

Something something Higgs boson

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Abstract

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1 Introduction

After the theoretical discovery of the Higgs mechanism and the associated particle, the Higgs boson, experimental confirmation of the prediction became a much sought-after prize in High Energy Physics. Although the theory was more or less accepted by the theoretical community since its inception in the 1960s, experimental evidence was required for verification.

Once CERN's Large Hadron Collider (LHC) emerged on the scene, with the necessary power to probe higher energy regimes of collisions, vast amounts of data started to be collected. In order to verify the theoretical predictions of particle physics, this data needed to be analysed carefully to try and distinguish signal of as-yet unknown particle decays vs effectively “proven” existing decays. The accepted threshold for a discovery in particle physics is the “gold-standard” 5-sigma rule. As a result, any announcement of a new experimental discovery (such as the Higgs boson) was required to meet this threshold.

For the Higgs boson in particular, various decay mechanisms were proposed that could be used to demonstrate its existence by comparing the rate of the by-products of those decays vs background mechanisms (that do not involve the Higgs). In 2013, when the Higgs was experimentally discovered by CERN, the evidence was provided by *bosonic* decays from the Higgs to the following pairs, $\gamma\gamma$, WW and ZZ .

The goal here is to instead examine the coupling of the Higgs to fermions, to verify that their mass can likewise be explained by the Higgs mechanism. Specifically, here the aim is to analyse the decay $H \rightarrow \tau\tau$; i.e. the decay of the Higgs to a pair of tau-flavour leptons. (The other candidates, electrons and muons fall outside the energy range of the LHC due to their lighter mass).

2 Theoretical Background

2.1 Physics

Not sure how much we actually need to include.

2.2 Statistics

Likewise.

3 Method

3.1 Missing Data

Early examinations of the dataset revealed the presence of a significant quantity of missing data. Elements of the data matrix recorded as missing were labelled by a value of -999. Before embarking on any model building or even feature engineering we analysed the missing values to try and ascertain whether there was any pattern to their locations. From the definitions of covariates in the dataset it was possible to infer potential causes of missing data. In particular, the estimated mass of the Higgs (*DER_mass_MMC*) “(may be undefined if the topology of the event is too far from the expected topology)”, indicates that this feature may be expected to be labelled as missing. Similarly, a set of the features mention a dependence on the number of jets measured in the interaction:

DER_deltaeta_jet_jet undefined if $PRI_jet_num \leq 1$

DER_mass_jet_jet undefined if $PRI_jet_num \leq 1$

DER_prodelta_jet_jet undefined if $PRI_jet_num \leq 1$

PRI_jet_subleading_pt undefined if $PRI_jet_num \leq 1$

PRI_jet_subleading_phi undefined if $PRI_jet_num \leq 1$

PRI_jet_subleading_eta undefined if $PRI_jet_num \leq 1$

DER_lep_eta_entrality undefined if $PRI_jet_num \leq 1$

PRI_jet_leading_pt undefined if $PRI_jet_num = 0$

PRI_jet_leading_eta undefined if $PRI_jet_num = 0$

PRI_jet_leading_phi undefined if $PRI_jet_num = 0$

So we can see that the missing data is explained by the combination of an unexpected topology and the number of jets observed. This implies that the physical data generating process is different in each of these regimes and that we ought to treat them differently in accordance. As such we chose to attempt to build a separate model for each regime.

3.2 Performance Metrics

3.2.1 AMS

The AMS metric is the Approximate Median discovery Significance; an approximation of the *significance* defined by

$$\begin{aligned} Z &= \Phi^{-1}(1 - p) \\ &= \sqrt{q_0} \\ &= \sqrt{2 \left(n \ln \frac{n}{\mu_b} - n + \mu_b \right)} \end{aligned}$$

where n is the (unknown) number of events in some search region \mathcal{G} , Φ^{-1} is the inverse Normal CDF, q_0 is a test statistic given by Wilks' Theorem and μ_b is the (unknown) expected number of background events. We replace $n \rightarrow s + b$ and $\mu_b \rightarrow b$ with s, b the estimator of the expected number of signal and background events respectively.

See [1] for more details.

Since this is the formal objective we aim to maximise, it might make sense to try and optimise it directly. A little analysis however reveals that the function is non-convex and therefore a poor choice. See A for the full calculation.

3.2.2 ROC and AUC

A ROC (Receiver Operating Characteristic) curve plots the true positive rate against the false positive rate for a binary classifier with a varying decision threshold. By doing this we can show the ability of the classifier to discriminate between the two classes and attain a metric for the performance that we can use to compare models. The worst possible model is given by a line at $y = x$ and any curve above this line is an improvement.

