Instructional workshop on OpenFOAM programming LECTURE # 3

Pavanakumar Mohanamuraly

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Outline

Recap of Week 1

1d Heat Equation

Finite difference

IduMatrix and fvMatrix

FOAM Finite volume - Operators (laplacian)

Recap of Week 1

- Code development compile/running
- Basic FOAM data-structures
- FOAM polyMesh and fvMesh
- ► FOAM fields

1d Heat equation

$$\frac{\partial \phi}{\partial t} - \kappa \frac{\partial^2 \phi}{\partial x^2} = f(x, t) \quad 0 \le x \le L, \quad t \ge 0$$
 (1)

with initial conditions,

$$\phi(x,0)=g_0(x)$$

and boundary conditions,

▶ Dirichlet type

$$\phi(0,t) = \phi_0$$
 and $\phi(L,t) = \phi_L$

Neumann type

$$rac{\partial \phi}{\partial x}(0,t) = \phi_0'$$
 and $rac{\partial \phi}{\partial x}(L,t) = \phi_L'$

Robins BC discussed on Day 2 (user defined BCs)



Finite difference (FD) basics ¹

▶ x discretized as uniformly spaced interval $0 \le x \le L$ such that

$$x_i = (i-1)\Delta x$$
 $i = 1, 2, ...N$ where, $\Delta x = \frac{L}{N-1}$

▶ Similarly, discretize t uniformly in $0 \le t \le T$

$$t_m = (m-1)\Delta t$$
 $i = 1, 2, ...M$ where, $\Delta t = \frac{T}{M-1}$ $\phi(x, t) o \phi(x_i, t_m) o \phi_i^m$

¹Source: http://www.nada.kth.se/~jjalap/numme/FDheat.pdf □ ▶ ∢ ♠ ▶ ∢ ⋛ ▶ ∢ ⋛ ▶ ↓ ⋛ → ○ ○ ○

2nd order central difference

Taylor series expansion

Forward

$$\phi_{i+1} = \phi_i + \Delta x \left. \frac{\partial \phi}{\partial x} \right|_{x_i} + \left. \frac{\Delta x^2}{2} \left. \frac{\partial^2 \phi}{\partial x^2} \right|_{x_i} + \dots$$
 (2)

Backward

$$\phi_{i-1} = \phi_i - \Delta x \left. \frac{\partial \phi}{\partial x} \right|_{x_i} + \left. \frac{\Delta x^2}{2} \left. \frac{\partial^2 \phi}{\partial x^2} \right|_{x_i} - \dots \right.$$
 (3)

Adding (2) and (3) we get

$$\phi_{i+1} + \phi_{i-1} = 2\phi_i + \Delta x^2 \left. \frac{\partial^2 \phi}{\partial x^2} \right|_{x_i} + O(\Delta x)^4 \tag{4}$$

$$\left. \frac{\partial^2 \phi}{\partial x^2} \right|_{x} = \frac{\phi_{i+1} - 2\phi_i + \phi_{i-1}}{\Delta x^2} + O(\Delta x)^2 \tag{5}$$



2nd order central difference

Subtracting (2) and (3) we get

$$\phi_{i+1} - \phi_{i-1} = 2\Delta x \left. \frac{\partial \phi}{\partial x} \right|_{x_i} + (\Delta x)^3 \frac{\partial^3 \phi}{\partial x^3} + \dots$$
 (6)

$$\left. \frac{\partial \phi}{\partial x} \right|_{x_i} = \frac{\phi_{i+1} - \phi_{i-1}}{2\Delta x} + O(\Delta x)^2 \tag{7}$$

Alternatively, using equation (2),

$$\left. \frac{\partial \phi}{\partial x} \right|_{x_i} = \frac{\phi_{i+1} - \phi_i}{\Delta x} + O(\Delta x)$$
 (8)

and equation (3),

$$\left. \frac{\partial \phi}{\partial x} \right|_{x} = \frac{\phi_i - \phi_{i-1}}{\Delta x} + O(\Delta x)$$
 (9)

Boundary condition

Dirichlet type

$$\phi_{i=1} = \phi_0 \quad \phi_{i=N} = \phi_L$$

Neumann type (1st order)

$$\phi_1 = \phi_2 - \phi_0' \Delta x$$
 and $\phi_N = \phi_{N-1} + \phi_L' \Delta x$

