PH-107

Quantum Physics and Applications

Elements of Statistical Physics-Ill

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Recap

We are interested in determining the number of particles *dN(E)* in the system with energy between *E* and *E+dE*.

$$N(E) dE = g(E) f(E) dE$$

- $N(E) \ dE$ Number of particles per unit volume in the energy range E and E+dE.
- $g(E) \ dE$ Density of states (number of energy states per unit volume) in the energy range E and E+dE.
- f(E) Probability Distribution function: Dependence on Particle Characteristics

Classical Particles and Quantum (Bosons and Fermions) Particles

Recap

We need to optimise $Q(\{N_i\})$ with two constraints

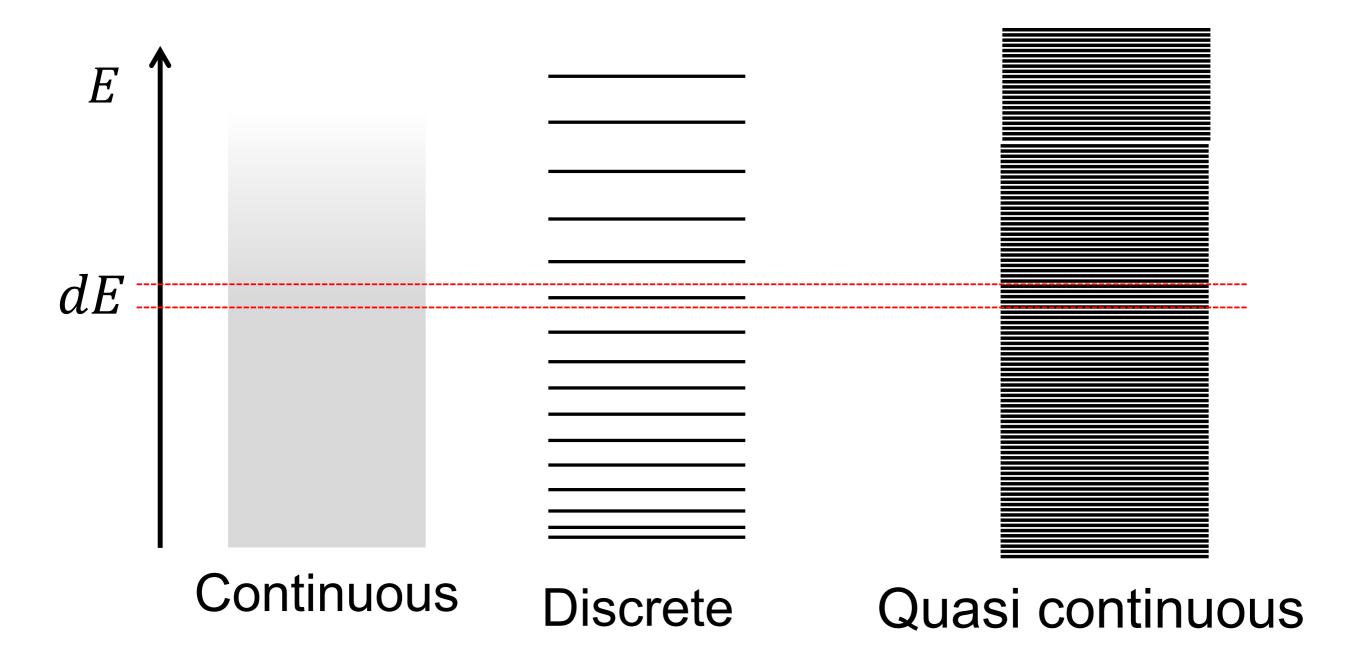
$$\sum_{i=1}^{\infty} N_i = N$$
 and $\sum_{i=1}^{\infty} E_i N_i = E$

$$f_{\rm MB}(E) = A e^{-\frac{E}{k_B T}}$$

$$f_{\rm FD}(E) = \frac{1}{1 + e^{(E - E_F)/k_B T}}$$

$$f_{\rm BE}(E) = \frac{1}{e^{(E/k_BT)} - 1}$$

Energy Spectrum



g(E)dE: No. of states in the interval dE (about E)

