

# Misner, Thorne and Wheeler's Gravitation Problems

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# Chapter 1

## Chapter 2

## Chapter 3

# The Electromagnetic Field

1. Derive equations:

$$|| F_{\beta}^{\alpha} || = \begin{vmatrix} 0 & E_x & E_y & E_z \\ E_x & 0 & B_z & -B_y \\ E_y & -B_z & 0 & B_x \\ E_z & B_y & -B_x & 0 \end{vmatrix} \quad (3.1)$$

and

$$|| F_{\alpha\beta} || = \begin{vmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & B_z & -B_y \\ E_y & -B_z & 0 & B_x \\ E_z & B_y & -B_x & 0 \end{vmatrix} \quad (3.2)$$

for the components of Faraday by comparing:

$$dp^{\alpha}/d\tau = eF_{\beta}^{\alpha}u^{\beta} \quad (3.3)$$

with

$$\frac{d\mathbf{p}}{d\tau} = \frac{1}{\sqrt{1-\mathbf{v}^2}} \frac{d\mathbf{p}}{dt} = \frac{e}{\sqrt{1-\mathbf{v}^2}} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) = e(u^0 \mathbf{E} + \mathbf{u} \times \mathbf{B}) \quad (3.4)$$

$$\frac{dp^0}{d\tau} = \frac{1}{\sqrt{1-\mathbf{v}^2}} \frac{dE}{dt} = \frac{1}{\sqrt{1-\mathbf{v}^2}} e \mathbf{E} \cdot \mathbf{v} = e \mathbf{E} \cdot \mathbf{u} \quad (3.5)$$

and by using definition:

$$F_{\alpha\beta} = \eta_{\alpha\gamma} F_{\beta}^{\gamma} \quad (3.6)$$

Consider equation 3.3 for the index  $\alpha = 0$ :

$$\frac{dp^0}{d\tau} = e[F_0^0 u^0 + F_1^0 u^1 + F_2^0 u^2 + F_3^0 u^3]$$

Equate this with 3.5:

$$e[F_0^0 u^0 + F_1^0 u^1 + F_2^0 u^2 + F_3^0 u^3] = e \mathbf{E} \cdot \mathbf{u} = e[E_1 u^1 + E_2 u^2 + E_3 u^3]$$

It is clear that:

$$\begin{aligned}
F_0^0 &= 0 \\
F_1^0 u^1 &= E_1 u^1 \Rightarrow F_1^0 = E_1 \\
F_2^0 u^2 &= E_2 u^2 \Rightarrow F_2^0 = E_2 \\
F_3^0 u^3 &= E_3 u^3 \Rightarrow F_3^0 = E_3
\end{aligned}$$

Now equating equation 3.4 with the remaining components of equation 3.3:

$$\begin{aligned}
\frac{dp^1}{d\tau} &= e[F_0^1 u^0 + F_1^1 u^1 + F_2^1 u^2 + F_3^1 u^3] = e[E_1 u^0 + B_3 u^2 - B_2 u^3] \\
\frac{dp^2}{d\tau} &= e[F_0^2 u^0 + F_1^2 u^1 + F_2^2 u^2 + F_3^2 u^3] = e[E_2 u^0 + B_1 u^3 - B_3 u^2] \\
\frac{dp^3}{d\tau} &= e[F_0^3 u^0 + F_1^3 u^1 + F_2^3 u^2 + F_3^3 u^3] = e[E_3 u^0 + B_2 u^1 - B_1 u^2]
\end{aligned}$$

and equating components as before:

