THE ROLE OF ICE COMPOSITIONS AND MORPHOLOGY FOR SNOWLINES AND THE C/N/O RATIOS IN ACTIVE DISKS

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ABSTRACT

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1. INTRODUCTION

Background info. Importance of volatiles in disks and planetary atmospheres, detections of snowlines in disks, C/O ratios etc. State again the importance of radial drift and gas accretion on the snowline locations, and that a systematic study of the combination of these two particular effects across the disk has not been done before. Then transition to the fact that we provide such a systematic study in Paper I and in this paper. Here, we expand the model of Paper I by making three additions: (1) we add N and CH4 in the static chemistry model, and explore how different abundances of CH4 and of the N main carriers (N2 and NH3) affect the C/O and N/O ratios, (2) we quantify the effect of radial drift and gas accretion on the N2, CH4 and NH3 snowline locations, and (3) we explore how different binding energies of CO and N2 affect their snowline locations.

2. COUPLED DRIFT-DESORPTION MODEL REVIEW

We begin with a brief review of Paper I's model for the effect of radial drift and viscous gas accretion on volatile snowline locations. We review our disk model assumptions and relevant timescales in Section 2.1, and summarize our findings in Section ??.

2.1. *Disk Model and Important Timescales*

We first assume a static disk, which is only irradiated by the central star and does not experience redistribution of solids or radial movement of the nebular gas. To quantify the effects of radial drift and gas accretion, we use a viscous disk with a spatially and temporally constant mass flux, $\dot{M}$. The viscous disk takes into account radial drift, gas accretion onto the central star, as well as accretion heating. We prefer this disk model to an irradiated or evolving disk (see Paper I) because it includes all the dynamical and thermal processes we are interested in for the scope of this paper, and therefore it is the most realistic one.

We model the static disk as a minimum mass solar nebula (MMSN), using a prescription for the gas surface density, Σ, and disk midplane temperature, T, similar to that of Chiang & Youdin (2010):

$$\Sigma = 2000\,(r/\mathrm{AU})^{-1}\ \mathrm{g\ cm^{-2}} \tag{1a}$$

$$T = 120\,(r/\mathrm{AU})^{-3/7}\ \mathrm{K}, \tag{1b}$$

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where r is the semimajor axis. Based on observations of protoplanetary disks (Andrews et al. 2010), we choose a flatter surface density than the one assumed by Chiang & Youdin (2010), where $\Sigma \propto r^{-3/2}$.

We use the Shakura & Sunyaev (1973) steady-state disk solution to model the viscous disk. Solving the Equation set of Appendix A in Paper I yields an expression for the temperature profile in a steady-state disk:

$$T_{\mathrm{act}} = \frac{1}{4r}\left(\frac{3G\kappa_0 \dot{M}^2 M_* \mu m_{\mathrm{p}} \Omega_{\mathrm{k}}}{\pi^2 \alpha k_{\mathrm{B}} \sigma}\right)^{1/3}. \tag{2}$$

Here G is the gravitational constant, $\kappa_0 = 2 \times 10^{-6}$ is a dimensionless coefficient in the opacity law $\kappa = \kappa_0 T_{\mathrm{act}}^2$, $M_* = M_\odot$ is the mass of the central star, $\mu = 2.35$ is the mean molecular weight of the nebular gas, m_{p} is the proton mass, $\Omega_{\mathrm{k}} = \sqrt{GM_\odot/r^3}$ is the Keplerian angular velocity, $\alpha = 0.01$ is a dimensionless coefficient (see below for details), k_{B} is the Boltzmann constant, and σ is the Stefan-Boltzmann constant. The final midplane temperature profile is computed as

$$T^4 = T_{\mathrm{act}}^4 + T_{\mathrm{irr}}^4, \tag{3}$$

where $T_{\mathrm{irr}} = T$ from Equation (1b). We use this expression because in addition to accretion heating, stellar irradiation also contributes to the disk thermal structure.

The steady-state disk has an α-viscosity prescription, where the kinematic viscosity is $\mu = \alpha c H$. Here $c \equiv \sqrt{k_{\mathrm{B}}T/(\mu m_{\mathrm{p}})}$ is the isothermal sound speed (with T from Equation 3), and $H \equiv c/\Omega_{\mathrm{k}}$ is the disk scale height. We can then determine the gas surface density for a viscous disk as (Shakura & Sunyaev 1973; see also Paper I for a more detailed explanation of these calculations):

$$\Sigma = \frac{\dot{M}}{3\pi\nu}. \tag{4}$$

We choose $\dot{M} = 10^{-8} M_\odot\ \mathrm{yr^{-1}}$, consistent with mass flux observations in disks (e.g., Andrews et al. 2010).

In what follows, we summarize the relevant timescales that affect volatile snowline locations.

