth material field $\mathcal{F}(\mathbf{X}, t)$. Material time derivative or *total time derivative* of the field is defined as,

$$\dot{\mathcal{F}}(\mathbf{X}, t) = \frac{D\mathcal{F}(\mathbf{X}, t)}{Dt} = \left(\frac{\partial \mathcal{F}(\mathbf{X}, t)}{\partial t} \right)_{\mathbf{X}} \text{ or } \dot{\mathcal{F}}(X_A, t) = \left(\frac{\partial \mathcal{F}(X_A, t)}{\partial t} \right)_{X_A}.$$

Material gradient of a material field:

Material gradient of the material field $\mathcal{F}(\mathbf{X}, t)$ is given as,

$$\text{Grad } \mathcal{F}(\mathbf{X}, t) = \left(\frac{\partial \mathcal{F}(\mathbf{X}, t)}{\partial \mathbf{X}} \right)_t$$

Spatial time derivative of a spatial field:

Consider a smooth spatial field $f(\mathbf{x}, t)$. The spatial time derivative of the field is the derivative of \mathbf{x} w.r.t time t holding the current position \mathbf{x} . fixed. It is also referred as *local time derivative*. It is simply denoted as,

$$\left(\frac{\partial f}{\partial t} \right)_{\mathbf{x}}.$$

It represents the change in f with time t as seen by an observer currently situated at \mathbf{x} .

Spatial gradient of a spatial field:

Spatial gradient of the spatial field $f(\mathbf{x}, t)$, $\text{grad}(\mathbf{x})$ is the derivative w.r.t. the current position \mathbf{x} holding time t fixed.

$$\text{grad } f(\mathbf{x}, t) = \left(\frac{\partial f(\mathbf{x}, t)}{\partial \mathbf{x}} \right)_t$$

Material time derivative of a spatial field:

The material time derivative of the spatial field $f(\mathbf{x}, t)$ is the time derivative of f holding \mathbf{X} fixed.

$$\dot{f}(\mathbf{x}, t) = \frac{Df(\mathbf{x}, t)}{Dt} = \left(\frac{\partial f[\boldsymbol{\chi}(X, t), t]}{\partial t} \right)_{\mathbf{X}=\boldsymbol{\chi}^{-1}(\mathbf{x}, t)}.$$

Now, consider a smooth scalar valued spatial field $\phi(\mathbf{x}, t)$, by chain rule

$$\dot{\phi}(\mathbf{x}, t) = \left(\frac{\partial \phi(\mathbf{x}, t)}{\partial t} \right)_{\mathbf{x}} + \left(\frac{\partial \phi(\mathbf{x}, t)}{\partial \mathbf{x}} \right)_{\mathbf{t}} \cdot \left(\frac{\partial \boldsymbol{\chi}(X, t)}{\partial t} \right)_{\mathbf{X}=\boldsymbol{\chi}^{-1}(\mathbf{x}, t)}.$$

$$\dot{\phi}(\mathbf{x}, t) = \frac{D\phi(\mathbf{x}, t)}{Dt} = \boxed{\frac{\partial \phi(\mathbf{x}, t)}{\partial t}} + \boxed{\text{grad } \phi(\mathbf{x}, t) \cdot \mathbf{v}(\mathbf{x}, t)} \text{ or } \dot{\phi}(\mathbf{x}, t) = \frac{D\phi}{Dt} = \boxed{\frac{\partial \phi}{\partial t}} + \boxed{\frac{\partial \phi}{\partial x_i} v_i}.$$

Here, the first term on the right hand side is the *spatial or local time derivative* of spatial scalar field ϕ , while the second term is called the *convective rate of change* of ϕ , which describes the changing position of a partial.