$$\begin{aligned}
F_0^1 u^0 &= E_1 u^0 \Rightarrow F_0^1 = E_1 \\
F_1^1 u^1 &= 0 \Rightarrow F_1^1 = 0 \\
F_2^1 u^2 &= B_3 u^2 \Rightarrow F_2^1 = B_3 \\
F_3^1 u^3 &= -B_2 u^3 \Rightarrow F_3^1 = -B_2 \\
F_0^2 u^0 &= E_2 u^0 \Rightarrow F_0^2 = E_2 \\
F_1^2 u^1 &= -B_3 u^1 \Rightarrow F_1^2 = -B_3 \\
F_2^2 u^2 &= 0 \Rightarrow F_2^2 = 0 \\
F_3^2 u^3 &= B_1 u^3 \Rightarrow F_3^2 = B_1 \\
F_0^3 u^0 &= E_3 u^0 \Rightarrow F_0^3 = E_3 \\
F_1^3 u^1 &= B_2 u^1 \Rightarrow F_1^3 = B_2 \\
F_2^3 u^2 &= -B_1 u^2 \Rightarrow F_2^3 = -B_1 \\
F_3^3 u^3 &= 0 \Rightarrow F_3^3 = 0
\end{aligned}$$

Collecting all of these components in matrix form and relabelling indices with the following mapping:

$$\begin{aligned}
\alpha = 1 &\rightarrow x \\
\alpha = 2 &\rightarrow y \\
\alpha = 3 &\rightarrow z
\end{aligned}$$

gives equation 3.1:

$$\| F_\beta^\alpha \| = \begin{vmatrix} F_0^0 & F_1^0 & F_2^0 & F_3^0 \\ F_0^1 & F_1^1 & F_2^1 & F_3^1 \\ F_0^2 & F_1^2 & F_2^2 & F_3^2 \\ F_0^3 & F_1^3 & F_2^3 & F_3^3 \end{vmatrix} = \begin{vmatrix} 0 & E_x & E_y & E_z \\ E_x & 0 & B_z & -B_y \\ E_y & -B_z & 0 & B_x \\ E_z & B_y & -B_x & 0 \end{vmatrix}$$

Now equation 3.6 can be used to convert the mixed Faraday tensor to the fully covariant one. Remember that for all components  $\alpha \neq \beta$  the Minkowski metric is zero. Therefore the only non-zero components in the sums created by the summation convention in equation 3.6 are:

$$\begin{aligned}
F_{00} &= \eta_{00}F_0^0 \Rightarrow F_{00} = -F_0^0 \\
F_{01} &= \eta_{00}F_1^0 \Rightarrow F_{01} = -F_1^0 \\
F_{02} &= \eta_{00}F_2^0 \Rightarrow F_{02} = -F_2^0 \\
F_{03} &= \eta_{00}F_3^0 \Rightarrow F_{03} = -F_3^0 \\
F_{10} &= \eta_{11}F_0^1 \Rightarrow F_{10} = F_0^1 \\
F_{11} &= \eta_{11}F_1^1 \Rightarrow F_{11} = F_1^1 \\
F_{12} &= \eta_{11}F_2^1 \Rightarrow F_{12} = F_2^1 \\
F_{13} &= \eta_{11}F_3^1 \Rightarrow F_{13} = F_3^1 \\
F_{20} &= \eta_{22}F_0^2 \Rightarrow F_{02} = F_0^2 \\
F_{21} &= \eta_{22}F_1^2 \Rightarrow F_{12} = F_1^2 \\
F_{22} &= \eta_{22}F_2^2 \Rightarrow F_{22} = F_2^2 \\
F_{23} &= \eta_{22}F_3^2 \Rightarrow F_{23} = F_2^3 \\
F_{30} &= \eta_{33}F_0^3 \Rightarrow F_{30} = F_0^3 \\
F_{31} &= \eta_{33}F_1^3 \Rightarrow F_{31} = F_1^3 \\
F_{32} &= \eta_{33}F_2^3 \Rightarrow F_{32} = F_2^3 \\
F_{33} &= \eta_{33}F_3^3 \Rightarrow F_{33} = F_3^3
\end{aligned}$$

Collecting the components into matrix form recovers the fully covariant Faraday tensor:

$$|| F_{\alpha\beta} || = \begin{vmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & B_z & -B_y \\ E_y & -B_z & 0 & B_x \\ E_z & B_y & -B_x & 0 \end{vmatrix} \quad (3.7)$$

## 2. From the transformation laws for components of vectors and 1-forms, derive the transformation law:

$$S^{\mu'\nu'}_{\lambda'} = S^{\alpha\beta}_{\gamma} \Lambda^{\mu'}_{\alpha} \Lambda^{\nu'}_{\beta} \Lambda^{\gamma}_{\lambda'}$$

