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Exercise 4: Hodgkin-Huxley Model

Introduction

In Hodgkin-Huxley Model, two new elements are introduced based on the single-compartment model (SCM) for simulating the generation of action potential, namely a transient voltage-dependent sodium conductance (g_{Na}) and a persistent voltage-dependent potassium conductance (g_{Na}) (Fig. 1).

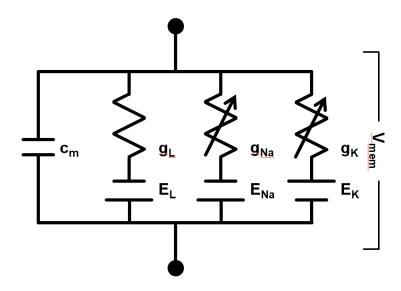


Fig. 1 Modified electrical circuit for Hodgkin-Huxley Model. Based on the SCM, two elements are added, namely a voltage-dependent sodium conductance (g_{Na}) and a voltage-dependent potassium conductance (g_K) . All components are connected in parallel. The battery in serial with each resistor annotates the reversal potential of the corresponding conductance $(E_L, E_{Na} \text{ or } E_K)$.

These conductances depend on the maximum conductances \bar{g}_{Na} and \bar{g}_{Na} and also the activation variables m(V), h(V) and n(V) as below:

$$g_{Na}(V) = \bar{g}_{Na}m(V)^3h(V)$$
 $g_K(V) = \bar{g}_Kn(V)^4$ $m(V), h(V), n(V) \in [0,1]$

Interestingly, the three activation variables showed different responses with changes in voltage (Fig. 2). On one hand, with an increasing membrane voltage, h decreases whereas m and n increase. On the other hand, the time constant of m is particularly small, which reveals its quick rate of change voltage. Only when membrane voltage reaches above -30mV do the time constants of h and h decrease to comparable levels. Base on the following equation,

$$\frac{dV}{dt} = \frac{1}{\tau_{eff}(V)} [V_{\infty}^{eff}(V) - V]$$

$$\tau_{eff}(V) = \frac{c_m}{\bar{g}_L + \bar{g}_{Na} m(V)^3 \mathrm{h}(V) + \bar{g}_K n(V)^4} \quad V_{\infty}^{eff}(V) = \frac{\bar{g}_L E_L + \bar{g}_{Na} m(V)^3 \mathrm{h}(V) E_{Na} + \bar{g}_K n(V)^4 E_K + i_e}{\bar{g}_L + \bar{g}_{Na} m(V)^3 \mathrm{h}(V) + \bar{g}_K n(V)^4}$$

one can trace the kinetics of membrane voltage with a current injected. However, the challenge will be that the values of m, h and n also change along with membrane voltage.

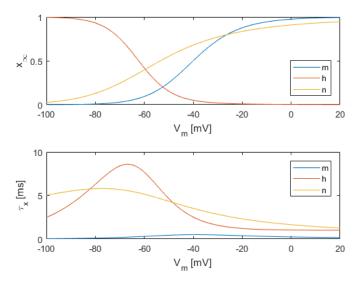


Fig. 2 Equilibrium states and time constants of the activation variables against membrane voltage. In the upper panel, h shows a decreasing trend when the membrane voltage increases; whereas m and n show increasing ones. In the lower panel, the time constant of m is particularly small, which reveals the quick rate of change in value whenever there is a change in voltage. That of h and h becomes comparable only when membrane voltage is above -30mV.

In this exercise, based on the SCM, the changes of membrane voltage (V_{mem}) and the activation variables (m, h and n) across the time were computed under different sets of their initial values and also different currents injected into the cell (I_e) using Matlab.

