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Supporting internal notes

Correlations between jets and charged particles in Pb+Pb Collisions at 5.02 TeV:
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Measurement of angular and momentum distributions of charged particles within and around jets in Pb+Pb and pp collisions at $\sqrt{s_{NN}} = 5.02$ TeV with ATLAS at the LHC

Studies of the fragmentation of jets into charged particles in heavy-ion collisions can provide information about the mechanism of jet quenching by the hot and dense QCD matter created in such collisions, the quark-gluon plasma. This paper presents a measurement of the angular distribution of charged particles around the jet axis in $\sqrt{s_{NN}} = 5.02$ TeV Pb+Pb and pp collisions, done using the ATLAS detector at the LHC. The measurement is performed for jets reconstructed with the anti- k_t algorithm with radius parameter $R = 0.4$, and is extended to a distance of $r = 0.8$ from the jet axis. The charged particles in this analysis have transverse momenta in the 1–63 GeV range while jets have a transverse momenta in the 126–316 GeV range and are restricted to have an absolute value of the jet rapidity be less than 1.7. Results are presented as a function of Pb+Pb collision centrality and distance from the jet axis for different jet and charged-particle transverse momenta ranges. The previously observed enhancement of charged particles with transverse momenta below 4 GeV is measured to increase for larger angular distances from the jet axis achieving a maximum at $r = 0.6$. Charged particles with transverse momenta above 4 GeV show an enhancement in Pb+Pb collisions only in the jet core for distances up to $r = 0.05$ and a suppression at all larger distances from the jet axis.

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The ATLAS Collaboration

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1 Introduction

Ultra-relativistic nuclear collisions at the Large Hadron Collider (LHC) produce hot, dense matter called the quark-gluon plasma, QGP (see Refs. [1, 2] for recent reviews). Hard-scattering processes occurring in these collisions produce jets that traverse and interact with the QGP, losing energy via a process called jet-quenching. The rates and characteristics of these jets in heavy-ion collisions can be compared to the same quantities in pp collisions, where we do not expect the production of QGP. This comparison can provide information on the properties of the QGP and how it interacts with partons from the hard scattering.

Jets with large transverse momenta are observed to be produced in central lead-lead (Pb+Pb) collisions at the LHC at a rate that is reduced by a factor of two with respect to the expectation from these cross sections measured in pp interactions scaled by the nuclear overlap function of Pb+Pb collisions [3–7]. Similarly, back-to-back dijet [8–10] and photon-jet pairs [11, 12] are observed to have less balanced transverse momenta in Pb+Pb collisions compared to pp collisions. These observations suggest that some of the energy from the hard-scattered parton may be transferred outside of the jet through its interaction with the QGP.

Complementary measurements look at how the structure of jets is different between Pb+Pb and pp collisions. Jet shape measurements have shown a broadening of the jets from pp collisions to Pb+Pb collisions [13, 14]. Additionally, longitudinal fragmentation functions have been measured in 2.76 TeV and 5.02 TeV Pb+Pb collisions; these indicate that Pb+Pb collisions have an excess of low and high

momentum particles along with a depletion of intermediate momentum particles inside the jet compared to pp collisions [15–18]. Particles carrying a large fraction of the jet momentum are generally closely aligned with the jet axis, whereas low momentum particles are observed to have a much broader angular distribution extending outside the jet [9, 19–21]. These observations suggest that the energy lost via jet-quenching is being transferred to soft particles within and around the jet through emission of soft gluons [22–28]. Measurements of yields of these particles as a function of transverse momentum and angular distance between the particle and the jet axis have a potential to provide more detailed information on the structure of jets in the QGP as well as provide constraints on the medium response to the jet.

This paper presents a measurement of charged particle p_T distributions inside and around jets. The measured yields are defined as:

$$D(p_T, r) = \frac{1}{N_{\text{jet}}} \frac{1}{A} \frac{dn_{\text{ch}}(p_T, r)}{dp_T},$$

where N_{jet} is the number of jets in consideration, $A = \pi(r_{\text{max}}^2 - r_{\text{min}}^2)$ is the area of an annulus around the jet with its inner and outer radii r_{min} and r_{max} respectively. The angular distance from the jet axis is given by $r = \sqrt{\Delta\eta^2 + \Delta\phi^2}$ ¹, and $n_{\text{ch}}(p_T, r)$ is the number of charged particles with a given p_T within the annulus. The ratios of the charged-particle yields measured in Pb+Pb and pp collisions,

$$R_{D(p_T, r)} = \frac{D(p_T, r)_{\text{Pb+Pb}}}{D(p_T, r)_{pp}},$$

allow evaluation of modifications of the yields from the QGP medium. The differences between the $D(p_T, r)$ distributions in Pb+Pb and pp collisions,

$$\Delta D(p_T, r) = D(p_T, r)_{\text{Pb+Pb}} - D(p_T, r)_{pp},$$

are also presented and allow for quantifying the absolute differences between the two collision systems.

The analysis is done using 0.49 nb⁻¹ of Pb+Pb collisions and 25 pb⁻¹ of pp collisions at center-of-mass energy of 5.02 TeV collected in 2015 by ATLAS. Jets are reconstructed with the anti- k_t algorithm [29] using a radius parameter $R = 0.4$ over a rapidity range of $|y^{\text{jet}}| < 1.7$. The measurement is presented for jets with transverse momenta p_T^{jet} , in the 126–316 GeV range, for charged particles with $p_T > 1.0$ GeV. It is done in successive intervals of r around the jet in annuli with inner and outer radii r_{min} and r_{max} : 0.0, 0.05, 0.1, 0.15, 0.2, 0.25, 0.3, 0.4, 0.5, 0.6, 0.7, 0.8.

2 Experimental setup

The measurements presented here are performed using the ATLAS calorimeter, inner detector, trigger, and data acquisition systems [30].

¹ ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector, and the z -axis along the beam pipe. The x -axis points from the IP to the center of the LHC ring, and the y -axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. The rapidity is defined as $y = 0.5 \ln[(E + p_z)/(E - p_z)]$ where E and p_z are the energy and the component of the momentum along the beam direction. Transverse momentum and transverse energy are defined as $p_T = p \sin \theta$ and $E_T = E \sin \theta$, respectively. $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$ gives the angular distance between two objects with relative differences $\Delta\eta$ and $\Delta\phi$ in pseudorapidity and azimuth respectively.

The calorimeter system consists of a sampling liquid-argon (LAr) electromagnetic (EM) calorimeter covering $|\eta| < 3.2$, a steel–scintillator sampling hadronic calorimeter covering $|\eta| < 1.7$, LAr hadronic calorimeters covering $1.5 < |\eta| < 3.2$, and two LAr forward calorimeters (FCal) covering $3.1 < |\eta| < 4.9$. The EM calorimeters are segmented longitudinally in shower depth into three layers with an additional pre-sampler layer. They have segmentation that varies with layer and pseudorapidity. The hadronic calorimeters have three sampling layers longitudinal in shower depth.

The inner detector measures charged particles within the pseudorapidity interval $|\eta| < 2.5$ using a combination of silicon pixel detectors, silicon microstrip detectors (SCT), and a straw-tube transition radiation tracker (TRT), all immersed in a 2 T axial magnetic field [30]. Each of the three detectors is composed of a barrel and two symmetric end-cap sections. The pixel detector is composed of four layers including the Insertable B-Layer" [31, 32]. The SCT barrel section contains four layers of modules with sensors on both sides, and each end-cap consists of nine layers of double-sided modules with radial strips. The TRT contains layers of staggered straws interleaved with fibers in the barrel and end-cap.

The zero-degree calorimeters (ZDCs) are located symmetrically at $z = \pm 140$ m and cover $|\eta| > 8.3$. The ZDCs use tungsten plates as absorbers, and quartz rods sandwiched between the tungsten plates as the active medium. In Pb+Pb collisions the ZDCs primarily measure “spectator” neutrons, that is neutrons that do not interact hadronically when the incident nuclei collide. A ZDC coincidence trigger is implemented by requiring the pulse height from both ZDCs to be above a threshold to accept the single-neutron peak.

A two-level trigger system is used to select the Pb+Pb and pp collisions. The first level is based on custom electronics while the second level, the High Level Trigger (HLT), is based on software [33]. Minimum-bias (MB) events are recorded using a logical OR of two triggers: 1) total energy Level-1 trigger; 2) veto on the total energy trigger and ZDC coincidence trigger at Level-1 with the additional requirement of least one track in the HLT. The total-energy trigger required a total transverse energy measured in the calorimeter system to be greater than 50 GeV. Jet events are selected by the HLT, seeded by a jet identified by the Level-1 jet trigger in pp collisions or by the total-energy trigger with a threshold of 50 GeV in Pb+Pb collisions. The Level-1 jet trigger utilized in pp collisions required a jet with transverse momentum greater than 20 GeV. The HLT jet trigger used a jet reconstruction procedure similar to that in the offline analysis as discussed in Section 4. It selected events containing jets with a transverse energy of at least 75 GeV in Pb+Pb collisions and at least 85 GeV in pp collisions. The measurement is performed in the jet transverse momentum range where the trigger is fully efficient.

3 Data sets and event selection

The Pb+Pb and pp data used in this analysis were recorded in 2015. The data samples consist of 25 pb^{-1} of $\sqrt{s} = 5.02 \text{ TeV } pp$ and 0.49 nb^{-1} of $\sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$ Pb+Pb data. In both samples, events are required to have a reconstructed vertex within 150 mm of the nominal IP along the beam axis. Only events taken during stable beam conditions and satisfying detector and data-quality requirements that include the calorimeters and inner tracking detectors being in nominal operating conditions are considered.

In Pb+Pb collisions, the event centrality reflects the overlap area of the two colliding nuclei and is characterized by $\Sigma E_{\text{T}}^{\text{FCal}}$, the total transverse energy deposited in the FCal [34]. The six centrality intervals used in this analysis are defined according to successive percentiles of the $\Sigma E_{\text{T}}^{\text{FCal}}$ distribution obtained in MB collisions, ordered from the most central (highest $\Sigma E_{\text{T}}^{\text{FCal}}$) to the most peripheral (lowest $\Sigma E_{\text{T}}^{\text{FCal}}$) collisions: 0–10%, 10–20%, 20–30%, 30–40%, 40–60%, 60–80%.

In addition to the jet-triggered sample, a separate MB Pb+Pb data sample was recorded with three trigger selections: the minimum-bias trigger and two total transverse-energy triggers with thresholds of 1.5 TeV and 6.5 TeV used to enhance the rate of central Pb+Pb events. This sample was combined with a set of 1.8×10^7 5.02 TeV hard-scattering dijet pp events generated with POWHEG+PYTHIA8 [35, 36] using the A14 tune of parameters [37] and the NNPDF23LO PDF set [38] to produce the Monte Carlo Overlay sample. This sample is reweighed on an event-by-event basis such that it has the same centrality distribution as the jet triggered sample. A separate set of 1.8×10^7 5.02 TeV hard-scattering dijet pp events generated with the same tune and PDFs was used as the pp MC. The detector response was simulated in both MC samples using GEANT4 [39, 40] and was used to evaluate the performance of the detector and analysis procedure. Another sample of MB Pb+Pb events was generated using HIJING (version 1.38b) [40] and was only used to evaluate the track reconstruction performance.

4 Jet and track selection

The jet reconstruction procedure is identical to that used in Ref. [7]. The anti- k_t algorithm is first run in four-momentum recombination mode, on $\Delta\eta \times \Delta\phi = 0.1 \times 0.1$ calorimeter towers with two anti- k_t distance parameter values ($R = 0.2$ and $R = 0.4$). The energies in the towers are obtained by summing the energies of calorimeter cells at the electromagnetic energy scale within the tower boundaries. Then, an iterative procedure is used to estimate the η -dependent underlying event (UE) transverse energy density, while excluding the regions populated by jets. The estimate of the UE contribution is performed on an event-by-event basis. Furthermore, the background is modulated to account for the azimuthal anisotropy in particle production [41]. The modulation accounts for the contribution of the second, third, and fourth order azimuthal anisotropy harmonics. The UE transverse energy is subtracted from calorimeter towers included in the jet and the four-momentum of the jet is updated accordingly. Then, a jet η - and p_T -dependent correction factor to the p_T^{jet} derived from the simulation samples is applied to correct for the calorimeter energy response [42]. The same calibration factors are applied both in pp and Pb+Pb collisions. An additional correction based on *in-situ* studies of jets recoiling against photons, Z bosons, and jets in other regions of the calorimeter is applied [43, 44]. The same jet reconstruction procedure without the azimuthal modulation of the UE is also applied to pp collisions. In this analysis, jets are required to have p_T^{jet} in the 126–316 GeV range, with rapidity $|y^{\text{jet}}| < 1.7$. To prevent nearby jets from distorting the measurement of $D(p_T, r)$ distributions, jets are rejected if there is another jet with a higher p_T^{jet} than the considered jet anywhere within an angular distance of $\Delta R < 1.0$. The isolation requirement removes approximately 0.01% of jets, and has a negligible impact on the final measurement.

Charged-particle tracks in Pb+Pb collisions are reconstructed from hits in the inner detector using the track reconstruction algorithm that has been optimized for the high hit density in heavy-ion collisions [45]. Tracks used in this analysis are required to have at least 9 (11) total silicon hits for charged particles with pseudorapidity, $|\eta^{\text{ch}}| \leq 1.65$ ($|\eta^{\text{ch}}| > 1.65$). At least one hit is required in one of the two innermost pixel layers. If the track trajectory passes through an active module in the innermost layer, then a hit in this layer is required. Additionally, a track must have no more than two holes in the pixel and SCT detectors together, where a hole is defined by the absence of a hit predicted by the track trajectory. All charged-particle tracks used in this analysis are required to have reconstructed transverse momentum $p_T^{\text{ch}} > 1.0$ GeV. In order to suppress a contribution from secondary particles², the distance of closest approach of the track to the

² Primary particles are defined as particles with a mean lifetime $\tau > 0.3 \times 10^{-10}$ s either directly produced in pp interactions or from subsequent decays of particles with a shorter lifetime. All other particles are considered to be secondary.

primary vertex is required to be less than a value that varies from 0.45 mm at $p_T^{\text{ch}} = 4$ GeV to 0.2 mm at $p_T^{\text{ch}} = 20$ GeV in the transverse plane and is less than 1.0 mm in the longitudinal direction.