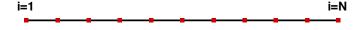
Neumann type (2nd order)

$$\phi_{i=0} = \phi_2 - 2\phi_0' \Delta x$$
 and $\phi_{N+1} = \phi_N - 2\phi_L' \Delta x$

▶ Halo nodes i = 0 and i = N + 1



- ▶ Need *cell-to-cell* connectivity for FOAM FD
- Avoid this by modifying FD implementation
- Original domain



Modified domain (half points - faces)



Modified domain



Split into face contribution (interior faces)

$$\frac{\partial^{2} \phi}{\partial x^{2}}\Big|_{x_{i}} = \frac{\phi_{i+1} - 2\phi_{i} + \phi_{i-1}}{\Delta x^{2}} = \frac{1}{\Delta x} \left(\underbrace{\frac{\phi_{i+1} - \phi_{i}}{\Delta x}}_{\mathbf{L}} - \underbrace{\frac{\phi_{i} - \phi_{i-1}}{\Delta x}}_{\mathbf{R}} \right) \tag{10}$$

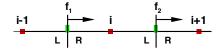
$$\frac{\phi_{i+1} - \phi_{i}}{\Delta x} = \frac{1}{\Delta x} \left(\underbrace{\frac{\phi_{i+1} - \phi_{i}}{\Delta x}}_{\mathbf{R}} - \underbrace{\frac{\phi_{i} - \phi_{i-1}}{\Delta x}}_{\mathbf{R}} \right) \tag{11}$$

$$\frac{\phi_{i} - \phi_{i-1}}{\Delta x} = \frac{1}{\Delta x} \left(\underbrace{\frac{\phi_{i} - \phi_{i-1}}{\Delta x}}_{\mathbf{R}} - \underbrace{\frac{\phi_{i} - \phi_{i-1}}{\Delta x}}_{\mathbf{R}} \right) \tag{12}$$

4□ ► 4□ ► 4 = ► 4 = ► 9 < 0</p>

Sign consistency examples (ignore Δx for clarity)

► Case (i)

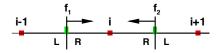


$$\mathbf{f_1} = \phi_{\mathsf{L}} - \phi_{\mathsf{R}} = \phi_{i-1} - \phi_i \qquad (13)$$

$$\mathbf{f_2} = \phi_{\mathsf{L}} - \phi_{\mathsf{R}} = \phi_i - \phi_{i+1}$$
 (14)

$$\underbrace{-\mathbf{f_1}}_{\mathbf{R}} + \underbrace{\mathbf{f_2}}_{\mathbf{I}} = -(\phi_{i+1} - 2\phi_i + \phi_{i-1}) = -\left. \frac{\partial^2 \phi}{\partial x^2} \right|_{x_i}$$
(15)

Case (ii)



$$\mathbf{f_1} = \phi_{\mathbf{L}} - \phi_{\mathbf{R}} = \phi_{i-1} - \phi_i \qquad (16)$$

$$\mathbf{f_2} = \phi_{\mathsf{L}} - \phi_{\mathsf{R}} = \phi_{i+1} - \phi_i \qquad (17)$$

$$\underbrace{-\mathbf{f_1}}_{\mathbf{R}} - \underbrace{\mathbf{f_2}}_{\mathbf{R}} = -(\phi_{i+1} - 2\phi_i + \phi_{i-1}) = -\left. \frac{\partial^2 \phi}{\partial x^2} \right|_{x_i}$$
(18)

As long as L/R is consistently defined we get the correct expression

Neumann boundary conditions

Remember the fact that

$$\phi_0' = \frac{\phi_1 - \phi_0}{\Delta x} = \phi_1^f \text{ and } \phi_L' = \frac{\phi_{N+1} - \phi_N}{\Delta x} = \phi_N^f$$
 (19)

where, f denotes the face value

lacktriangle Simply enforce values ϕ_0' and ϕ_I' for the Neumann patch face

Dirichlet boundary conditions

- Patch face value set to enforces correct values (weak)
- Introduce halo nodes at both ends

At
$$i = N$$

$$\frac{\phi_{N+1} - \phi_N}{\Delta x} = \frac{1}{\Delta x} (\phi_L - \phi_N)$$
(20)

