

Introduction to Tensors

Eigenvalues and Eigenvectors

For a tensor \mathbf{A} , if there exists scalars λ_i and corresponding normalized vectors $\hat{\mathbf{n}}_i$ such that

$$\mathbf{A}\hat{\mathbf{n}}_i = \lambda_i\hat{\mathbf{n}}_i, \quad (i = 1, 2, 3; \text{ no summation})$$

then λ_i are called *eigenvalues* (or principal values) and $\hat{\mathbf{n}}_i$ are called *eigenvectors* (or principal directions or principal axes) of tensor \mathbf{A} .

A set of homogeneous algebraic equation to determine unknown eigenvalues λ_i (i=1,2,3) and unknown eigenvectors $\hat{\mathbf{n}}_i$ (i=1,2,3) are,

$$(\mathbf{A} - \lambda_i\mathbf{I})\hat{\mathbf{n}}_i = \mathbf{o}, \quad (i = 1, 2, 3; \text{ no summation})$$

Eigenvalues characterize the physical nature of a tensor. They do not depend upon the coordinates. For a *positive definite tensor* all eigenvalues are *real* (and *positive*). Also, the set of eigenvectors of a symmetric tensor form a *mutually orthogonal* (or *orthonormal*) basis $\{\hat{\mathbf{n}}_i\}$.

The trivial solution of the system given by $(\mathbf{A} - \lambda_i \mathbf{I})\hat{\mathbf{n}}_i = \mathbf{o}$ is $\hat{\mathbf{n}}_i = \mathbf{o}$.

For system to have solutions $\hat{\mathbf{n}}_i \neq \mathbf{o}$ following condition should be satisfied,

$$\det(\mathbf{A} - \lambda_i \mathbf{I}) = 0$$

where $\det(\mathbf{A} - \lambda_i \mathbf{I}) = -\lambda_i^3 + I_1 \lambda_i^2 - I_2 \lambda_i + I_3$. This requires to solve a cubic equation

$$\lambda^3 - I_1 \lambda^2 + I_2 \lambda - I_3 = 0,$$

which is known as characteristic equation of \mathbf{A} . The solution of the equations are the eigenvalues λ_i ($i=1,2,3$). *Scaler Invariants* (or *principal scalar invariants*) I_1 , I_2 and, I_3 in terms of eigenvalues are,

$$I_1 = \text{tr} \mathbf{A} = A_{ii} = \lambda_1 + \lambda_2 + \lambda_3$$

$$I_2 = \frac{1}{2} [I_1^2 - \text{tr}(\mathbf{A}^2)] = \frac{1}{2} [A_{ii}A_{jj} - A_{ij}A_{ji}] = \lambda_1 \lambda_2 + \lambda_2 \lambda_3 + \lambda_3 \lambda_1$$

$$I_3 = \det \mathbf{A} = \lambda_1 \lambda_2 \lambda_3$$

Application of \mathbf{A} to the equation $\mathbf{A}\hat{\mathbf{n}}_i = \lambda_i\hat{\mathbf{n}}_i$ results,

$$\mathbf{A}^2\hat{\mathbf{n}}_i = \lambda_i\mathbf{A}\hat{\mathbf{n}}_i = \lambda_i^2\hat{\mathbf{n}}_i.$$

Operating the above equation with tensor \mathbf{A} again gives, $\mathbf{A}^3\hat{\mathbf{n}}_i = \lambda_i^3\hat{\mathbf{n}}_i$.

Similarly, after repeated application of \mathbf{A} , we can write a general relation as,

$$\mathbf{A}^a\hat{\mathbf{n}}_i = \lambda_i^a\hat{\mathbf{n}}_i, \text{ where } a \text{ is a positive integer.}$$

By multiplying the characteristic equation by $\hat{\mathbf{n}}_i$ and using the above relation, we obtain the *Caley-Hamilton equation*.

$$\mathbf{A}^3 - I_1\mathbf{A}^2 + I_2\mathbf{A} - I_3 = 0.$$

Spectral decomposition of a tensor

Any symmetric tensor \mathbf{A} can be represented using its eigenvalue and eigenvectors as basis vectors $\{\hat{\mathbf{n}}_i\}$. A unit tensor in $\{\hat{\mathbf{n}}_i\}$ basis vectors is represented as,

$$\mathbf{I} = \hat{\mathbf{n}}_i \otimes \hat{\mathbf{n}}_i.$$

Tensor \mathbf{A} can be written as,

$$\mathbf{A} = \mathbf{A}\mathbf{I} = (\mathbf{A}\hat{\mathbf{n}}_i) \otimes \hat{\mathbf{n}}_i = \sum_{i=1}^3 \lambda_i \hat{\mathbf{n}}_i \otimes \hat{\mathbf{n}}_i$$

Now, let us find the $(ij)^{\text{th}}$ component of tensor \mathbf{A} relative to the a basis of eigenvectors $\{\hat{\mathbf{n}}_i\}$,

$$A_{ij} = \hat{\mathbf{n}}_i \cdot (\mathbf{A}\hat{\mathbf{n}}_j) = \hat{\mathbf{n}}_i \cdot (\lambda_j \hat{\mathbf{n}}_j) = \lambda_j \delta_{ij}. (j = 1, 2, 3; \text{no summation})$$