Method

In this exercise, different vectors were generated for illustrating Hodgkin-Huxley Model. Membrane capacity (c_m , also c_m in the code), maximal conductances (g in the code) and corresponding reversal potentials (E in the code) were saved at the very beginning, as they do not change in the whole exercise. g and E are both vectors in the code, each of which contains the values for leaky conductance ($\bar{g}_L \& E_L$), voltage-dependent sodium conductance ($\bar{g}_{Na} \& E_{Na}$) and voltage-dependent potassium conductance ($\bar{g}_K \& E_K$). They are saved as vectors as this design will speed up the calculation in the following parts. The values of these seven parameters are as below:

$$c_m = 10nF/mm^2; \ \bar{g}_L = 3\mu S/mm^2; \ \bar{g}_{Na} = 1200\mu S/mm^2; \ \bar{g}_K = 360\mu S/mm^2;$$

 $E_L = -54.402mV; \ E_L = 50mV; \ E_L = -77mV.$

Since the aim of this exercise is to observe the changes of variables across the time, a time vector (t in the code) from 0 to 40ms with step-size (dt in the code) of 0.1ms was generated. Besides, because, in this time, all the situations of current injected l_e were constant only, a single variable for current (i_e in the code) was generated and reused in all situations.

Next, a function with a for-loop (HH_time.m) was used for the generation of dynamic variables, including the activation variables (m, h and n in the code), effective time constant of membrane voltage (τ_{eff} , also V_{eff} in the code), effective equilibrium state of membrane voltage (V_{∞}^{eff} , also V_{eff} in the code) and, most importantly, membrane voltage (V_{eff}). For E_{eff} in the code) and, most importantly, membrane voltage (V_{eff}). For E_{eff} in the corresponding of the for-loop with the initial value assigned in each situation. The corresponding values of T_{eff} and T_{eff} were then calculated, again, based on the following equations:

$$\tau_{eff}(t) = \frac{c_m}{\bar{g}_L + \bar{g}_{Na} m^3 h + \bar{g}_K n^4} \qquad V_{\infty}^{eff}(t) = \frac{\bar{g}_L E_L + \bar{g}_{Na} m^3 h E_{Na} + \bar{g}_K n^4 E_K + i_e}{\bar{g}_L + \bar{g}_{Na} m^3 h + \bar{g}_K n^4}$$

In each round of the for-loop, the new equilibrium states (x_infty in the code, where x is m, h or n) and time constants (tau_x in the code, where x is m, h or n) for m, h and n were firstly generated using the function files provided (HH_equi_tau_m.m, HH_equi_tau_h.m and HH_equi_tau_n.m). Namely, by providing the old voltage value in these functions, the rates of activating ($\alpha(V)$) and inactivating ($\beta(t)$) of each of the activation variables were calculated in the sub-function (e.g. HH_alpha_beta_m) in each of the files as below:

$$\alpha_m(V) = \frac{0.1/ms(V + 40mV)}{1 - e^{-0.1(V + 40mV)}} \qquad \alpha_h(V) = \frac{0.07}{ms}e^{-0.05(V + 65mV)} \qquad \alpha_n(V) = \frac{0.01/ms(V + 55mV)}{1 - e^{-0.1(V + 55mV)}}$$

$$\beta_m(V) = \frac{4}{ms} e^{-0.0556(V + 65mV)} \qquad \beta_h(V) = \frac{1/ms}{1 + e^{-0.1(V + 35mV)}} \qquad \beta_n(V) = \frac{0.125}{ms} e^{-0.0125(V + 65mV)}$$

Then new *x_infty* and *tau_x* were calculated:

$$x_{\infty}(V) = \frac{\alpha_{x}(V)}{\alpha_{x}(V) + \beta_{x}(V)} \qquad \qquad \tau_{x}(V) = \frac{1}{\alpha_{x}(V) + \beta_{x}(V)}$$

Using these new values, the new value of m, h and n were calculated and added to the vectors as followed:

$$x(i+1) = x_{\infty}(V) + [x(i) - x_{\infty}(V)]e^{\frac{-dt}{\tau_{\chi}(V)}}$$