Desorption timescale. Following **?** and Paper I,

Review disk models, desorption model, relevant timescales. State that we use a steady-state disk for the coupled drift-desorption evolution, since it is the most realistic, therefore only summarize the static and steady-state (viscous) disk. Summarize the findings of Paper I, i.e. particles of certain sizes desorb instantaneously and at a fixed particle size dependent location.
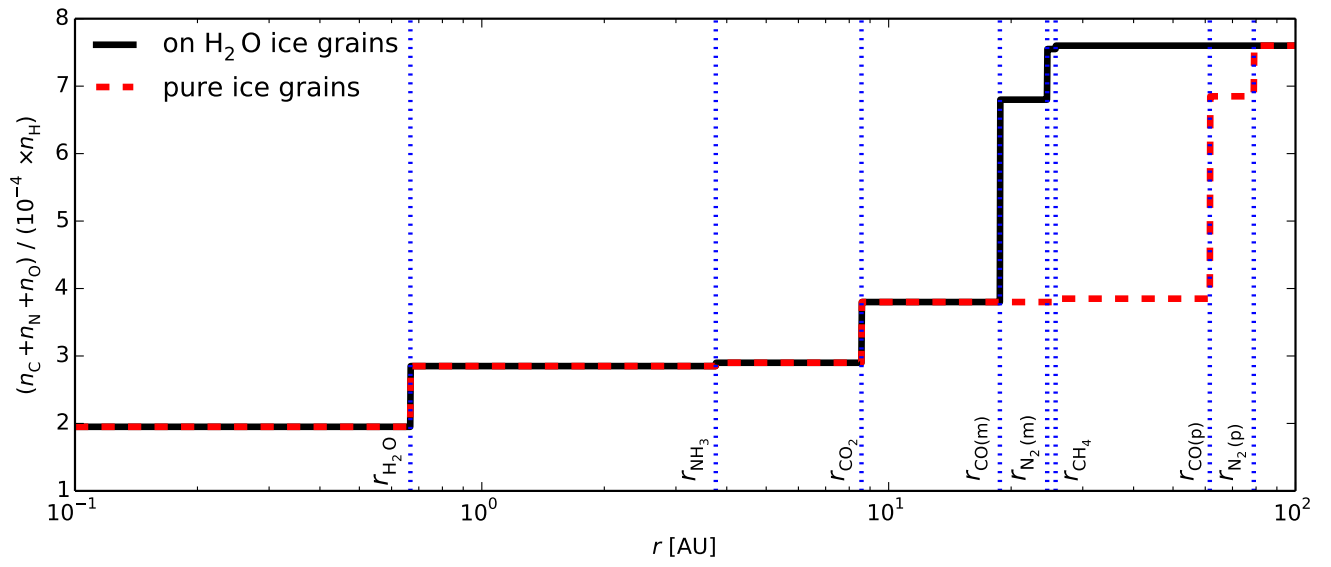
3. CH4 AND C/O RATIOS

FIG. 1.— CNO abundance in grains...

Discuss observed abundances for CH4 and the choices that we make (no CH4, median value, maximum value). State that desorption energies for H2O, CO2 and CH4 are well constrained experimentally, and that the CO2 and CH4 binding energies are only weakly dependent on whether it's pure CO2/CH4 or combined with H2O, but that is not the case for CO (and N2 as we will show in the next section). Present new binding energies for CO as pure ice and mixed with water. Show Figure 1 and discuss how different CH4 abundances and binding energies affect snowline locations and C/O ratio: CO-H2O mixture (though I think it's rather CO layered on top of H2O) moves the CO snowline inward by ~40 AU (will calculate percentages too); the maximum reasonable abundance of CH4 changes the C/O ratio by less than 10%. Show Figure 2 and quantify the effect of drift and accretion on the CH4 snowline compared to a static disk. While CH4 has only a modest effect on the C/O ratio in a static disk, this effect may be larger in a viscous disk, as the C gas abundances inside the CH4 snowline may be enhanced due to the differential motion of the desorbed ices and overall nebular gas (refer to Paper I). In this study, however, we neglect these effects and therefore do not include CH4 in estimating the C/O ratio (as an aside, the figures that include CH4 in the C/O ratio with drift and desorption are quite messy due to snowlines overlapping). Show Figure 3 and estimate the difference between CO-H2O and CO pure ice snowlines in the case of drift and accretion, as well as the comparisons for the static disk for the CO snowline.

