

# ME232: Dynamics

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Room # 106

The shape of the path followed by  $m$  may be obtained by solving the first of (55), and (56) and eliminating the time  $t$ .

To that end substitute  $r = 1/u$ . Thus,  $\dot{r} = -(1/u^2)\dot{u}$ , which from (56) becomes  $\dot{r} = -h(\dot{u}/\dot{\theta})$  or  $\dot{r} = -h(du/d\theta)$ . The second time derivative is  $\ddot{r} = -h(d^2u/d\theta^2)\dot{\theta}$  , which by combining with (56), become  $\ddot{r} = -h^2u^2(d^2u/d\theta^2)$ .

Substitution into the first of (55) now gives,

$$-Gm_0u^2 = -h^2u^2\frac{d^2u}{d\theta^2} - \frac{1}{u}h^2u^4, \quad \text{or} \quad \frac{d^2u}{d\theta^2} + u = \frac{Gm_0}{h^2}, \quad \text{.....(57)}$$

which is a nonhomogeneous linear differential equation.

The solution of (57) may be verified by direct substitution and is

$$u = \frac{1}{r} = C \cos(\theta + \delta) + \frac{Gm_0}{h^2}, \quad \text{.....(58)}$$

where  $C$  and  $\delta$  are the two integration constants. The phase angle  $\delta$  may be eliminated by choosing the  $x$ -axis so that  $r$  is a minimum when  $\theta = 0$ . Thus,

$$\frac{1}{r} = C \cos \theta + \frac{Gm_0}{h^2}. \quad \text{.....(59)} \tag{77}$$

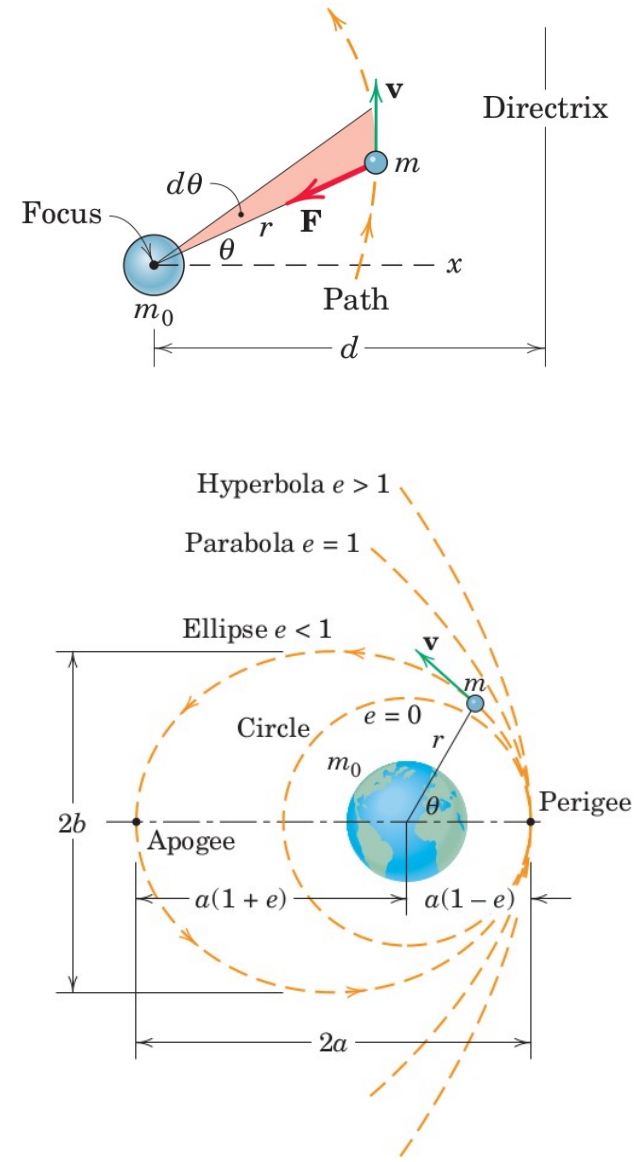
The interpretation of (59) requires a knowledge of the equations for conic sections. A conic section is formed by the locus of a point which moves so that the ratio  $e$  of its distance from a point (focus) to a line (directrix) is constant. Thus,  $e = r / (d - r \cos \theta)$ , which may be rewritten as

$$\frac{1}{r} = \frac{1}{d} \cos \theta + \frac{1}{ed}. \quad \dots\dots\dots(60)$$

It can be observed that (59) and (60) have same forms. Thus, we see that the motion of  $m$  is along a conic section with  $d = 1/C$  and  $ed = h^2/(Gm_0)$ , or

$$e = h^2 C / Gm_0. \quad \dots\dots\dots(61)$$

The three cases to be investigated correspond to  $e < 1$  (ellipse),  $e = 1$  (parabola), and  $e > 1$  (hyperbola). The trajectory for each of these cases is shown in figure.



### Case 1: ellipse ( $e < 1$ ):

From (60) we deduce that  $r$  is a minimum when  $\theta = 0$  and is a maximum when  $\theta = \pi$ . Thus,

$$2a = r_{\min} + r_{\max} = \frac{ed}{1+e} + \frac{ed}{1-e} \quad \text{or} \quad a = \frac{ed}{1-e^2}.$$

With the distance  $d$  expressed in terms of  $a$ , (60) and the maximum and minimum values of  $r$  may be written as

$$\frac{1}{r} = \frac{1 + e \cos \theta}{a(1 - e^2)}, \quad r_{\min} = a(1 - e), \quad r_{\max} = a(1 + e). \quad \dots\dots\dots(62)$$

Equation (62) is an expression of **Kepler's first law**, which says that **the planets move in elliptical orbits around the sun as a focus**. The period  $\tau$  for the elliptical orbit is the total area  $A$  of the ellipse divided by the constant rate  $\dot{A}$  at which the area is swept through.

Thus, from (56),

$$\tau = \frac{A}{\dot{A}} = \frac{\pi ab}{\frac{1}{2}r^2\dot{\theta}} \quad \text{or} \quad \tau = \frac{2\pi ab}{h}.$$

Eliminate  $\dot{\theta}$  or  $h$  in the expression for  $\tau$  by substituting (61), the identity  $d = 1/C$ , the geometric relationships  $a = ed/(1-e^2)$  and  $b = a(1-e^2)^{1/2}$  for the ellipse, and the equivalence  $Gm_0 = gR^2$ . The result after simplification is

$$\tau = 2\pi \frac{a^{3/2}}{R\sqrt{g}}. \quad \text{.....(63)}$$

In this equation note that  $R$  is the mean radius of the central attracting body and  $g$  is the absolute value of the acceleration due to gravity at the surface of the attracting body. Equation (63) expresses **Kepler's third law** of planetary motion which states that **the square of the period of motion is proportional to the cube of the semi-major axis of the orbit**.

### Case 2: parabola ( $e = 1$ ):

Equations (60) and (61) become,

$$\frac{1}{r} = \frac{1}{d} (1 + \cos \theta) \quad \text{and} \quad h^2 C = G m_0. \quad \dots\dots\dots(64)$$

The radius vector becomes infinite as  $\theta$  approaches  $\pi$ , so the dimension  $a$  is infinite.

### Case 3: hyperbola ( $e > 1$ ):

From (60) we see that the radial distance  $r$  becomes infinite for the two values of the polar angle  $\theta_1$  and  $-\theta_1$  defined by  $\cos \theta_1 = -1/e$ . Only branch I corresponding to  $-\theta_1 < \theta < \theta_1$  represents a physically possible motion.

