Thinfilm

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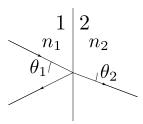
1 Matrice de transfert

$$\begin{array}{c|c} 1 & 2 \\ t_{12} & & \\ \hline & t_{21} \\ \hline & \\ r_{12} & & \\ \hline & \\ & & \\ \end{array}$$

$$\begin{pmatrix} \rightarrow \\ \leftarrow \end{pmatrix}_1 = \frac{1}{t_{12}} \begin{pmatrix} 1 & -r_{21} \\ r_{12} & t_{12}t_{21} - r_{12}r_{21} \end{pmatrix} \begin{pmatrix} \rightarrow \\ \leftarrow \end{pmatrix}_2 \tag{2}$$

Dans l'équation (2), \rightarrow_1 est la phase et amplitude de l'onde incidente depuis la gauche. \leftarrow_1 est la phase et amplitude de l'onde réfléchie vers la gauche. \rightarrow_2 est celle de l'onde transmise et finalement \leftarrow_2 est nul.

2 Fresnel



Polarisation perpendiculaire

$$\begin{cases}
 r_s = r_{12} = \frac{n_1 \cos \theta_1 - n_2 \cos \theta_2}{n_1 \cos \theta_1 + n_2 \cos \theta_2} = \frac{\eta_1 - \eta_2}{\eta_1 + \eta_2} \\
 t_s = t_{12} = \frac{2n_1 \cos \theta_1}{n_1 \cos \theta_1 + n_2 \cos \theta_2} = \frac{2\eta_1}{\eta_1 + \eta_2}
\end{cases}$$
(3)

Avec $\eta = n \cos \theta$.

Polarisation parallèle

$$\begin{cases}
 r_p = \frac{n_1 \cos \theta_2 - n_2 \cos \theta_1}{n_2 \cos \theta_1 + n_1 \cos \theta_2} = \frac{\eta_1 - \eta_2}{\eta_1 + \eta_2} \\
 t_p = \frac{2n_1 \cos \theta_1}{n_2 \cos \theta_1 + n_1 \cos \theta_2} = \frac{2n_1/\cos \theta_2}{\eta_1 + \eta_2} = \frac{2\frac{n_1}{n_2} \eta_2}{\eta_1 + \eta_2}
\end{cases} (4)$$

Avec $\eta = n/\cos\theta$.

3 Matrice de transfert de Fresnel

Polarisation perpendiculaire

$$\frac{1}{t_{12}} \begin{pmatrix} 1 & -r_{21} \\ r_{12} & t_{12}t_{21} - r_{12}r_{21} \end{pmatrix} = \frac{1}{2\eta_{1}} \begin{pmatrix} \eta_{1} + \eta_{2} & \eta_{1} - \eta_{2} \\ \eta_{1} - \eta_{2} & \frac{2\eta_{1}2\eta_{2}}{\eta_{1} + \eta_{2}} - \frac{(\eta_{1} - \eta_{2})(\eta_{2} - \eta_{1})}{\eta_{1} + \eta_{2}} \end{pmatrix}$$

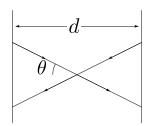
$$= \frac{1}{2\eta_{1}} \begin{pmatrix} \eta_{1} + \eta_{2} & \eta_{1} - \eta_{2} \\ \eta_{1} - \eta_{2} & \eta_{1} + \eta_{2} \end{pmatrix}$$

$$= \frac{1}{2\eta_{1}} \begin{pmatrix} \eta_{1} & 1 \\ \eta_{1} & -1 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ \eta_{2} & -\eta_{2} \end{pmatrix}$$
(5)

Polarisation parallèle

$$\frac{1}{t_{12}} \begin{pmatrix} 1 & -r_{21} \\ r_{12} & t_{12}t_{21} - r_{12}r_{21} \end{pmatrix} = \frac{1}{2\frac{n_1}{n_2}\eta_2} \begin{pmatrix} \eta_1 + \eta_2 & \eta_1 - \eta_2 \\ \eta_1 - \eta_2 & \frac{2\frac{n_1}{n_2}\eta_22\frac{n_2}{n_1}\eta_1}{\eta_1 + \eta_2} - \frac{(\eta_1 - \eta_2)(\eta_2 - \eta_1)}{\eta_1 + \eta_2} \end{pmatrix}
= \frac{1}{2\frac{n_1}{n_2}\eta_2} \begin{pmatrix} \eta_1 + \eta_2 & \eta_1 - \eta_2 \\ \eta_1 - \eta_2 & \eta_1 + \eta_2 \end{pmatrix}
= \frac{1}{2n_1} \begin{pmatrix} \eta_1 & 1 \\ \eta_1 & -1 \end{pmatrix} \frac{n_2}{\eta_2} \begin{pmatrix} 1 & 1 \\ \eta_2 & -\eta_2 \end{pmatrix}$$
(6)

