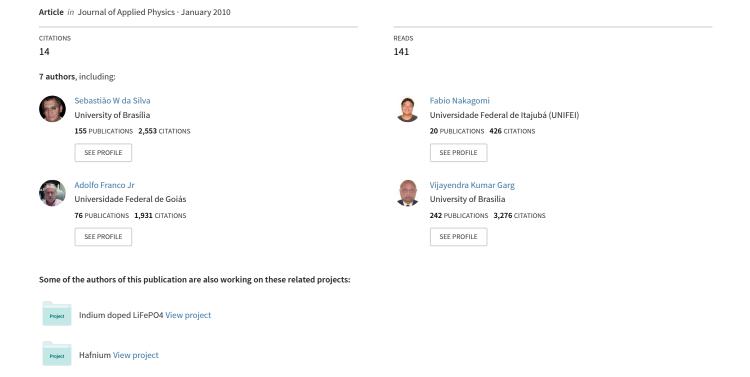
Effect of the Zn content in the structural and magnetic properties of Zn x Mg 1-x Fe 2 O 4 mixed ferrites monitored by Raman and Mössbauer spectroscopies



Effect of the Zn content in the structural and magnetic properties of $Zn_xMg_{1-x}Fe_2O_4$ mixed ferrites monitored by Raman and Mössbauer spectroscopies

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Samples of $Zn_xMg_{1-x}Fe_2O_4$ ($0 \le x \le 1$) synthesized by the combustion reaction method were investigated by x-ray diffraction, Mössbauer spectroscopy, and Raman spectroscopy. All the samples are found to have a cubic spinel structure and the lattice parameter increases linearly with increasing Zn-content (x). The Mössbauer data showed that the replacement of Mg^{2+} ions for Zn^{2+} ions changes substantially the hyperfine parameter. Moreover, it was verified the presence of Fe^{3+} ions both in A and B sites. The Raman spectra showed five predicted Raman bands for the spinel structure and it was observed the splitting of the A_{1g} Raman mode into tree branches, where each one have been attributed to peaks belonging to each ion (Zn, Mg, and Fe) in the tetrahedral positions. © 2010 American Institute of Physics. [doi:10.1063/1.3350903]

I. INTRODUCTION

The magnetic properties of ferrites are strongly dependent on the cation distributions. Spinel magnetic oxides can be represented by MFe₂O₄ formula (M is a divalent metal ion and Fe is a trivalent iron ion). Occupation of the tetrahedral sites (A sites) entirely with divalent cations produce a normal spinel structure, while occupation of the octahedral sites (B sites) with divalent cations yield an inverse spinal structure. In bulk form, ZnFe₂O₄ is usually assumed to be a completely normal spinel with all Fe^{3+} ions on B sites and all Zn^{2+} ions on A sites. On the other hand, it has been showed that in the MgFe₂O₄ ferrite, the site preference of the divalent ions leads to a predominantly inverse structure with Mg^{2+} ions mainly on B sites and Fe^{3+} ions distributed almost equally among A and B sites. Ferrite spinels may also contain mixture of two divalent metal ions such as Zn_xMg_{1-x}Fe₂O₄ in which Mg²⁺ and Zn²⁺ ratio may be varied. In this case, it is expected that zinc ions replace magnesium ions, between x=0 and 1, Zn^{2+} ions appear to enter preferentially tetrahedral positions while the Fe³⁺ ions should be displaced from these sites for the octahedral sites. However, in spite of the fact that for nanocrystalline systems such behavior is not always observed and there are few studies on this subject.^{2,3} In present work, the effect of Zn²⁺ nonmagnetic ions on structural and magnetic properties of cubic $Zn_rMg_{1-r}Fe_2O_4$ ($0 \le x \le 1$) nanoparticles have been investigated by means of x-ray diffraction (XRD), Mössbauer spectroscopy, and Raman spectroscopy.