In matrix form,

$$[\mathbf{A}] = \begin{bmatrix} \lambda_1 & 0 & 0 \\ 0 & \lambda_2 & 0 \\ 0 & 0 & \lambda_3 \end{bmatrix}.$$

Example

For the given tensor \mathbf{A} , find the eigenvalues and corresponding eigenvectors.

$$\mathbf{A} = \alpha(\mathbf{I} - \mathbf{e}_1 \otimes \mathbf{e}_1) + \beta(\mathbf{e}_1 \otimes \mathbf{e}_2 + \mathbf{e}_2 \otimes \mathbf{e}_1)$$

First let us find the invariants of the tensor,

$$I_1 = 2a, \quad I_2 = a^2 - \beta^2, \quad I_3 = -a\beta^2$$

Characteristic equation now become,

$$\lambda^3 - 2a\lambda^2 + (a^2 - \beta^2)\lambda - (-a\beta^2) = 0.$$

Roots of the equation are,

$$\lambda_1 = \alpha, \quad \lambda_2 = \frac{\sqrt{4\beta^2 + \alpha^2} + \alpha}{2}, \quad \lambda_3 = \frac{-\sqrt{4\beta^2 + \alpha^2} - \alpha}{2}.$$

To find the eigenvectors we use the following relation and substitute the value of λ in it,

$$\mathbf{A}\hat{\mathbf{n}}_i = \lambda_i\hat{\mathbf{n}}_i, \quad (i = 1, 2, 3; \text{ no summation})$$

Substituting $\lambda_1 = a$ in equation we get following set of equations,

$$\begin{aligned} -\alpha \hat{n}_1^1 + \beta \hat{n}_1^2 &= 0 \\ \beta \hat{n}_1^1 &= 0 \end{aligned}$$

From above two equations we get, $\hat{n}_1^1 = \hat{n}_1^2 = 0$.

To determine the third component, we have another equation as,

$$(\hat{n}_1^1)^2 + (\hat{n}_1^2)^2 + (\hat{n}_1^3)^2 = 1,$$

which gives us $\hat{n}_1^3 = 1$. Thus, we have the eigenvector corresponding to eigenvalue $\lambda_1 = a$ is $\hat{\mathbf{n}}_1 = \{0, 0, 1\}$.

Similarly, eigenvectors $\hat{\mathbf{n}}_2$ and $\hat{\mathbf{n}}_3$ corresponding to λ_2 and λ_3 can also be determined.

Problem Set

Problem 1: Given that $T_{ij} = 2\mu E_{ij} + \delta_{ij} E_{kk}$. Find $T_{mn} E_{mn}$.

Problem 2: Show that (i) $e_{ijk} e_{ijk} = 6$, (ii) $e_{ijp} e_{ijq} = 2\delta_{pq}$.

Problem 3: Using the properties of ∇ operator, prove that

(i) $\nabla \cdot (\mathbf{A}^T \mathbf{u}) = \nabla \cdot \mathbf{A} \cdot \mathbf{u} + \mathbf{A} : \nabla \mathbf{u}$

(ii) $\nabla (\phi \mathbf{u}) = \mathbf{u} \otimes \nabla \phi + \phi \nabla \mathbf{u}$

Problem 4: Show that

$$\int_S \mathbf{u} \cdot \mathbf{A} \mathbf{n} \, dS = \int_V \nabla \cdot \mathbf{A}^T \mathbf{u} \, dV$$

Problem 5: The most general form of a fourth-order isotropic tensor can be expressed by

$$\alpha \delta_{ij} \delta_{kl} + \beta \delta_{ik} \delta_{jl} + \gamma \delta_{il} \delta_{jk}$$

where α , β , and γ are arbitrary constants. Verify that this form remains the same under the general transformation.

Problem Set

Problem 6:

Provided that \mathbf{T} is symmetric, show that $\text{tr}(\nabla \times \mathbf{T}) = 0$.

Problem 7:

Let a new right-handed Cartesian coordinate system be represented by the set $\{\bar{\mathbf{e}}_i\}$ of basis vectors with transformation law, $\bar{\mathbf{e}}_2 = -\sin \theta \mathbf{e}_1 + \cos \theta \mathbf{e}_2$ and $\bar{\mathbf{e}}_3 = \mathbf{e}_3$.

The origin of the new coordinate system coincides with the old origin.

(a) Find $\bar{\mathbf{e}}_1$ in terms of the old set $\{\mathbf{e}_i\}$ of basis vectors.

(b) Find the orthogonal matrix $[\mathbf{Q}]$ and express the new coordinates in terms of the old one.

(c) Express the vector $\mathbf{u} = -6\mathbf{e}_1 - 3\mathbf{e}_2 + \mathbf{e}_3$ in terms of the new set $\{\bar{\mathbf{e}}_i\}$ of basis vectors.

Problem 8:

Proof that $\nabla \times (\mathbf{u} \times \mathbf{v}) = u(\nabla \cdot \mathbf{v}) - v(\nabla \cdot \mathbf{u}) + (\nabla u)\mathbf{v} - (\nabla v)\mathbf{u}$.