

THE ROLE OF ICE COMPOSITIONS FOR SNOWLINES AND THE C/N/O RATIOS IN ACTIVE DISKS

ANA-MARIA A. PISO^{1,2}, JAMILA PEGUES^{1,3}, KARIN I. ÖBERG¹

Draft version October 10, 2016

ABSTRACT

The elemental compositions of planets define their chemistry, and could potentially be used as beacons for their formation location if the elemental gas and grain ratios of planet birth environments, i.e. protoplanetary disks, are well understood. In disks, the ratios of volatile elements, such as C/O and N/O, are regulated by the abundance of the main C, N, O carriers, their ice binding environment, and the presence of snowlines of major volatiles at different distances from the central star. We explore the effects of disk dynamical processes, molecular compositions and abundances, and ice compositions on the snowline locations of the main C, O and N carriers, and the C/N/O ratios in gas and dust throughout the disk. The gas-phase N/O ratio enhancement in the outer disk (exterior to the H₂O snowline) exceeds the C/O ratio enhancement for all reasonable volatile compositions. Ice compositions and disk dynamics individually change the snowline location of N₂, the main nitrogen carrier, by a factor of 2-3, and when considered together the range of possible N₂ snowline locations is $\sim 11\text{--}79$ AU in a standard disk model. Observations that anchor snowline locations at different stages of planet formation are therefore key to develop C/N/O ratios as a probe of planet formation zones.

1. INTRODUCTION

The chemical composition of protoplanetary disks is largely dictated by the freeze-out of volatile species. The snowline locations of volatile molecules are therefore crucial in determining disk chemical abundances in gas and dust, as well as planet compositions.

Carbon and oxygen bearing molecules, such as H₂O, CO₂ and CO, as well as the carbon-to-oxygen (C/O) ratio in protoplanetary disks and in giant planet atmospheres have been extensively studied from a theoretical standpoint (Öberg et al. 2011b, Ali-Dib et al. 2014, Madhusudhan et al. 2014, Mollière et al. 2015), and snowlines of volatiles such as H₂O and CO have been detected (Zhang et al. 2013, Qi et al. 2013). However, disk chemistry involves many other molecular compounds (Henning & Semenov 2013) including nitrogen bearing species and hydrocarbons (e.g., Mandell et al. 2012), which may affect the compositions of nascent planets.

Both in Solar system comets and in protoplanetary disks, volatile carbon and oxygen are primarily contained in H₂O, CO₂ and CO (e.g., Lodders 2003, Mumma & Charnley 2011, Öberg et al. 2011b, Boogert et al. 2015). However, some fraction of carbon may also be carried by CH₄ (e.g., Öberg et al. 2008), which may change the C/O ratio in gas and in dust at some disk locations. In the case of nitrogen, chemical models of the protostellar nebula (e.g., Owen et al. 2001) and of protoplanetary disks (e.g., Rodgers & Charnley 2002) suggest that N₂ was the dominant form of nitrogen, and that giant planets have accreted their nitrogen content primarily as N₂ (Mousis et al. 2014). Due to the high volatility of N₂, the gas phase nitrogen-to-oxygen (N/O) ratio in

the outer disk is expected to be high, perhaps more enhanced than the C/O ratio compared to the Solar value. The C/O ratio in gas cannot exceed unity (i.e., a factor of ~ 2 enhancement compared to the Solar value) since the major volatile carbon carrier is CO. In contrast, the N/O ratio mainly depends on the relative depletion of N₂ and oxygen carriers, and it will increase as each of the oxygen carrier snowlines (H₂O, CO₂, CO) is crossed. Beyond the CO snowline, there is no strict upper limit to the N/O ratio. The spatial extent of this latter region depends on the relative bond strengths of CO and N₂ to ice, but may be quite large (see Section 3). Giant planets that form at wide separations should thus have an excess of nitrogen in their atmospheres, which could be used to trace their formation origin. In addition to N₂, a fraction of the nitrogen abundance may also be carried by less volatile species such as NH₃ (Bottinelli et al. 2010, Mumma & Charnley 2011). The present day N₂ in Titan's atmosphere, for example, is thought to originate from accretion of primordial NH₃ (Atreya et al. 1978, Mandt et al. 2014).

The snowline locations of the main carbon, oxygen and nitrogen carriers strongly depend on the ice grain composition. Very volatile species, such as CO and N₂, present binding energies, and therefore snowline locations, that are sensitive to the details of the composition of the icy grain mantles. Spectroscopic observations suggest that up to 90% of the CO ice is frozen in a layer that is thick enough and separated from the H₂O ice layer underneath so that it can be considered pure ice (e.g., Pontoppidan et al. 2003). However, if the CO (or N₂) ice layer is thin enough (\sim monolayer coverage), then it will interact with the H₂O ice substrate (e.g., Collings et al. 2003). The ice binding energy is significantly larger in this water dominated environment than in the pure ice case, as shown by laboratory experiments (Collings et al. 2003, Öberg et al. 2005, Bisschop et al. 2006, Fayolle et al. 2016). This implies that ices in different environments

¹ Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138

² UCLA, 595 Charles E. Young Drive East, Los Angeles, CA 90095

³ Department of Astrophysical Sciences, Princeton University, Princeton, NJ 08544

will sublime at different radii, which will substantially change the disk regions where these volatiles are present in gaseous or solid form (see Section 3.2). In protostellar cores, H₂O ice is primarily amorphous (e.g., Williams & Herbst 2002, van Dishoeck et al. 2014). When temperatures exceed ~ 80 -90 K, H₂O ice acquires a crystalline structure (Schegerer & Wolf 2010). The CO binding energy is larger in an amorphous porous H₂O ice environment than in the amorphous compact or crystalline cases (Noble et al. 2012, Fayolle et al. 2016). The N₂ binding energy is also larger for a porous versus compact H₂O ice substrate. No equivalent studies have yet been performed for the deposition of N₂ on crystalline H₂O ice, but we expect the N₂ binding energy in this environment to follow the same trend as the CO binding energy, as CO and N₂ display a similar desorption behavior. To explore the range of distances at which CO and N₂ in different environments desorb, we consider the limiting cases: pure ices (lowest binding energy) and ices residing on an amorphous porous H₂O ice substrate (highest binding energy). We refer to the latter simply as water dominated ices unless noted otherwise.

In this work, we expand the coupled drift-desorption model developed in (Piso et al. 2015; hereafter Paper I) by considering additional volatile molecules and abundances, ice compositions, as well as nitrogen-to-oxygen (N/O) ratios. This paper is organized as follows. In Section 2, we review the drift-desorption model developed in Paper I. We discuss the effect of different abundances of the main carbon, oxygen and nitrogen carriers, grain compositions and disk dynamics on snowline locations and the C/N/O ratios in Section 3. We address the implications of our results in Section 4 and summarize our findings in Section 5.

2. COUPLED DRIFT-DESORPTION MODEL

We begin with a brief review of Paper I’s model for the effect of radial drift and viscous gas accretion on volatile snowline locations. We review our disk model in Section 2.1, and summarize our numerical method and results in Section 2.2.