The efficiency, ε , for reconstructing charged particles in Pb+Pb and pp collisions is determined using the MC samples described above. It is evaluated as a function of the generator-level primary particle transverse momentum, p_T^{truth} , and pseudorapidity, η^{truth} , by associating tracks to generator-level primary particles [40]. For Pb+Pb collisions the efficiency is also evaluated separately in each centrality interval used in the measurement. In both collision systems the efficiency increases slowly with p_T^{truth} and is seen to be independent of p_T^{jet} in the measurement phase space. It is approximately 80% at $p_T = 1$ GeV and rises to approximately 90% at $p_T = 63$ GeV. The variation in efficiency between the most central and peripheral Pb+Pb collisions is approximately 3%.

The contribution of reconstructed tracks that cannot be matched to a generated primary particle in the pp MC samples (“fake” tracks), along with the residual contribution of tracks matched to secondary particles, are together less than 2% in the entire p_T^{ch} range under study in both pp and Pb+Pb collisions.

5 Analysis procedure

The analysis procedure is similar to the procedure in Ref. [18] with the additional requirement of being done differentially in r . Measured tracks are associated with a reconstructed jet if they fall within $\Delta R < 0.8$ of the jet axis and are constructed as:

$$\frac{d^2 n_{\text{ch}}^{\text{meas}}}{dp_T^{\text{ch}} dr} = \frac{1}{\varepsilon(p_T^{\text{ch}}, \eta^{\text{ch}})} \frac{\Delta n_{\text{ch}}(p_T^{\text{ch}}, r)}{\Delta p_T^{\text{ch}} \Delta r}$$

where $\Delta n_{\text{ch}}(p_T^{\text{ch}}, r)$ represents the number of tracks within a given p_T^{ch} and r range. The efficiency correction is applied as a $1/\varepsilon(p_T^{\text{ch}}, \eta^{\text{ch}})$ weight on a track-by-track basis, assuming $p_T^{\text{ch}} = p_T^{\text{truth}}$. While that assumption is not strictly valid, the efficiency varies sufficiently slowly with p_T^{truth} that the error introduced by this assumption is less than 1%.

The measured track yields need to be corrected for the UE, fake tracks and secondaries. In pp collisions, the UE contribution from hard scatterings not associated with jet production is negligible. The contributions from fake tracks and secondary charged particles are estimated from MC samples and subtracted. This procedure is similar to previous measurements [18, 46].

For Pb+Pb collisions, the UE, fake track, and secondary contributions are estimated together in a two step process: first, MC overlay is used to generate $\eta^{\text{jet}}-\phi^{\text{jet}}$ maps of the average number of charged particles in a given annulus around a reconstructed jet. This is done for charged particles without a truth match and as a function of p_T^{jet} , η^{jet} , ϕ^{jet} , angle of the jet to the second order event plane³ $d\Psi_{\text{jet}}$, r , p_T^{ch} , and centrality. In the second step, the $\eta^{\text{jet}}-\phi^{\text{jet}}$ maps are used to generate the UE distribution for jets with a given η^{jet} , ϕ^{jet} , and $d\Psi_{\text{jet}}$. This distribution includes fakes, and is given by $d^2 n_{\text{ch}}^{\text{UE+Fake}}(p_T^{\text{ch}}, r)/dp_T^{\text{ch}} dr$. The yields decrease with decreasing collision centrality, increasing p_T^{ch} , and increasing azimuthal distance from Ψ_2 . The subtracted distributions are then given by

³ The second order event plane angle Ψ_2 is determined on an event-by-event basis by a standard method using the ϕ variation of transverse energy in the FCal [41]

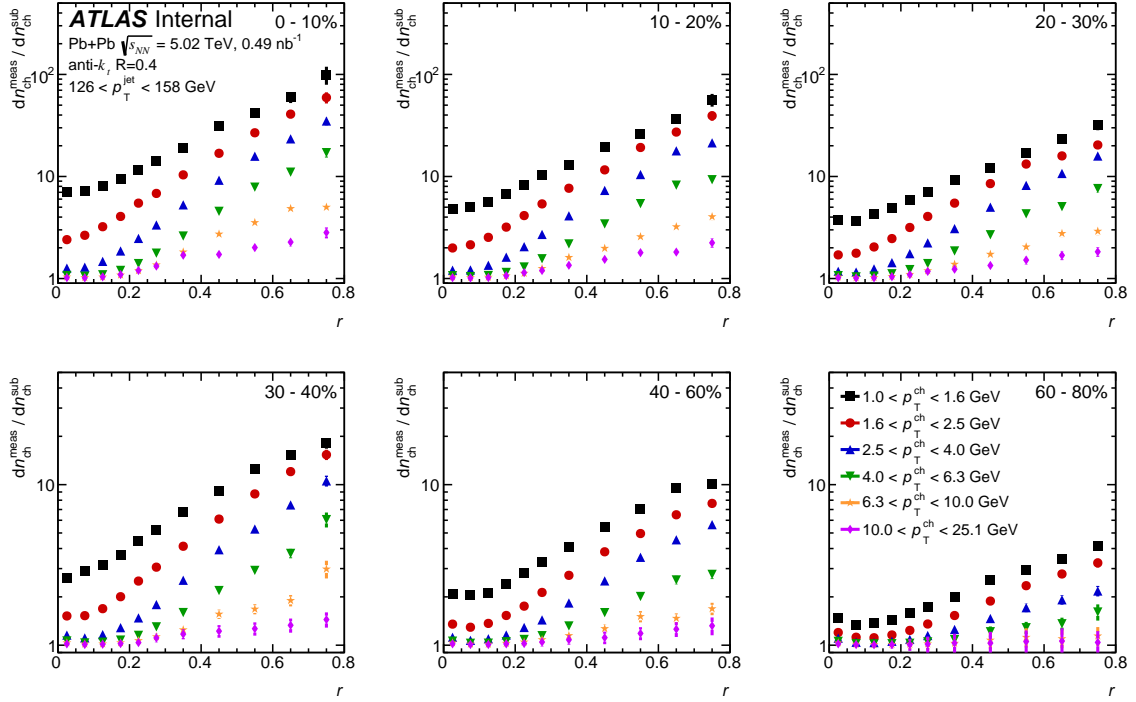


Figure 1: Ratio of the raw charged particle distributions to those after the subtraction of the UE and fake tracks as a function of r for different p_T^{ch} intervals, six centrality selections and for p_T^{jet} between 126–158 GeV.

$$\frac{d^2 n_{\text{ch}}^{\text{sub}}}{dp_T^{\text{ch}} dr} = \frac{d^2 n_{\text{ch}}^{\text{meas}}}{dp_T^{\text{ch}} dr} - \frac{d^2 n_{\text{ch}}^{\text{UE+Fake}}(r)}{dp_T^{\text{ch}} dr}$$

Figure 1 shows the ratio of the charged-particle distributions before and after the UE and fake track subtraction, as a function of r for different p_T^{ch} intervals and $126 < p_T^{\text{jet}} < 158$ GeV for six centrality selections. The UE contribution is the highest for 1.0 GeV charged particles at large values of r and in central collisions, with the background being approximately 100 times the signal, and slowly decreasing with increasing p_T^{jet} . It rapidly decreases for more peripheral collisions, larger p_T^{ch} and smaller r . In addition, due to the steeply falling nature of the jet p_T spectra, the smearing due to jet energy resolution leads to a net migration of jets from lower p_T^{jet} to higher p_T^{jet} values, such that a jet reconstructed with a given transverse momentum will correspond, on average, to a lower truth jet p_T . This “up-feeding” induces a difference between the UE yields determined using the MC overlay events and the actual UE contribution to reconstructed jets. The magnitude of this difference is centrality dependent, and can be seen clearly for charged particles with higher p_T .

To remove the effects of the bin migration due to the jet energy and track momentum resolution, the subtracted $d^2 n_{\text{ch}}^{\text{sub}}/dp_T^{\text{ch}} dr$ distributions are corrected by a two-dimensional Bayesian unfolding [47] in p_T^{ch} and p_T^{jet} as implemented in the RooUnfold package [48]. Two-dimensional unfolding is used because the calorimetric jet energy response depends on the fragmentation pattern of the jet [49]. Four-dimensional response matrices are created from the pp and Pb+Pb MC samples using the generator-level and reconstructed p_T^{jet} , and the generator-level and reconstructed charged-particle p_T^{ch} . They are corrected for tracking efficiencies

and are evaluated in bins of r and centrality. The Bayesian procedure requires a choice in the number of iterations. Additional iterations reduce the sensitivity to the choice of prior, but may amplify statistical fluctuations in the distributions. After four iterations the charged particle distributions are found to be stable for both the Pb+Pb and pp data. A separate one-dimensional Bayesian unfolding is used to correct the measured p_T^{jet} spectra that are used to normalize the unfolded charged particle distributions. To achieve better correspondence with the data, the response matrices for both the one and two dimensional unfolding are reweighted so that the distributions match the shapes in the reconstructed data.

An independent bin-by-bin unfolding procedure is also used to correct for migrations originating from the finite jet and track angular resolutions. Two corresponding $D(p_T, r)$ distributions are evaluated in MC samples, one using generator-level jets and primary particles and the other using reconstructed jets and charged particles with their reconstructed p_T replaced by generator-level transverse momentum, p_T^{truth} . The ratio of these two MC distributions provides a correction factor which is then applied to the data.

The final particle-level corrected distributions, normalized by the area of the annulus under question are defined as:

$$D(p_T, r) = \frac{1}{N_{\text{jet}}^{\text{unfolded}}} \frac{1}{A(r)} \frac{d^2 n_{\text{ch}}^{\text{unfolded}}(r)}{dp_T dr},$$

where $N_{\text{jet}}^{\text{unfolded}}$ is the unfolded number of jets in a given p_T^{jet} interval.

The performance of the full analysis procedure is validated in the MC samples by comparing the fully corrected charged particle distributions to the generator-level distributions. Good closure ($< 3\%$) is seen for charged particles with $p_T < 10$ GeV. It is to be noted that adding or removing particles carrying a large fraction of the jet momentum near the edge of the jet can significantly alter its reconstructed momentum and position; this instability leads to some non-closure in the analysis procedure. Results are presented where the non-closure in the pp MC sample is less than 5%.

6 Systematic uncertainties

The following sources of systematic uncertainty are considered: the jet energy scale (JES), the jet energy resolution (JER), the sensitivity of the unfolding to the prior, the UE contribution, the residual non-closure of the analysis procedure, and tracking-related uncertainties. For each systematic variation, the $D(p_T, r)$ distributions along with their ratios and differences are re-evaluated. The difference between the varied and nominal distributions is used as an estimate of the uncertainty.

The systematic uncertainty due to the JES in Pb+Pb collisions is composed of two parts: a centrality-independent baseline component and a centrality-dependent component. Only the centrality-independent baseline component is used in pp collisions; it is determined from *in-situ* studies of the calorimeter response [42, 49, 50] and the relative energy scale difference between the jet reconstruction procedures in heavy-ion [50] and pp collisions [51]. The centrality-dependent uncertainty reflects a modification of parton showers by the Pb+Pb environment. It is evaluated by comparing calorimeter p_T^{jet} and the sum of the transverse momentum of charged particles within the jet in data and MC. The size of the centrality-dependent uncertainty on the JES reaches 0.5% in the most central collisions. Each component that contributes to the JES uncertainty is varied separately by ± 1 standard deviation for each interval in p_T^{jet} and the response matrix is recomputed accordingly. The data are then unfolded with the modified

matrices. The resulting uncertainty from the JES increases with increasing charged-particle p_T at fixed p_T^{jet} and decreases with increasing p_T^{jet} , and is at the level of 2–4%.

The uncertainty on the $D(p_T, r)$ distributions due to the JER is evaluated by repeating the unfolding procedure with modified response matrices, where an additional contribution is added to the resolution of the reconstructed p_T^{jet} using a Gaussian smearing procedure. The smearing factor is evaluated using an *in-situ* technique in 13 TeV pp data that involves studies of dijet energy balance [52, 53]. An additional uncertainty is included to account for differences between the tower-based jet reconstruction and that used in analyses of 13 TeV pp data. The resulting uncertainty from the JER is symmetrized to account for negative variations of the JER. The size of the resulting uncertainty on the $D(p_T, r)$ distributions due to the JER typically reaches 4–5% for the highest charged-particle p_T intervals and decreases to 2–3% with decreasing charged-particle p_T at fixed p_T^{jet} .

The uncertainties related to track reconstruction and selection originate from several sources. Uncertainties related to the material description in simulation and the track transverse momentum resolution are obtained from studies in data and simulation described in Ref. [54]. The sensitivity of the tracking efficiency to the description of the inactive material in the MC samples is evaluated by varying the material description. This resulting uncertainty in the track reconstruction efficiency is between 0.5% and 2% in the track p_T range used in the analysis. The systematic uncertainty on the fakes and secondaries is 30% in both collision systems [54]. The contamination of fake tracks is less than 2% and the resulting uncertainty in the $D(p_T, r)$ distributions is at most 5%. An additional uncertainty takes into account a possible residual misalignment of the tracking detectors in pp and Pb+Pb data-taking. The alignment in these datasets is checked *in-situ* with $Z \rightarrow \mu^+ \mu^-$ events, and the track- p_T dependent uncertainty arises from the finite size of this sample. The resulting uncertainties in the $D(p_T, r)$ distributions are typically less than 0.1%. An additional uncertainty in the tracking efficiency due to the high local track density in the core of jets is 0.4% [55] for all p_T^{jet} ranges in this analysis. The uncertainty due to the track selection is evaluated by repeating the analysis with an additional requirement on the significance of the distance of closest approach of the track to the primary vertex. This uncertainty affects the track reconstruction efficiencies, track momentum resolution, and rate of fake tracks. The resulting uncertainty typically varies between 1–2%. Finally, the track-to-particle association requirements are varied. This variation affects the track reconstruction efficiency, track momentum resolution, and rate of fake tracks. The resulting systematic uncertainty is $\leq 0.1\%$ on the $D(p_T, r)$ distributions. All track-related systematic uncertainties are added in quadrature and presented as the total tracking uncertainty.