4 Matrice de transfert de propagation



On a $t_{12}=t_{21}=\psi$ et aucune réflexion.

$$\frac{1}{t_{12}} \begin{pmatrix} 1 & -r_{21} \\ r_{12} & t_{12}t_{21} - r_{12}r_{21} \end{pmatrix} = \begin{pmatrix} \psi^{-1} & 0 \\ 0 & \psi \end{pmatrix}$$
 (7)

 $\psi=\exp(ik\frac{d}{\cos\theta})=\exp(ink_0\frac{d}{\cos\theta})=\exp(in\frac{2\pi}{\lambda_0}\frac{d}{\cos\theta}) \text{ où } k_0 \text{ est le nombre d'onde dans le vide et } \lambda_0 \text{ est la longueur d'onde dans le vide.}$

5 Multi couches

$$\left(\begin{array}{c} \rightarrow \\ \leftarrow \\ \end{array} \right)_{inc} \quad n_{inc}, \theta_{inc} \quad n_{1}, d_{1} \quad n_{2}, d_{2} \quad n_{3}, d_{3} \quad n_{exit} \quad \left(\begin{array}{c} \rightarrow \\ \leftarrow \\ \end{array} \right)_{exit}$$

Polarisation perpendiculaire

$$\begin{pmatrix}
\rightarrow \\
\leftarrow
\end{pmatrix}_{inc} = \frac{1}{2\eta_{inc}} \begin{pmatrix} \eta_{inc} & 1 \\ \eta_{inc} & -1 \end{pmatrix} \underbrace{\begin{pmatrix} 1 & 1 \\ \eta_{1} & -\eta_{1} \end{pmatrix}}_{1} \underbrace{\begin{pmatrix} \psi_{1}^{-1} & 0 \\ 0 & \psi_{1} \end{pmatrix}}_{L_{3}} \underbrace{\frac{1}{2\eta_{1}}}_{1} \begin{pmatrix} \eta_{1} & 1 \\ \eta_{1} & -1 \end{pmatrix}}_{1} \begin{pmatrix} 1 & 1 \\ \eta_{2} & -\eta_{2} \end{pmatrix} \underbrace{\begin{pmatrix} \psi_{2}^{-1} & 0 \\ 0 & \psi_{2} \end{pmatrix}}_{U_{2}}$$

$$\frac{1}{2\eta_{2}} \begin{pmatrix} \eta_{2} & 1 \\ \eta_{2} & -1 \end{pmatrix} \underbrace{\begin{pmatrix} 1 & 1 \\ \eta_{3} & -\eta_{3} \end{pmatrix}}_{1} \underbrace{\begin{pmatrix} \psi_{3}^{-1} & 0 \\ 0 & \psi_{3} \end{pmatrix}}_{1} \underbrace{\frac{1}{2\eta_{3}}}_{1} \begin{pmatrix} \eta_{3} & 1 \\ \eta_{3} & -1 \end{pmatrix}}_{1} \underbrace{\begin{pmatrix} 1 & 1 \\ \eta_{exit} & -\eta_{exit} \end{pmatrix}}_{exit} \underbrace{\begin{pmatrix} \rightarrow \\ \leftarrow \end{pmatrix}}_{exit}$$

$$= \frac{1}{2\eta_{inc}} \begin{pmatrix} \eta_{inc} & 1 \\ \eta_{inc} & -1 \end{pmatrix} L_{1}L_{2}L_{3} \begin{pmatrix} 1 & 1 \\ \eta_{exit} & -\eta_{exit} \end{pmatrix} \underbrace{\begin{pmatrix} \rightarrow \\ \leftarrow \end{pmatrix}}_{exit}$$

$$(8)$$

$$L = \begin{pmatrix} 1 & 1 \\ \eta & -\eta \end{pmatrix} \begin{pmatrix} \psi^{-1} & 0 \\ 0 & \psi \end{pmatrix} \frac{1}{2\eta} \begin{pmatrix} \eta & 1 \\ \eta & -1 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{\psi + \psi^{-1}}{2} & -\frac{\psi - \psi^{-1}}{2\eta} \\ -\eta \frac{\psi - \psi^{-1}}{2} & \frac{\psi + \psi^{-1}}{2} \end{pmatrix}$$