4. NITROGEN CARRIERS AND N/O RATIOS

Similar to the previous section, but with more details. Discuss that nitrogen is abundant in the solar system and disks and primarily found as N2. Due to the high volatility of N2, the gas phase N/O ratio in the outer disk may be even more enhanced than the C/O ratio. A fraction of the nitrogen abundance may be also carried by NH3. Discuss NH3 observed abundances and the choices that we make (no NH3, median, maximum). State that the NH3 desorption energy is only weakly dependent on whether it's pure NH3 or combined with H2O, but that is not the case for N2. Present new binding energies for N2 as pure ice and combined with water. Show Figure 4 and discuss how different nitrogen abundances and binding energies affect snowline locations and N/O ratio: N2 combined with H2O moves the N2 snowline inward by ~50 AU (will calculate percentages too); the maximum reasonable abundance of NH3 changes the N/O ratio by ~15%. In the outer disk, the N/O ratio is enhanced by a factor of ~4 compared to the solar value, twice as much as the C/O enhancement. Show Figure 5 and quantify the effect of drift and accretion on the NH3 snowline compared to a static disk. While NH3 does not have a significant effect on the N/O ratio in a static disk, this effect may be larger in a viscous disk, as the N gas abundance inside the NH3 snowline may be enhanced due to the differential motion of the desorbed ices and overall nebular gas (refer to Paper I). In this study, however, we neglect these effects and therefore do not include NH3 in estimating the N/O ratio (again, N/O ratio figure with drift is quite messy when including NH3 and does not add any information that is not already shown in Figures 4 and 5). Show Figure 6 and estimate the difference between N2-H2O and N2 pure
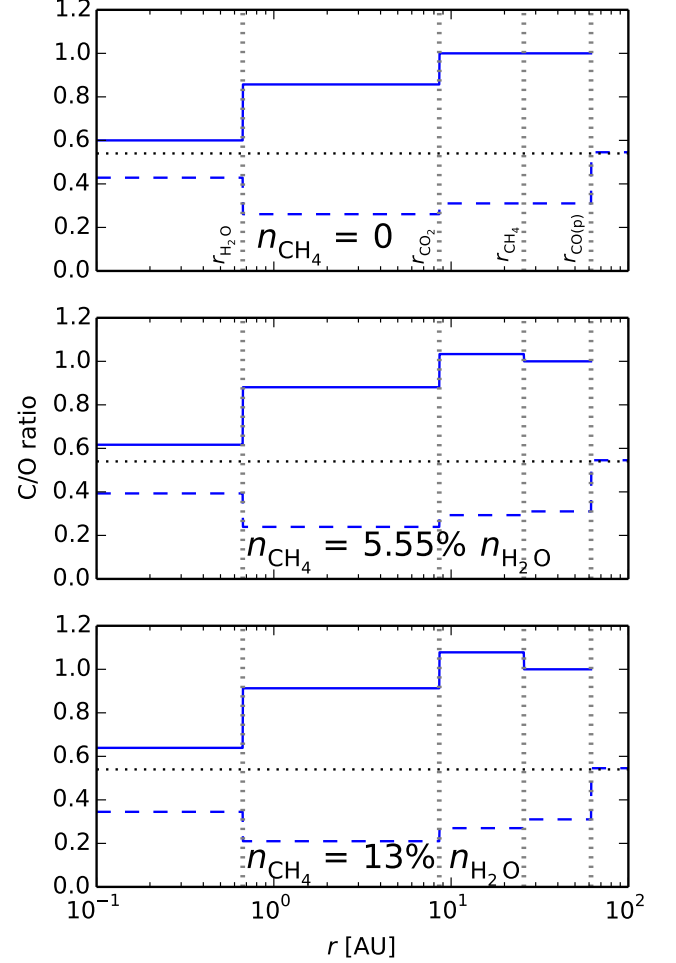

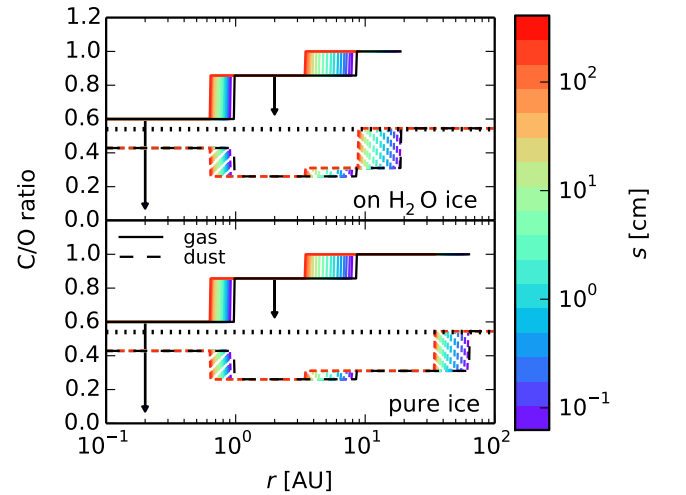FIG. 2.— C/O ratio in a static disk for different CH4 abundances and CO binding energies...