2.1. Disk Model

In this work we consider both a static and a viscous disk. The static disk is irradiated by the central star and does not experience redistribution of solids or radial movement of the nebular gas. To quantify the effects of radial drift and gas accretion, we use a viscous disk with a spatially and temporally constant mass flux, \dot{M} . The viscous disk takes into account radial drift, gas accretion onto the central star, as well as accretion heating. We focus on this disk model which includes all the dynamical and thermal processes we are interested in for the scope of this paper, and do not further consider the other disk models presented in Paper I.

Following Chiang & Youdin (2010), the temperature profile for a static disk is

$$T_{\text{irr}} = 120 (r/\text{AU})^{-3/7} \text{ K}, \quad (1)$$

where r is the semimajor axis. We use the Shakura & Sunyaev (1973) steady-state disk solution to model the viscous disk. From Paper I, the viscous disk temperature

profile is computed as

$$T_{\text{visc}}^4 = \left[\frac{1}{4r} \left(\frac{3G\kappa_0\dot{M}M_*\mu m_p\Omega_k}{\pi^2\alpha k_B\sigma} \right)^{1/3} \right]^4 + T_{\text{irr}}^4. \quad (2)$$

Here G is the gravitational constant, $\kappa_0 = 2 \times 10^{-6}$ is a dimensionless opacity coefficient, $M_* = M_\odot$ is the mass of the central star, $\mu = 2.35$ is the mean molecular weight of the nebular gas, m_p is the proton mass, $\Omega_k = \sqrt{GM_\odot/r^3}$ is the Keplerian angular velocity, $\alpha = 0.01$ is a dimensionless coefficient (see below for details), k_B is the Boltzmann constant, and σ is the Stefan-Boltzmann constant.

The steady-state disk has an α -viscosity prescription, where the kinematic viscosity is $\nu = \alpha cH$. Here $c \equiv \sqrt{k_B T_{\text{visc}}/(\mu m_p)}$ is the isothermal sound speed and $H \equiv c/\Omega_k$ is the disk scale height. We can then determine the gas surface density for a viscous disk as (Shakura & Sunyaev 1973; see also Paper I for a more detailed explanation of these calculations):

$$\Sigma = \frac{\dot{M}}{3\pi\nu}. \quad (3)$$

We choose $\dot{M} = 10^{-8} M_\odot \text{ yr}^{-1}$, consistent with mass flux observations in disks (e.g., Andrews et al. 2010). As described in Paper I, the mass flux rate \dot{M} and stellar luminosity L_* will vary throughout the disk lifetime (Kennedy et al. 2006, Chambers 2009), in contrast with our simplified model which assumes that both quantities are constant. This effect will be most pronounced in the inner disk (\lesssim few AU), where accretion heating dominates. We thus acknowledge that the location of the H₂O snowline may be determined by the decline in \dot{M} or L_* with time, rather than radial drift (see Paper I, Section 2.1 for a more detailed explanation).

2.2. Desorption-Drift Equations and Results

The model is described in full in Paper I, here we review and summarize key concepts and results. For a range of initial icy grain sizes composed of a single volatile, we showed in Paper I that the timescale on which these particles desorb is comparable to their radial drift time, as well as to the accretion timescale of the nebular gas onto the central star. We thus have to take into account both drift and gas accretion when we calculate the disk location at which a particle desorbs, since that location may be different from the snowline position in a static disk for a given volatile (see Figure 1 and Öberg et al. 2011b). We determine a particle’s final location in the disk by solving the following coupled differential equations:

$$\frac{ds}{dt} = -\frac{3\mu_x m_p}{\rho_s} N_x R_{\text{des},x} \quad (4a)$$

$$\frac{dr}{dt} = \dot{r}, \quad (4b)$$

where s is the particle size, t is time, μ_x is the mean molecular weight of volatile x , $\rho_s = 2 \text{ g cm}^{-3}$ is the density of an icy particle, $N_x \approx 10^{15} \text{ sites cm}^{-2}$ is the number of adsorption sites of molecule x per cm^{-2} , $R_{\text{des},x}$ is the desorption rate of species x , and \dot{r} is the particle’s radial drift velocity. We calculate R_{des} and \dot{r} as follows.

The desorption rate $R_{\text{des},x}$ (per molecule) is (Hollenbach et al. 2009)

$$R_{\text{des},x} = \nu_x \exp(-E_x/T_{\text{grain}}), \quad (5)$$

where E_x is the adsorption binding energy in units of Kelvin, T_{grain} is the grain temperature (assumed to be the same as the disk temperature, see Paper I), and $\nu_x = 1.6 \times 10^{11} \sqrt{(E_x/\mu_x)} \text{ s}^{-1}$ is the molecule's vibrational frequency in the surface potential well. We discuss our choices for E_x for the different volatile species in Section 3.1.

Following Chiang & Youdin (2010) and Birnstiel et al. (2012), a particle's radial drift velocity can be approximated as

$$\dot{r} \approx -2\eta\Omega_k r \left(\frac{\tau_s}{1 + \tau_s^2} \right) + \frac{\dot{r}_{\text{gas}}}{1 + \tau_s^2}, \quad (6)$$

where the first term is the drift velocity in a non-accreting disk and the second term accounts for the radial movement of the gas. Here $\eta \approx c^2/(2v_k^2)$, where v_k is the Keplerian velocity, and $\tau_s \equiv \Omega_k t_s$ is the dimensionless stopping time:

$$t_s = \begin{cases} \rho_s s / (\rho c), & s < 9\lambda/4 \text{ Epstein drag} \\ 4\rho_s s^2 / (9\rho c \lambda), & s < 9\lambda/4, \text{Re} \lesssim 1 \text{ Stokes drag,} \end{cases} \quad (7)$$

where ρ is the disk mid-plane density, λ is the mean free path and Re is the Reynolds number. The gas accretion velocity \dot{r}_{gas} is determined from $\dot{M} = -2\pi r \dot{r}_{\text{gas}} \Sigma$, for a fixed \dot{M} and with Σ given by Equation (3).

For a particle of initial size s_0 , we solve the Equation set (4) with the initial conditions $s(t_0) = s_0$ and $r(t_0) = r_0$, where t_0 is the time at which we start the integration and r_0 is the particle's initial location. We stop our simulation after $t_d = 3 \text{ Myr}$, the disk lifetime, since this is roughly the timescale on which planets form, and determine the desorption timescale t_{des} from $s(t_{\text{des}}) = 0$, and thus a particle's desorption distance $r_{\text{des}} = r(t_{\text{des}})$. Our results are insensitive to our choice of t_0 as long as $t_0 \ll t_d$. We note that a particle's size is initially fixed and only changes due to desorption. We thus do not take into account processes such as grain coagulation or fragmentation, which nonetheless occur in disks (e.g., Birnstiel et al. 2012, Pérez et al. 2012). We discuss the effect of these processes on snowline locations in Paper I.