The systematic uncertainty associated with the UE subtraction has two components: limited statistics of charged particles associated with a jet without a corresponding generator particle in the Pb+Pb MC, and a comparison to an alternative UE estimation done using the cone method. The cone method uses jet triggered events to estimate the background and is adapted from [17, 18]. A regular grid of 9 cones of size $R = 0.8$ is used to cover the inner detector region. Cones are excluded if they are within an angular distance of 1.6 to a reconstructed jet with $p_T^{\text{jet}} > 90$ GeV or if they contain a charged particle with $p_T > 10$ GeV. This exclusion reduces biases from any hard processes. The resulting UE charged particle yields $dn_{\text{ch}}^{\text{UECone}}/dp_T^{\text{ch}}$ are evaluated over the 1–10 GeV range as a function of p_T^{ch} , p_T^{jet} , centrality, and r , and are subsequently averaged over all cones. The UE uncertainty on the $D(p_T, r)$ distributions is less than 10% for $r < 0.4$ and sharply decreases with increasing charged-particle p_T . It reaches a maximum of 40% at the largest angular distances from the jet axis and is the dominant source of the systematic uncertainty for low charged-particle p_T at large r . In particular, the component from the limited statistics dominates in the most central collisions, while the component from the alternative estimation method dominates elsewhere.

The systematic uncertainty on the unfolding procedure is estimated by generating the response matrices from the MC distributions without any re-weighting to match shapes in data. Conservatively, the difference between the nominal $D(p_T, r)$ distribution and $D(p_T, r)$ unfolded with the un-reweighted response matrices is taken as the systematic uncertainty, and is at the level of 5–7%.

Conservatively, an additional uncertainty to account for possible residual limitations in the analysis procedure is assigned by evaluating the non-closure of the unfolded distributions in simulations. This is typically at the level of 3–4% and is described in Section 5.

The correlations between the various systematic components are considered in evaluating the $R_{D(p_T, r)}$ and $\Delta D(p_T, r)$ distributions. The unfolding and non-closure uncertainties are taken to be uncorrelated between pp and Pb+Pb collisions, while all others are taken to be correlated. For these, the $R_{D(p_T, r)}$ and $\Delta D(p_T, r)$ distributions are re-evaluated by applying the variation to both collision systems; the resulting variations of the ratios from their central values are used as the correlated systematic uncertainty.

Examples of systematic uncertainties in the $D(p_T, r)$ distributions for jets in the 126–158 GeV p_T^{jet} range measured in pp and Pb+Pb collision systems are shown in Figure 2. The uncertainties on the $R_{D(p_T, r)}$ distributions are shown in Figure 3. It can be seen that the dominant systematic uncertainty on the Pb+Pb and the $R_{D(p_T, r)}$ distributions is from the UE estimation. While it is less than 5% for $r < 0.3$, it is approximately 40% for charged particles with $p_T = 1$ GeV at $r = 0.8$. The uncertainties in the pp system are smaller, with the dominant systematic uncertainty coming from the tracking. This uncertainty is approximately 10% for $r < 0.1$ and decreases to less than 5% at larger distances.

7 Results

The $D(p_T, r)$ distributions are studied as a function of p_T^{jet} for pp data and Pb+Pb collisions with different centralities. Ratios and differences between $D(p_T, r)$ distributions in Pb+Pb and pp collisions are evaluated to explore the interplay between the hot and dense matter and the parton shower.

The $D(p_T, r)$ distributions evaluated in pp and Pb+Pb collisions for $126 < p_T^{\text{jet}} < 158$ GeV are shown in Figure 4. The distributions exhibit a difference in shape between Pb+Pb and pp collisions, with the Pb+Pb distributions being broader at low p_T ($p_T < 4$ GeV) and narrower at high p_T ($p_T > 4$ GeV) in 0–10% central collisions. This modification is centrality dependent and is smaller for peripheral Pb+Pb collisions.

In order to quantify the differences seen in Figure 4, ratios of the $D(p_T, r)$ distributions in Pb+Pb collisions to those measured in pp collisions for $126 < p_T^{\text{jet}} < 158$ GeV and $200 < p_T^{\text{jet}} < 251$ GeV jets are presented in Figure 5. They are shown as a function of r for different p_T and centrality selections. In 0–10% central collisions, $R_{D(p_T, r)}$ is greater than unity for $r < 0.8$ for charged particles with p_T less than 4.0 GeV in both jet selections. For these particles, the enhancement of yields in Pb+Pb collisions compared to those in pp collisions grows with increasing r up to approximately $r = 0.3$, with $R_{D(p_T, r)}$ reaching up to two for $1.0 < p_T < 2.5$ GeV. The value of $R_{D(p_T, r)}$ is approximately constant for r in the interval 0.3–0.6 and decreases for $r > 0.6$. For charged particles with $p_T > 4.0$ GeV, $R_{D(p_T, r)}$ shows a depletion outside the jet core for $r > 0.05$. The magnitude of this depletion increases with increasing r up to $r = 0.3$ and is approximately constant thereafter. The observed behavior inside the jet cone, $r < 0.4$, agrees with the measurement of the inclusive jet fragmentation functions [10, 17, 18], where yields of fragments with $p_T < 4$ GeV are observed to be enhanced and yields of charged particles with intermediate p_T are suppressed in Pb+Pb collisions compared to those in pp collisions. For 30–40% mid-central collisions, the enhancement of particles with $p_T < 4.0$ GeV is similar to that in the most central collisions, however the

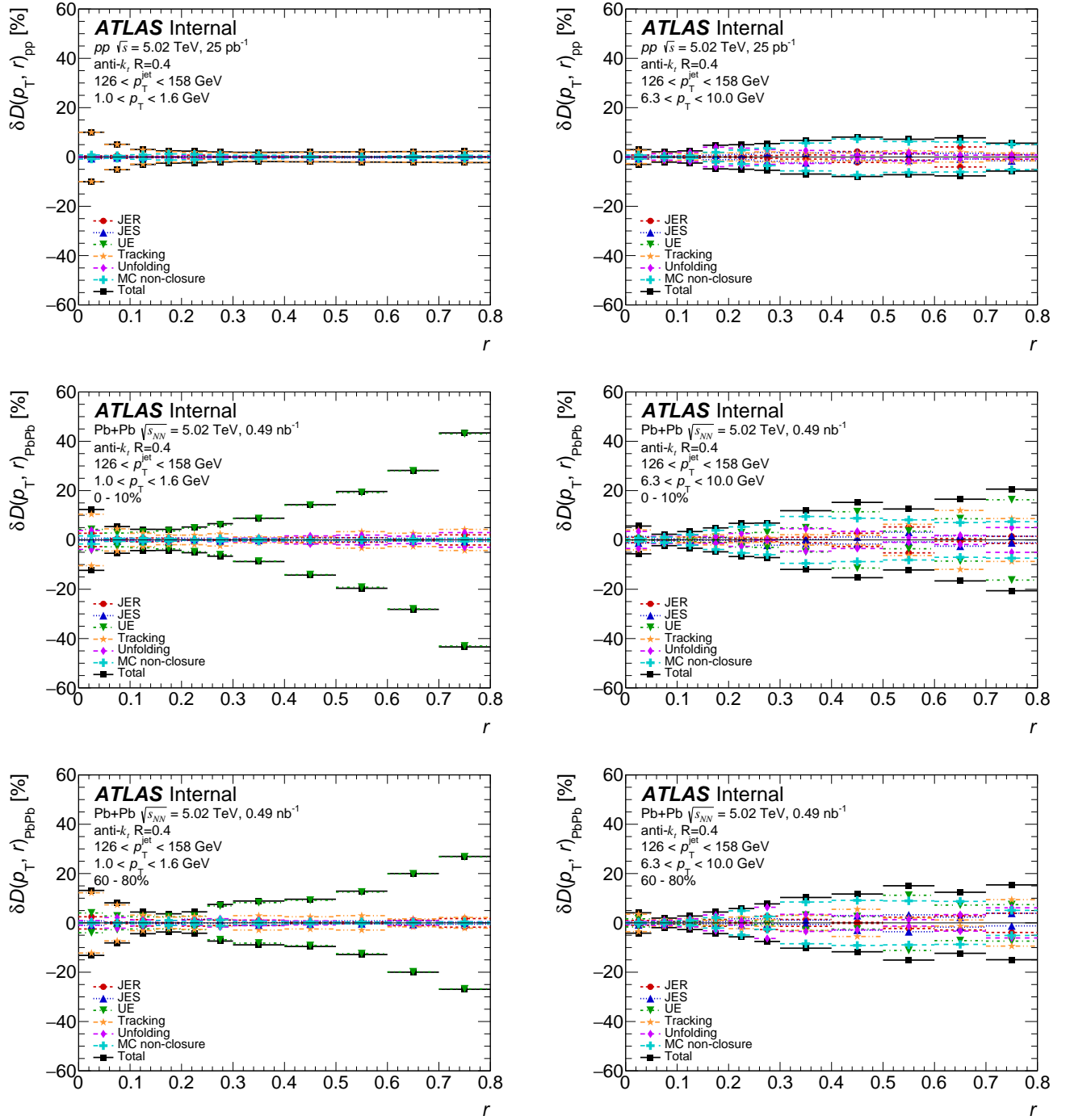


Figure 2: Relative size of the systematic uncertainties for $D(p_T, r)$ distributions in pp (top), central 0–10% Pb+Pb (middle), and peripheral 60–80% Pb+Pb (bottom) collisions for tracks with $1.0 < p_T < 1.6$ GeV (left) and $6.3 < p_T < 10$ GeV (right) in jets with $126 < p_T^{\text{jet}} < 158$ GeV. The systematic uncertainties due to JES, JER, unfolding, UE contribution, MC non-closure, and tracking are shown along with the total systematic uncertainty from all sources.

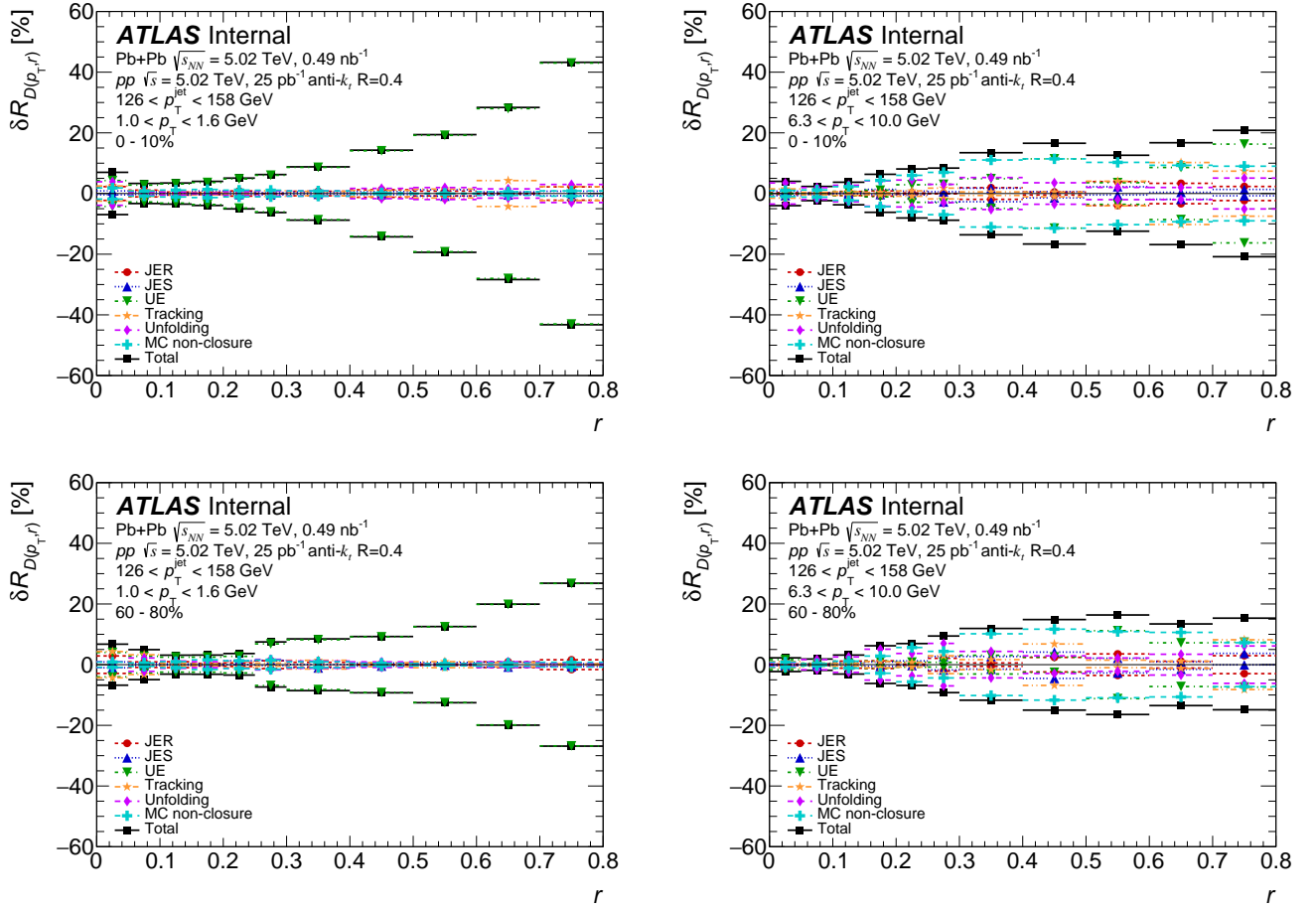


Figure 3: Relative size of the systematic uncertainties for $R_{D(p_T, r)}$ distributions for 0–10% (top) and 60–80% (bottom) Pb+Pb collisions, for tracks with $1.0 < p_T < 1.6$ GeV (left) and $6.3 < p_T < 10.0$ GeV (right), in jets with $126 < p_T^{\text{jet}} < 158$ GeV. The systematic uncertainties due to JES, JER, unfolding, UE contribution, MC non-closure, and tracking are shown along with the total systematic uncertainty from all sources.

