

PH-107

Quantum Physics and Applications

Elements of Statistical Physics-I

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Statistical Physics

What is it and who needs it?

We now know how to deal with the dynamics of a **single particle**, both in the classical and the quantum world.

We use either Newton's laws (classical) or the TDSE/TISE (quantum) to understand the time-evolution of the state of the particle.

But, how good is this knowledge in a macroscopic world consisting of many particles (interacting simultaneously with each other)?

Statistical Physics

What is it and who needs it?

As a simple example, description of the motion of air molecules in a closed room needs tracking at least 10^{23} particles.

Attempt to describe the motion of each molecule will clearly exhaust the computational power of any conceivable machine, not to mention that of the investigator!

Therefore, we need a means to assess the **collective behavior** of a large system, without having to track the evolution of each particle in the system.

Statistical Physics

What is it and who needs it?

The purpose of **statistical physics** is exactly this:

To explain the macroscopic (or collective) properties (for example, energy, momentum etc.) or thermodynamic parameters (pressure, volume, temperature, magnetization etc.) of a large collection of particles.

Specifically, we are interested in determining the number of particles $dN(E)$ in the system with energy between E and $E+dE$.

Henceforth, the term we will use for the large collection of particles is **system**.

Statistical Physics

Number of particles between E and $E+dE$

Obviously, $dN(E)$ depends on the total number of particles in the system, N .

A little less obviously, $dN(E)$ depends on two other factors:

1. The number of available energy states within the interval E and $E+dE$.

2. The probability that a particle can actually occupy a particular state.

Statistical Physics

Number of particles between E and $E+dE$

We may write the number of available energy states within the interval E and $E+dE$ in terms of the **density of states**, $g(E)$.

And we write the probability of a particle occupying an available state in the interval E and $E+dE$ in terms of the **probability distribution function**, $f(E)$.

The product $g(E)f(E)dE$ thus gives $dN(E)$.

Statistical Physics

Closer look at $g(E)$: Example of many particles in a 3D infinite potential box

Assume that we have a **large number of electrons** (\sim Avogadro's Number, N_A), in a macroscopic 3D potential box ($L = 0.1$ m).

The available energy states (assuming no interaction between the electrons) are of the form:

$$E = \frac{\pi^2 \hbar^2}{2mL^2} (n_x^2 + n_y^2 + n_z^2) = 3.76 \times 10^{-17} (n_x^2 + n_y^2 + n_z^2) \text{ eV}$$

Thus, the spacing between the energy states $\Delta E \sim 10^{-17}$ eV.

Statistical Physics

Closer look at $g(E)$: Example of many particles in a 3D infinite potential box

The available energy levels are thus very closely spaced, defining a **quasi-continuum** of states

This explains the use of a density-of-states (DOS) function, **$g(E)$** .

But still, the number of states within a certain energy range is finite (though quite large).

Statistical Physics

Closer look at $f(E)$:

The need for $f(E)$, or a probabilistic description of an ensemble arises because

(a) We have a (classical) uncertainty due to the fact that **we cannot determine** the state of every particle at all times.

(b) Additionally (for quantum mechanical particles), we have an inherent uncertainty of **each individual particle's** quantum state.

Statistical Physics

$f(E)$: Dependence on Particle Characteristics

Accordingly, $f(E)$ depends on the **identity of the particles** in the system

i.e., whether the particles are

Classical Particles: Distinguishable, which need not obey **Pauli's Exclusion Principle**.

Bosons (Quantum Particles): Indistinguishable, which do not obey Pauli's Exclusion Principle.

Fermions (Quantum Particles): Indistinguishable, which obey Pauli's Exclusion Principle.

We know that quantum particles, Bosons and Fermions, obey integer and half-integer spin statistics.

Are there particles, which obey fractional quantum statistics? What are they called?

Statistical Physics

Pauli's Exclusion Principle

Based on Empirical Evidence, i.e., The Periodic table of Elements:

It is impossible for two **electrons** in an atom or molecule to have the same values of the four quantum numbers: n as the principal quantum number, ℓ as the angular-momentum quantum number, m_ℓ as the magnetic quantum number, and m_s as the spin quantum number.