As we show in Paper I, a particle of initial size s_0 can experience three outcomes after $t_d = 3 \text{ Myr}$: (1) it can remain at its initial location, (2) it can drift towards the host star, then stop without evaporating significantly, and (3) it can completely desorb on a timescale shorter than 3 Myr. Particles in scenarios (1) and (2) are thus not affected by radial drift or gas accretion, and the snowline locations are those for a static disk. In contrast, the grains in case (3) desorb practically *instantaneously* and *at a fixed particle-size dependent location* in the disk, regardless of their initial position. The snowline locations for these particles will thus be fixed for a given initial particle size and disk model. We have found that grains with sizes $\sim 0.001 \text{ cm} \lesssim s \lesssim 7 \text{ m}$ satisfy this condition for our fiducial disk.

3. RESULTS

3.1. Snowlines in a Static Disk: The Importance of Ice Compositions

As we note in Section 1, the disk volatile composition and the ice composition determine the location of important snowlines. In this work we focus on the primary carbon, oxygen and nitrogen carriers, i.e. H_2O , CO_2 , CO , N_2 , and to a lesser extent, CH_4 and NH_3 . Our standard model is based on the median ice abundances observed toward Solar-type protostars (Öberg et al. 2011a), which are $n_{\text{CO}_2} = 0.29 \times n_{\text{H}_2\text{O}}$, $n_{\text{CO}} = 0.38 \times n_{\text{H}_2\text{O}}$, $n_{\text{CH}_4} = 0.0555 \times n_{\text{H}_2\text{O}}$ (hereafter CH_4 -mid) and $n_{\text{NH}_3} = 0.055 \times n_{\text{H}_2\text{O}}$ (hereafter NH_3 -mid). Here $n_{\text{H}_2\text{O}} \approx 10^{-4} \times n_{\text{H}}$ is the total water abundance (van Dishoeck 2006), with n_{H} the hydrogen abundance in the disk midplane. For CO , we also take into account that the observed CO ice only traces some of the CO reservoir due to its high volatility, and similarly to Öberg et al. (2011b) and Paper I we set the total CO abundance to $0.9 \times 10^{-4} n_{\text{H}}$. Finally, we assume that all nitrogen not found in NH_3 is in N_2 and assume a Solar nitrogen abundance, $n_{\text{N}} = 8 \times 10^{-5} n_{\text{H}}$ (Lodders 2003). In effect, this model assumes no chemical evolution between the protostellar and disk midplane stages. This is reasonable for material that accretes onto the disk at large radii (Visser et al. 2009), but may overestimate the contribution of the original volatiles to the total volatile budget in the innermost disk.

We determine the location of the H_2O , CO_2 , CO , CH_4 , N_2 and NH_3 snowlines in our static disk by balancing desorption with readsorption, following Hollenbach et al. (2009). The binding energies of H_2O , CO_2 , CO , CH_4 , N_2 and NH_3 as pure ices are 5800 K, 2000 K, 834 K, 1300 K, 767 K and 2965 K, respectively (Fraser et al. 2001, Collings et al. 2004, Fayolle et al. 2016, Garrod & Herbst 2006, Martín-Doménech et al. 2014). For our fiducial disk model, these energies correspond to disk temperatures of 143 K, 48 K, 21 K, 30 K, 18 K and 68 K, respectively. For CO and N_2 as water dominated ices, the binding energies are 1388 K and 1266 K, respectively (Fayolle et al. 2016), corresponding to disk temperatures of 34 K and 31 K, respectively. Figure 1 shows the resulting snowline locations, assuming CO and N_2 pure ices (top panel), and CO and N_2 in water dominated ices (bottom panel). The ordinate displays the total carbon, oxygen and nitrogen abundance in solids as a function of the hydrogen total abundance. As expected, the total grain abundance increases with semimajor axis, as more and more species freeze out. Freeze-out at the CO_2 and CO snowlines pulls more heavy elements into the grains than in the case of the H_2O snowline.

Figure 1, bottom panel, displays the snowline locations when CO and N_2 are in an amorphous porous water environment (see Section 1). The snowlines move outward by a factor of ~ 3.3 if the ices are pure, and by up to a factor of ~ 2 for an amorphous compact water substrate (not shown; our estimates are based on the results of Fayolle et al. 2016). The CO snowline moves outward by a factor of ~ 1.5 if CO is in a crystalline rather than an amorphous porous water environment (not shown; our estimates are based on the results of Noble et al. 2012), and a similar trend is expected for the N_2 snowline (see Section 1). **Based on recent observations (Luspay-Kuti et al. 2016) and theoretical models (Lectez**