The AUC (Area Under the [ROC] Curve) provides a compression of the assessment the ROC curve provides by integrating the ROC curve and giving a value in $[0, 1]$ that represents an overall view of the discriminatory power of the model. A value of 0.5 corresponds to the worst model ($y = x$) whilst 1.0 would represent a perfect model.

Note that for our problem, the actual ratio of signal:background events is very small, which would cause a problem in interpreting the AUC reliably. In the actual dataset, this problem is mitigated by providing comparable numbers of rows with signal labels as to rows with background labels, each with an associated weight, such that the overall weighted events of signal and background match the actual observed data.

3.3 Logistic Regression

3.4 SVM

3.5 Feature Engineering

In order to improve the performance of our model, we attempted to carry out some feature engineering to extract as much information as possible from the dataset.

3.5.1 Redundancy

By considering the basic physics of the beam it is possible to see that there is some redundancy amongst the (primitive) features in the dataset. We can exploit this fact and reduce our initial feature space slightly, by transforming the redundant features into a set of new derived features that contain the same information. More specifically, the redundancy comes from the consideration that the physical phenomena should be invariant to certain symmetries; a rotation about the beam (z) axis and reflections in the xy -plane.

3.5.2 Higher-Order Effects

RBF Centroids To attempt to capture generic higher order interaction-type effects, we implemented a set of RBF centroid features. The implementation works as

Algorithm 1: Augment covariate matrix with RBF centroid features

follows

1. Generate heuristic estimate of RBF hyperparameter s .
2. Select n_c points from \mathbf{X}_{train} , $\mathbf{X}_c = \{\mathbf{x}_i^{train}, \dots, \mathbf{x}_{n_c}^{train}\}$.
3. **for** $i=1:n_c$ **do**
 - Calculate $\mathbf{k}_{RBF}(\mathbf{x}_i^{train}, \mathbf{X}_{train}; s)$
 - $\mathbf{X}_{train} \leftarrow \mathbf{X}_{train} \oplus \mathbf{k}_{RBF}(\mathbf{x}_i^{train}, \mathbf{X}_{train}; s)$
4. Fit model.
5. **for** $i=1:n_c$ **do**
 - Calculate $\mathbf{k}_{RBF}(\mathbf{x}_i^{train}, \mathbf{X}_{test}; s)$
 - $\mathbf{X}_{test} \leftarrow \mathbf{X}_{test} \oplus \mathbf{k}_{RBF}(\mathbf{x}_i^{train}, \mathbf{X}_{test}; s)$

3.5.3 Transformations

Polynomial

3.5.4 Interactions

4 Results

4.1 Feature Engineering

4.2 Model Selection

4.3 Predictions

To actually calculate AMS scores, we need to convert out probabilistic outputs into binary labels which means we need to pick a decision rule, i.e. a threshold over which we assign a label 1 (or s) vs 0 (or b).

5 Conclusion

Appendices

A AMS

Is the AMS metric convex?

$$\text{AMS} = \sqrt{2 \left((s+b) \log\left(1 + \frac{s}{b}\right) - s \right)}$$

In the documentation, the implication is that in general, we expect $b \gg s$ giving the approximate AMS as $\text{AMS} \sim \frac{s}{\sqrt{b}}(1 + \mathcal{O}(\frac{s}{b}))$.

If we directly check the derivatives of the approximate AMS in the large b regime:

$$\begin{aligned} \partial_s^2 \frac{s}{\sqrt{b}} &= 0 \\ \partial_b^2 \frac{s}{\sqrt{b}} &= \frac{3}{4} \frac{s}{b^{5/2}} > 0 \\ \partial_{bs}^2 &= -\frac{1}{2b^{3/2}} < 0 \end{aligned}$$

which implies non-convexity (and non-concavity).

To check more carefully though, calculate the terms of the Hessian of the original and then apply the approximation:

$$\begin{aligned} \partial_s \text{AMS} &= \text{AMS}^{-1} \log\left(1 + \frac{s}{b}\right) \\ \partial_b \text{AMS} &= \text{AMS}^{-1} \left(\log\left(1 + \frac{s}{b}\right) - \frac{s}{b} \right) \\ \partial_s^2 \text{AMS} &= \text{AMS}^{-1} \left(\frac{1}{b+s} - \text{AMS}^{-2} \log\left(1 + \frac{s}{b}\right)^2 \right) \\ \partial_b^2 \text{AMS} &= \text{AMS}^{-1} \left(\frac{s^2}{b^2(s+b)} - \text{AMS}^{-2} \left(\log\left(1 + \frac{s}{b}\right) - \frac{s}{b} \right)^2 \right) \\ \partial_{sb}^2 \text{AMS} &= -\text{AMS}^{-1} \left(\frac{s}{b(s+b)} - \text{AMS}^{-2} \log\left(1 + \frac{s}{b}\right) \left(\log\left(1 + \frac{s}{b}\right) - \frac{s}{b} \right) \right) \end{aligned}$$

and then verify whether any of them are ever < 0 .