326 depletion of particles with $p_T > 4.0$ GeV is not as strong. For 60–80% peripheral collisions, $R_{D(p_T, r)}$ has
 327 no significant r dependence and the values of $R_{D(p_T, r)}$ are within approximately 50% of unity.

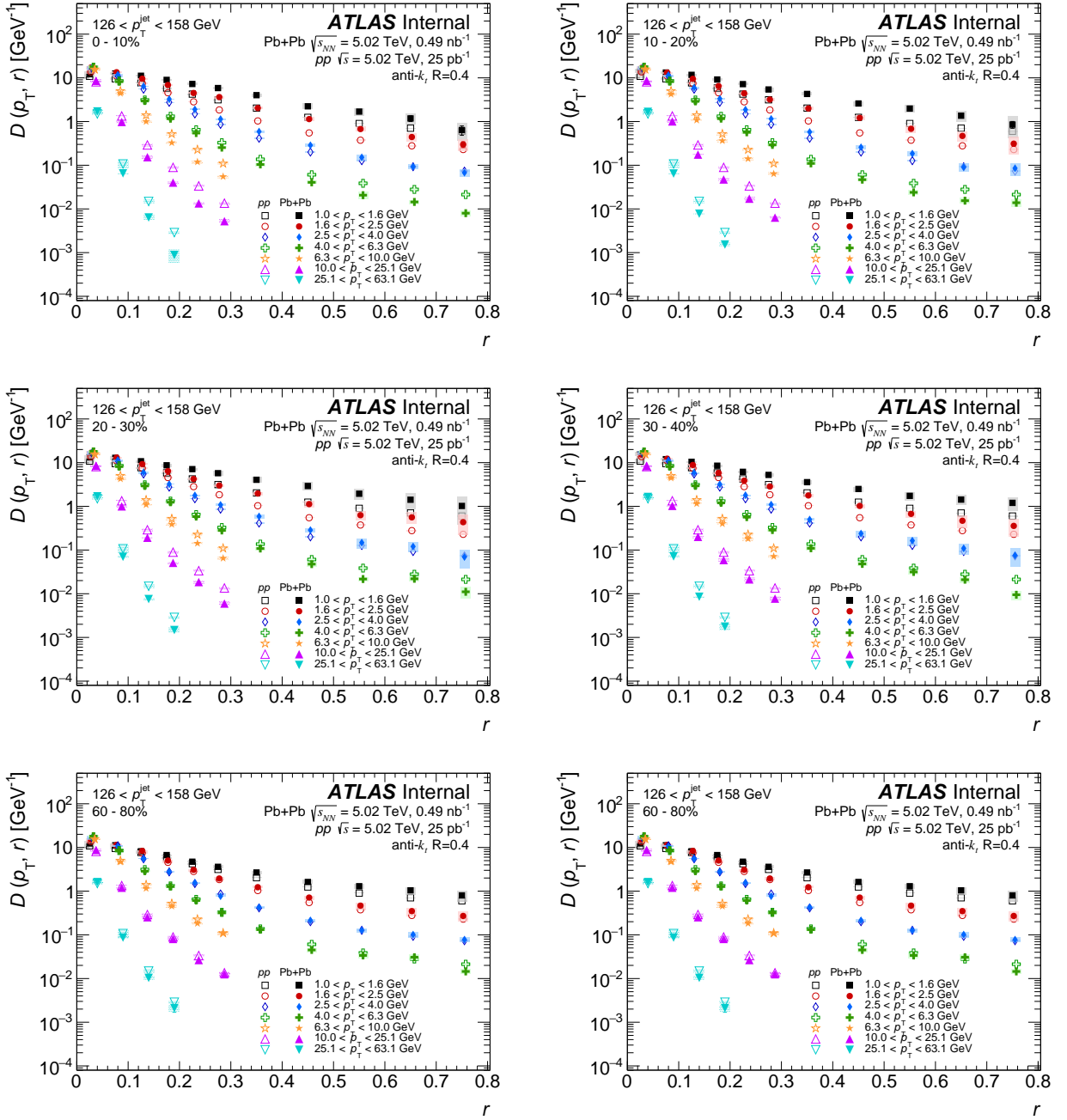


Figure 4: The $D(p_T, r)$ distributions in pp (open symbols) and $Pb+Pb$ (closed symbols) as a function of angular distance r for p_T^{jet} of 126 to 158 GeV. The colors represent different track p_T ranges, and each panel is a different centrality selection. The vertical bars on the data points indicate statistical uncertainties while the shaded boxes indicate systematic uncertainties. The widths of the boxes are not indicative of the bin size and the points are shifted horizontally for better visibility.

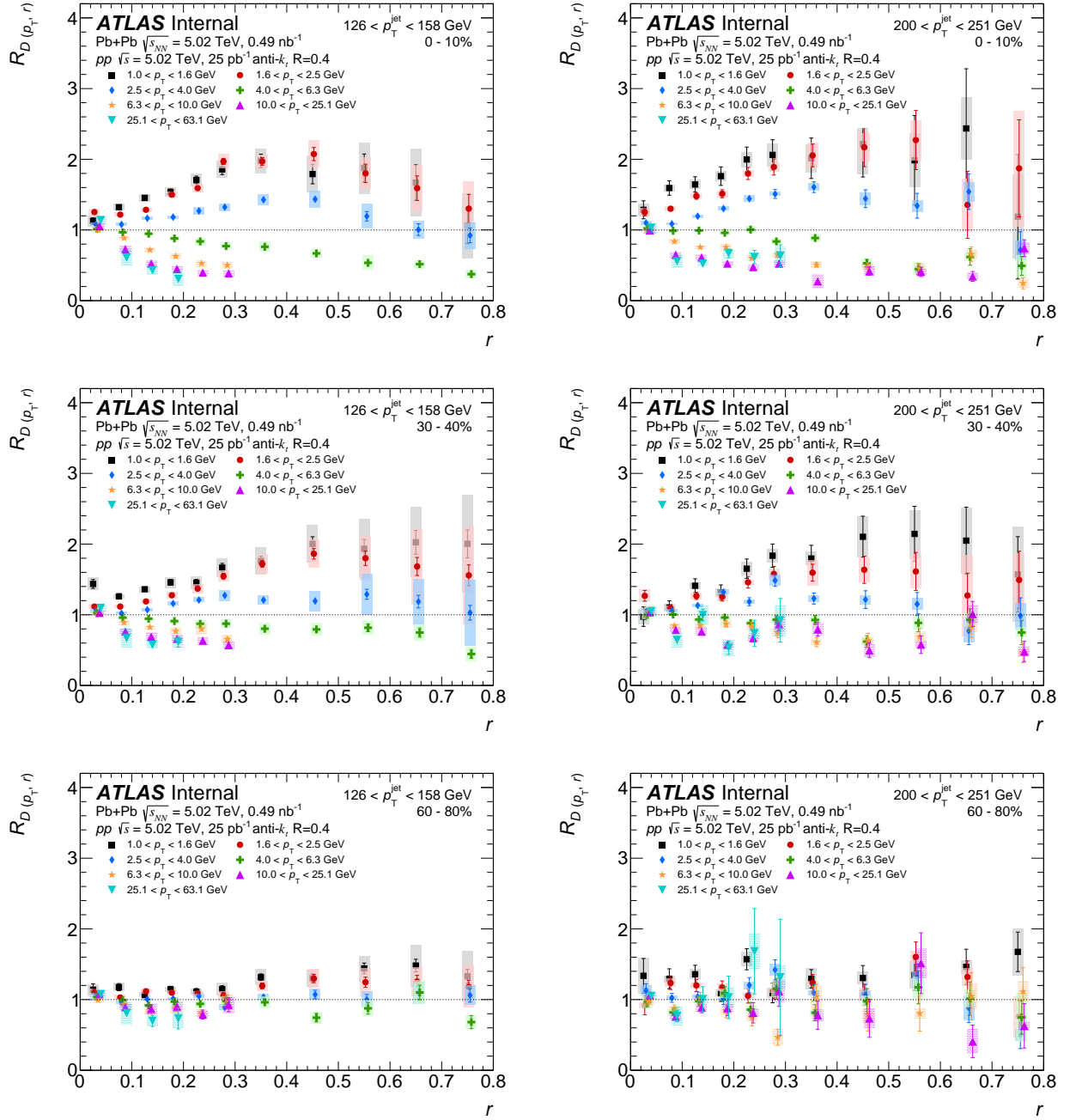


Figure 5: Ratios of $D(p_T, r)$ distributions in Pb+Pb and pp collisions as a function of angular distance r for p_T^{jet} of 126 to 158 GeV (left) and of 200 to 251 GeV (right) for six p_T selections. Different centrality selections are shown: 0–10% (top), 30–40% (middle), 60–80% (bottom). The vertical bars on the data points indicate statistical uncertainties while the shaded boxes indicate systematic uncertainties. The widths of the boxes are not indicative of the bin size and the points are shifted horizontally for better visibility.

8 Discussion

8.1 $R_{D(p_T, r)}$ distributions

Here the centrality, p_T^{jet} and the charged-particle p_T dependence of the $R_{D(p_T, r)}$ distributions introduced in Section 7 are discussed.

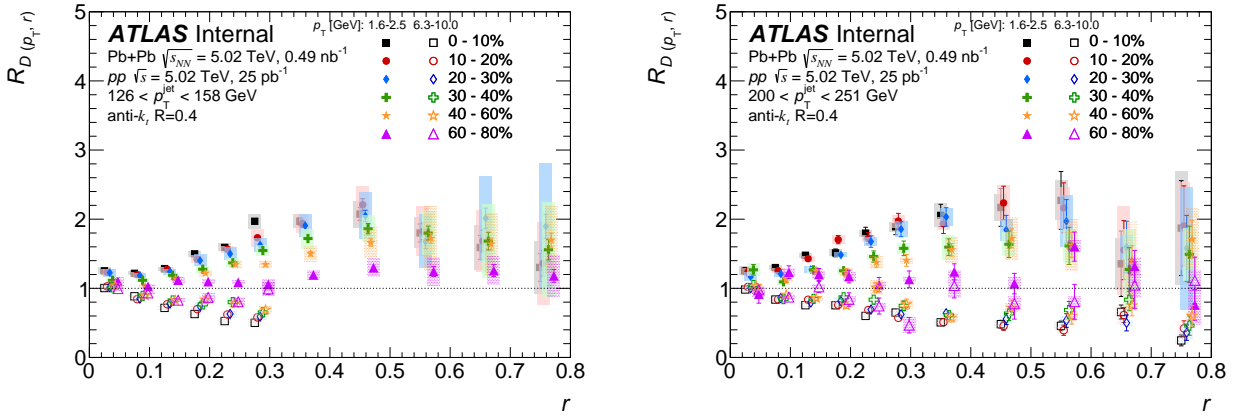


Figure 6: The $R_{D(p_T, r)}$ distributions for p_T^{jet} of 126–158 GeV and 200–251 GeV as a function of angular distance r for two p_T selections, 1.6–2.5 GeV (closed symbols) and 6.3–10.0 GeV (open symbols), and six centrality intervals. The vertical bars on the data points indicate statistical uncertainties while the shaded boxes indicate systematic uncertainties. The widths of the boxes are not indicative of the bin size and the points are shifted horizontally for better visibility.

The centrality dependence of $R_{D(p_T, r)}$ for two charged-particle p_T intervals: 1.6–2.5 GeV and 6.3–10.0 GeV, and two different p_T^{jet} ranges: 126–158 GeV and 200–251 GeV, is presented in Figure 6. For both p_T^{jet} selections and 1.6–2.5 GeV charged particles, the magnitude of the excess increases with increasing collision centrality and r for $r < 0.3$. The magnitude of the excess is approximately a factor of two in the most central collisions for $r > 0.3$. A continuous centrality dependent suppression of yields of charged-particles with $6.3 < p_T < 10.0$ GeV is observed. The magnitude of the modifications decreases with decreasing collision centrality for both p_T intervals and p_T^{jet} selections.

The $R_{D(p_T, r)}$ distributions for low and high p_T particles in the different p_T^{jet} selections are directly overlaid in Figure 7. These distributions are for the 0–10% most central collisions, and suggest an enhancement in $R_{D(p_T, r)}$ with increasing p_T^{jet} for $r < 0.25$ for low p_T charged particles. No significant p_T^{jet} dependence is seen at larger r values. Furthermore, the $R_{D(p_T, r)}$ distributions for high- p_T charged particles do not seem to show any significant p_T^{jet} dependence at any r .