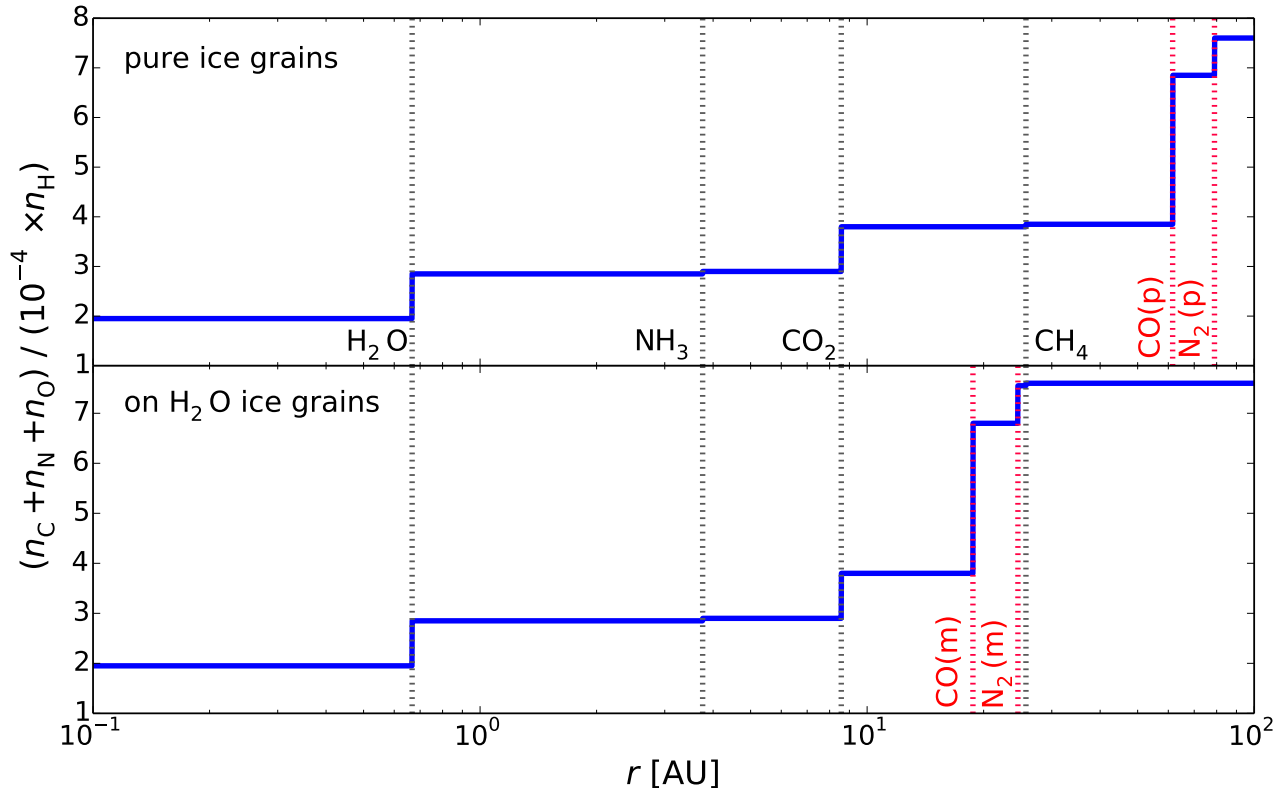


FIG. 1.— The total carbon, nitrogen and oxygen abundance in solids as a function of semimajor axis in a static disk, for CO and N₂ as pure ices (top panel) and water dominated ices (bottom panel). Relevant volatile snowlines are marked by the vertical dashed lines. The grain abundances are calculated as a function of the observed median CH₄ and NH₃ abundances in protostellar cores. The total grain abundance increases with semimajor axis as more and more species freeze out.

et al. 2015, Mousis et al. 2016) for the composition of comet 67P/Churyumov-Gerasimenko, a fraction of the disk CO (and possibly N₂) content may be trapped in clathrates. This would imply CO and N₂ snowline locations roughly corresponding to CO and N₂ in a crystalline H₂O ice environment as described above. The results of Figure 1 thus represent the limiting cases for the positions of the CO and N₂ snowlines, for different compositions of the icy grains and morphology of the water ice substrate. This variation in snowline locations changes the chemical abundances both in gas and dust throughout the disk, directly affecting the compositions of nascent giant planets forming in situ.

3.2. C/N/O Ratios in Static Disks

In this section we determine the C/O and N/O ratios in gas and dust throughout our static disk, and to what extent they are affected by the presence of CH₄ and NH₃ over the full range of observed CH₄ and NH₃ abundances toward low-mass protostars. In this section we only consider pure ices, i.e. ices that are layered on a silicate mantle. In reality, the icy grain will have a layered structure with volatiles residing on top of a H₂O ice substrate. However, if the volatile ice layer is thick enough and separated from the H₂O ice layer, the interaction between H₂O and the other volatile species will be minimal and thus the ices can be considered pure rather than water dominated (see also Section 1).

We explore the parameter space of possible CH₄ abundances by assuming three different scenarios: (1) no CH₄, (2) CH₄-mid, and (3) the maximum CH₄ observed abundance (hereafter CH₄-max), $n_{\text{CH}_4\text{-max}} = 0.13 \times n_{\text{H}_2\text{O}}$ (Öberg et al. 2008). Since the abundance of carbon grains is uncertain, we assume that all the carbon that is not in the form of CH₄, CO and CO₂ is in carbon grains, so that we reproduce the Solar C/O ratio (gas+dust) of 0.54.

Figure 2 shows the C/O ratio in gas and dust as a function of semimajor axis in a static disk: no CH₄ (top panel), CH₄-mid (middle panel) and CH₄-max (bottom panel). As in Öberg et al. (2011b) and Paper I, a gaseous C/O ratio of unity can be achieved between the CO₂ and CO snowlines, where oxygen gas is significantly depleted. The gas-phase C/O ratio may be further enhanced between the CO₂ and CH₄ snowlines due to the presence of additional carbon gas from CH₄. In this region, the C/O ratio increases by 3% for CH₄-mid and by 8% for CH₄-max, as displayed in the middle and bottom panels of Figure 2. Based on the range of observed CH₄ protostellar abundances, its presence in the disk only modestly affects the C/O ratio.

We assume that the main nitrogen-bearing species are N₂ and NH₃, since other volatiles that contain nitrogen have significantly lower abundances in comparison (e.g., Mumma & Charnley 2011). Similarly to the case of CH₄, we explore the parameter space of possible NH₃ abundances using observations toward low-mass protostars,

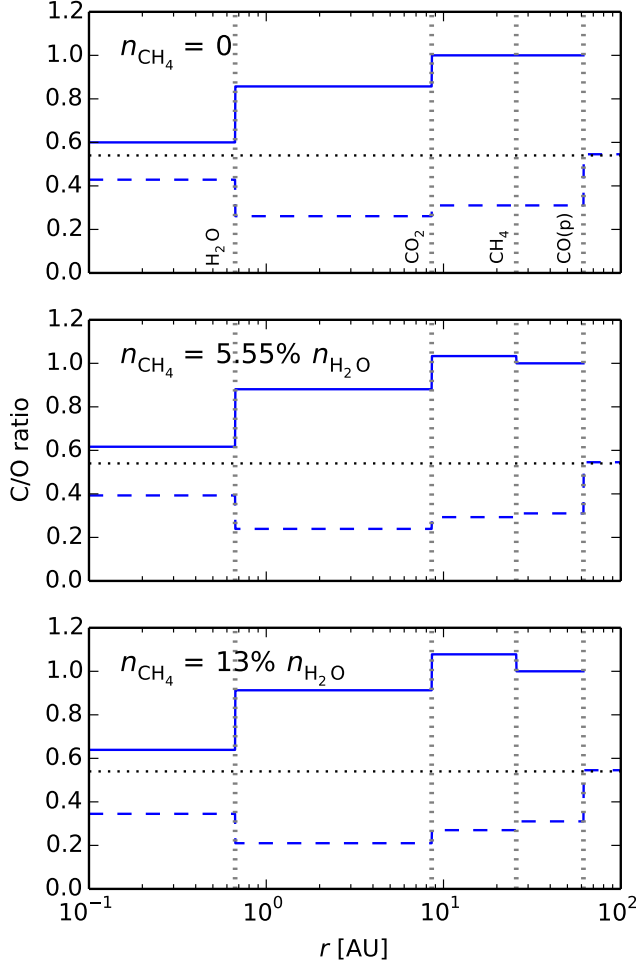


FIG. 2.— The C/O ratio in gas (solid lines) and dust (dashed lines) as a function of semimajor axis in a static disk, assuming no carbon is present in the form of CH_4 (top panel), the median observed CH_4 abundance is assumed (middle panel), and the maximum observed CH_4 abundance is assumed (bottom panel). The C/O estimates are performed assuming that the CO ices are in pure form. The vertical dotted lines mark the snowline locations of the main C and O carriers. The horizontal dotted lines represent the stellar C/O value. The presence of methane only modestly increases the C/O ratio in gas between the CO_2 and CH_4 snowlines.

as follows: (1) no NH_3 , (2) NH_3 -mid, and (3) the maximum observed NH_3 abundance $n_{\text{NH}_3-\text{max}} = 0.15 \times n_{\text{H}_2\text{O}}$ (Bottinelli et al. 2010). In each case, the N_2 abundance then simply follows as $n_{\text{N}_2} = (n_{\text{N}} - n_{\text{NH}_3})/2$.