Afterwards, the new τ_{eff} and V_{∞}^{eff} were calculated with the new value of m, h and n as above, and saved. At the end, the new V was computed as below and saved:

$$V(i+1) = V_{\infty}^{eff}(t) + [V(i) - V_{\infty}^{eff}(t)]e^{\frac{-dt}{\tau_{eff}(t)}}$$

In the three situations, the i_e and initial values of V, m, h and n were as below:

	1 st situation	2 nd situation	3 rd situation
i _e (nA/mm²)	0	100	62
V (mV)	-70	Ending of 1 st situation	-65
m	0	Ending of 1 st situation	0.08
h	0	Ending of 1 st situation	0.5
n	0.3	Ending of 1 st situation	0.37

Results & Discussion

Calculation of τ_{eff} and V_{∞}^{eff}

In Hodgkin-Huxley model, the dynamic equation becomes:

$$c_m \frac{dV(t)}{dt} = -i_L - i_{Na} - i_K + i_e$$

in which:

$$i_L = \bar{g}_L(V(t) - E_L)$$
 $i_{Na} = \bar{g}_{Na}m^3h(V(t) - E_{Na})$ $i_K = \bar{g}_K n^4(V(t) - E_K)$

Therefore,

$$c_m \frac{dV(t)}{dt} = (\bar{g}_L E_L + \bar{g}_{Na} m^3 h E_{Na} + \bar{g}_K n^4 E_K + i_e) - (\bar{g}_L + \bar{g}_{Na} m^3 h + \bar{g}_K n^4) [V(t)]$$

$$\frac{dV(t)}{dt} = \frac{\bar{g}_L + \bar{g}_{Na}m^3h + \bar{g}_Kn^4}{c_m} \left\{ \frac{\bar{g}_L E_L + \bar{g}_{Na}m^3hE_{Na} + \bar{g}_Kn^4E_K + i_e}{\bar{g}_L + \bar{g}_{Na}m^3h + \bar{g}_Kn^4} - [V(t)] \right\}$$

$$\dot{\tau}_{eff}(t) = \frac{c_m}{\bar{g}_L + \bar{g}_{Na} m^3 h + \bar{g}_K n^4} \qquad V_{\infty}^{eff}(t) = \frac{\bar{g}_L E_L + \bar{g}_{Na} m^3 h E_{Na} + \bar{g}_K n^4 E_K + i_e}{\bar{g}_L + \bar{g}_{Na} m^3 h + \bar{g}_K n^4}$$

Activation variables and voltage reached steady state when no current is injected

Without any current input, the membrane voltage and activation variables (m, h and n) eventually achieved their respective plateaus (Fig. 3). In the upper panel, with the initial values of activation variables, the $V_{\infty}^{eff}(0)$ was -65.5mV, which was higher than V(0). This difference caused the increase in V. Similar situations happened to the three activation variables in the lower panel of Fig. 3. Refer to Fig. 2, the value of $m_{\infty}(-65.5)$, $h_{\infty}(-65.5)$ and $n_{\infty}(-65.5)$ were 0.03, 0.75 and 0.25, which were different than the initial value. Among the three, h had the most significant different and time constant at t=0, which therefore showed the greatest rate of change.

Despite so, it is interesting that it is actually n, which has the least difference between n(0) and $n_{\infty}(0)$, that promoted the change in V_{∞}^{eff} (Data not shown). This may because \bar{g}_K was relatively smaller that \bar{g}_{Na} , and hence a slight change in the n caused a more significant effect on the denominator in the equation of V_{∞}^{eff} (see above). With an almost unnoticeable decrease in n, driven by the difference with the equilibrium state, V_{∞}^{eff} showed an increase. However, with a sustained increase in V, the equilibrium state of n also increased, and thus V_{∞}^{eff} decreased. Eventually, the difference between V and V_{∞}^{eff} gradually diminished, causing the achievement of equilibrium state of all the variables, which are also the values in a resting neuron.