In Figure 5, it was shown that for central and mid-central collisions, there is an enhancement of charged particles with $p_T < 4.0$ GeV and a suppression of charged particles with $p_T > 4.0$ GeV. In Figure 8 the p_T dependance for selections in r is directly investigated for 126–158 GeV and 200–251 GeV jets, in the following centrality intervals: 0–10%, 30–40% and 60–80%. Interestingly, there is no significant suppression of the yields in Pb+Pb collisions for $r < 0.05$ at all measured p_T . For larger r values the yields are enhanced for charged-particles with $p_T < 4$ GeV and suppressed for higher p_T charged-particles in both the 0–10% and 30–40% centrality selections and both p_T^{jet} ranges presented here. The magnitude of the

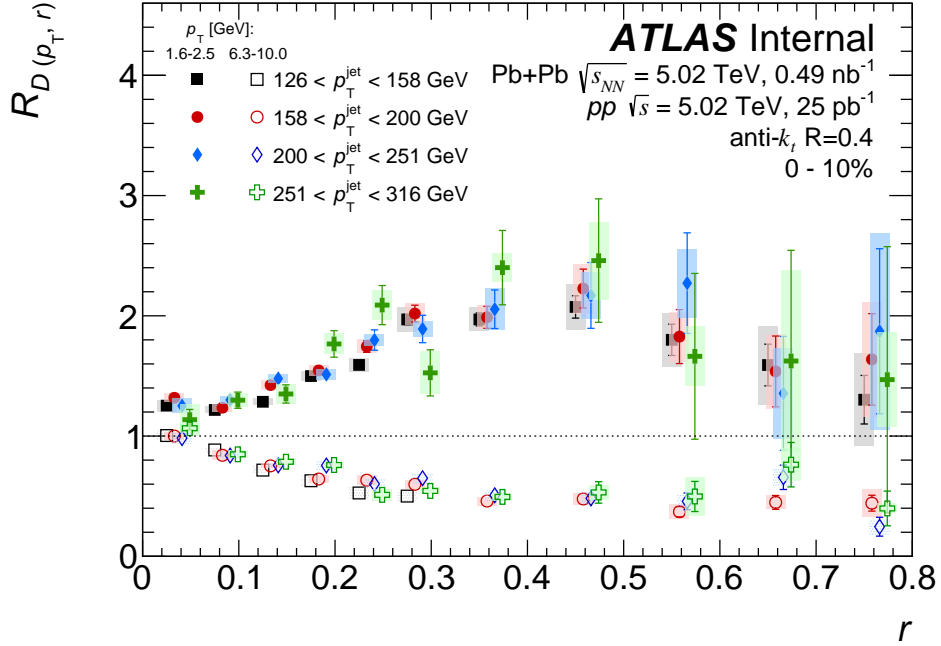


Figure 7: $R_{D(p_T, r)}$ as a function of r for 0–10% collisions for charged particles with $1.0 < p_T < 1.6$ GeV (closed symbols) and $6.3 < p_T < 10.0$ GeV (open symbols) for different p_T^{jet} selections. The vertical bars on the data points indicate statistical uncertainties while the shaded boxes indicate systematic uncertainties. The widths of the boxes are not indicative of the bin size and the points are shifted horizontally for better visibility.

enhancement increases for decreasing p_T at in the $p_T < 4$ GeV p_T , while the suppression is enhanced with increasing p_T for $4 < p_T < 10$ GeV, after which it is approximately constant. At fixed p_T the magnitude of the deviation from unity is largest for $0.3 < r < 0.4$ and $0.5 < r < 0.6$. In the 60–80% central collisions, the same trend remains true (but with smaller magnitude modifications) for $126 < p_T^{\text{jet}} < 158$ GeV; for the higher p_T^{jet} selection the larger uncertainties do not allow a clear conclusion to be drawn for peripheral collisions.

One possible explanation of the modification of the jet fragmentation in the kinematic region of $p_T > 4$ GeV range is the larger expected energy loss of gluon-initiated jets leading to a relative enhancement of quark jets in Pb+Pb collisions compared to pp collisions at a given p_T^{jet} value [18, 56]. Since gluon jets have a broader distribution of particle transverse momentum with respect to the jet direction compared to quark-initiated jets [57], such an effect could potentially describe the narrowing of particle distribution around the jet direction for particles with $p_T > 4.0$ GeV observed here, though no calculations of this are available.

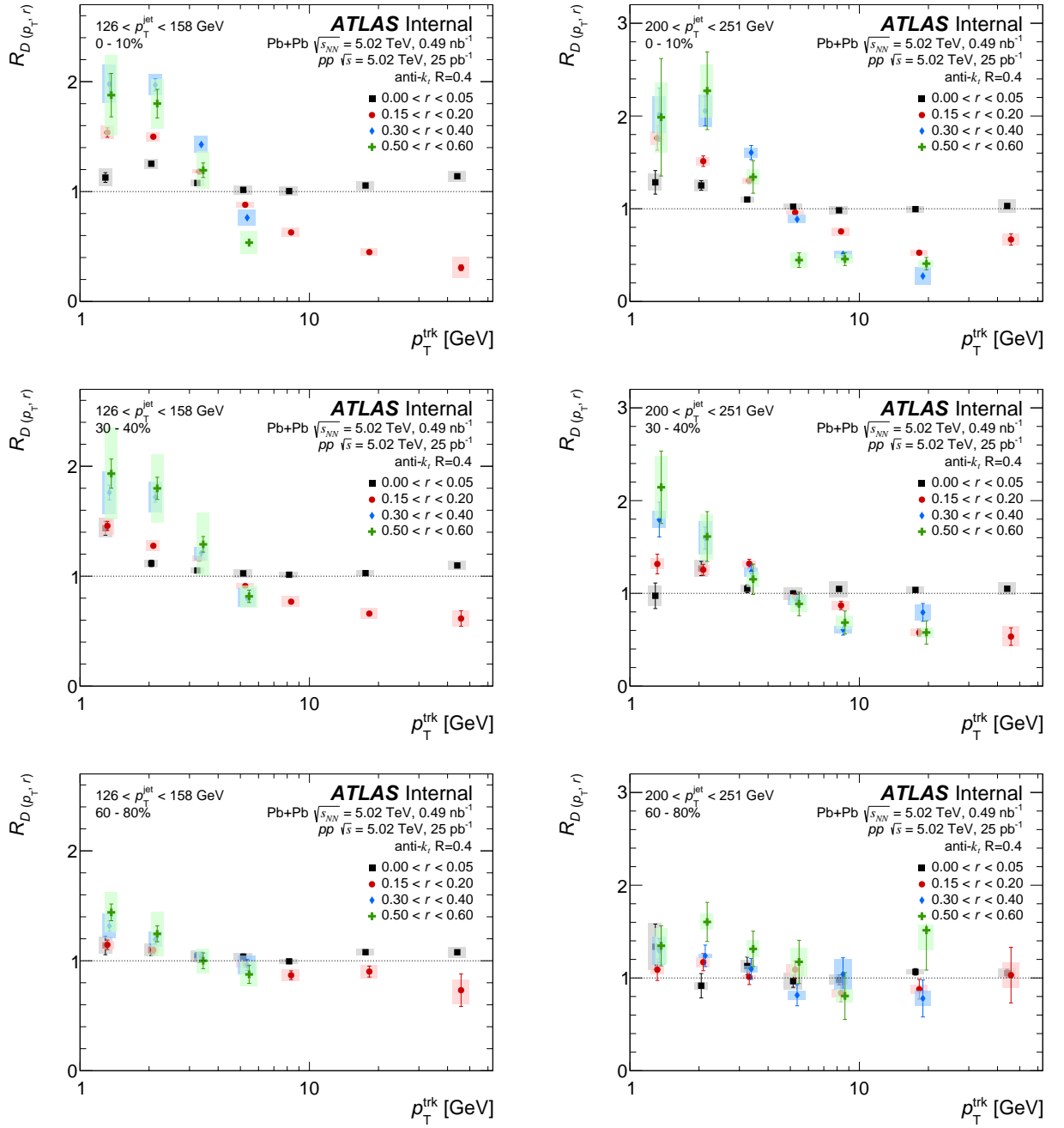


Figure 8: $R_{D(p_T, r)}$ as a function of p_T for 0–10% (top), 30–40% (middle), and 60–80% (bottom) Pb+Pb collisions in two different p_T^{jet} selections: 126–158 GeV (left) and 200–251 GeV (right). The different colors indicate different angular distances from the jet axis. The vertical bars on the data points indicate statistical uncertainties while the shaded boxes indicate systematic uncertainties. The widths of the boxes are not indicative of the bin size and the points are shifted horizontally for better visibility.

8.2 Differences of $D(p_T, r)$ distributions

In addition to the ratios of the $D(p_T, r)$ distributions, differences between the charged particle yields are also evaluated to quantify the modification in terms of the particle density. These are given as:

$$\Delta D(p_T, r) = D(p_T, r)_{\text{Pb+Pb}} - D(p_T, r)_{pp}$$

These differences are presented as a function of r for different p_T selections in 0–10% central collisions in Figure 9. These distributions show an excess in the charged-particle yield density for Pb+Pb collisions compared to pp collisions for charged particles with $p_T < 4.0$ GeV. This excess ranges from 0.5 to 4 particles per unit area per GeV for 1 GeV charged particles in 126–158 GeV jets for 0–10% central Pb+Pb collisions and increases with increasing p_T^{jet} . The largest excesses for charged particles with $p_T < 4.0$ GeV is within the jet cone. For large r values, the density decreases, but remains positive. A depletion for higher p_T particles of approximately 0.5 particles per unit area per GeV is seen for 126–158 GeV jets in 0–10% central Pb+Pb collisions. The magnitude of this depletion increases for higher p_T^{jet} . There is a minimum in the $\Delta D(p_T, r)$ distributions of charged particles with $4.0 < p_T < 25.1$ GeV at $0.05 < r < 0.10$ that is seen at many p_T^{jet} ranges under investigation. The magnitudes of the excesses and deficits discussed here are dependent on the sizes of the charged-particle p_T selections chosen. In order to remove that dependence, Section 8.3 provides similar quantities in which a wider charged-particle p_T range is integrated over.

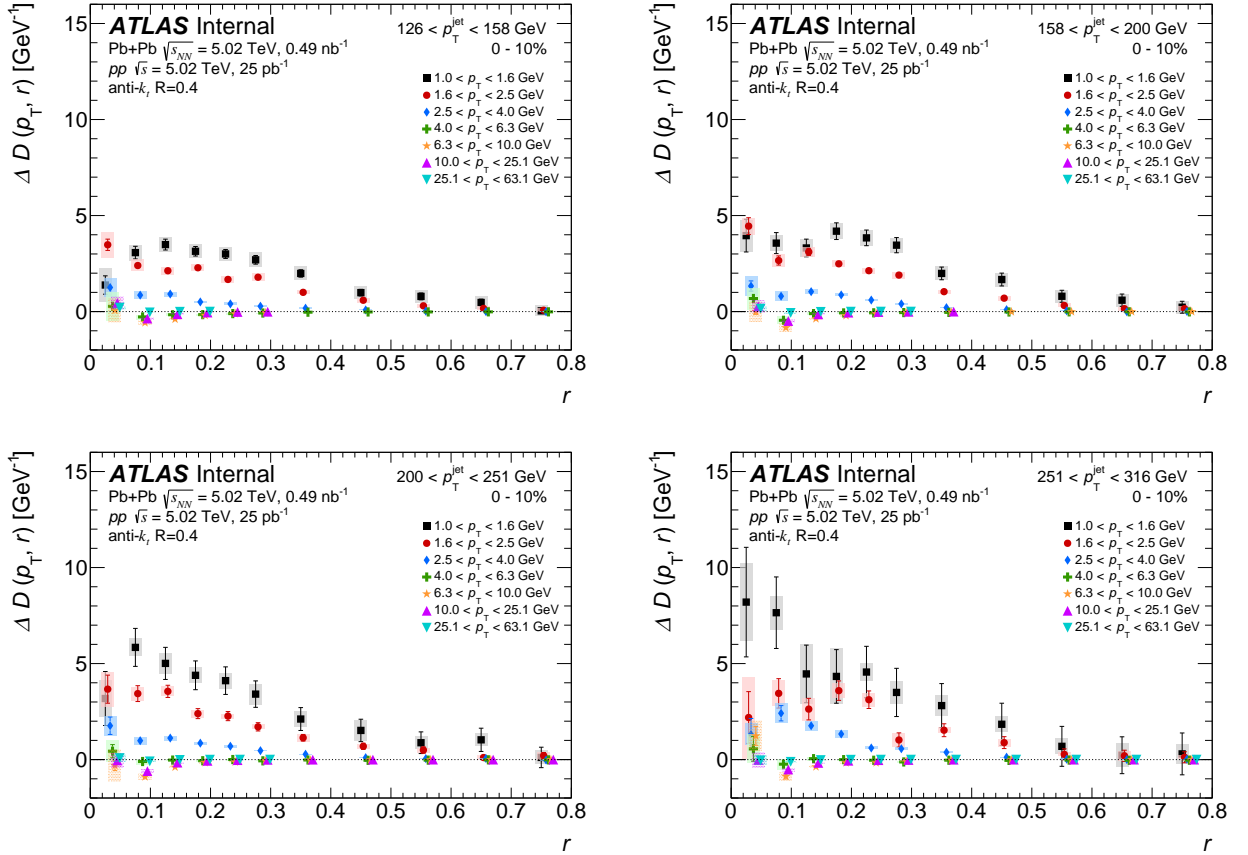


Figure 9: $\Delta D(p_T, r)$ as a function of r in central collisions for all p_T ranges in four p_T^{jet} selections: 126–158 GeV, 158–200 GeV, 200–251 GeV, and 251–316 GeV. The vertical bars on the data points indicate statistical uncertainties while the shaded boxes indicate systematic uncertainties. The widths of the boxes are not indicative of the bin size and the points are shifted horizontally for better visibility.

8.3 p_T integrated distributions

Motivated by similar studies of the enhancement of soft fragments in jet fragmentation functions in Pb+Pb compared to pp collisions from Ref. [18], the $D(p_T, r)$ distributions can be integrated for charged particles with $p_T < 4$ GeV to construct the quantities $\Theta(r)$ and $P(r)$ defined as:

$$\Theta(r) = \int_1^4 D(p_T, r) dp_T$$

$$P(r) = \int_0^r \int_1^4 D(p_T, r') dp_T dr'$$

The $\Theta(r)$ values are integrated between 1–4 GeV in p_T of charged particles to provide a summary look at the p_T region of enhancement discussed above. The $P(r)$ values further add a running integral over r and provide information about the jet shape. Both of these quantities can be compared between the pp and Pb+Pb systems to give the following distributions:

$$\begin{aligned}\Delta_{\Theta(r)} &= \Theta(r)_{\text{Pb+Pb}} - \Theta(r)_{pp} \\ R_{\Theta(r)} &= \frac{\Theta(r)_{\text{Pb+Pb}}}{\Theta(r)_{pp}} \\ R_{P(r)} &= \frac{P(r)_{\text{Pb+Pb}}}{P(r)_{pp}}\end{aligned}$$

(the quantity $\Delta_{P(r)}$ can also be analogously defined, but is omitted from the present discussion). These aggregate quantities are intended to provide some summary information about the location with respect to the jet axis, magnitude, and $p_{\text{T}}^{\text{jet}}$ dependence of the low- p_{T} charged-particle excess discussed above. The ratio quantities are useful for comparisons to other Pb+Pb measurements; $\Delta_{\Theta(r)}$ is very similar to $\Delta D(p_{\text{T}}, r)$, however it is integrated over charged-particle p_{T} from 1–4 GeV.