At
$$i = 1$$

$$\frac{\phi_{i=1} - \phi_{i=0}}{\Delta x} = \frac{1}{\Delta x} (\phi_1 - \phi_0)$$
(21)

▶ Remember that if dy and dz is = 1 then $\Delta x = \Delta V$, where ΔV is cell volume

Sign consistency for boundary (ignore Δx for clarity)

▶ Case (i) At x = 0

$$i=0$$
 $T_{x=0}$
 $i=1$
 R
 L

Dirichlet type

$$\mathbf{f_{x=0}} = \frac{\phi_{\mathsf{L}} - \phi_{\mathsf{R}}}{\Delta x} = \frac{\phi_{1} - \phi_{\mathsf{X=0}}}{\Delta x} \tag{22}$$

Neumann type

$$\mathbf{f}_{\mathbf{x}=\mathbf{0}} = \phi_0' \tag{23}$$



Sign consistency for boundary (ignore Δx for clarity)

▶ Case (i) At x = 0

Dirichlet type

$$\mathbf{f_{x=L}} = \frac{\phi_{L} - \phi_{R}}{\Delta x} = \frac{\phi_{N} - \phi_{N+1}}{\Delta x}$$
 (24)

Neumann type

$$\mathbf{f}_{\mathbf{x}=\mathbf{L}} = \phi_L' \tag{25}$$



How to access field boundary values?

fixedValue patch type

```
const tmp<scalarField> sf;
scalarField some( x.boundaryField()[ipatch].
    valueBoundaryCoeffs(sf) );
```

fixedGradient patch type

```
scalarField some( x.boundaryField()[ipatch].
    gradientBoundaryCoeffs() );
```

Hands on - FD solver I

- ▶ For convenience set dx = dy = dz = 1 in the 1d grid
- Ignore time term for now and solve for steady-state
- Create a volScalarField phi and read from input file
- Create a surfaceScalarField named flux
- ▶ Loop over all internal faces of *flux* and set face value

$$\phi_f = \phi_{owner} - \phi_{neighbour} \tag{26}$$

For Neumann patches set face value to the gradient

$$\phi_f = \phi'(x = 0 \text{ or } x = L) \tag{27}$$

For Dirichlet patches set face value using

$$\phi_f = \phi_{owner} - \phi(x = 0 \text{ or } x = L)$$
 (28)



Hands on - FD solver I

- Create a volScalarField named residue and initialize to zero
- Loop over all faces of flux
- Add the face value to the owner cell and subtract from neighbour cell of residue
- ▶ Do the same for boundary faces (no *neighbour* cell)
- Now residue contains the laplacian

Matrix forms

Possible to construct matrix version of the discrete operator

Matrix form of discrete system of Poisson Equations

- ▶ BC enforced separately using boundaryCoeff and internalCoeff
- ► For zeroGradient BC this matrix becomes the full system
- ▶ Most of the entries of A are zeros \implies A is sparse

Sparse Matrix Storage in FOAM

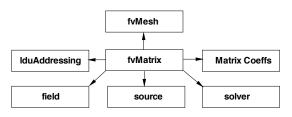
- ► FOAM uses the *LDU* sparse matrix storage
- ▶ Uses the owner/neighbour data for addressing non-zero L/U entries
- Optimized for symmetric matrices also handles asymmetric matrix
- Key limitation "cannot store entries beyond first level of cell neighbours"
 - Big issue for hybrid and highly skewed meshes
 - ► Higher order FVM (> 2nd order)

IduMatrix

- ► Abstract base class implementing the *ludMatix* and solvers
- Uses the *IduMatrix* to construct the *owner/neighbour* addressing scheme

- ▶ *IduMatrix* is not usable, need non-abstract class *fvMatrix*
- ► FOAM operators like grad, div, etc, return fvMatrix

fvMatrix



Solves the system of equations

$$\mathbf{A}\mathbf{x} = \mathbf{b} \tag{29}$$

where, \mathbf{x} is the *field*, \mathbf{b} is the *source* and sparse matrix A entries are *Matrix Coeffs*.