II. EXPERIMENTS

A series of mixed magnesium-zinc ferrite with the general formula $Zn_xMg_{1-x}Fe_2O_4$, with $0.0 \le x \le 1.0$ was synthe-

sized by the combustion reaction method⁴ without subsequent calcination steps. All reagents, iron nitrate Fe(NO₃)₃.9H₂O, zinc nitrate Zn(NO₃)₂.6H₂O, magnesium nitrate Mg(NO₃)₂.6H₂O, and urea CO(NH₂)₂ as fuel, were analytical grade and manipulated in air without the protection of nitrogen or inert. The chemical composition (total zinc and magnesium contents) of all ferrite samples was determinate by atomic absorption spectrophotometry using the commercial Perkin-Elmer 5000 system. The calculated composition is in good agreement with the core chemical composition determinate by atomic spectroscopy (see Table I). The as-prepared powder were characterized by XRD using a Shimadzu diffractometer (model 6000) with Cu-K_a radiation $(\lambda = 1.5418 \text{ Å})$ in a wide range of Bragg angles $(15^{\circ} < 2\theta)$ $< 80^{\circ}$) with a scanning rate of 2°/min at room temperature. Liquid nitrogen-temperature Mössbauer spectra of the asproduced samples were recorded in the transmission geometry using the ⁵⁷Co source in Rh matrix. The system velocity was calibrated with a thin iron foil whereas the spectra were

TABLE I. The calculated composition, based on the proportions of the starting materials, and chemical composition determinate by atomic spectroscopy.

Calculated $Zn_xMg_{1-x}Fe_2O_4$	Determined		
	x (Zn)	(1-x) (Mg)	
0.0	0.000	0.999	
0.2	0.192	0.808	
0.4	0.390	0.610	
0.5	0.487	0.513	
0.6	0.617	0.383	
0.7	0.699	0.300	
0.8	0.812	0.188	
1.0	0.998	0.000	

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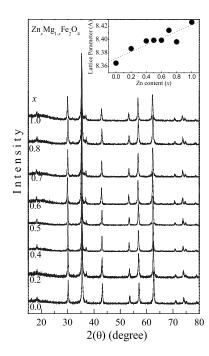


FIG. 1. XRD patterns of $Zn_xMg_{1-x}Fe_2O_4$ ($0 \le x \le 1$) samples. The inset shows a plot of the deduced lattice parameter as a function of x.

least-square fitted to a combination of Lorentzian-like lines. The Raman system used to record the spectra of the samples was a commercial triple spectrometer (Jobin Yvon Model T64000) equipped with a charge-coupled device detector. The 514 nm line from an argon-ion laser was used to illuminate the samples at an optical power around 0.2 mW. All Raman measurements were performed at room temperature.

III. RESULTS AND DISCUSSION

Figure 1 shows all XRD patterns of the spinel ferrite prepared by combustion reaction method at room temperature. XRD patterns revealed that all specimens exhibited sharp and intense peaks that correspond to the cubic inverse spinel structure of MgFe₂O₄.⁵ The absence of extra reflections in the diffraction patterns ensures the phase purity. The mean particle sizes were calculated from x-ray line broadening of the (311) diffraction peak using Scherrer equation.⁶ The mean particle diameter is nearly the same for all specimens; being ca. 40 and 42 nm for x=0.0 and x=1.0, respectively. The XRD data show that the lattice parameter of cubic cell increases linearly with increasing Zn content (x) following Vegard's law approximately,⁷ see the inset on Fig. 1.

All synthesized $Zn_xMg_{1-x}Fe_2O_4$ samples reported in this study were examined by Mössbauer spectroscopy at 77 K, as shown in Fig. 2. At low Zn content $(0 \le x \le 0.4)$ the Mössbauer spectra was analyzed by using the assumption of Lorentzian line shapes and resolved into two subspectra: one corresponding to A site Fe ions in tetrahedral sites and the other to B site Fe ion in octahedral sites. At intermediate Zn content, the Mössbauer spectra are broader and this can be explained considering the breaking of local magnetic ordering due to increase in Zn content. For samples of x=0.8 and 1.0 the Mössbauer spectra exhibits a doublet, as expected for a paramagnetic phase.