Figure 3 shows the snowline locations of the main oxygen and nitrogen carriers and the N/O ratio in gas and dust as a function of semimajor axis in a static disk, for our three choices of the NH_3 abundance: no NH_3 (top panel), NH_3 -mid (middle panel) and NH_3 -max (bottom panel). For comparison, the horizontal dotted lines show the average N/O ratio in the disk. As expected, the gaseous N/O ratio generally exhibits an increasing trend towards the outer disk as more oxygen gas is depleted, with small decreases between the NH_3 and CO_2 snowlines (by 6% for NH_3 -mid and by 18% for NH_3 -max, respectively) due to NH_3 freeze-out. While the presence of NH_3 only moderately affects our results for the N/O ratio, NH_3 is important since otherwise the nitro-

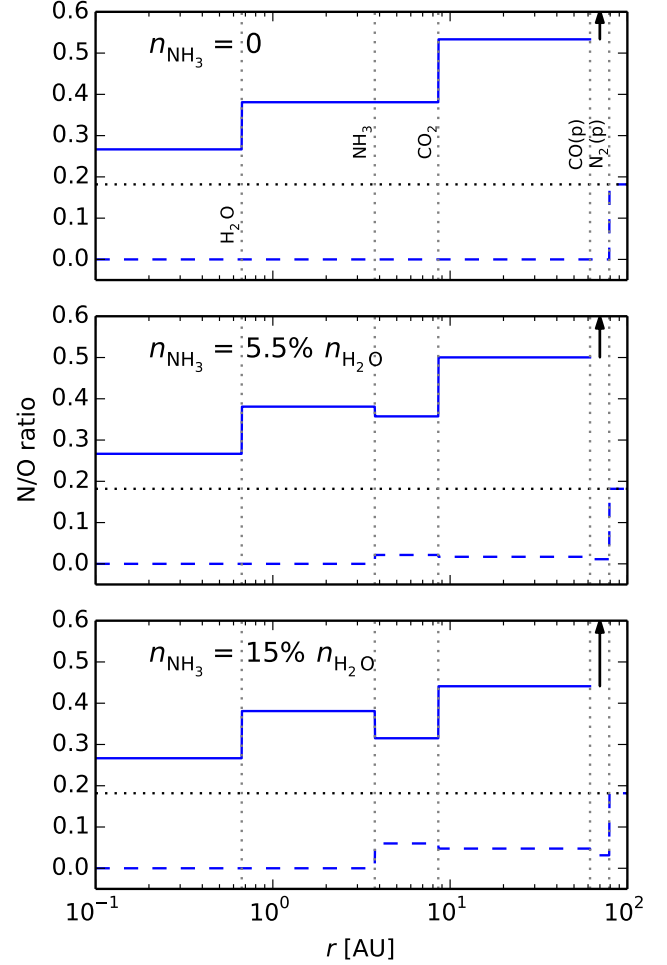


FIG. 3.— The N/O ratio in gas (solid lines) and dust (dashed lines) as a function of semimajor axis in a static disk, assuming no nitrogen is present in the form of NH_3 (top panel), the median observed NH_3 abundance is assumed (middle panel), and the maximum observed NH_3 abundance is assumed (bottom panel). The N/O estimates are performed assuming that the CO and N_2 ices are in pure form. The vertical dotted lines mark the snowline locations of the main C, O and N carriers. The horizontal dotted lines represent the average N/O value in the disk. The gas-phase N/O ratio is enhanced by a factor of two between the H_2O and CO_2 snowlines compared to its average value, and by a factor of three between the CO_2 and CO snowlines. The arrows mark a highly elevated N/O ratio in gas between the CO and N_2 snowlines due to the depletion of oxygen gas in this region. The presence of NH_3 moderately decreases the N/O ratio in gas between the NH_3 and CO_2 snowlines.

gen content in solid bodies would be more depleted than is observed for comets and asteroids (Wyckoff et al. 1991, Mumma & Charnley 2011, Bergin et al. 2015).

The gas-phase N/O ratio is enhanced by a factor of two outside the H_2O snowline compared to its average value, by more than a factor of three between the CO_2 and CO snowlines, and by orders of magnitude between the CO and N_2 snowlines. This latter region can span tens of AU depending on disk parameters and the relative CO and N_2 ice binding environment. This N/O enhancement is more pronounced than the C/O gas phase enhancement of a factor of two in the outer disk (see Figure 2).

3.3. C/N/O Ratios in Dynamic Disks

Here we use the model of Section 2 to estimate the movement of the CO and N₂ snowlines for different grain morphologies in a viscous disk. Figure 4 shows the H₂O, CO₂ and CO snowline locations for particles with initial sizes $\sim 0.05 \text{ cm} \lesssim s \lesssim 7 \text{ m}$ as well as estimates for the C/O ratio in gas and dust in a viscous disk, with the CO snowline calculated under different grain morphologies as noted above. We assume there is no carbon in the form of CH₄. The true snowline for particles that desorb outside the static snowline is the static snowline itself, hence desorbing particles with $s < 0.05 \text{ cm}$ do not form true snowlines. If the CO binding environment is known, the CO snowline moves inward by up to $\sim 50\%$ compared to a static disk for each case (pure and water dominated ices) due to disk dynamics. The full range of potential CO snowlines taking into account both ice compositions and disk dynamics span $\sim 8.7 \text{ AU}$ to $\sim 61 \text{ AU}$, which is a factor of ~ 7 difference. This implies that gas phase C/O ratios of order unity may be reached in the giant planet forming zone, and the CO snowline may be inside 10 AU for certain disk parameters.

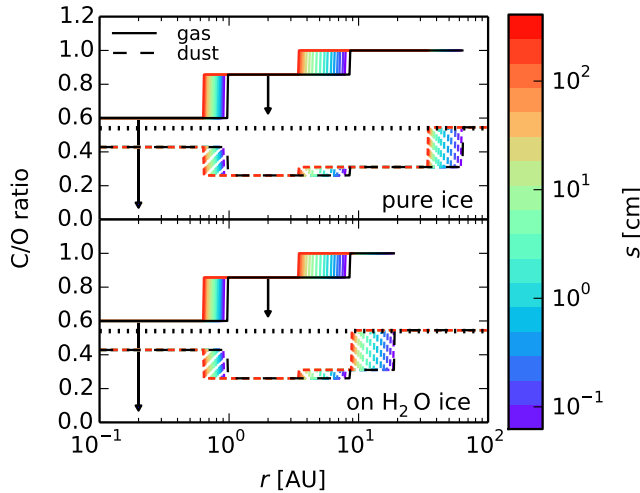


FIG. 4.— C/O ratio estimates in gas (solid lines) and dust (dashed lines) as function of semimajor axis in a viscous disk, for CO as pure ice (top panel) or as water dominated ices (bottom panel). The H₂O, CO₂ and CO snowlines are shown for particles with initial sizes $\sim 0.05 \text{ cm} \lesssim s \lesssim 7 \text{ m}$ as indicated by the color bar. The C/O ratio in a static disk (black lines) is shown for comparison. The arrows show that the C/O ratio in gas will decrease inside the H₂O and CO₂ snowlines in the viscous disk, as the relative fluxes of the desorbed icy particles and the overall nebular gas will cause an excess of oxygen gas inside these snowlines (see Paper I for details). The presence of CO in a water ice environment rather than as pure ice moves the CO snowline significantly inward by $\sim 70\%$. Taken together, disk dynamics and ice compositions move the CO snowline inward by a factor of ~ 7 .