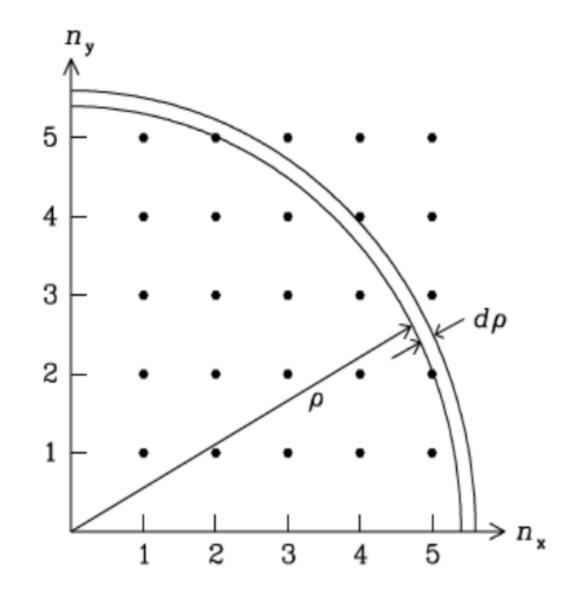
Calculation of g(E): Infinite 3D Box

$$E = \frac{\pi^2 \hbar^2}{2mL^2} (n_x^2 + n_y^2 + n_z^2)$$

Note that we can label the states as (n_x, n_y, n_z) and arrange them in a (imaginary) 3D space defined by the axes, n_x , n_y , and n_z .

Integer Space

This is how we invoke the concept of a integer space (space defined by axes, n_x , n_y , and n_z)



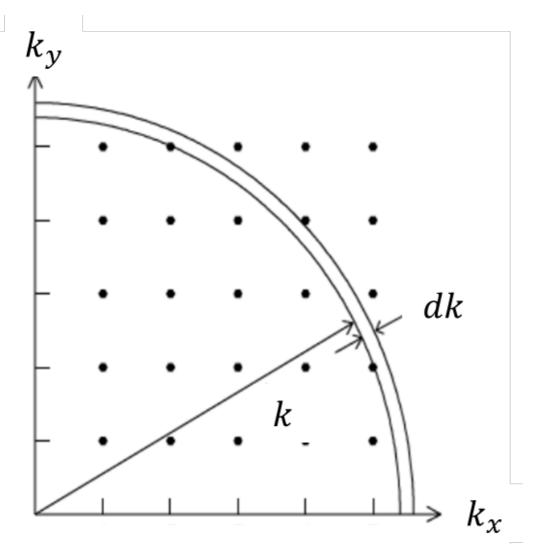
$$\rho_{\rm 2D} = \sqrt{n_x^2 + n_y^2}$$

$$\rho_{3D} = \sqrt{n_x^2 + n_y^2 + n_z^2}$$

All states which lie within the shell of thickness dp about p are being assigned the energy

