Notes on Cordier-Dubuc Paper

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1 Weak formulation

The solution consists of two stages:

- (A) Solve for the magnetic scalar potential Ψ , imposing the desired fluxoid quantization condition through the inhomogeneous boundary term B_{QS}
- (B) Solve for the shifted vector potential \mathbf{A}^* and the phase θ .

Because both vector- and scalar-valued fields are involved, the finite-element function space will be a direct product of $Q_p(\Omega_{\rm in} \cup \Omega_{\rm out})$ (for the scalar fields) and $Q_p^3(\Omega_{\rm in} \cup \Omega_{\rm out})$ (for the vector fields).

TODO. Investigate possible issues with convergence related to choice of finite-element spaces and the LBB condition for saddle-point problems. I believe the above (naive) choices should be OK, but there may be better choices / I may have misunderstood the convergence criteria.

Definition 1 (Basis functions). Let

- φ_i denote the scalar basis elements
- ullet Φ_i denote the vector basis elements

Definition 2 (System matrices). The following matrices will be used to assemble the linear systems for Items (A)

and (B):

$$M_{ij}^{(\text{int})} := \int_{\Omega_{(\text{int})}} \nabla \varphi_i \cdot \nabla \varphi_j \, d\Omega \tag{0.1}$$

$$M_{ij}^{(\text{ext})} := \int_{\Omega_{(\text{ext})}} \nabla \varphi_i \cdot \nabla \varphi_j \, d\Omega \tag{0.2}$$

$$C_{ij}^{(\text{int/ext})} := \int_{\Omega_{(\text{int/ext})}}^{(\text{ext})} (\mathbf{\nabla} \times \mathbf{\Phi}_i) \cdot (\mathbf{\nabla} \times \mathbf{\Phi}_j) \, d\Omega$$
 (0.3)

$$D_{ij}^{(\text{int/ext})} := \int_{\Omega_{(\text{int/ext})}} (\mathbf{\nabla} \cdot \mathbf{\Phi}_i) (\mathbf{\nabla} \cdot \mathbf{\Phi}_j) \, d\Omega$$
 (0.4)

$$S_{ij}^{(\text{int/ext})} := \int_{\Omega_{(\text{int/ext})}} \mathbf{\Phi}_i \cdot \mathbf{\Phi}_j \, d\Omega \tag{0.5}$$

$$T_{ij}^{(\text{int/ext})} := \int_{\Omega_{(\text{int/ext})}} \mathbf{\Phi}_i \cdot \mathbf{\nabla} \varphi_j \, d\Omega$$
 (0.6)

$$E_{ij}^{(\text{ext})} := \int_{\Omega_{(\text{ext})}} (\nabla \times \mathbf{\Phi}_i) \cdot \nabla \varphi_j \, d\Omega$$
 (0.7)

$$H_i^{(\text{ext})} := \int_{\Gamma_H} \mathbf{\Phi}_i \cdot \mathbf{H}_{(\text{ext})} \, d\Omega \tag{0.8}$$

$$B_i^{(QS)} := \int_{S_a} \varphi_i B_{QS} \, \mathrm{d}S \tag{0.9}$$

The weak formulation of the problem is then

$$M_{ij}^{(\text{ext})} \Psi_j = B_i^{(\text{QS})} \tag{A}$$

$$\begin{cases}
 \left[C_{ij}^{(\text{int})} + p D_{ij}^{(\text{int})} + \frac{1}{\lambda^2} S_{ij}^{(\text{int})} \right] A_j^* = \frac{\Phi_0}{2\pi \lambda^2} T_{ij}^{(\text{int})} \theta_i \\
 \frac{1}{\mu_0 \lambda^2} \widetilde{T}_{ij}^{(\text{int})} A_j^* = \frac{\Phi_0}{2\pi \mu_0 \lambda^2} M_{ij}^{(\text{int})} \theta_i \\
 \left[C_{ij}^{(\text{ext})} + p D_{ij}^{(\text{ext})} \right] A_i^* = -E_{ij}^{(\text{ext})} \Psi_j + \mathbf{H}_i^{(\text{ext})}
\end{cases} \tag{B}$$

• Macroscopic complex wavefunction of Cooper pairs (isn't this equivalent to considering a GL equation?):

$$J_s = \frac{1}{\mu_0 \lambda_L^2} \left(-A + \frac{\Phi_0}{2\pi} \nabla \phi \right) \tag{0.10}$$

1.1 Discretization

Following deal.ii's step-46 tutorial, we look for solutions in the following function space:

$$(\Psi, \mathbf{A}^*, \theta) \in \underbrace{[Z^1(\Omega_{\mathrm{int}} \times Q_{p+1}^1(\Omega_{\mathrm{ext}})]}_{\Psi} \times \underbrace{[Q_p^d(\Omega_{\mathrm{int} \cup \Omega_{\mathrm{ext}}})]}_{\mathbf{A}^*} \times \underbrace{[Z^1(\Omega_{\mathrm{ext}}) \times Q_{p+1}^1(\Omega_{\mathrm{int}})]}_{\theta}$$
(0.11)

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2 References (forward)

- C. Cordier, S Flament, and C. Dubuc, "A 3D finite element formulation for calculating meissner currents in superconductors," IEEE Trans. Appl. Superconduct., vol. 9, pp. 2–6, Mar. 1999.
 - Only allows calculations with multiply-connected regions if the fluxoid is zero

section References (backward) None found. (Searched: Science Web)

3 Questions

- 1. Values of B_{QS} ? (Seems arbitrary except for the requirement that the flux through any "holes" should be the desired value)
- 2. Boundary conditions on Γ_{12} are apparently enforced implicitly without the need to discretize eqs. 26 29. This is very convenient but somewhat mysterious...

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