Figure 10 shows the $\Delta_{\Theta(r)}$ distributions as a function of r in centrality intervals: 0–10%, 30–40%, 60–80%. In the most central collisions, a significant $p_{\text{T}}^{\text{jet}}$ dependence to $\Delta_{\Theta(r)}$ is observed; for $r < 0.4$ (particles within the jet cone) $\Delta_{\Theta(r)}$ increases with increasing $p_{\text{T}}^{\text{jet}}$. The value of $\Delta_{\Theta(r)}$ decreases in more peripheral collisions and the $p_{\text{T}}^{\text{jet}}$ dependence is no longer significant.

Figure 11 shows $R_{\Theta(r)}$ and $R_{P(r)}$ for the following centrality intervals: 0–10%, 30–40% and 60–80%. The $R_{\Theta(r)}$ distributions of the most central collisions show a maximum for $r \sim 0.4$ and a flattening or a decrease for larger r . However, since $R_{\Theta(r)}$ remains at or above unity for the full range of r values presented $R_{P(r)}$ continues to slowly increase with increasing r over the full measured range. In more peripheral collisions, the magnitude of the excess is reduced and the trends in $R_{\Theta(r)}$ are less clear, however the slow increase of $R_{P(r)}$ is clearly seen for the 30–40% central collisions.

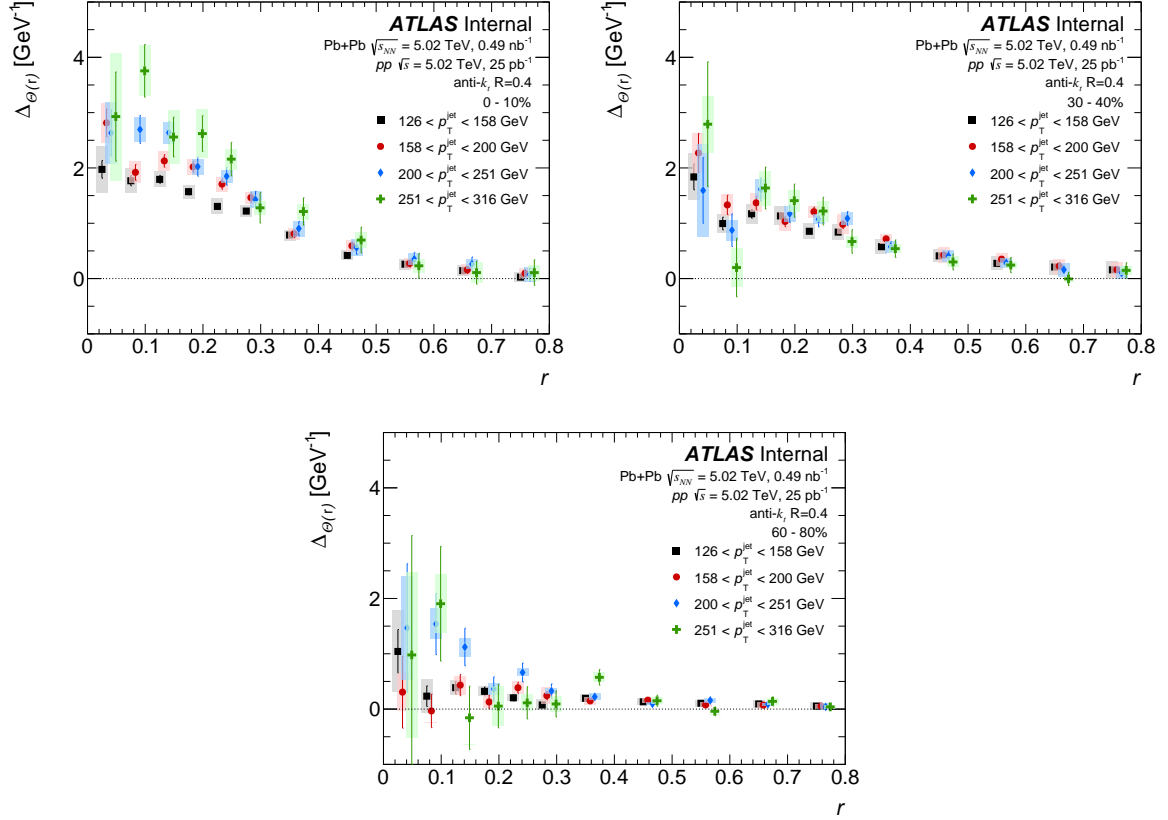


Figure 10: $\Delta\theta(r)$ as a function of r for charged-particles with $p_T < 4$ GeV in four p_T^{jet} selections: 126–158 GeV, 158–200 GeV, 200–251 GeV, and 251–316 GeV and three centrality selections: 0–10% (left), 30–40% (middle) and 60–80% (right). The vertical bars on the data points indicate statistical uncertainties while the shaded boxes indicate systematic uncertainties. The widths of the boxes are not indicative of the bin size and the points are shifted horizontally for better visibility.

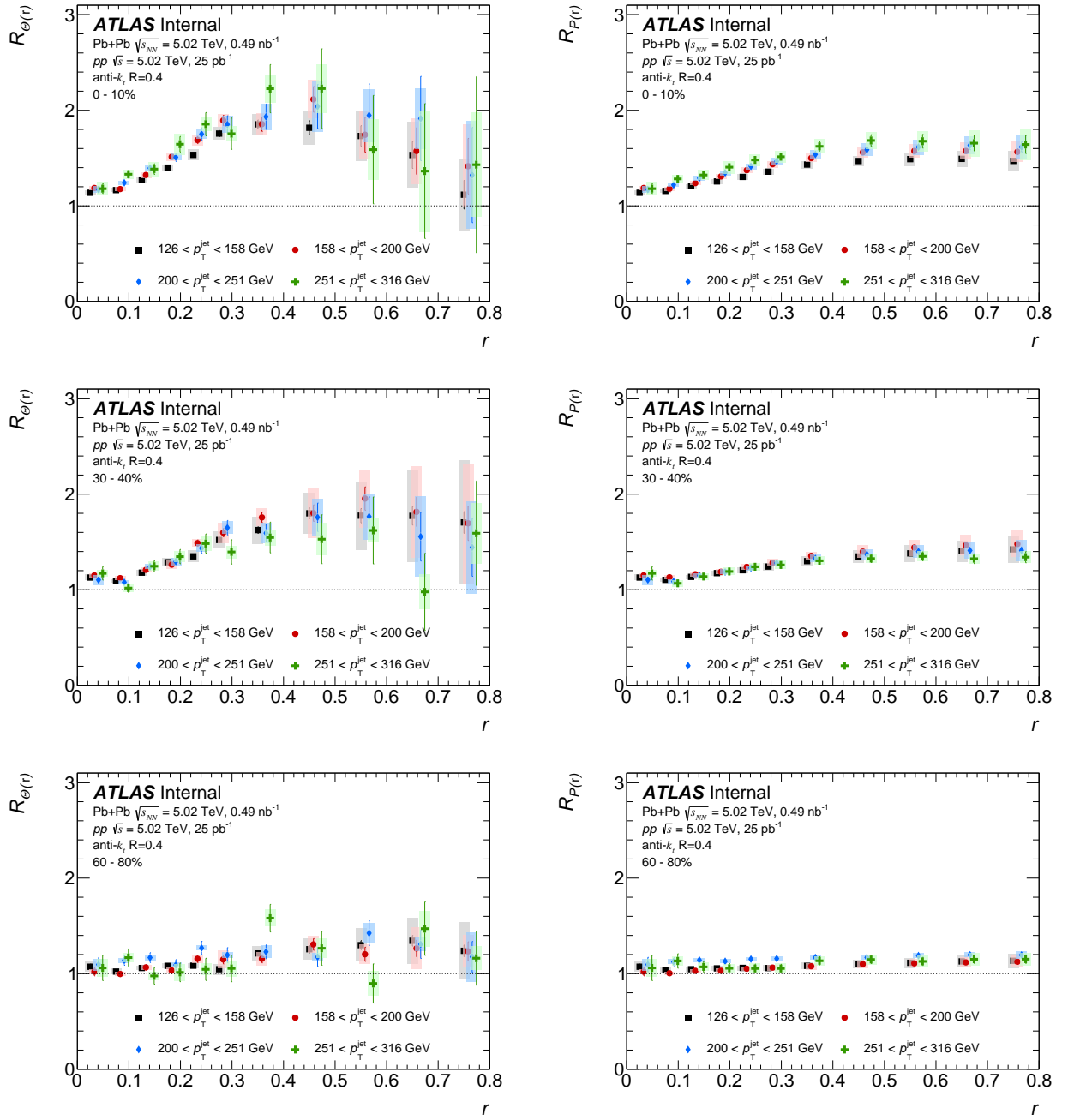


Figure 11: $R_{\Theta(r)}$ (left) and $R_{P(r)}$ (right) as a function of r for charged-particles with $p_T < 4$ GeV ranges in four p_T^{jet} selections: 126–158 GeV, 158–200 GeV, 200–251 GeV, and 251–316 GeV and three centrality selections: 0–10% (top), 30–40% (middle) and 60–80% (bottom). The vertical bars on the data points indicate statistical uncertainties while the shaded boxes indicate systematic uncertainties. The widths of the boxes are not indicative of the bin size and the points are shifted horizontally for better visibility.

9 Summary

This paper presents a measurement of the yields of charged particles, $D(p_T, r)$, inside and around $R = 0.4$ anti- k_t jets with $|y^{\text{jet}}| < 1.7$ up to a distance of $r = 0.8$ from the jet axis. The yields are measured in intervals of p_T^{jet} from 126 to 316 GeV in Pb+Pb and pp collisions at 5.02 TeV as a function of charged particle p_T and the angular distance r between the jet axis and charged particle.

These results show a broadening of the $D(p_T, r)$ distribution for low p_T particles inside the jet in central Pb+Pb collisions compared to those in pp collisions while for higher p_T particles angular distributions are narrower in Pb+Pb collisions compared to pp collisions. These modifications are centrality dependent, and decrease with decreasing centrality. The $R_{D(p_T, r)}$ distributions for charged particles with $p_T < 4$ GeV are above unity and grow with increasing angular separation up to $r \sim 0.3$, showing weak to no dependence on r in the interval $0.3 < r < 0.6$ followed with a small decrease in the enhancement for $0.6 < r < 0.8$. For charged particles with $p_T > 4$ GeV, a suppression in $R_{D(p_T, r)}$ is observed, and the distributions decrease with increasing r for $0.05 < r < 0.3$, with no r dependence for $r > 0.3$. For all charged-particle p_T values, the $R_{D(p_T, r)}$ values are greater than or equal to unity for $r < 0.05$. Between $0.1 < r < 0.25$, a statistically significant trend of increasing $R_{D(p_T, r)}$ with increasing p_T^{jet} is observed for low- p_T particles. No significant p_T^{jet} dependence is seen for particles with $p_T > 4$ GeV. These new measurements will help distinguish modification of the jet itself and any response of the hot QCD matter to the jet.