$$E = \frac{\pi^2 \hbar^2}{2mL^2} \rho^2$$

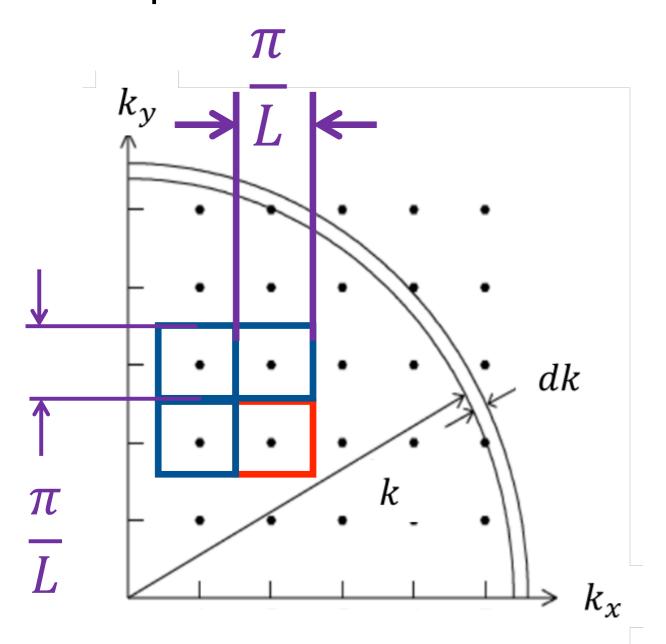
Since $n_{x,y,z}=\frac{k_{x,y,z}\ L}{\pi}$, we could define an equivalent space with axes $k_x=\frac{n_x\pi}{L}, k_y=\frac{n_y\pi}{L}, \ {\rm and}\ k_z=\frac{n_z\pi}{L}$



$$k_{\rm 2D} = \sqrt{k_x^2 + k_y^2}$$

$$k_{3D} = \sqrt{k_x^2 + k_y^2 + k_z^2}$$

What is the area (volume) of the 2D (3D) k-space occupied by each point?



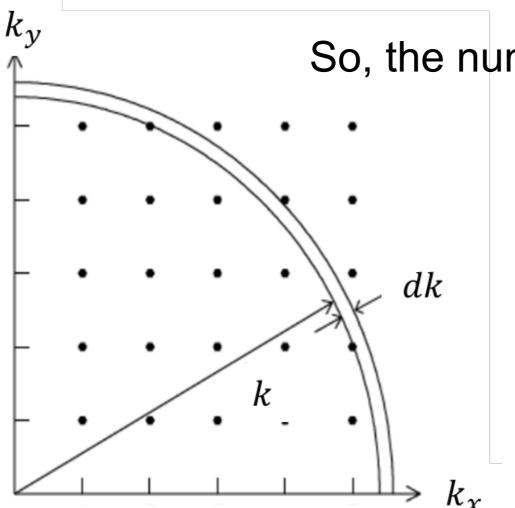
This area is $\Delta k_x \times \Delta k_y = \left(\frac{\pi}{L}\right)^2$ In 3D, this volume is $\left(\frac{\pi}{L}\right)^3$

So, the density of points in the k-space is 1 state per $\left(\frac{\pi}{L}\right)^2$ area

or 1 state per $\left(\frac{\pi}{L}\right)^3$ volume

We get the density of states in the k space, F(k), as

$$F(k) = \frac{1}{\left(\frac{\pi}{L}\right)^3} = \frac{V}{\pi^3}$$



So, the number of states in the octant of the shell is

$$F(k) d\mathbf{k} = \frac{V}{\pi^3} \left(\frac{1}{8}\right) 4\pi k^2 dk$$

Here $d\mathbf{k} = dk_x dk_y dk_z$

Reciprocal Space

But the number of states between k and k+dk should be the same as the number of states between E and E+dE.

$$F(k)d\mathbf{k} = G(E)dE$$

$$F(k)d\mathbf{k} = \left(\frac{L}{\pi}\right)^3 d\mathbf{k} = G(E)dE$$

This implies

$$\frac{F(k)d\mathbf{k}}{L^3} = \frac{F(k)d\mathbf{k}}{V} = \left(\frac{1}{\pi}\right)^3 d\mathbf{k} = \frac{G(E)}{V} dE = g(E)dE$$

i.e.,
$$f(k)d\mathbf{k} = \frac{1}{\pi^3} \left(\frac{1}{8}\right) 4\pi k^2 dk = g(E)dE$$

Reciprocal Space

Now, let us recall
$$E = \frac{\hbar^2 k^2}{2m}$$

And make use of it in

$$\frac{1}{\pi^3} \left(\frac{1}{8}\right) 4\pi k^2 dk = g(E)dE$$

to get the expression for the density of states in terms of energy as

$$g(E) = \frac{1}{4\pi^2} \left(\frac{2m}{\hbar^2}\right)^{\frac{3}{2}} \sqrt{E}$$

Reciprocal Space

If there are other degeneracies of each k-state (for example the spin-degeneracy of Fermions), we need to multiply g(E) with the degeneracy factor, (let's say) λ .