Using the AMS approximation, and $\log\left(1 + \frac{s}{b}\right) \sim \frac{s}{b} - \frac{1}{2} \left(\frac{s}{b}\right)^2$ we get the following

$$\begin{aligned} \partial_s^2 \text{AMS} &\sim \frac{\sqrt{b}}{s} \left[\frac{1}{(b+s)} - \frac{b}{s^2} \left(\frac{s^2}{b^2} - \frac{s^3}{b^3} \right) \right] \\ &= \frac{\sqrt{b}}{s(s+b)} - \frac{1}{s\sqrt{b}} + \frac{1}{b^{3/2}} \\ &\gtrsim \epsilon \end{aligned}$$

Since in this regime of large b this term is negative, the Hessian cannot be positive definite and hence the metric is non-convex. If the Hessian is *concave* though, we can simply optimize $-\text{AMS}$, so we need to check the other terms.

$$\begin{aligned}
\partial_b^2 \text{AMS} &\sim \frac{\sqrt{b}}{s} \left[\frac{s^2}{b^2(s+b)} - \frac{b}{s^2} \left(\frac{1}{4} \frac{s^4}{b^4} \right) \right] \\
&= \frac{s}{b^{3/2}(s+b)} - \frac{1}{4} \frac{s}{b^{5/2}} \\
&\gtrsim \epsilon
\end{aligned}$$

$$\begin{aligned}
\partial_{bs}^2 \text{AMS} &\sim -\frac{\sqrt{b}}{s} \left[\frac{s}{b(s+b)} + \frac{b}{s^2} \left(\frac{s}{b} - \frac{1}{2} \frac{s}{b^2} \right) \left(-\frac{1}{2} \frac{s^2}{b^2} \right) \right] \\
&= \frac{1}{2b^{3/2}} \left(1 - \frac{1}{2} s \right) - \frac{1}{\sqrt{b}(s+b)} \\
&< 0
\end{aligned}$$

where we take ϵ to be some suitably small positive number.

If we decide to take some tolerance, $\epsilon \sim 10^{-6}$ at which to ignore higher order terms in $\frac{s}{b}$ we can check specifically that the claims on the Hessian above hold in that regime.

So, let's assume we can ignore terms of $O(\frac{s}{b})^3$ then for $\epsilon \sim 10^{-6}$ a ratio of $\frac{s}{b} \sim 10^{-2}$ would suffice; so to test this, let us choose $s = 1$ and $b = 100$ for simplicity. With this, we find that $\partial_s^2 \text{AMS} \sim +10^{-6}$, $\partial_b^2 \text{AMS} \sim +10^{-6}$ and $\partial_{bs}^2 \text{AMS} \sim -10^{-4}$ which satisfies the claims above. Finally, to verify that this is a reasonable assertion, if we check the ratio of $s : b$ in the data (using the weight vector) we find a ratio of approximately 10^{-3} , which means our assumption was conservative.

B Modelling

B.1 Newton's Method for Logistic Regression

From the conditional distribution under the logistic regression model for a single point $\mathbf{x} \in R^d$, $y \in \{+1, -1\}$ and coefficients $\beta \in R^d$.

$$\begin{aligned}
p(y | \mathbf{x}, \beta) &= \sigma(f(\mathbf{x}; \beta) \cdot y) \\
&= (1 + \exp(y\langle \mathbf{x}, \beta \rangle))^{-1}
\end{aligned}$$

We want to minimise the negative loglikelihood:

$$\begin{aligned}
l(\mathbf{x}, y, \beta) &= \log(1 + e^{y\langle \mathbf{x}, \beta \rangle}) \\
\partial_{\beta_a} l &= \frac{yx_a e^{y\langle \mathbf{x}, \beta \rangle}}{1 + e^{y\langle \mathbf{x}, \beta \rangle}} \\
\partial_{\beta_a \beta_b}^2 l &= \frac{y^2 x_a x_b e^{y\langle \mathbf{x}, \beta \rangle}}{(1 + e^{y\langle \mathbf{x}, \beta \rangle})^2} \left(1 + e^{y\langle \mathbf{x}, \beta \rangle} - e^{y\langle \mathbf{x}, \beta \rangle} \right) \\
&= \frac{x_a x_b e^{y\langle \mathbf{x}, \beta \rangle}}{(1 + e^{y\langle \mathbf{x}, \beta \rangle})^2} \\
&= x_a x_b w(x)
\end{aligned}$$

where in the last line we define a weight function $w(x) = \frac{e^{-x}}{(1+e^{-x})^2}$ (and the element-wise vector version, $\mathbf{w}(\mathbf{x}) = \frac{e^{-\mathbf{x}}}{(1+e^{-\mathbf{x}})^2}$).