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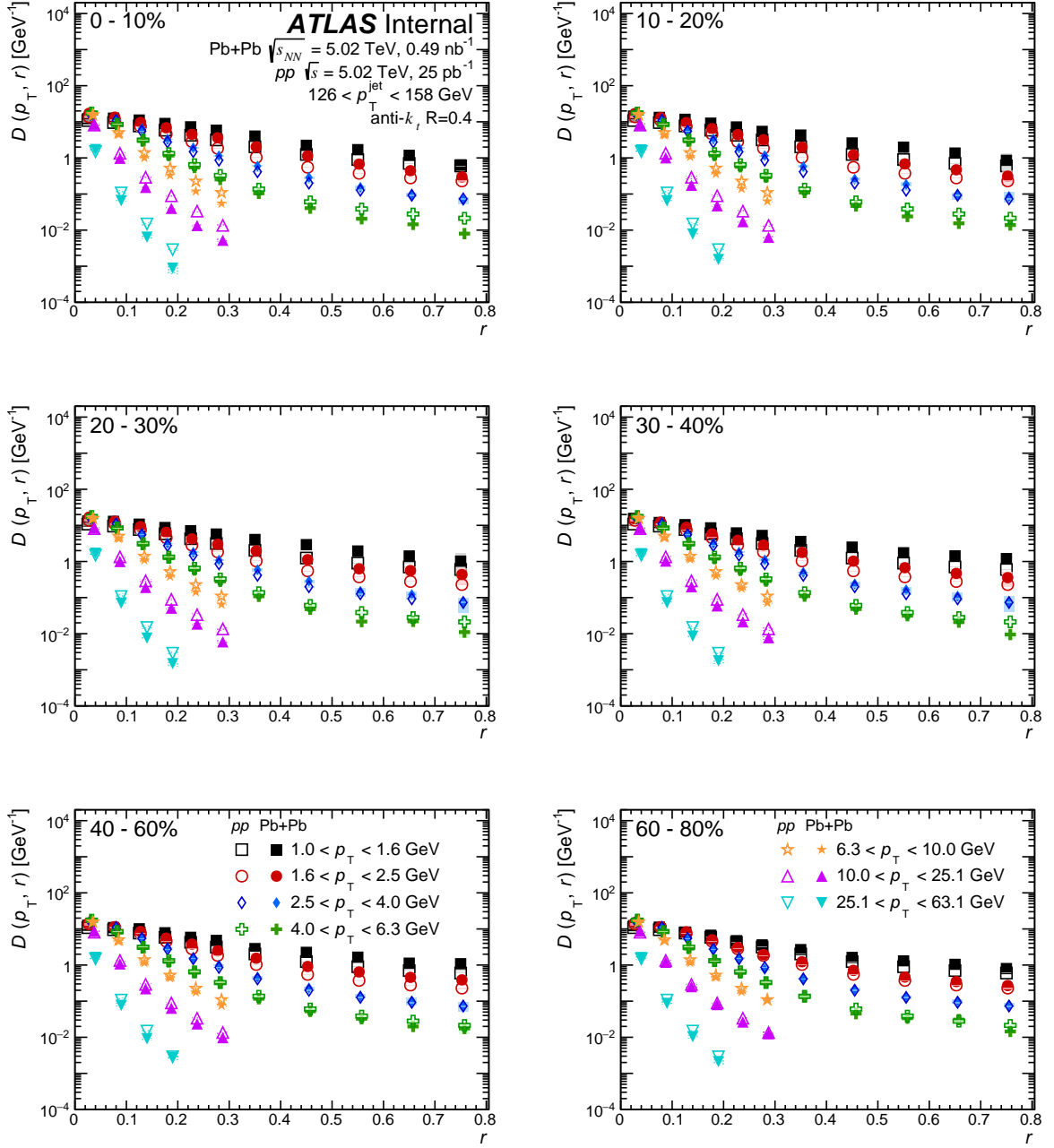
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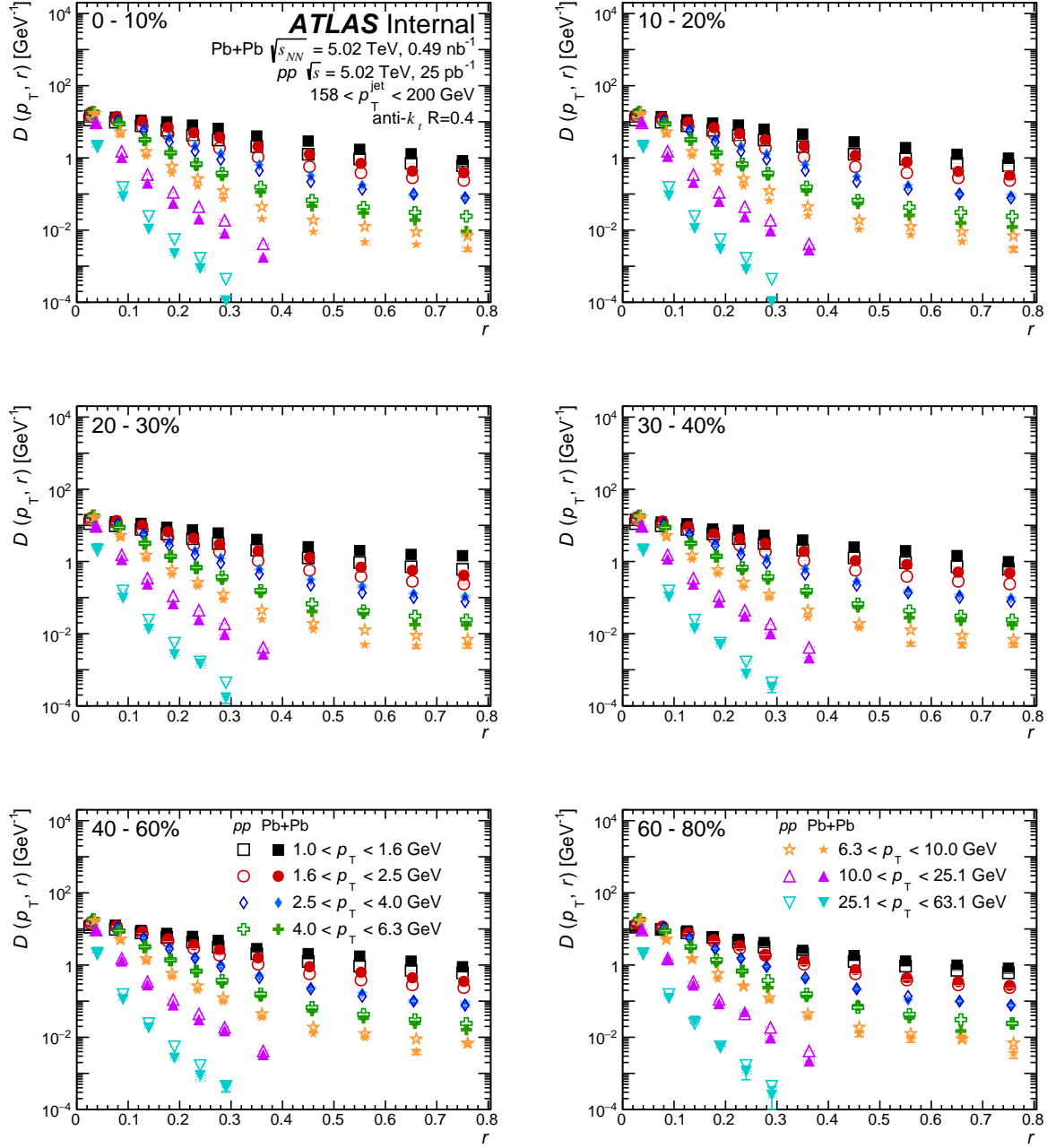
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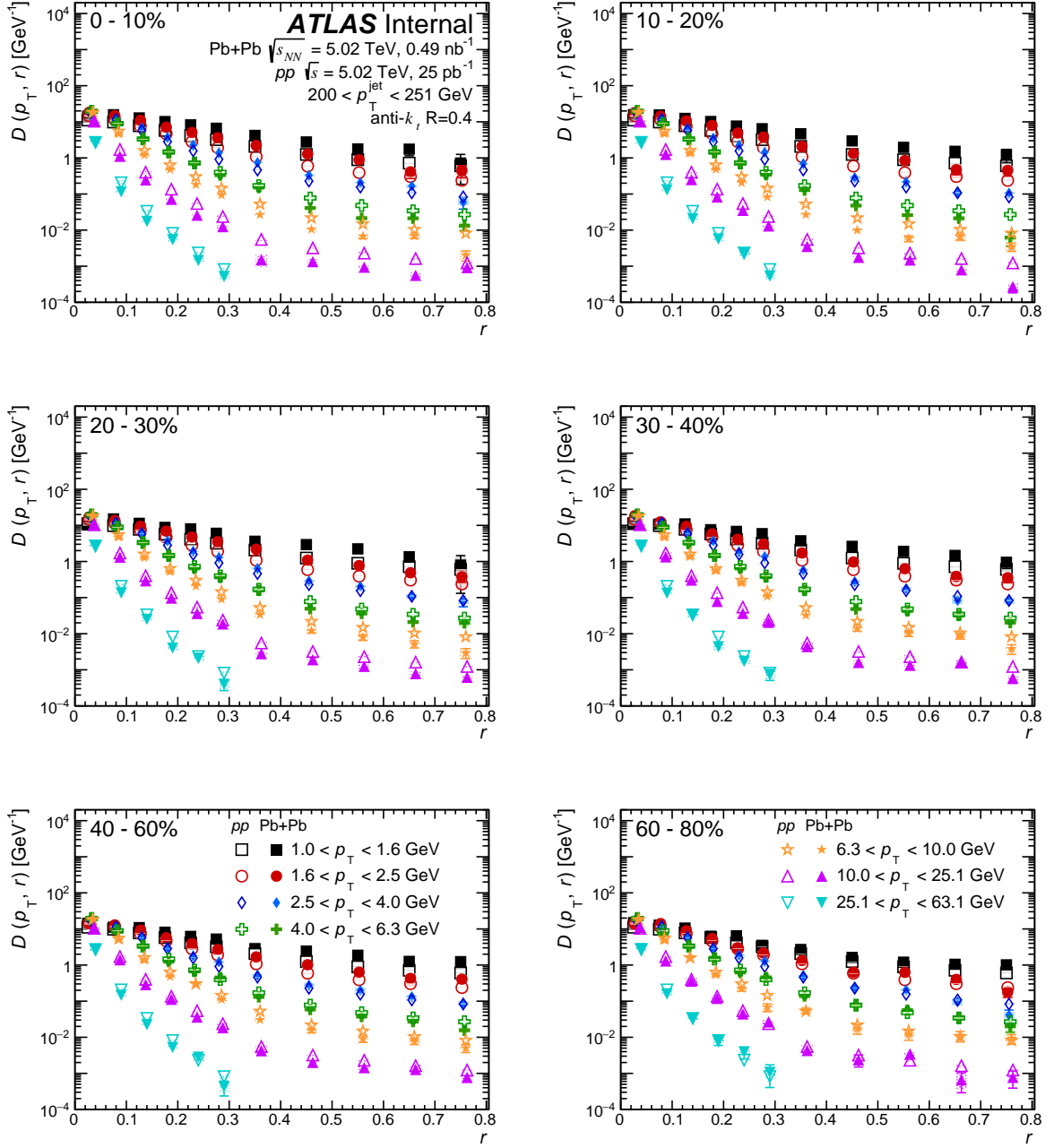
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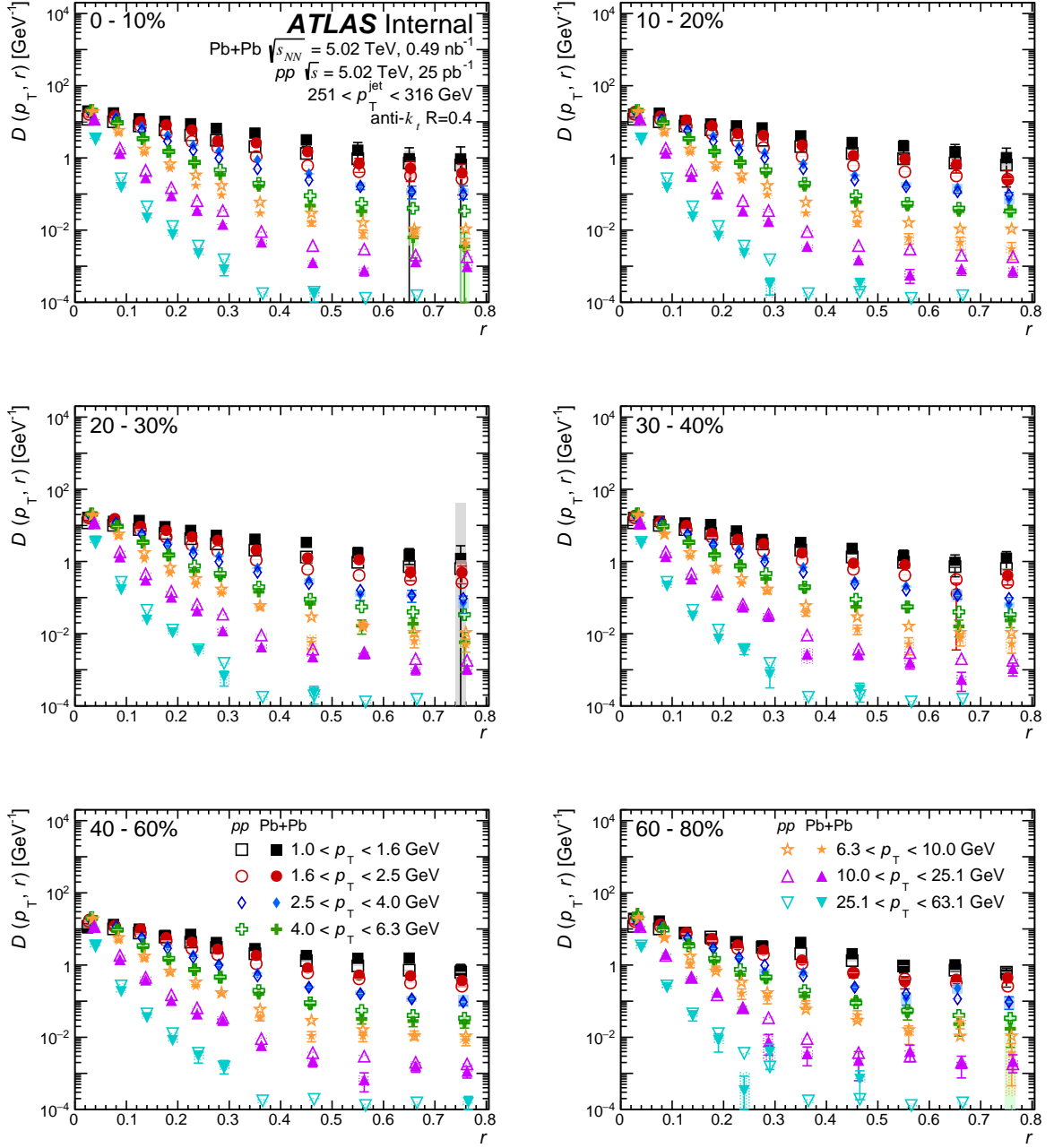
Auxiliary material

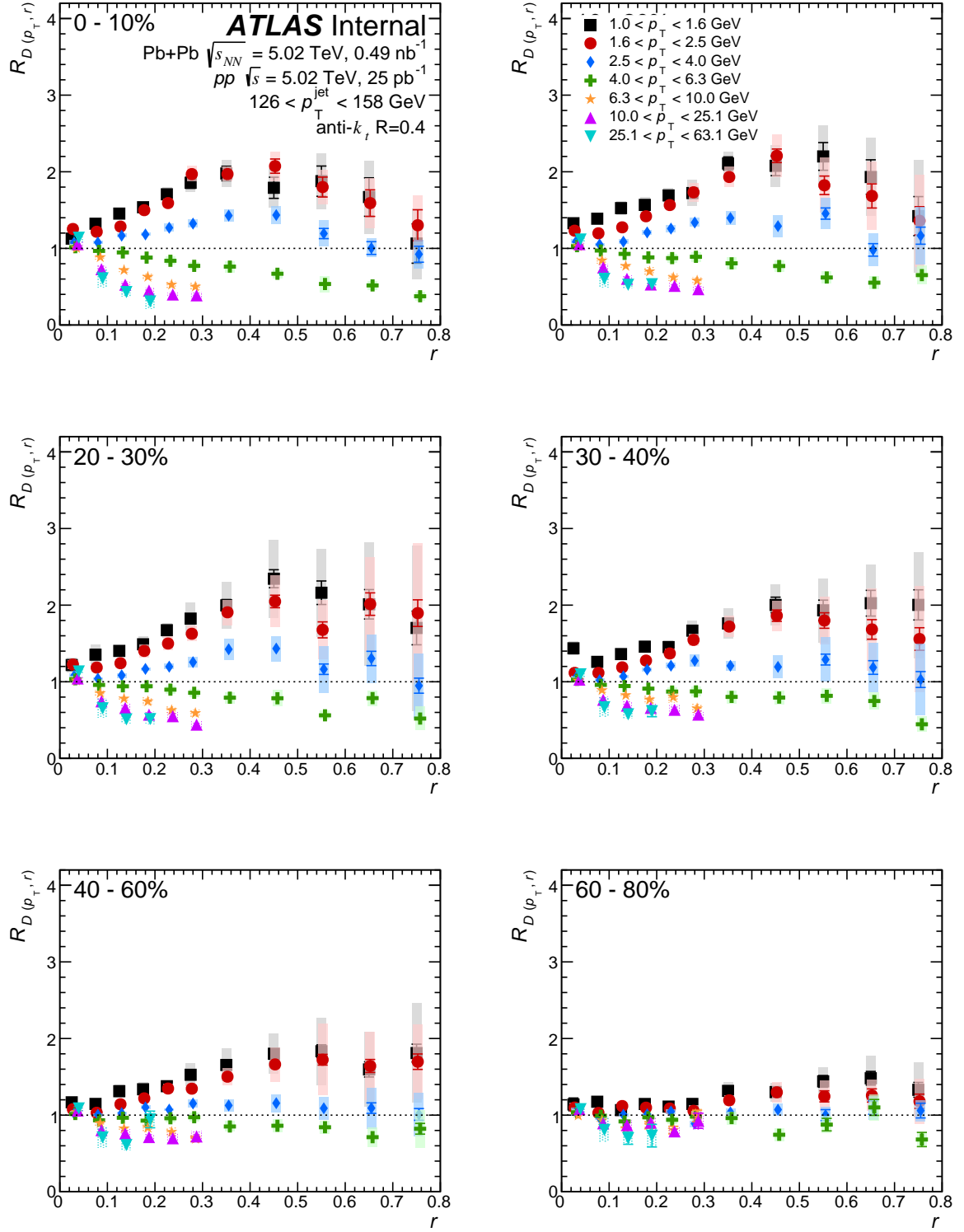
In an ATLAS paper, auxiliary plots and tables that are supposed to be made public should be collected in an appendix that has the title ‘Auxiliary material’. This information will appear on the public webpage, but will not be included in the document submitted to arXiv and to the journal.