So,
$$g(E) = \boxed{\lambda} \times \frac{1}{4\pi^2} \left(\frac{2m}{\hbar^2}\right)^{\frac{3}{2}} \sqrt{E}$$

So, in case of electrons (Fermions) bound to the 3D potential well, we have

$$g(E) = \frac{1}{2\pi^2} \left(\frac{2m}{\hbar^2}\right)^{\frac{3}{2}} \sqrt{E}$$

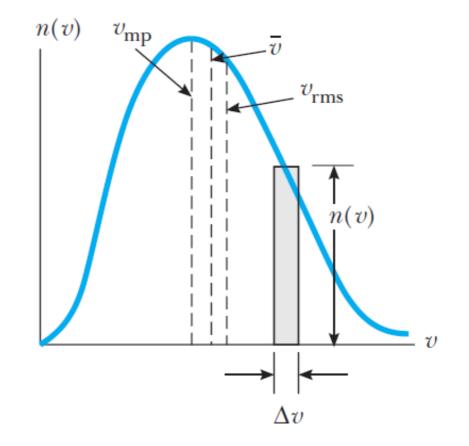
Application of M-B Distribution

Maxwell-Boltzmann Distribution

Maxwell speed distribution function for gas molecules at thermal equilibrium at temperature T

$$n(v) \ dv = \frac{4\pi N}{V} \left(\frac{m}{2\pi k_B T}\right)^{\frac{3}{2}} v^2 e^{-\frac{mv^2}{2k_B T}} \ dv$$

Here, n(v) dv is the number of (ideal) gas molecules (of mass m) per unit volume with speed between v and v + dv.



We know
$$n(E) dE = g(E) f(E) dE$$

In the present case:
$$f_{\rm MB}(E) = A \ e^{-\frac{E}{k_B T}}$$

Also, we know
$$E = \frac{1}{2}mv^2$$

So, every v corresponds to unique E

and therefore,
$$g(E) dE = g(v) dv$$

Density of states in the energy interval E and E+dE

Density of states in the velocity interval v and v+dv

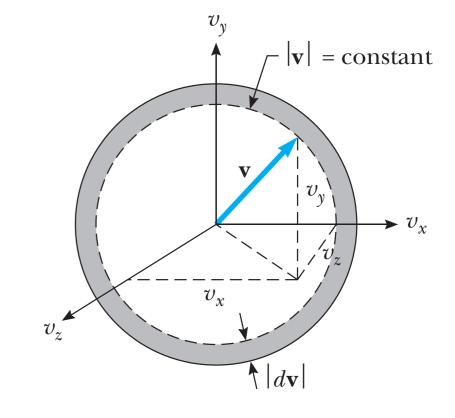
So, let us estimate number of states with velocity between v and v+dv

$$g(v) \ dv = C4\pi v^2 \ dv$$

$$\downarrow$$
 constant

$$n(E) dE = n(v) dv = g(v) Ae^{-\frac{mv^2}{2k_BT}} dv$$

= $A'4\pi v^2 e^{-\frac{mv^2}{2k_BT}} dv$



Total number of particles per unit volume =
$$\frac{N}{V}$$

$$A' = AC$$

Total number of particles per unit volume = $\frac{N}{V}$

$$\frac{N}{V} = \int_0^\infty n(v) \ dv = 4\pi A' \int_0^\infty v^2 \ e^{-\frac{mv^2}{2k_B T}} \ dv$$