Consider the tensor  $\mathbf{S}$  of rank  $(2, 1)$ , in geometric notation the transformation between two sets basis vectors and 1-forms reads:

$$\mathbf{S}(\sigma, \rho, \nu) = \mathbf{S}(\sigma', \rho', \nu')$$

In component form this reads:

$$S^{\alpha\beta}_{\gamma} \sigma_{\alpha} \rho_{\beta} \nu^{\gamma} = S^{\mu'\nu'}_{\lambda'} \sigma_{\mu'} \rho_{\nu'} \nu^{\lambda'} \quad (3.8)$$

Using the Lorentz transformation laws to transform one basis into the other for  $\sigma$ ,  $\rho$ ,  $\nu$  gives:

$$\begin{aligned}\sigma_\alpha &= \Lambda_\alpha^{\mu'} \sigma_{\mu'} \\ \rho_\beta &= \Lambda_\beta^{\nu'} \rho_{\nu'} \\ \nu^\beta &= \Lambda_{\lambda'}^\gamma \nu^{\lambda'}\end{aligned}$$

and substituting these transformations into equation 3.8:

$$\begin{aligned}S^{\mu'\nu'}_{\lambda'} \sigma_{\mu'} \rho_{\nu'} \nu^{\lambda'} &= S^{\alpha\beta}_\gamma (\Lambda_\alpha^{\mu'} \sigma_{\mu'}) (\Lambda_\beta^{\nu'} \rho_{\nu'}) (\Lambda_{\lambda'}^\gamma \nu^{\lambda'}) \\ S^{\mu'\nu'}_{\lambda'} \sigma_{\mu'} \rho_{\nu'} \nu^{\lambda'} &= S^{\alpha\beta}_\gamma \Lambda_\alpha^{\mu'} \Lambda_\beta^{\nu'} \Lambda_{\lambda'}^\gamma \sigma_{\mu'} \rho_{\nu'} \nu^{\lambda'}\end{aligned}$$

Equating the components gives the desired transformation law:

$$S^{\mu'\nu'}_{\lambda'} = S^{\alpha\beta}_\gamma \Lambda_\alpha^{\mu'} \Lambda_\beta^{\nu'} \Lambda_{\lambda'}^\gamma$$

**3. Raising and lowering indices. Derive:**

$$S^\alpha_{\beta\gamma} = \eta_{\beta\mu} S^{\alpha\mu}_\gamma \quad (3.9)$$

**and:**

$$S^{\alpha\mu}_\gamma = \eta^{\mu\beta} S^\alpha_{\beta\gamma} \quad (3.10)$$

**from:**

4.

5.

6.

**7. Maxwell's Equations. Show, by explicit examination of components, that the geometric laws**

$$F_{\alpha\beta,\gamma} + F_{\beta\gamma,\alpha} + F_{\gamma\alpha,\beta} = 0 \quad (3.11)$$

$$F^{\alpha\beta}_{,\beta} = 4\pi J^\alpha \quad (3.12)$$

**do reduce to Maxwell's equations**

$$\nabla \cdot \mathbf{B} = 0 \quad (3.13)$$

$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \times \mathbf{E} = 0 \quad (3.14)$$

$$\nabla \cdot \mathbf{E} = 4\pi \rho \quad (3.15)$$

$$\frac{\partial \mathbf{E}}{\partial t} - \nabla \times \mathbf{B} = -4\pi \mathbf{J} \quad (3.16)$$

8.

9.