► Subject to BC specified on the *field* using a specified *solver*



fvm and fvc namespace

- Two possible ways to construct discrete operators
 - ▶ Matrix-free scoped in *fvc* namespace
 - Matrix-based scoped in fvm namespace
- Explicit matrix construction impossible for non-linear operations like limiting
- Operators defined in fvc return field types
- ▶ Operators defined in *fvm* return *fvMatrix* types

Krylov solvers ²

- For large sparse systems it only possible to perform matrix-vector products (Ax)
- Even explicit construction of matrix A is impossible
- ▶ But it is possible to construct a *Krylov sequence* $\{x, Ax, A^2x, A^3x, ...\}$
- Krylov sub-space methods build up on the sequence and look for
 - Eigenvectors and
 - Invariant sub-spaces

within the Krylov sub-space

- A large class of LA algorithms like CG, BICG, GMRES are based on Krylov sub-spaces
- Preconditioning is normally performed to speed up convergence

²Source: "The Matrix Eigenvalue Problem" by David S. Watkins

Krylov sub-space solvers in *fvMatrix*

- Krylov sub-space solvers available for fvMatrix
 - Symmetric matrix
 - ► GAMG Geometric/Algebraic Multi-Grid solver
 - ► ICCG Incomplete Cholesky Conjugate Gradient
 - ▶ PCG Preconditioned Conjugate Gradient solver
 - Asymmetric matrix
 - BICCG Bi-Conjugate Gradient
 - GAMG Geometric/Algebraic Multi-Grid solver
 - PBiCG Preconditioned Bi-Conjugate Gradient
 - Matrix preconditioners
 - DIC Diagonal Incomplete Cholesky
 - DILU Diagonal Incomplete LU
 - GAMG Geometric/Algebraic Multi-Grid solver

Multigrid and Krylov solver beyond the scope of the workshop

How to use *fvMatrix*?

Will look at the specific constructor using field variable

```
fvScalarMatrix A( x, x.dimensions() );
```

- Access non-zero upper diagonal entries using A.upper()
- Access diagonal entries using A.diag()
- Access non-zero lower diagonal entries using A.lower()
- First access to A.upper() after construction makes A symmetric
 - ▶ lowerPtr_ = upperPtr_;
- Since fvMatrix entries are based on mesh.owner()/neighbour() following holds true
 - Size of diag().size() = number of cells
 - Size of upper()/lower().size() = number of internal faces

negDiagSum() function

Negated row-wise sum of non-zero off diagonal coeffs

```
void Foam::lduMatrix::negSumDiag()
  const scalarField& Lower
    = const_cast<const lduMatrix &>(*this).lower();
  const scalarField& Upper
    = const_cast<const lduMatrix &>(*this).upper();
  scalarField& Diag = diag();
  const labelUList& l = lduAddr().lowerAddr();
  const labelUList& u = lduAddr().upperAddr();
  for (register label face=0; face<1.size(); face++)</pre>
    Diag[l[face]] -= Lower[face];
    Diag[u[face]] -= Upper[face];
```

Hands on - Create FD matrix A from example

- ▶ Ignore BC patches for now
- Access A.upper() and assign that to 1
- Calculate the diagonal term using the negSumDiag() function
- Now check if the matrix is symmetric using A.symmetric()
- ▶ It should return true or false?

Hints

- ► A.upper()/diag() is a scalarField so forAll works
- ▶ Lazy yet efficient option is to use assignment operator =

Specifying the BC

The BC is specified using two fields

A.boundaryCoeffs()

- Source term that goes to the RHS
- ► For Neumann condition the *-fixedGradient* value is specified
- ▶ For Dirichlet condition the -fixedValue divided by Δx

A.internalCoeffs()

- ► Term that goes to the *LHS* matrix coefficient
- ▶ For Neumann condition it is zero
- ▶ For Dirichlet condition it is -1 divided by Δx

Show on black-board how this calculated

Access boundary values (fixedValue)

Access field connected to A using A.psi()

```
forAll( A.psi().boundaryField(), ipatch ) {
  i f
    A.psi().boundaryField()[ipatch].type()
    == "fixedValue"
  ) {
    const tmp<scalarField> sf;
    scalarField value
      A.psi().boundaryField()[ipatch].
          valueBoundaryCoeffs(sf)
    );
    A.boundaryCoeffs()[ipatch] = -1.0 * value;
    A.internalCoeffs()[ipatch] = -1.0;
```

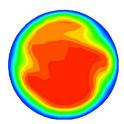
Access boundary values (fixedGradient)

```
if
  A.psi().boundaryField()[ipatch].type()
  == "fixedGradient."
  Info << ipatch << " fixedGradient ";</pre>
  scalarField gradient
    A.psi().boundaryField()[ipatch].
        gradientBoundaryCoeffs()
  );
  A.boundaryCoeffs()[ipatch] = -1.0 * gradient;
  Info << " gradient = " << gradient << "\n";</pre>
```

Access boundary values (zeroGradient)

How to access field boundary values?