Similarly, material time derivative of a smooth vector-valued spatial field $\mathbf{v}(\mathbf{x}, t)$ will be,

$$\dot{\mathbf{v}}(\mathbf{x}, t) = \frac{D\mathbf{v}(\mathbf{x}, t)}{Dt} = \frac{\partial \mathbf{v}(\mathbf{x}, t)}{\partial t} + \text{grad } \mathbf{v} \cdot \mathbf{v}(\mathbf{x}, t) \quad \text{or} \quad \dot{v}_i(\mathbf{x}, t) = \frac{Dv_i}{Dt} = \frac{\partial v_i}{\partial t} + \frac{\partial v_i}{\partial x_j} v_j.$$

If $\mathbf{v}(\mathbf{x}, t)$ is considered as spatial description of velocity then $\dot{\mathbf{v}}(\mathbf{x}, t)$ describes the spatial acceleration field (\mathbf{a}). Hence,

$$\mathbf{a} = \frac{D\mathbf{v}}{Dt} = \frac{\partial \mathbf{v}}{\partial t} + \text{grad } \mathbf{v} \cdot \mathbf{v} \quad \text{or} \quad a_i = \frac{Dv_i}{Dt} = \frac{\partial v_i}{\partial t} + \frac{\partial v_i}{\partial x_j} v_j.$$

Again, the first term on the right hand side $\partial v_i / \partial t$ describes the local acceleration (local rate of change of the velocity field). The second term describes the *convective acceleration* (convective rate of change of the velocity field).

Spatial velocity field $\mathbf{v}(\mathbf{x}, t)$ is defined as,

$$\mathbf{v}(\mathbf{x}, t) = \frac{D\mathbf{x}(\mathbf{X}, t)}{Dt} = \left(\frac{\partial \mathbf{x}(\mathbf{X}, t)}{\partial t} \right)_{\mathbf{X}} = \frac{\partial \mathbf{x}}{\partial t}$$

Rate of deformation tensor

Consider the spatial derivative of spatial velocity field $\mathbf{v}(\mathbf{x}, t)$ defined by

$$\mathbf{l}(\mathbf{x}, t) = \frac{\partial \mathbf{v}(\mathbf{x}, t)}{\partial \mathbf{x}} = \text{grad } \mathbf{v}(\mathbf{x}, t) \quad \text{or} \quad l_{ij} = \frac{\partial v_i}{\partial x_j},$$

where spatial field $\mathbf{l}(\mathbf{x}, t)$ is known as *spatial velocity gradient*, which is in general a non-symmetric second order tensor.

Material velocity gradient of a material velocity field $\mathbf{V}(\mathbf{X}, t)$ is defined as

$$\frac{\partial \mathbf{V}(\mathbf{X}, t)}{\partial \mathbf{X}} = \frac{\partial}{\partial \mathbf{X}} \left(\frac{\partial \boldsymbol{\chi}(\mathbf{X}, t)}{\partial t} \right) = \frac{\partial}{\partial t} \left(\frac{\partial \boldsymbol{\chi}(\mathbf{X}, t)}{\partial \mathbf{X}} \right) = \dot{\mathbf{F}}(\mathbf{X}, t)$$

$$\dot{\mathbf{F}}(\mathbf{X}, t) = \frac{\partial \mathbf{V}(\mathbf{X}, t)}{\partial \mathbf{X}} = \text{Grad } \mathbf{V}$$

Using the chain rule of differentiation, spatial velocity gradient can be written as,

$$\begin{aligned} \boldsymbol{l}(\boldsymbol{x}, t) &= \frac{\partial \boldsymbol{v}(\boldsymbol{x}, t)}{\partial \boldsymbol{x}} = \frac{\partial \boldsymbol{v}(\boldsymbol{x}, t)}{\partial \boldsymbol{X}} \frac{\partial \boldsymbol{X}(\boldsymbol{x}, t)}{\partial \boldsymbol{x}} = \frac{\partial \dot{\boldsymbol{\chi}}(\boldsymbol{X}, t)}{\partial \boldsymbol{X}} \frac{\partial \boldsymbol{X}(\boldsymbol{x}, t)}{\partial \boldsymbol{x}} \\ &= \frac{\partial}{\partial t} \left(\frac{\partial \boldsymbol{\chi}(\boldsymbol{X}, t)}{\partial \boldsymbol{X}} \right) \frac{\partial \boldsymbol{X}(\boldsymbol{x}, t)}{\partial \boldsymbol{x}} = \frac{\partial \boldsymbol{F}}{\partial t} \boldsymbol{F}^{-1} \end{aligned}$$

Thus, $\boldsymbol{l} = \dot{\boldsymbol{F}} \boldsymbol{F}^{-1}$, hence $\dot{\boldsymbol{F}} = \boldsymbol{l} \boldsymbol{F}$, or $l_{ij} = \dot{F}_{ik} F_{kj}^{-1}$, hence $\dot{F}_{ij} = l_{ik} F_{kj}$.