10. **More differentiation. (a) Justify the formula,**

$$d(u^\mu u_\mu)/d\tau = 2u_\mu(du^\mu/d\tau),$$

**by writing out the summation**  $u^\mu u_\mu = \eta_{\mu\nu} u^\mu u^\nu$  **explicitly** Writing out the components explicitly yields:

$$\begin{aligned} u^\mu u_\mu &= \eta_{00} u^0 u^0 + \eta_{11} u^1 u^1 + \eta_{22} u^2 u^2 + \eta_{33} u^3 u^3 \\ &= \eta_{00} (u^0)^2 + \eta_{11} (u^1)^2 + \eta_{22} (u^2)^2 + \eta_{33} (u^3)^2 \\ &= -(u^0)^2 + (u^1)^2 + (u^2)^2 + (u^3)^2 \end{aligned}$$

Taking a total derivative of the above with respect to  $\tau$ :

$$\begin{aligned} \frac{d}{d\tau}(u^\mu u_\mu) &= \frac{d}{d\tau}(-(u^0)^2 + (u^1)^2 + (u^2)^2 + (u^3)^2) \\ &= -2u^0 \frac{du^0}{d\tau} + 2u^1 \frac{du^1}{d\tau} + 2u^2 \frac{du^2}{d\tau} + 2u^3 \frac{du^3}{d\tau} \\ &= 2[-u^0 \frac{du^0}{d\tau} + u^1 \frac{du^1}{d\tau} + u^2 \frac{du^2}{d\tau} + u^3 \frac{du^3}{d\tau}] \\ &= 2\eta_{\mu\nu} u^\mu \frac{du^\nu}{d\tau} \\ &= 2u_\mu \frac{du^\mu}{d\tau} \end{aligned}$$

Therefore the formula is justified.

**b) Let  $\delta$  indicate a variation or small change, and justify the formula:**

$$\delta(F_{\alpha\beta} F^{\alpha\beta}) = 2F_{\alpha\beta} \delta F^{\alpha\beta}$$

The variation will obey the product rule as follows and remembering that the variation of the metric components would be zero:

$$\begin{aligned} \delta(F_{\alpha\beta} F^{\alpha\beta}) &= (\delta F_{\alpha\beta}) F^{\alpha\beta} + F_{\alpha\beta} (\delta F^{\alpha\beta}) \\ &= (\delta(\eta_{\alpha\gamma} \eta_{\beta\nu} F^{\alpha\beta}) F^{\gamma\nu}) + F_{\alpha\beta} (\delta F^{\alpha\beta}) \\ &= (\delta(\eta_{\alpha\gamma}) \eta_{\beta\nu} F^{\alpha\beta} + \eta_{\alpha\gamma} \delta(\eta_{\beta\nu}) F^{\alpha\beta} + \eta_{\alpha\gamma} \eta_{\beta\nu} \delta F^{\alpha\beta}) F^{\gamma\nu} + F_{\alpha\beta} (\delta F^{\alpha\beta}) \\ &= \eta_{\alpha\gamma} \eta_{\beta\nu} \delta F^{\alpha\beta} F^{\gamma\nu} + F_{\alpha\beta} (\delta F^{\alpha\beta}) \\ &= \delta F^{\alpha\beta} (\eta_{\alpha\gamma} \eta_{\beta\nu} F^{\gamma\nu}) + F_{\alpha\beta} (\delta F^{\alpha\beta}) \\ &= \delta F^{\alpha\beta} F_{\alpha\beta} + F_{\alpha\beta} \delta F^{\alpha\beta} \\ &= 2F_{\alpha\beta} \delta F^{\alpha\beta} \end{aligned}$$

Therefore the formula is justified.



**c) Compute**  $(F_{\alpha\beta}F^{\alpha\beta})_{,\mu}=?$

$$\begin{aligned}
 (F_{\alpha\beta}F^{\alpha\beta})_{,\mu} &= F_{\alpha\beta,\mu}F^{\alpha\beta} + F_{\alpha\beta}F^{\alpha\beta}_{,\mu} \\
 &= (\eta_{\alpha\gamma}\eta_{\beta\nu}F^{\alpha\beta})_{,\mu}F^{\gamma\nu} + F_{\alpha\beta}F^{\alpha\beta}_{,\mu} \\
 &= F^{\alpha\beta}_{,\mu}F_{\alpha\beta} + F_{\alpha\beta}F^{\alpha\beta}_{,\mu} \\
 &= 2F_{\alpha\beta}F^{\alpha\beta}_{,\mu}
 \end{aligned}$$