$$= \begin{pmatrix} \cos \delta & -\frac{i \sin \delta}{\eta} \\ -\eta i \sin \delta & \cos \delta \end{pmatrix}$$

$$(9)$$

où
$$\delta = n \frac{2\pi}{\lambda_0} \frac{d}{\cos \theta}$$

Si on prend $\begin{pmatrix} \rightarrow \\ \leftarrow \end{pmatrix}_{exit} = \begin{pmatrix} t \\ 0 \end{pmatrix}$ et $\begin{pmatrix} \rightarrow \\ \leftarrow \end{pmatrix}_{inc} = \begin{pmatrix} 1 \\ r \end{pmatrix}$ on a alors:
$$\begin{pmatrix} 1/t \\ r/t \end{pmatrix} = \frac{1}{2\eta_{inc}} \begin{pmatrix} \eta_{inc} & 1 \\ \eta_{inc} & -1 \end{pmatrix} L_1 L_2 L_3 \begin{pmatrix} 1 \\ \eta_{exit} \end{pmatrix}$$
(10)

Polarisation parallèle

$$\begin{pmatrix}
\rightarrow \\
\leftarrow
\end{pmatrix}_{inc} = \frac{1}{2n_{inc}} \begin{pmatrix} \eta_{inc} & 1 \\ \eta_{inc} & -1 \end{pmatrix} \underbrace{\begin{pmatrix} \eta_{1} & 1 \\ \eta_{1} & -\eta_{1} \end{pmatrix}}_{\eta_{1}} \begin{pmatrix} \eta_{1}^{-1} & 0 \\ 0 & \psi_{1} \end{pmatrix} \underbrace{\frac{1}{2n_{1}} \begin{pmatrix} \eta_{1} & 1 \\ \eta_{1} & -1 \end{pmatrix}}_{\eta_{2}} \underbrace{\begin{pmatrix} \eta_{2} & 1 \\ \eta_{2} & -\eta_{2} \end{pmatrix}}_{\eta_{2}} \begin{pmatrix} \psi_{2}^{-1} & 0 \\ 0 & \psi_{2} \end{pmatrix} \\
\frac{1}{2n_{2}} \begin{pmatrix} \eta_{2} & 1 \\ \eta_{2} & -1 \end{pmatrix} \underbrace{\frac{n_{3}}{\eta_{3}} \begin{pmatrix} 1 & 1 \\ \eta_{3} & -\eta_{3} \end{pmatrix}}_{\eta_{3}} \underbrace{\begin{pmatrix} \psi_{3}^{-1} & 0 \\ 0 & \psi_{3} \end{pmatrix}}_{2n_{3}} \underbrace{\frac{1}{\eta_{3}} \begin{pmatrix} \eta_{3} & 1 \\ \eta_{3} & -1 \end{pmatrix}}_{qexit} \underbrace{\begin{pmatrix} 1 & 1 \\ \eta_{exit} & -\eta_{exit} \end{pmatrix}}_{qexit} \underbrace{\begin{pmatrix} \rightarrow \\ \leftarrow \end{pmatrix}}_{exit} \\
= \frac{1}{2\eta_{inc}} \begin{pmatrix} \eta_{inc} & 1 \\ \eta_{inc} & -1 \end{pmatrix} L_{1}L_{2}L_{3} \underbrace{\frac{n_{exit}}{\eta_{exit}} \begin{pmatrix} 1 & 1 \\ \eta_{exit} & -\eta_{exit} \end{pmatrix}}_{qexit} \underbrace{\begin{pmatrix} \rightarrow \\ \leftarrow \end{pmatrix}}_{exit}$$

$$(11)$$

Les matrices L ont la même forme que dans le cas perpendiculaire.

Si on prend
$$\begin{pmatrix} \rightarrow \\ \leftarrow \end{pmatrix}_{exit} = \begin{pmatrix} t \\ 0 \end{pmatrix}$$
 et $\begin{pmatrix} \rightarrow \\ \leftarrow \end{pmatrix}_{inc} = \begin{pmatrix} 1 \\ r \end{pmatrix}$ on obtient :
$$\begin{pmatrix} 1/t \\ r/t \end{pmatrix} = \frac{1}{2\eta_{inc}} \begin{pmatrix} \eta_{inc} & 1 \\ \eta_{inc} & -1 \end{pmatrix} L_1 L_2 L_3 \begin{pmatrix} n_{exit}/\eta_{exit} \\ n_{exit} \end{pmatrix}$$
(12)