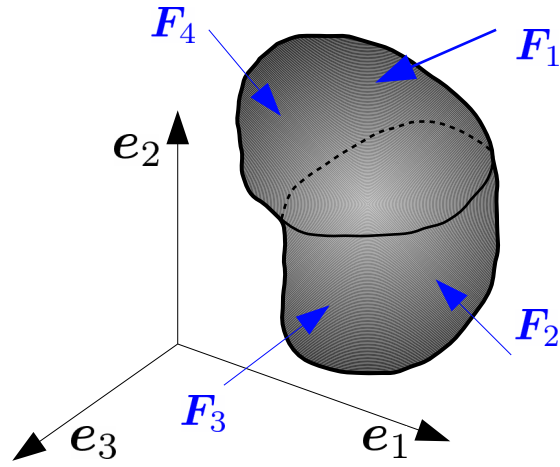
Spatial velocity gradient can be additively decomposed into a symmetric and anti-symmetric part as,

$$\boldsymbol{l}(\boldsymbol{x}, t) = \boldsymbol{d}(\boldsymbol{x}, t) + \boldsymbol{w}(\boldsymbol{x}, t),$$

where $\boldsymbol{d}(\boldsymbol{x}, t)$ is called the rate of deformation tensor and $\boldsymbol{w}(\boldsymbol{x}, t)$ is called the spin tensor, and they are defined as

$$\boldsymbol{d} = \frac{1}{2} (\boldsymbol{l} + \boldsymbol{l}^T), \text{ and } \boldsymbol{w} = \frac{1}{2} (\boldsymbol{l} - \boldsymbol{l}^T).$$

Stress

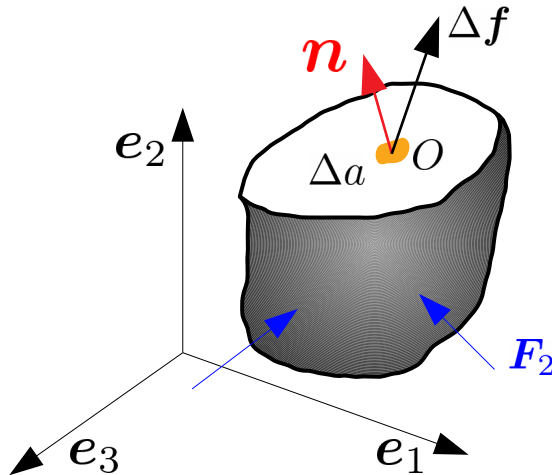


Consider a body in the deformed configuration acted upon by several external forces. Let the body be cut by an imaginary plane passing through point O .

Now, consider an elemental area ΔA in the neighborhood of point O . The reaction force on the area is $\Delta \mathbf{F}$, and normal to the area is \mathbf{n} .

We define a traction vector \mathbf{t} corresponding to normal \mathbf{n} and point O is defined as,

$$\mathbf{t}(\mathbf{x}, \mathbf{n}) = \lim_{\Delta a \rightarrow 0} \frac{\Delta \mathbf{f}}{\Delta a} = \frac{d\mathbf{f}}{da}$$



where $\mathbf{t}(\mathbf{x}, \mathbf{n})$ is called *Cauchy-traction vector* (force per unit surface area in the current configuration), exerted on the area da having normal \mathbf{n} . Traction vectors are also referred to as *surface traction*, *contact forces*, *stress vectors* or *loads*.

Cauchy's stress theorem

There exist a unique second-order tensor fields so that,

$$\mathbf{t}(\mathbf{x}, \mathbf{n}) = \boldsymbol{\sigma}^T(\mathbf{x}) \cdot \mathbf{n}, \quad \text{or} \quad t_i = \sigma_{ji} n_j.$$

where $\boldsymbol{\sigma}$ is known as *Cauchy stress tensor*. This relation between the stress tensor and the traction vector is known as *Cauchy's stress theorem*. It states that if traction vector \mathbf{t} depends upon the outward unit normal \mathbf{n} then it must be linear in \mathbf{n} . It immediately follows that

$$\mathbf{t}(\mathbf{x}, \mathbf{n}) = -\mathbf{t}(\mathbf{x}, -\mathbf{n}),$$

for all unit normal vectors \mathbf{n} . This is known as *Newton's (third) law of action and reaction*.