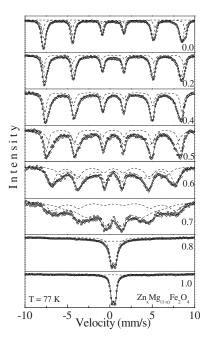


FIG. 2. Mössbauer spectra of $Zn_xMg_{1-x}Fe_2O_4$ ($0 \le x \le 1$) samples recorded at 77 K.

Mössbauer parameters such as relative area (I), hyperfine field (HF), isomer shift (IS), and quadrupole splitting (QS) for all samples are listed in Table II. It is observed that the isomer shift values present opposite behavior with increasing Zn content, while A site values decrease, B site values increase up to x=0.4. For x>0.4, the IS have close values. This can be explained through the bonding nature of iron in both sites. With the increasing x, the Fe population in the A site increases while in the B site, it decreases. As ionic radius values of Fe³⁺ (0.63 Å-A sites and 0.78 Å-B sites) are smaller than ionic radius values for Zn²⁺ (0.74 Å-A site and 0.88 Å-B site) and Mg²⁺ (0.71 Å-A site and 0.80 Å-B site), there is an expected increase in the orbital's overlapping of the ions in the A sites and a decrease in the B sites, resulting in IS changing.

TABLE II. Fitted Mössbauer parameters for $Zn_xMg_{1-x}Fe_2O_4$ ($0 \le x \le 1$) samples recorded at 77 K.

Samples (x) Zn	Site	I (%)	HF (T)	IS (mm/s)	QS (mm/s)
0.0	A	26.7	523.5	0.473	0.058
	B	73.3	503.5	0.380	-0.016
0.2	A	20.5	513.8	0.441	0.059
	B	79.5	495.5	0.404	-0.041
0.4	A	33.1	501.6	0.430	0.041
	B	66.9	477.2	0.426	-0.099
0.5	A	35.0	492.7	0.420	0.023
	B	65.0	454.7	0.447	-0.069
0.6	A	46.7	452.6	0.422	0.020
	B	53.3	395.0	0.451	-0.121
0.7	A	46.5	450.3	0.430	-0.046
	B	53.5	350.0	0.440	-0.197
0.8	A	Dub	• • •	0.430	0.470
	B				
1.0	A	Dub	•••	0.440	0.378
	В				

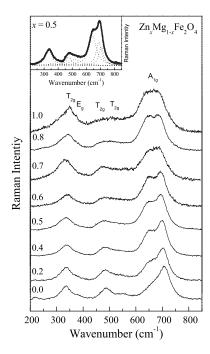


FIG. 3. The Raman spectra of $Zn_xMg_{1-x}Fe_2O_4$ ($0 \le x \le 1$) samples recorded at room temperature. The inset shows the Lorentzian-like assignment to the Raman modes for the sample x=0.5 in accordance with the five active optical modes.

The HF associated to both crystallographic sites (A and B) of samples $Zn_xMg_{1-x}Fe_2O_4$ ($0 \le x \le 0.7$) decreases as the Zn-content increases. This can be explained considering the superexchange interaction among the neighboring metal ions in the spinel structure. With increasing Zn content the probability of the Fe ions to find a Zn ion as nearest neighbor increases, thus HF decreases.

The relative areas of the fitted Mössbauer spectra gave the Fe ion distribution for each site. However, in the case of quaternary ferrites this information is not enough to characterize the site occupation, since it can be present in different divalent ions. In addition, for large Zn-content ($x \ge 0.8$), the presence of the paramagnetic doublet makes it impossible to determine the cation distribution to each site. So, the Raman spectroscopy was employed for this study.