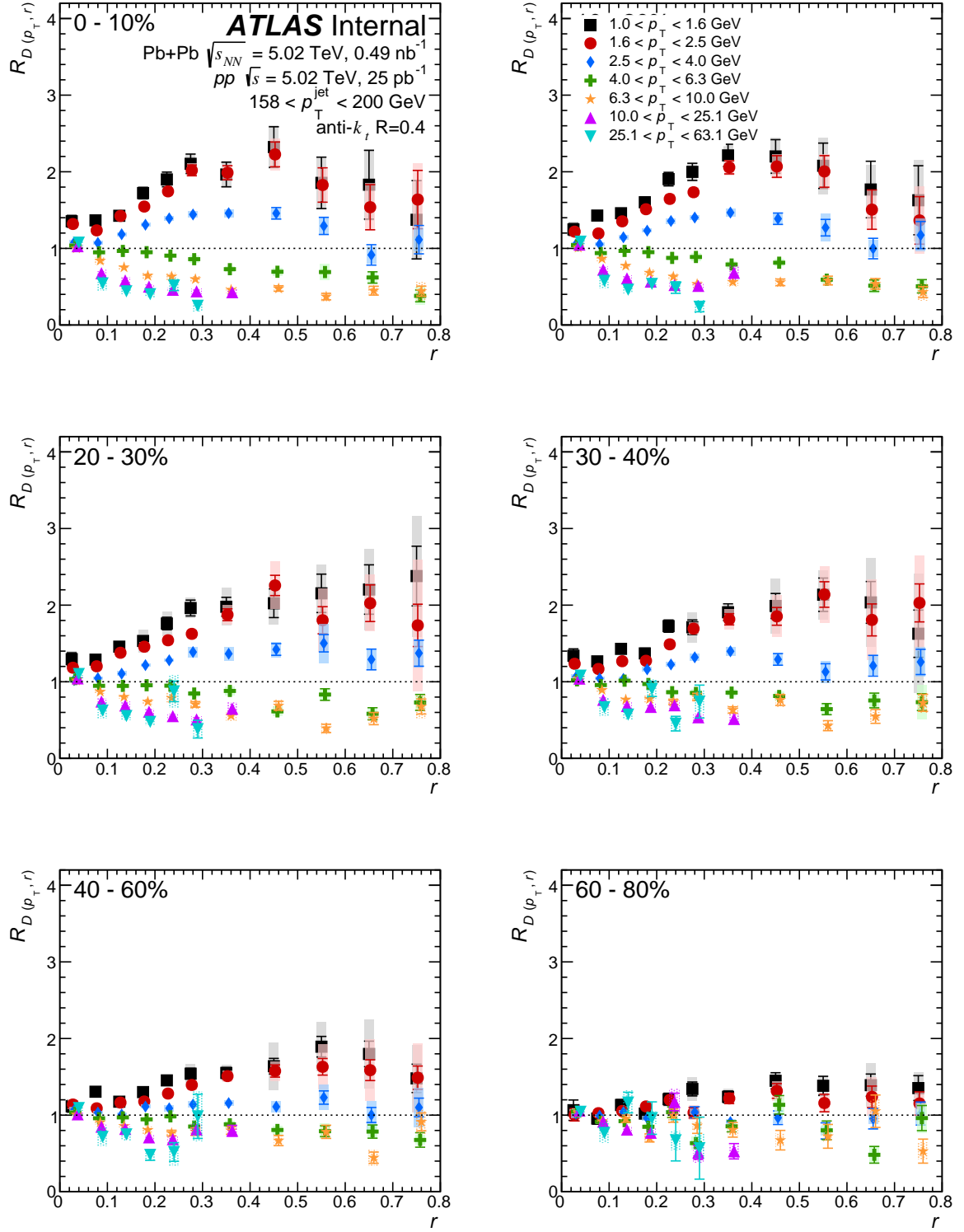
Figure 12: Full set of $D(p_T, r)$ distributions for 126–158 GeV jets.

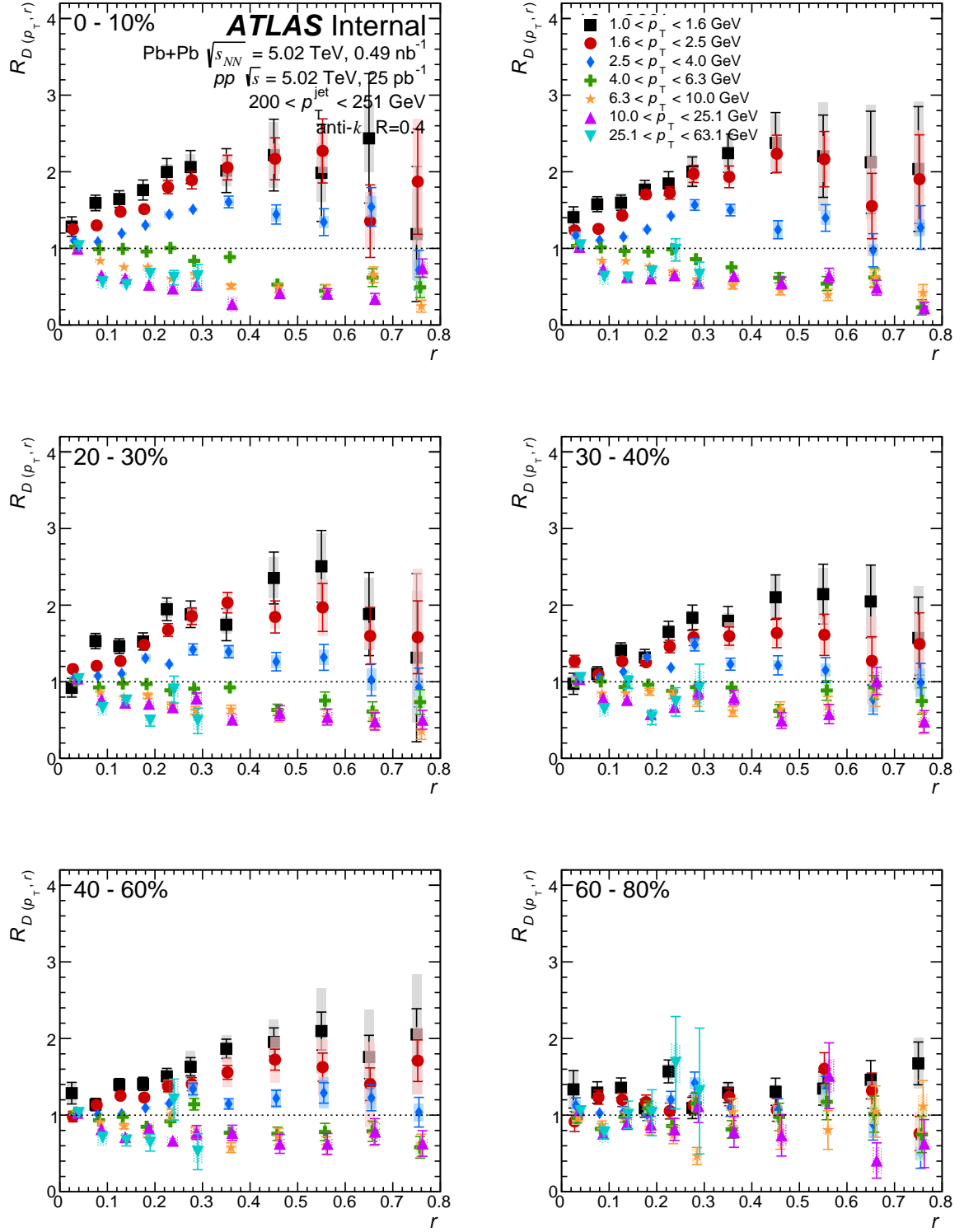
Figure 13: Full set of $D(p_T, r)$ distributions for 158–200 GeV jets.

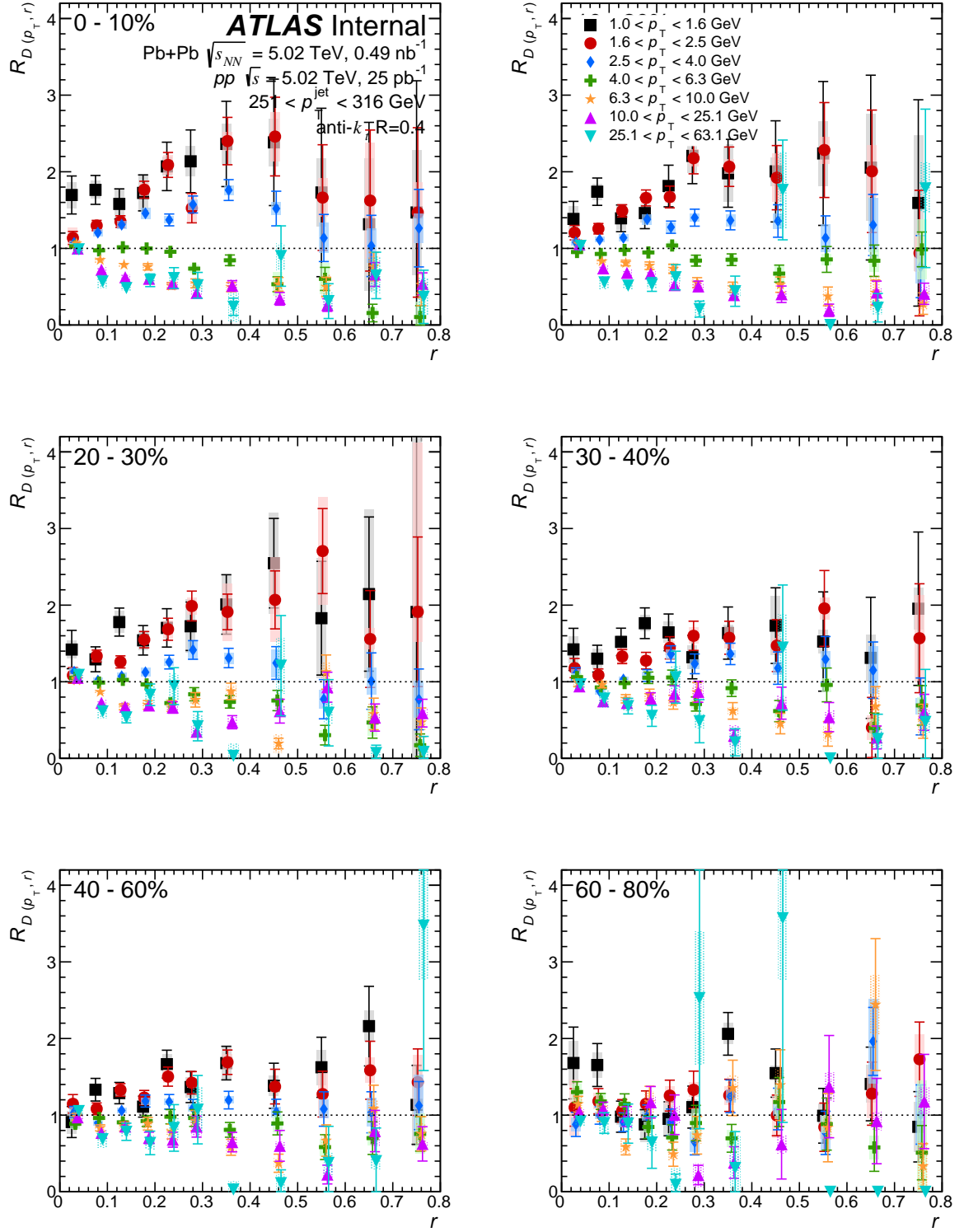
Figure 14: Full set of $D(p_T, r)$ distributions for 200–251 GeV jets.

Figure 15: Full set of $D(p_T, r)$ distributions for 251–316 GeV jets.

Figure 16: Full set of $R_{D(p_T, r)}$ distributions for 126–158 GeV jets.

Figure 17: Full set of $R_{D(p_T, r)}$ distributions for 158–200 GeV jets.

Figure 18: Full set of $R_{D(p_T, r)}$ distributions for 200–251 GeV jets.

Figure 19: Full set of $R_D(p_T, r)$ distributions for 251–316 GeV jets.