

Quantum theories for hydrogen atom

(83)

PAGE No.	
DATE	

- Assume that: electron is stationary
- Symmetry suggests polar coordinates.
- For one electron in three dimensions,

$$\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} + \frac{2m}{\hbar^2} (E - V) \psi = 0 \quad \text{--- (1)}$$

$\nabla^2 \psi$

where E = total energy, $V = P.E. = \left(\frac{-e^2}{4\pi\epsilon_0 r} \right)$

- Let's convert to (r, θ, ϕ) , θ = zenith angle $\in [0, \pi]$
 ϕ = azimuth angle $\in [0, 2\pi]$

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} + \frac{2m^2 \sin^2 \theta}{\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0 r^2} + E \right) \psi = 0 \quad \text{--- (2)}$$

- We take here:

(a) $\psi \rightarrow$ Normalisable

(b) ψ and its derivatives are continuous and single-valued functions.

Solving the eq. (2), will give three set of quantum numbers.

To define a particle's motion in three-dimensional box, we need three quantum numbers.

↑ [Bohr only considers one quantum model].

(because ψ must be zero at the walls of the box in the x, y, z -directions.)

~~##~~ and in hydrogen-atom, electron's motion is guided by inverse square electric field of the nucleus and not by the walls.

Separation of variables:-

$$\Rightarrow \Psi(r, \theta, \phi) = R(r) \cdot \Theta(\theta) \cdot \Phi(\phi)$$

\uparrow only depend on r \uparrow only on θ \uparrow only on ϕ .

• Substitute the values in eqn: (2):

$$\Rightarrow \frac{\sin^2 \theta}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \frac{\sin \theta}{\Theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) + \frac{1}{\phi} \frac{d^2 \Phi}{d\phi^2} + \frac{2mr^2 \sin^2 \theta}{\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) = 0.$$

$$\Rightarrow \frac{\sin^2 \theta}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \frac{\sin \theta}{\Theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) = 1 \text{ (say)};$$

$$\text{hence } \boxed{1 + \frac{2mr^2 \sin^2 \theta}{\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) = -\frac{1}{\phi} \frac{d^2 \Phi}{d\phi^2}} \quad (3)$$

This equation can be correct only if both sides of it are equal to the same constant, since they are functions of different variables,

$$\boxed{-\frac{1}{\phi} \frac{d^2 \Phi}{d\phi^2} = m_\ell^2} \quad (4)$$

$$\Rightarrow \frac{\sin^2 \theta}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \frac{\sin \theta}{\Theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) + \frac{2mr^2 \sin^2 \theta}{\hbar^2} (\dots) = m_\ell^2.$$

⇒ Divide by $\sin^2 \theta$:

$$\frac{1}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \frac{2mr^2}{\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) = m_\ell^2 - \frac{1}{\Theta \sin \theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right)$$

Again same thing, both should be equal to some common constant, here, that constant is $l(l+1)$.

$$\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \left[\frac{2m}{\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0 r} - E \right) - \frac{l(l+1)}{r^2} \right] R = 0. \quad (5)$$

and

$$\frac{1}{\sin\theta} \frac{d}{d\theta} \left(\sin\theta \frac{d\Theta}{d\theta} \right) + \left(l(l+1) - \frac{m^2}{\sin^2\theta} \right) \Theta = 0 \quad (6)$$

Eq. 5 & 6 gives the respective separation of variable equations.