- Number of coefficients equal the number of faces in the patch
- Consequence of the fact that it is possible to have different values for each face in the patch
- ▶ Achieved using the *nonuniform* key word in the field dictionary
- Inlet velocity profiles for pipe flows



Hands on: Print boundary field type and value

- Loop over all patches of the field x
- Print the type of boundary and boundary value
- Set correct values for A.boundaryCoeffs()/internalCoeffs()

Solving the LA problem

- ► So far we have only discussed setting up **Ax**
- ▶ What about **b** ?

Solving the LA problem

- So far we have only discussed setting up Ax
- What about b?
- It is accessed using A.source()
- ▶ If everything is set then we are ready to solve the LA problem

$$\mathbf{A}\mathbf{x} = \mathbf{b} \tag{30}$$

The *solution* dictionary

system/fvSolution dictionary

Hands on - Post-process 1d data

- ▶ 1*d* case is actually 3*d* with once cell thickness
- ▶ Print the result to a file *profile.dat*

```
#include <fstream>
....
std::ofstream fout("initial_sol.dat");
forAll( mesh.C(), i )
  fout << mesh.C()[i][0] << " " << x[i] << "\n";</pre>
```

Visualize using gnuplot

```
gnuplot> plot "initial_sol.dat" using 1:2 with lines
```

Hands on - Using fvMatrix for solution

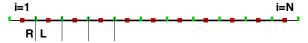
- Solve the 1d FD problem without time derivative using fvMatrix
- Proper boundary condition and value should be set
- Choose appropriate solver (write solver dictionary)
- ► Set the *A.source*() to *zero*
- Essentially solving the equation

$$\frac{\partial^2 \phi}{\partial x^2} = 0 \tag{31}$$

 Should give a straight line profile if Dirichlet BC is specified on both ends (use the previous hand-on to visualize results in gnuplot)

Finite Volume (FV) discretization of 1d heat equation

► FV domain



Split into face contribution (interior faces)

$$\frac{\partial^{2} \phi}{\partial x^{2}}\Big|_{x_{i}} = \frac{\phi_{i+1} - 2\phi_{i} + \phi_{i-1}}{\Delta x^{2}} = \frac{1}{\Delta x} \left(\underbrace{\frac{\phi_{i+1} - \phi_{i}}{\Delta x}}_{\mathbf{L}} - \underbrace{\frac{\phi_{i} - \phi_{i-1}}{\Delta x}}_{\mathbf{R}} \right) \tag{32}$$

$$\frac{\phi_{i+1} - \phi_i}{\Delta x} = \frac{1}{\Delta x} \left(\underbrace{\phi_{i+1}}_{\mathbf{p}} - \underbrace{\phi_i}_{\mathbf{p}} \right) \tag{33}$$

$$\frac{\phi_i - \phi_{i-1}}{\Delta x} = \frac{1}{\Delta x} \left(\underbrace{\phi_i} - \underbrace{\phi_{i-1}}_{2} \right) \tag{34}$$



Finite Volume (FV) discretization of 1d heat equation

- It is not a surprise that we end up with the same set of equations
- ► FV and FD give the same discretization for equally spaced grids in this case
- But this time we make use of FOAM operators to make the job simple

FOAM operators - laplacianScheme

- fvm :: laplacian is an in-built operator which yields correct fvMatrix
- ► Need to define the *laplacianScheme* in the *fvSchemes* dictionary

```
laplacianSchemes
{
    /// Green-gauss type with
    /// non-orthogonality correction
    default Gauss linear corrected;
}
```

Constructor requires only the field as argument

```
fvScalarMatrix A( fvm::laplacian(x) );
A.source() = 0.0;
A.solve();
```

Hands on - FOAM'ish solver

- ► Implement the FV solver with similar inputs and equations as the FD hands on
- ▶ Plot 1*d* solution using *gnuplot*

End of Week 2 Day 1