FIG. 3.— C/O ratio as function of semimajor axis for CO combined with H2O (top panel) and pure CO ice (bottom panel).... Drift and gas accretion move the CO snowlines inward by x% and y%, respectively.

ice snowlines in the case of drift, as well as the comparisons for the static disk for the N2 snowline. State that there will be an overabundance of gas-phase N/O between the CO and N2 snowlines, as there is no oxygen gas in this region.
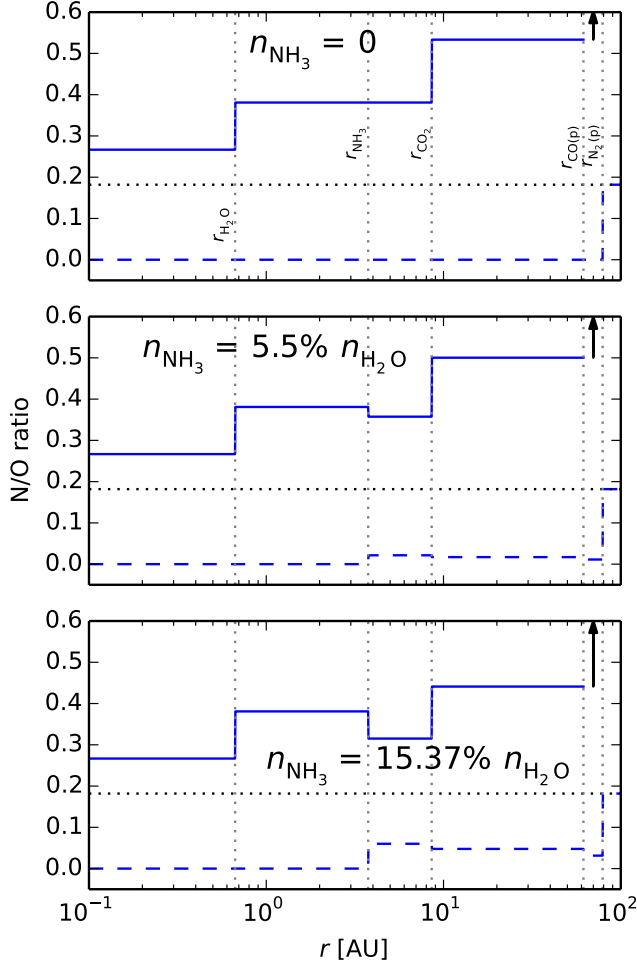
5. DISCUSSION

Discuss how entrapment of volatiles by H2O affects volatile abundances and C/O ratios. Re-emphasize the fact that the C/O and N/O ratios are upper estimates, and that CH4 and NH3 might matter in a viscous disk. State that we plan to address this in a future paper. More TBD.

FIG. 4.— N/O ratio in a static disk for different NH3 abundances and N2 binding energies...

6. SUMMARY

Maybe we can include the summary in the discussion section?

REFERENCES

Andrews, S. M., Wilner, D. J., Hughes, A. M., Qi, C., & Dullemond, C. P. 2010, ApJ, 723, 1241

Chiang, E. & Youdin, A. N. 2010, Annual Review of Earth and Planetary Sciences, 38, 493

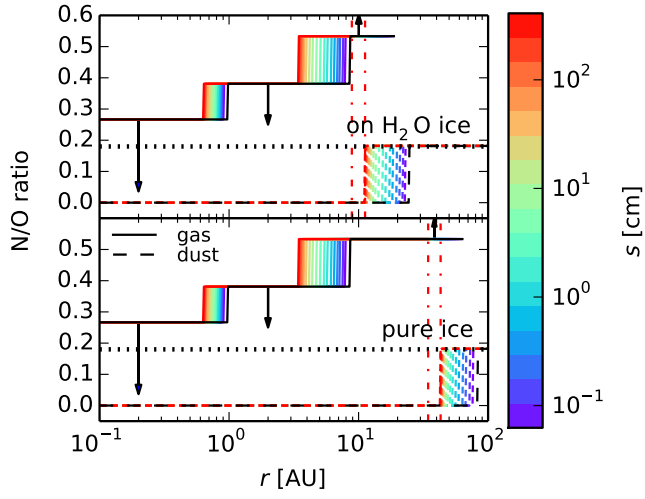Shakura, N. I. & Sunyaev, R. A. 1973, A&A, 24, 337

FIG. 5.— N/O ratio as function of semimajor axis for N2 combined with H2O (top panel) and pure N2 ice (bottom panel)....
Drift and gas accretion move the N2 snowlines inward by x% and y%, respectively. Overabundance of gas-phase N/O between the CO and N2 snowlines, marked by the vertical red dash-dotted lines for the largest drifting particles in our model.