

Viscous stresses

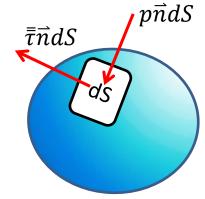
Complete system of Navier-Stokes equation



Conservation laws

•
$$\frac{\partial}{\partial t} \int_{\Omega} \rho \, d\Omega + \int_{\partial \Omega} \rho (v \cdot \vec{n}) \, dS = 0$$

•
$$\frac{\partial}{\partial t} \int_{\Omega} \rho \vec{v} \, d\Omega + \oint_{\partial \Omega} \rho \vec{v} (\vec{v} \cdot \vec{n}) dS$$



$$= \int_{\Omega} \rho \vec{f_e} d\Omega - \oint_{\partial \Omega} p \cdot \vec{n} dS + \oint_{\partial \Omega} (\bar{\tau} \cdot \vec{n}) dS$$
(Moment equation)



Conservation laws

•
$$\frac{\partial}{\partial t} \int_{\Omega} \rho E \, d\Omega + \oint_{\partial \Omega} \rho e(\vec{v} \cdot \vec{n}) dS =$$

$$\oint_{\partial \Omega} k (\nabla T \cdot \vec{n}) dS + \int_{\Omega} (\rho \vec{f_e} \cdot \vec{v}) d\Omega +$$

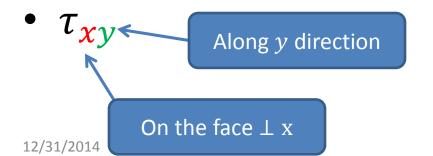
$$\oint_{\partial \Omega} (\bar{\tau} \cdot \vec{v}) \cdot \vec{n} dS \qquad \text{(Energy equation)}$$

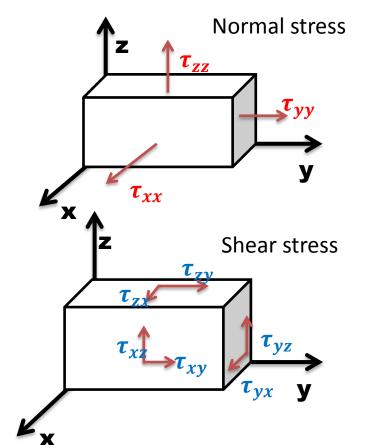


Viscous stress

Stress Tensor

$$\bullet \ \ \bar{\bar{\tau}} = \begin{bmatrix} \tau_{\chi\chi} & \tau_{\chi y} & \tau_{\chi z} \\ \tau_{y\chi} & \tau_{yy} & \tau_{yz} \\ \tau_{z\chi} & \tau_{zy} & \tau_{zz} \end{bmatrix}$$







Viscous stress

Proportional to the velocity gradient:

Newtonian fluid (water, air, thin oil)





• Otherwise:

Non-Newtonian fluid (butter, ketchup)







6

Viscosity

In the case of a Newtonian fluid

$$\tau_{xx} = \lambda \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) + 2\mu \frac{\partial u}{\partial x}$$
$$\tau_{xy} = \tau_{yx} = \mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right)$$

 λ is second viscosity coefficient μ is dynamic viscosity $\nu = \mu/\rho$ kinematic viscosity

Stokes hypothesis: $\lambda = -2/3\mu$

$$ar{ar{ au}} = egin{bmatrix} au_{\chi\chi} & au_{\chi y} & au_{\chi Z} \ au_{\chi\chi} & au_{\chi y} & au_{\chi Z} \ au_{\chi\chi} & au_{\chi y} & au_{\chi Z} \end{bmatrix}$$



- In the case of a Newtonian fluid,
 conservation law gives Navier-Stokes equations
- Take the momentum equation as an example



•
$$\frac{\partial}{\partial t} \int_{\Omega} \rho \vec{v} \, d\Omega + \int_{\partial \Omega} \rho \vec{v} (\vec{v} \cdot \vec{n}) + p \vec{n} dS$$

$$= \int_{\partial\Omega} \overline{\overline{t}} \cdot \overrightarrow{n} \, dS + \int_{\Omega} \rho \overrightarrow{f_e} \, d\Omega$$