Figure 5 shows the H₂O, CO₂, CO and N₂ snowline locations in a viscous disk for the same range of initial particle sizes as in Figure 4, and with the CO and N₂ snowlines calculated assuming different grain morphologies as explained above, as well as estimates for the N/O ratio throughout the disk. For simplicity, we assume that all nitrogen is the form of N₂. This choice is justified since the presence of some NH₃ only moderately changes

the N/O ratio (see Figure 3), and since we are primarily interested in the N₂ snowline locations rather than exact values for the N/O ratio. The innermost N₂ snowlines in the viscous disk, created by particles with $s \sim 7 \text{ m}$ for our fiducial model, are located at $r_{\text{N}_2, \text{pure}} \approx 42 \text{ AU}$ for N₂ as pure ice and at $r_{\text{N}_2, \text{water}} \approx 11 \text{ AU}$ for N₂ in water dominated ices. Thus for each case (pure versus water dominated ices), the N₂ snowline moves inward by up to 50% due to disk dynamics. By taking into account both ice compositions and disk dynamics, the full range of potential N₂ snowlines span ~ 11 to $\sim 79 \text{ AU}$, which is a factor of ~ 7 difference. Similarly to the case for CO, the N₂ snowline may be close to 10 AU for certain disk models.

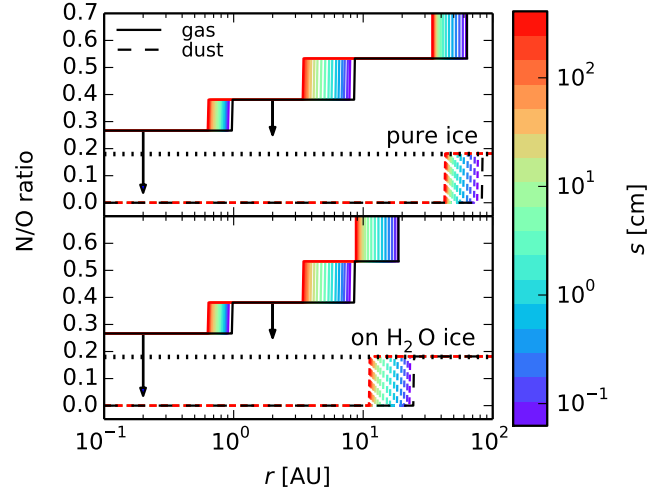


FIG. 5.— N/O ratio estimates in gas (solid lines) and dust (dashed lines) as function of semimajor axis in a viscous disk, for CO and N₂ as pure ices (top panel) or as water dominated ices (bottom panel). The H₂O, CO₂, CO and N₂ snowlines are shown for particles with initial sizes $\sim 0.05 \text{ cm} \lesssim s \lesssim 7 \text{ m}$ as indicated by the color bar. The N/O ratio in a static disk (black lines) is shown for comparison. The arrows show that the N/O ratio in gas will decrease inside the H₂O and CO₂ snowlines in the viscous disk, as the relative fluxes of the desorbed icy particles and the overall nebular gas will cause an excess of oxygen gas inside these snowlines (see Paper I for details). Radial drift and gas accretion move the N₂ snowline inward by up to $\sim 50\%$ compared to a static disk. The presence of N₂ in a water ice environment rather than as pure ice moves the N₂ snowline significantly inward by $\sim 70\%$. Taken together, disk dynamics and ice compositions move the N₂ snowline inward by a factor of ~ 7 . The results of an enhanced gas-phase N/O ratio between the H₂O and CO snowlines compared to its average value, and of highly elevated N/O ratios in gas between the CO and N₂ snowlines (see Figure 3), are preserved.

4. DISCUSSION

This study shows that the gas-phase N/O ratio in protoplanetary disks is considerably enhanced throughout most of the disk midplane compared to its average value. As demonstrated in Figure 6, the gaseous N/O ratio is enhanced by a factor of two beyond the H₂O snowline, by more than a factor of three between the CO₂ and CO snowlines, and by several orders of magnitude between the CO and N₂ snowlines. Thus constraining the N/O ratio in a giant planet atmosphere could be used to trace its formation origins.

Theoretical models of the magnitude and role of N/O (and N/C) ratios in exoplanet atmospheres are needed in order to use these ratios as probes for a planet’s formation location. Models that explore the effect of varying the C/O ratio in exoplanet atmospheres exist in literature, and they display a large and observable effect on gas giant envelope chemistry (Lodders 2009, Mollière et al. 2015). However, no similar model explorations exist for the effect of N/O and C/N/O ratios, and both are needed to exploit this potential constraint. Given the existence of such theoretical models, measurements of the N/O ratio in planetary envelopes may be possible to infer from atmospheric compositions of nitrogen versus carbon and oxygen bearing species. Nitrogen carriers have not been targeted so far due to lack of instrument sensitivity, but such observations and detections are likely in the near future with the advent of JWST (e.g., NH₃, Greene et al. 2016). The N/O ratio enhancement is larger than that of the gas phase C/O ratio throughout most of the disk. Thus measurements of an enhanced C/O ratio in an exoplanet atmosphere could be corroborated (disproved) by measurements of enhanced (non-enhanced) N/O ratios. Moreover, Figure 6 shows that giant planets that have formed in situ between the H₂O and CO snowlines are expected to present elevated both C/O and N/O ratios in their atmospheres, whereas planets between the CO and N₂ snowlines will have a highly enhanced N/O ratio in their atmospheres, but not C/O.

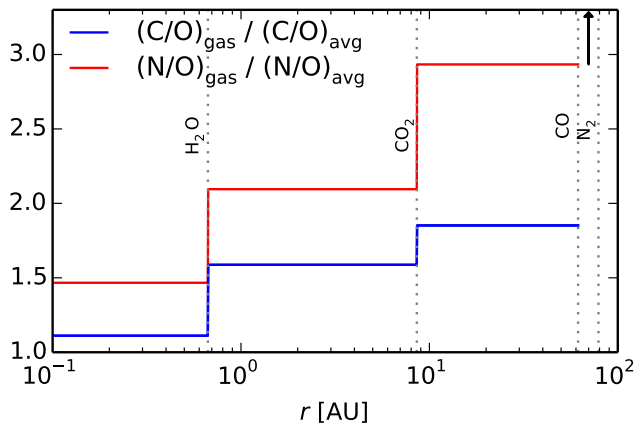


FIG. 6.— Gas phase C/O (blue curve) and N/O (red curve) ratios divided by the average C/O and N/O ratio in a static disk, assuming CO and N₂ are pure ices, and there is no CH₄ or NH₃. The dashed vertical lines mark the H₂O, CO₂, CO and N₂ snowlines. The arrow indicates that the N/O ratio is enhanced by orders of magnitude compared to its average value between the CO and N₂ snowlines. The gaseous N/O ratio is enhanced throughout most of the disk, and more enhanced than the C/O ratio.