$$\int_0^\infty v^2 e^{-\frac{mv^2}{2k_B T}} dv = \frac{\sqrt{\pi}}{4} \left(\frac{m}{2k_B T}\right)^{-3/2}$$

$$A' = \left(\frac{N}{V}\right) \left(\frac{m}{2\pi k_B T}\right)^{3/2}$$

$$n(v) dv = 4\pi \left(\frac{N}{V}\right) \left(\frac{m}{2\pi k_B T}\right)^{3/2} v^2 e^{-\frac{mv^2}{2k_B T}} dv$$

Maxwell Boltzmann speed distribution for gas molecules in thermal equilibrium at temperature *T*.

$$n(v) dv = 4\pi \left(\frac{N}{V}\right) \left(\frac{m}{2\pi k_B T}\right)^{3/2} v^2 e^{-\frac{mv^2}{2k_B T}} dv$$

No. of gas molecules per unit volume having velocity in v and v+dv

Total number of molecules per unit volume

$$\frac{n(v)}{N/V} = 4\pi \left(\frac{m}{2\pi k_B T}\right)^{3/2} v^2 e^{-\frac{mv^2}{2k_B T}} = f_{\text{MB}}(v)$$

Maxwell-Boltzmann speed distribution

Fraction of molecules having velocity in v and v+dv

$$n(v) dv = 4\pi \left(\frac{N}{V}\right) \left(\frac{m}{2\pi k_B T}\right)^{3/2} v^2 e^{-\frac{mv^2}{2k_B T}} dv$$

Average speed

$$\bar{v} = \frac{\int_0^\infty v \ n(v) \ dv}{\int_0^\infty n(v) \ dv} = \frac{\int_0^\infty v \ n(v) \ dv}{N/V} = 4\pi \left(\frac{m}{2\pi k_B T}\right)^{3/2} \int_0^\infty v^3 \ e^{-\frac{mv^2}{2k_B T}} dv$$

$$\int_0^\infty v^3 e^{-\frac{mv^2}{2k_B T}} dv = \frac{1}{2} \left(\frac{m}{2k_B T} \right)^{-2}$$

$$\bar{v} = \langle v \rangle = \sqrt{\frac{8k_BT}{\pi m}}$$

$$n(v) dv = 4\pi \left(\frac{N}{V}\right) \left(\frac{m}{2\pi k_B T}\right)^{3/2} v^2 e^{-\frac{mv^2}{2k_B T}} dv$$

Mean square speed

$$\bar{v}^2 = \langle v^2 \rangle = \frac{3k_B T}{m}$$

$$\int_0^\infty v^4 e^{-\frac{mv^2}{2k_B T}} dv = \frac{3\sqrt{\pi}}{8} \left(\frac{m}{2k_B T}\right)^{-5/2}$$

Root mean square speed

$$v_{\rm rms} = \sqrt{\langle v^2 \rangle} = \sqrt{\frac{3k_B T}{m}}$$

Most probable speed

We know

$$f_{\text{MB}}(v) = 4\pi \left(\frac{m}{2\pi k_B T}\right)^{3/2} v^2 e^{-\frac{mv^2}{2k_B T}}$$