Using divergence theorem:

•
$$\frac{\partial}{\partial t} \int_{\Omega} \rho \vec{v} \, d\Omega + \int_{\Omega} \rho \vec{v} (\nabla \cdot \vec{v}) + \nabla p d\Omega$$

$$= \int_{\Omega} \nabla \cdot \overline{t} d\Omega + \int_{\Omega} \rho \overrightarrow{f_e} d\Omega$$



•
$$\frac{\partial}{\partial t}(\rho\vec{v}) + \rho\vec{v}(\nabla \cdot \vec{v}) + \nabla p = \nabla \cdot \overline{\overline{t}} + \rho \overrightarrow{f_e}$$

•
$$\nabla \cdot \bar{t} = \nabla \cdot \begin{bmatrix} \tau_{xx} & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \tau_{yy} & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \tau_{zz} \end{bmatrix} = \begin{bmatrix} \frac{\partial}{\partial x} (\tau_{xx}) + \frac{\partial}{\partial x} (\tau_{xy}) + \frac{\partial}{\partial z} (\tau_{xz}) \\ \frac{\partial}{\partial x} (\tau_{yx}) + \frac{\partial}{\partial x} (\tau_{yy}) + \frac{\partial}{\partial z} (\tau_{yz}) \\ \frac{\partial}{\partial x} (\tau_{zx}) + \frac{\partial}{\partial x} (\tau_{zy}) + \frac{\partial}{\partial z} (\tau_{zz}) \end{bmatrix}$$

- For incompressible fluid, we can further simplify
- Take the first row as an example



• Take the first row as an example

$$abla \cdot \overline{\overline{t}} =
abla \cdot egin{array}{cccc} au_{xx} & au_{xy} & au_{xz} \ au_{yx} & au_{yy} & au_{yz} \ au_{zx} & au_{zy} & au_{zz} \ \end{array}$$

•
$$\frac{\partial}{\partial x}(\tau_{xx}) + \frac{\partial}{\partial y}(\tau_{xy}) + \frac{\partial}{\partial z}(\tau_{xz})$$

= $\frac{\partial}{\partial x}\left(2\mu\frac{\partial u}{\partial x}\right) + \frac{\partial}{\partial y}\left[\mu\left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}\right)\right] + \frac{\partial}{\partial z}\left[\mu\left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x}\right)\right]$
= $2\mu\frac{\partial^{2}u}{\partial x^{2}} + \mu\frac{\partial^{2}u}{\partial y^{2}} + \mu\frac{\partial^{2}v}{\partial xy} + \mu\frac{\partial^{2}u}{\partial z^{2}} + \mu\frac{\partial^{2}w}{\partial xz}$
= $\mu\left(\frac{\partial^{2}u}{\partial x^{2}} + \frac{\partial^{2}u}{\partial y^{2}} + \frac{\partial^{2}u}{\partial z^{2}}\right) + \mu\frac{\partial}{\partial x}\left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z}\right)$



•
$$\frac{\partial}{\partial t}(\rho \vec{v}) + \rho \vec{v}(\nabla \cdot \vec{v}) + \nabla p$$

= $\nabla \cdot \bar{t} + \rho \vec{f}_{e}$
= $\mu \Delta \vec{v} + \rho \vec{f}_{e}$

Assumption:

- 1. Viscosity is constant
- 2. Newtonian fluid
- 3. Simplified mass continuity equation

$$\nabla \cdot \vec{v} = 0$$



Reference

- J. Blazek. Computational Fluid Dynamics: Principles and Applications: 2001
- X. Li. Partial Differential Equations and Applications in Computational Fluid Dynamics: 2014
- Wikipidia. Derivation of the Navier–Stokes equations



- Five equations for five conservative variables ρ , ρu , ρv , ρw , ρE
- But they contain seven unknown variables:

 ρ , u, v, w, p, T, E

We have to supply two additional equations

We have to provide the viscosity coefficient μ and the thermal conductivity coefficient κ