Due to disk dynamics and ice compositions, the locations of the CO and N₂ snowlines, and thus the disk regions with highly elevated gas phase N/O and C/O ratios, are uncertain and may span tens of AU. Both ice morphologies discussed in this study, pure and water dominated ices, are plausible in protoplanetary disks and depend on whether H₂O and CO ices formed on similar timescales or successively (e.g., Garrod & Pauly 2011). Observations of protostellar cores show that a large fraction of CO is bound in a pure ice multilayer (Pontoppidan

et al. 2003), but theoretical models also suggest an icy mantle structure where CO resides on a H₂O ice layer (e.g., Collings et al. 2003). One can also imagine a scenario where CO is in a water binding environment and N₂ is not. This could be attributed to the fact that H₂O may bind preferentially to CO than N₂, since both H₂O and CO are polar molecules while N₂ is not. It is also possible for N₂ ices to form later than CO (e.g., Pagani et al. 2012), and thus be deposited on the outer layers of the icy mantles which are typically water poor (e.g., Garrod & Pauly 2011). The impact of the ice environment on the snowline location is much smaller in the case of CO₂ and NH₃, as their binding energies and behavior are closer to that of H₂O. No detailed measurements for the CH₄ binding energy in a water environment exist so far, but due to its low desorption temperature a similar behavior to that of CO and N₂ would be expected. While the presence of some carbon in the form of CH₄ only modestly affects our results, CH₄ may become important in disks where a large fraction of the CO abundance has been converted into hydrocarbons (e.g., Du et al. 2015).

Changes in stellar luminosity (e.g., Kennedy et al. 2006) and gas mass accretion rate (e.g., Chambers 2009), as well as the evolution of icy dust particles due to grain growth and fragmentation (e.g., Birnstiel et al. 2012), may introduce additional uncertainties in the snowline locations, and thus the C/N/O ratios. Moreover, the diffusion of vapor across the snowlines following the cold finger effect (Stevenson & Lunine 1988, Cyr et al. 1998) will change the shape of the C/O and N/O curves and therefore the magnitude of the C/N/O ratios between different snowlines. The effect of dynamical processes on snowline locations is discussed in more detail in Paper I, Section 5.2. Given the number of uncertainties in snowline locations, detections of snowlines in a sample of disks at different evolutionary stages are needed to provide observational constraints on the relative importance of ice compositions and disk dynamics in setting snowline locations. The uncertainties in snowline locations caused by disk dynamics, ice compositions, and other effects outlined above can be resolved in extreme cases, such as a detection of a CO snowline at a temperature corresponding to pure CO ice desorption in a static disk (e.g., Qi et al. 2013 at ~ 17 K) or CO desorption from a water dominated ice in a dynamic disk. In intermediate cases it is more difficult to resolve the relative importance of ice compositions and disk dynamics. For example, the CO snowline in HD 163296 is at a higher temperature of ~ 25 K (Qi et al. 2015), which could be caused either by CO being in a water dominated environment or by dynamical effects that push the CO snowline inward. Detections of multiple snowlines in the same disk could potentially break this degeneracy.

Uncertainties in snowline locations of this magnitude also affect interpretations of Solar system observations. Recent measurements of nitrogen abundance in comet 67P/Churyumov-Gerasimenko found a N₂/CO ratio ~ 10⁻³ (Rubin et al. 2015). A low N₂/CO ratio is consistent with comets having formed inside the N₂ snowline where N₂ is still in the gas phase. However, it is also possible that the measured N₂ abundance in 67P may be due to post-formation processes such as radiogenic heating (Rubin et al. 2015), and thus may not reflect

the comet's primordial composition. Theoretical models suggest that Jupiter-family comets, such as 67P, originate from the Kuiper belt (Duncan & Levison 1997; but see Rubin et al. 2015 for alternative formation scenarios for 67P). It is thus possible, in principle, to use measurements of the N_2 abundance in Jupiter-family comets to determine where the N_2 snowline was located in our Solar system. However due to the uncertainty in the calculated location of the N_2 snowline (see Section 3.3), as well as the uncertainty of the formation zone of Jupiter-family comets (anywhere between 5 and >30 AU; Pontoppidan et al. 2014), more detailed modeling is needed.

5. SUMMARY

In this paper we explore the role of icy grain compositions and disk dynamics on the snowline locations of major volatile carrier molecules and the C/N/O ratios in protoplanetary disks. We enhance the coupled drift-desorption model developed in Piso et al. (2015) by adding more carbon- and nitrogen-bearing species into our framework, and by considering different binding ice environments. Our results can be summarized as follows:

1. Due to the high volatility of N_2 , the gaseous N/O ratio outside the H_2O snowline is enhanced by a factor of two compared to its average value, by more than a factor of three between the CO_2 and CO snowlines, and by many orders of magnitude between the CO and N_2 snowlines due to the complete depletion of oxygen gas in this region. This enhancement is more pronounced than in the case of the gas-phase C/O ratio, which is increased by at most a factor of two compared to the stellar value.
2. The effect of CH_4 and NH_3 on the C/O and N/O ratios is small, even when we consider the maximum observed CH_4 and NH_3 abundances in proto-

stellar cores. In this scenario, the gas phase C/O ratio increases by 8% between the CO_2 and CH_4 snowlines, and the gaseous N/O ratio decreases by 18% between the NH_3 and CO_2 snowlines. In both cases, large gas phase C/O and N/O ratios in the outer disk are preserved.

3. Grain composition sensitively affects the CO and N_2 snowline locations. If CO and N_2 reside in water dominated rather than pure ices, their snowlines move inward by up to $\sim 70\%$. This effect is separate from that of radial drift and viscous gas accretion, which also cause an inward movement of the CO and N_2 snowlines by up to $\sim 50\%$.
4. The locations of the CO and N_2 snowlines are uncertain when we consider both viscous versus static disks, and pure versus water dominated ices. The snowlines in a viscous disk with CO or N_2 in a water environment are by up to a factor of ~ 7 closer to the host star than in a static disk with CO or N_2 as pure ices.

Our results have direct consequences for the composition of nascent giant planets. The considerable inward movement of the CO and N_2 snowlines due to the ice grains being water dominated rather than pure ices implies that giant planets with high C/O and/or N/O ratios in their atmospheres may form closer in than previously predicted by theoretical models. Moreover, our model shows that wide separation gas giants may have an excess of nitrogen in their envelopes, which may be used to trace their origins.