Spiked change generated in membrane voltage under sufficient current input

With a current of $100nA/mm^2$, spiked patterns were observed in membrane voltage (Fig. 4). In each cycle, a slow depolarisation was initially observed. Once the membrane voltage reached around -50mV, an overshooting pattern followed by a repolarisation phase (i.e. the spike), appeared. Correspondingly, in the depolarisation phase, m and h were generally increasing, whereas n was decreasing. In the overshooting phase, at the very beginning m increased drastically, causing the increasing of g_{Na} , and so as V_{∞}^{eff} . However, soon after so, h decreased with the rise of membrane voltage and therefore counteracted the increase of m. At the same time, n increased and caused an increase of g_{Na} . These two factors contributed to the slowdown of overshooting and repolarising phase of V.

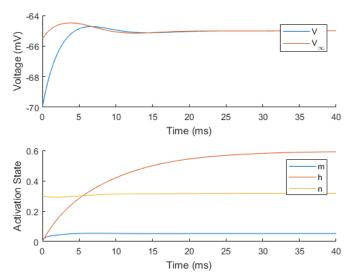


Fig. 3 Voltage and activation variables stabilised across time without any current input. In the both panel, the variables eventually reached plateau. Despite insignificant in terms of value, the change in n (lower panel) synchronised with that in V_{∞}^{eff} (upper panel).

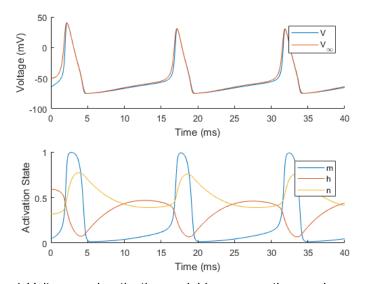


Fig. 4 Voltage and activation variables across time under a current input of 100nA/mm². In the upper panel, membrane voltage showed a spiky pattern similar to the sharp of action potentials.

Spike generation in membrane voltage depends on the amount of current input

With a current of $62nA/mm^2$, spiked patterns were observed in a similar manner as in the case of $100nA/mm^2$ (Fig. 5). However, the inter-spike interval was longer with a lower input current. Interestingly, with an even lower current input of $61nA/mm^2$, despite the successful generation of one spike, the membrane voltage and also the activation variables evenutally stabilised. This may due to the reduced difference between V and V_{∞}^{eff} , thus the value of m could not reach to its exponential increasing phase as in Fig. 2. These results showed that the generation of action potential depends on how much the current is given to a neuron. On contrary, an extreme high current (e.g $1000nA/mm^2$) also distorted the formation of spikes (Data not shown).

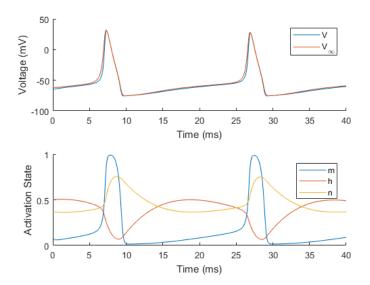


Fig. 5 Voltage and activation variables across time under a current input of 62nA/mm². In the upper panel, membrane voltage showed a similar result as Fig. 4, but in a slower rate, i.e. larger inter-spike intervals.

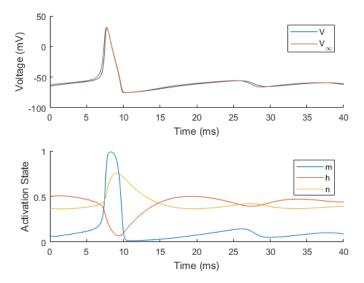


Fig. 6 Voltage and activation variables across time under a current input of 61nA/mm². In the upper panel, membrane voltage showed a spike, but afterwards stabilise. Similar pattern occurred to the activation variables.