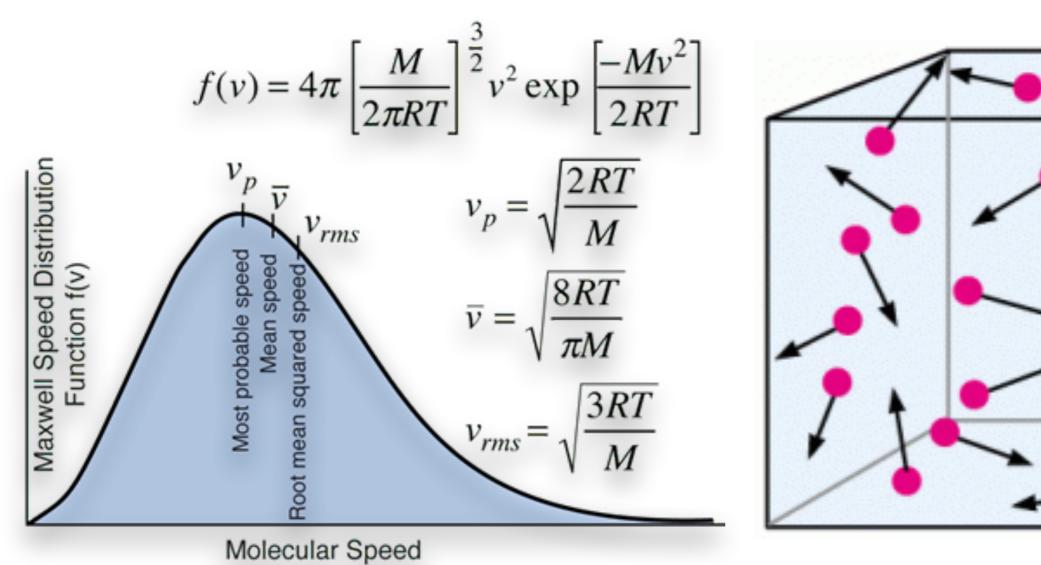
Let us differentiate this wrt v

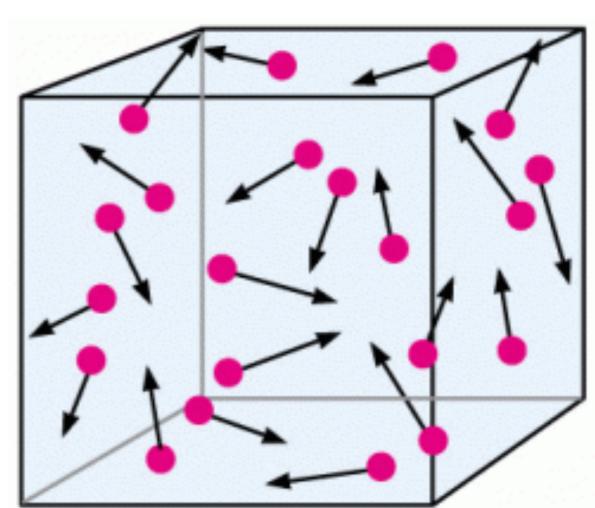
$$\frac{f_{\text{MB}}(v)}{dv} = 4\pi \left(\frac{m}{2\pi k_B T}\right)^{3/2} \left(2v - 2v^3 \frac{m}{2k_B T}\right) e^{-\frac{mv^2}{2k_B T}} = 0$$

$$v^2 = \frac{2k_BT}{m} \qquad v_{\rm mp} = \sqrt{\frac{2k_BT}{m}}$$

$$v_{\rm rms} > \bar{v} > v_{\rm mp} \Longrightarrow \sqrt{3} > \sqrt{8/\pi} > \sqrt{2}$$

From Kinetic Theory of Gases





Average Kinetic Energy

Kinetic Energy =
$$\frac{1}{2}mv^2$$

Average Kinetic Energy =
$$\langle K \rangle = \frac{1}{2} m \langle v^2 \rangle$$

$$\langle K \rangle = \frac{3}{2} k_B T$$

Mean square speed

$$\bar{v}^2 = \langle v^2 \rangle = \frac{3k_B T}{m}$$

This is the statement of equipartition theorem

A classical molecule in thermal equilibrium at temperature T has an average energy of k_BT/2 for each degree of freedom.

Show that standard deviation of the molecular speeds is given by

$$\sigma_v = \sqrt{3 - \frac{8}{\pi}} \cdot \sqrt{\frac{k_{\rm B}T}{m}}$$

Recommended Readings

Statistical Physics, Chapter 10

