

Mechanical Processing in Internally Coupled Ears

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Auditory Systems



Independent Ears

Eustachian tubes typically very narrow.

Effectively independent eardrum vibrations.



Coupled Ears

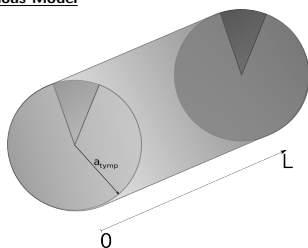
Eardrums connected through wide eustachian tubes and a large mouth cavity.

Eardrums vibrations influence each other.

Advantages of Low Frequency Hearing

Mouth Cavity

Previous Model

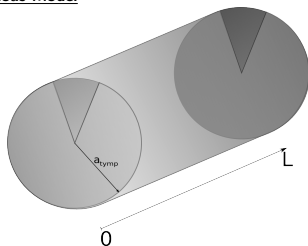


a_{tymp} fixed.

$$V_{\text{cyl}} = \pi a_{\text{tymp}}^2 L$$

Mouth Cavity

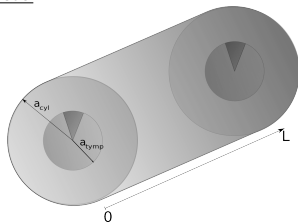
Previous Model



a_{tymp} fixed.

$$V_{\text{cyl}} = \pi a_{\text{tymp}}^2 L$$

Our Model

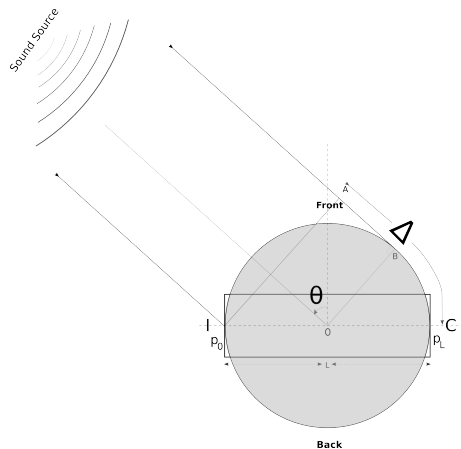


a_{tymp} , V_{cyl} fixed.

$$a_{\text{cyl}} = \sqrt{V_{\text{cyl}} / \pi L}$$

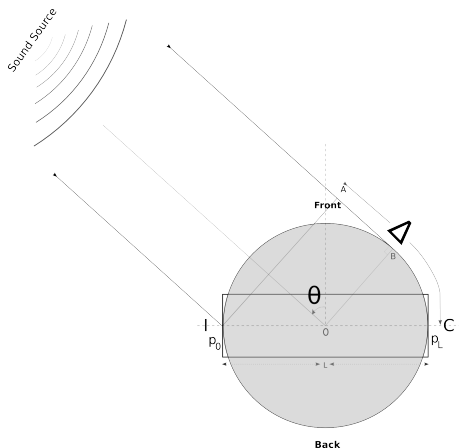
Acoustic Head Model

- **I** - Ipsilateral ear, **C** - Contralateral ear.
 p_0 , p_L - sound pressure on eardrums, θ - sound source direction.



Acoustic Head Model

- ▶ Sound source “far away”.
- ▶ Phase difference between sound at both ears - $\Delta = 1.5kL \sin \theta$.
- ▶ No appreciable amplitude difference, $|p_0| = |p_L|$.



Cavity Pressure

3D Wave Equation

$$\frac{1}{c^2} \partial_t^2 p(x, r, \phi, t) = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial p(x, r, \phi, t)}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 p(x, r, \phi, t)}{\partial \phi^2} + \frac{\partial^2 p(x, r, \phi, t)}{\partial x^2} \quad (1)$$

To be solved using the separation ansatz

$$p(x, r, \phi, t) = f(x)g(r)h(\phi)e^{j\omega t}$$

No-penetration at the cavity boundary, i.e. normal derivative vanishes

$$-j\rho\omega\mathbf{n} \cdot \nabla p(x, r, \phi; t)|_{r=a_{\text{cyl}}} = \left. \frac{\partial p}{\partial r} \right|_{r=a_{\text{cyl}}} = 0 \quad (2)$$

No-penetration at the cavity boundary, i.e. normal derivative vanishes

$$-j\rho\omega\mathbf{v} = \mathbf{n} \cdot \nabla p(x, r, \phi; t)|_{r=a_{\text{cyl}}} = \left. \frac{\partial p}{\partial r} \right|_{r=a_{\text{cyl}}} = 0 \quad (2)$$

Pressure Modes

$$p_{qs}(x, r, \phi; t) = \left[A_{qs} e^{j\zeta_{qs}x} + B_{qs} e^{-j\zeta_{qs}x} \right] \cos q\phi J_q(\nu_{qs}r) e^{j\omega t} \quad (3)$$

such that, $\left. \frac{\partial J_q(\nu_{qs}r)}{\partial r} \right|_{r=a_{\text{cyl}}} = 0$ and $\zeta_{qs} = \sqrt{k^2 - \nu_{qs}^2}$

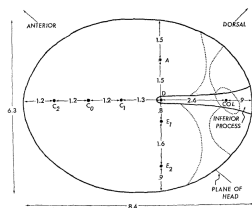
Plane Wave Mode

$$p_{00}(x, r, \phi; t) = \left[A e^{jkx} + B_{qs} e^{-jkx} \right] e^{j\omega t} \quad (4)$$

Trivially satisfies the no-penetration condition.

Eardrum

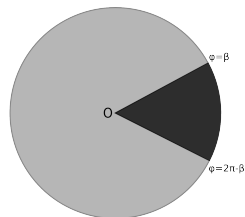
Sketch of a Tokay eardrum as seen from the outside^a.



COL - approximate position opposite the extracolumella insertion.

^aG. A. Manley, "The middle ear of the tokay gecko," *Journal of Comparative Physiology*, vol. 81, no. 3, pp. 239-250, 1972

The ICE eardrum.



Extracolumella (dark) - rigid, stationary.

Tympanum - assumed linear elastic.

Rigidly clamped at the boundaries ($r = a_{\text{typ}}$
and $\phi = \beta, 2\pi - \beta$)

Membrane Vibrations

Membrane EOM

$$-\partial_t^2 u(r, \phi; t) - 2\alpha \partial_t u(r, \phi; t) + c_M^2 \Delta_{(2)} u(r, \phi; t) = \frac{1}{\rho_m d} \psi(r, \phi; t) \quad (5)$$

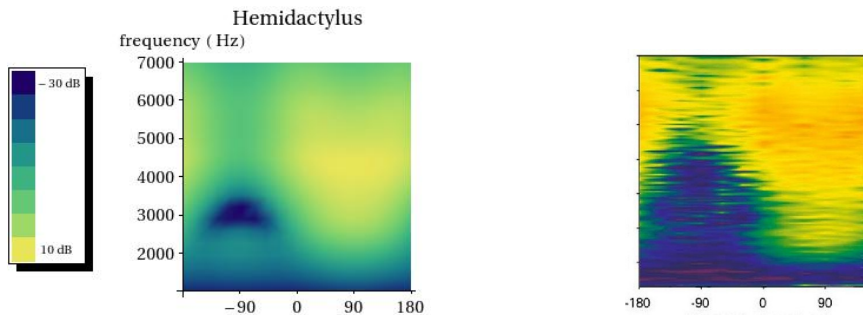
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Membrane EOM

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Membrane parameters

α - damping coefficient, c_M^2 - propagation velocity
 ρ_m - density, d - thickness.



Conclusion

Thank You