Figure 3 shows the room temperature Raman spectra for $\operatorname{Zn}_x \operatorname{Mg}_{(1-x)}\operatorname{Fe}_2\operatorname{O}_4$ $(0 \le x \le 1)$ samples, recorded in the range of 200–850 cm⁻¹. In this region, we have assigned the Raman modes in accordance with the five active optical modes $(A_{1g} + E_g + 3T_{2g})$ characteristic of the cubic spinel O_h^7 $(Fd\overline{3}m)$ space group; see the inset on Fig. 3. It is known from literature that Raman modes present in region 650–710 cm⁻¹ has A_{1g} symmetry and is related to the tetrahedral sublattice. 9 As

discussed in previous work¹⁰ the mass difference between the three ions $(Zn^{2+}, Mg^{2+}, and Fe^{3+})$ splits the A_{1g} mode in three different energy values; the lightest ion (Mg²⁺) responds for the Raman mode peaking at 706 cm⁻¹ (for x =0.5) whereas the heaviest one (Zn^{2+}) responds for the 645 cm⁻¹ mode. Because the iron possesses an intermediate mass, we have associated the 689 cm⁻¹ mode for the Fe³⁺ ion in the tetrahedral sublattice (see the inset on Fig. 3). From Fig. 3, it is clearly observed in the spectra changes in the relative intensities of the peaks belonging to each ion. Following Seong et al., 11 the integrated intensity of the Raman mode is proportional to the number of the corresponding oscillators. From this point of view, we can state that while the intensity of the $A_{1g}(Mg)$ mode decreases with the Mg-content decreases, an opposite behavior happens to $A_{1p}(Zn)$ mode when x increases. Thus it is clear that Mg, Fe, and Zn ions are present in both sites of the spinel structure.

IV. CONLUSION

The influence of Zn-content on the structural and magnetic properties of $Zn_xMg_{1-x}Fe_2O_4$ ($0 \le x \le 1$) quaternary spinel ferrites was investigated using Raman and Mössbauer spectroscopies. From Mössbauer data an increase was observed in the orbital's overlapping of the ions in the A sites and a decrease in the B sites, resulting in IS changing. Also, it was verified that the HF field decreases with increasing x which was due to the replacement of the ions in the spinel structure. The Raman results showed the presence of zinc, iron, and magnesium ions in tetrahedral and octahedral sites.

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¹V. Sepelák, D. Baabe, F. J. Litterst, and K. D. Becker, J. Appl. Phys. 88, 5884 (2000).

²S. A. Mazen, S. F. Mansour, and H. M. Zaki, Cryst. Res. Technol. **38**, 471 (2003).

C. Upadhyay, H. C. Verma, and S. Anand, J. Appl. Phys. 95, 5746 (2004).
A. Franco, Jr., T. E. P. Alves, E. C. O. Lima, E. S. Nunes, and V. Zapf, Appl. Phys. A: Mater. Sci. Process. 94, 131 (2009).

⁵JCPDS Card No. 73–2211.

⁶B. D. Cullity, *Elements of X-ray Diffraction* (Addison-Wesley, Reading, 1978).

⁷A. R. Denton and N. W. Ashcroft, Phys. Rev. A **43**, 3161 (1991).

⁸C. M. Srivastava, S. N. Shringi, and R. G. Srivastava, Phys. Rev. B 14, 2041 (1976).

⁹J. Kreisel, G. Lucazeu, and H. Vincent, J. Solid State Chem. **137**, 127 (1998)

¹⁰F. Nakagomi, S. W. da Silva, V. K. Garg, A. C. de Oliveira, P. C. Morais, and A. Franco, Jr., J. Solid State Chem. 182, 2423 (2009).

¹¹M. J. Seong, M. C. Hanna, and A. Mascarenhas, Appl. Phys. Lett. 79, 3974 (2001).