This work is supported by a Simons Collaboration on the Origins of Life (SCOL) investigator award to KIO. JP was supported by the Bankeer institute.

REFERENCES

- Ali-Dib, M., Mousis, O., Petit, J.-M., & Lunine, J. I. 2014, *ApJ*, 785, 125
- Andrews, S. M., Wilner, D. J., Hughes, A. M., Qi, C., & Dullemond, C. P. 2010, *ApJ*, 723, 1241
- Atreya, S. K., Donahue, T. M., & Kuhn, W. R. 1978, *Science*, 201, 611
- Bergin, E. A., Blake, G. A., Ciesla, F., Hirschmann, M. M., & Li, J. 2015, *Proceedings of the National Academy of Science*, 112, 8965
- Birnstiel, T., Klahr, H., & Ercolano, B. 2012, *A&A*, 539, A148
- Bisschop, S. E., Fraser, H. J., Öberg, K. I., van Dishoeck, E. F., & Schlemmer, S. 2006, *A&A*, 449, 1297
- Boogert, A. C. A., Gerakines, P. A., & Whittet, D. C. B. 2015, *ARA&A*, 53, 541
- Bottinelli, S., Boogert, A. C. A., Bouwman, J., et al. 2010, *ApJ*, 718, 1100
- Chambers, J. E. 2009, *ApJ*, 705, 1206
- Chiang, E., & Youdin, A. N. 2010, *Annual Review of Earth and Planetary Sciences*, 38, 493
- Collings, M. P., Anderson, M. A., Chen, R., et al. 2004, *MNRAS*, 354, 1133
- Collings, M. P., Dever, J. W., Fraser, H. J., & McCoustra, M. R. S. 2003, *Ap&SS*, 285, 633
- Cyr, K. E., Sears, W. D., & Lunine, J. I. 1998, *Icarus*, 135, 537
- Du, F., Bergin, E. A., & Hogerheijde, M. R. 2015, *ArXiv e-prints*, arXiv:1506.03510
- Duncan, M. J., & Levison, H. F. 1997, *Science*, 276, 1670
- Fayolle, E. C., Balle, J., Loomis, R., et al. 2016, *ApJ*, 816, L28
- Fraser, H. J., Collings, M. P., McCoustra, M. R. S., & Williams, D. A. 2001, *MNRAS*, 327, 1165
- Garrod, R. T., & Herbst, E. 2006, *A&A*, 457, 927
- Garrod, R. T., & Pauly, T. 2011, *ApJ*, 735, 15
- Greene, T. P., Line, M. R., Montero, C., et al. 2016, *ApJ*, 817, 17
- Henning, T., & Semenov, D. 2013, *Chemical Reviews*, 113, 9016
- Hollenbach, D., Kaufman, M. J., Bergin, E. A., & Melnick, G. J. 2009, *ApJ*, 690, 1497
- Kennedy, G. M., Kenyon, S. J., & Bromley, B. C. 2006, *ApJ*, 650, L139
- Lectez, S., Simon, J.-M., Mousis, O., et al. 2015, *ApJ*, 805, L1
- Lodders, K. 2003, *ApJ*, 591, 1220
- . 2009, *ArXiv e-prints*, arXiv:0910.0811
- Luspay-Kuti, A., Mousis, O., Hässig, M., et al. 2016, *Science Advances*, 2, 1501781
- Madhusudhan, N., Amin, M. A., & Kennedy, G. M. 2014, *ApJ*, 794, L12
- Mandell, A. M., Bast, J., van Dishoeck, E. F., et al. 2012, *ApJ*, 747, 92
- Mandt, K. E., Mousis, O., Lunine, J., & Gautier, D. 2014, *ApJ*, 788, L24
- Martín-Doménech, R., Muñoz Caro, G. M., Bueno, J., & Goesmann, F. 2014, *A&A*, 564, A8
- Mollière, P., van Boekel, R., Dullemond, C., Henning, T., & Mordasini, C. 2015, *ApJ*, 813, 47
- Mousis, O., Fletcher, L. N., Lebreton, J.-P., et al. 2014, *Planet. Space Sci.*, 104, 29
- Mousis, O., Lunine, J. I., Luspay-Kuti, A., et al. 2016, *ApJ*, 819, L33
- Mumma, M. J., & Charnley, S. B. 2011, *ARA&A*, 49, 471
- Noble, J. A., Congiu, E., Dulieu, F., & Fraser, H. J. 2012, *MNRAS*, 421, 768
- Öberg, K. I., Boogert, A. C. A., Pontoppidan, K. M., et al. 2008, *ApJ*, 678, 1032
- . 2011a, *ApJ*, 740, 109
- Öberg, K. I., Murray-Clay, R., & Bergin, E. A. 2011b, *ApJ*, 743, L16

- Öberg, K. I., van Broekhuizen, F., Fraser, H. J., et al. 2005, *ApJ*, 621, L33
- Owen, T., Mahaffy, P. R., Niemann, H. B., Atreya, S., & Wong, M. 2001, *ApJ*, 553, L77
- Pagani, L., Bourgoïn, A., & Lique, F. 2012, *A&A*, 548, L4
- Pérez, L. M., Carpenter, J. M., Chandler, C. J., et al. 2012, *ApJ*, 760, L17
- Piso, A.-M. A., Öberg, K. I., Birnstiel, T., & Murray-Clay, R. A. 2015, *ApJ*, 815, 109
- Pontoppidan, K. M., Salyk, C., Bergin, E. A., et al. 2014, *Protostars and Planets VI*, 363
- Pontoppidan, K. M., Fraser, H. J., Dartois, E., et al. 2003, *A&A*, 408, 981
- Qi, C., Öberg, K. I., Andrews, S. M., et al. 2015, *ApJ*, 813, 128
- Qi, C., Öberg, K. I., Wilner, D. J., et al. 2013, *Science*, 341, 630
- Rodgers, S. D., & Charnley, S. B. 2002, *MNRAS*, 330, 660
- Rubin, M., Altwegg, K., Balsiger, H., et al. 2015, *Science*, 348, 232
- Scheegerer, A. A., & Wolf, S. 2010, *A&A*, 517, A87
- Shakura, N. I., & Sunyaev, R. A. 1973, *A&A*, 24, 337
- Stevenson, D. J., & Lunine, J. I. 1988, *Icarus*, 75, 146
- van Dishoeck, E. F. 2006, *Proceedings of the National Academy of Science*, 103, 12249
- van Dishoeck, E. F., Bergin, E. A., Lis, D. C., & Lunine, J. I. 2014, *Protostars and Planets VI*, 835
- Visser, R., van Dishoeck, E. F., Doty, S. D., & Dullemond, C. P. 2009, *A&A*, 495, 881
- Williams, D., & Herbst, E. 2002, *Surface Science*, 500, 823
- Wyckoff, S., Tegler, S. C., & Engel, L. 1991, *ApJ*, 367, 641
- Zhang, K., Pontoppidan, K. M., Salyk, C., & Blake, G. A. 2013, *ApJ*, 766, 82