To implement Newton's method we want to try and simplify the expressions. First define the logistic and logit functions as

$$\begin{aligned}\gamma(x) &= (1 + e^{-x})^{-1} \\ \gamma^{-1}(p) &= \log\left(\frac{p}{1-p}\right)\end{aligned}$$

and then write out the derivatives in terms of the logistic function:

$$\begin{aligned}\partial_{\beta_a} l(y = +1) &= x_a \gamma(\langle \mathbf{x}, \boldsymbol{\beta} \rangle) \\ \partial_{\beta_a} l(y = -1) &= -x_a (1 - \gamma(\langle \mathbf{x}, \boldsymbol{\beta} \rangle))\end{aligned}$$

If we then map $y \in \{1, -1\}$ to $y' \in \{0, 1\}$ we can condense these into a single expression

$$\partial_{\beta_a} l(y) = -x_a (y - \gamma(\langle \mathbf{x}, \boldsymbol{\beta} \rangle))$$

and rewrite the loss as $l(y') = \log(1 + e^{(-1)^{y'} \langle \mathbf{x}, \boldsymbol{\beta} \rangle})$.

Now for our entire dataset D , assumed to be IID, we have a total loss, gradient and Hessian

$$\begin{aligned}L &= \sum_{i \in D} l_i = \sum_{i \in D} \log(1 + e^{y_i \langle \mathbf{x}_i, \boldsymbol{\beta} \rangle}) \\ \partial_{\boldsymbol{\beta}} L &= - \sum_{i \in D} \mathbf{x}_i (y_i - \gamma(\langle \mathbf{x}_i, \boldsymbol{\beta} \rangle)) \\ &= -\mathbf{X}^T (\mathbf{y} - \gamma(\mathbf{X}\boldsymbol{\beta})) \\ \mathbf{H} &= \sum_{i \in D} \mathbf{x}_i \mathbf{x}_i^T w(\mathbf{x}_i) \\ &= \mathbf{X}^T \mathbf{W} \mathbf{X}\end{aligned}$$

with $\mathbf{X} = \begin{pmatrix} \mathbf{x}_1^T \\ \vdots \\ \mathbf{x}_n^T \end{pmatrix} \in R^{n \times d}$, $\mathbf{W} = \text{diag}(\mathbf{w}(\mathbf{x})) \in R^{n \times n}$.

B.1.1 Regularisation

We can add constraints to our minimisation problem, to help constrain the coefficients and effectively implement L_2 regularisation by formulating the following convex optimisation problem.

$$\begin{aligned}\min. \quad & L \\ \text{subject to} \quad & \|\boldsymbol{\beta}\|^2 \leq c\end{aligned}$$

From this we can construct the Lagrangian $\mathcal{L} = L + \lambda(\|\boldsymbol{\beta}\|^2 - c)$ and hence augment the gradient and Hessian we calculated above as

$$\begin{aligned}\partial_{\boldsymbol{\beta}} \mathcal{L} &= \partial_{\boldsymbol{\beta}} L + 2\lambda \boldsymbol{\beta} \\ \tilde{\mathbf{H}} &= \mathbf{H} + 2\lambda \mathbf{I}\end{aligned}$$

Since we have a free parameter c that determines the optimal value $\lambda^*(c)$ of our dual variable, we can equivalently ignore c and treat λ as a parameter to be tuned.

B.1.2 Newton’s method

Algorithm 2: Newton’s method

1. Initialise parameters; $\beta_0, \lambda_0, \epsilon, L_0 = L(\beta_0), m$
 2. **for** $i=1:m$ **do**
 - Compute $\partial_{\beta}L$ and \mathbf{H}
 - Calculate Newton step $\Delta\beta = -\mathbf{H}^{-1}\partial_{\beta}L$
 - Compute new loss L
 - Test stopping criterion: $|L - L_0| < \epsilon$
 3. Return $\hat{\beta}$.
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References

- [1] Claire Adam-Bourdarios, Glen Cowan, Cécile Germain, Isabelle Guyon, Balázs Kégl, and David Rousseau. The Higgs boson machine learning challenge. In Glen Cowan, Cécile Germain, Isabelle Guyon, Balázs Kégl, and David Rousseau, editors, *Proceedings of the NIPS 2014 Workshop on High-energy Physics and Machine Learning*, volume 42 of *Proceedings of Machine Learning Research*, pages 19–55, Montreal, Canada, 13 Dec 2015. PMLR.