Statistical Physics

Pauli's Exclusion Principle

A More General Statement:

Two identical Fermions (particles with **half-integer spin**, of which electron is an example) cannot occupy the same quantum state simultaneously.

The principle does not apply to particles with **integer spin** (Bosons) so that any number of Bosons can occupy the same state.

Pauli's Exclusion Principle

A more rigorous statement:

The total wave function for two identical fermions is **antisymmetric** with respect to exchange of the particles.

This means that the wave function changes its sign if the space *and* spin co-ordinates of any two particles are interchanged.

If we consider a two-Fermion wave function (including spin) $\Psi(r_1\sigma_1, r_2\sigma_2)$ then

$$\Psi(r_1\sigma_1, r_2\sigma_2) = -\Psi(r_2\sigma_2, r_1\sigma_1)$$

Pauli's Exclusion Principle

A more rigorous statement:

The total wave function for two identical Bosons is **symmetric** with respect to exchange of the particles.

i.e., if we consider a two-Boson wave function (including spin) $\Psi(r_1\sigma_1, r_2\sigma_2)$ then

$$\Psi(r_1\sigma_1, r_2\sigma_2) = +\Psi(r_2\sigma_2, r_1\sigma_1)$$

We thus see, that the number of Quantum particles that can be placed in any state depends on the **spin degree of freedom** of the particles.

Statistical Physics

The calculation of $f(E)$

The purpose of the rest of the lecture(s) on this topic is the calculation of $f(E)$ for classical particles (Maxwell-Boltzmann), Bosons (Bose-Einstein) and Fermions (Fermi-Dirac).

How do we calculate $f(E)$?

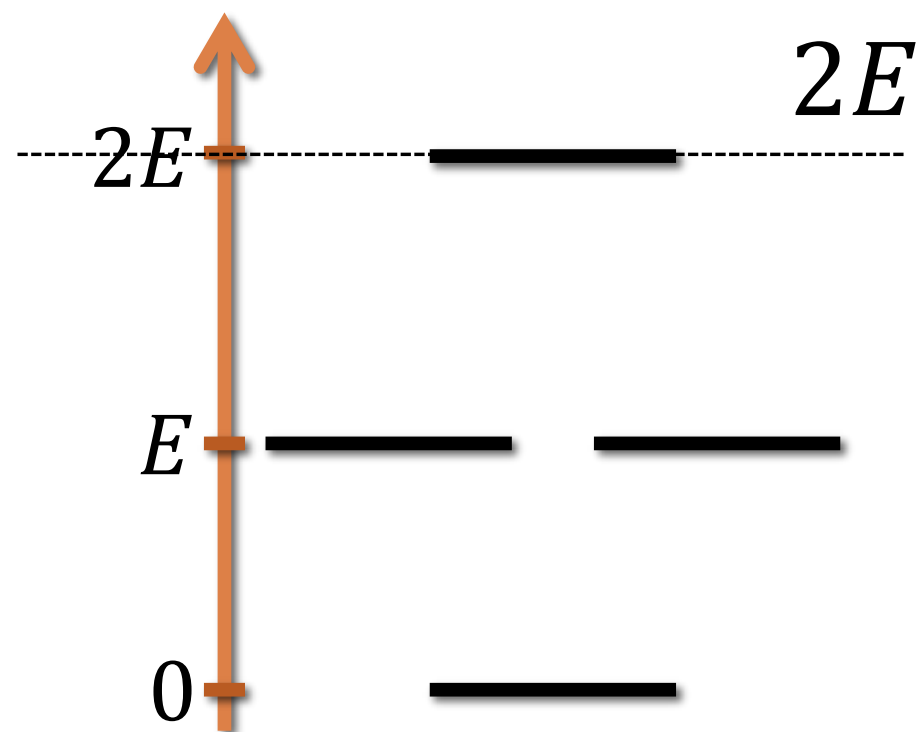
Let us analyze the question further: For a given fixed energy E of the system, how do I know what is $dN(E)$, or (to start with) $f(E)$?

Let's start answering this question with a simple example

Statistical Physics

2 particles occupying 3 states

Let's say that we have a system of **two particles** with total energy $2E$. The energy states available to the system are 0, E , and $2E$. The degeneracies of the states are 1, 2, and 1, respectively.

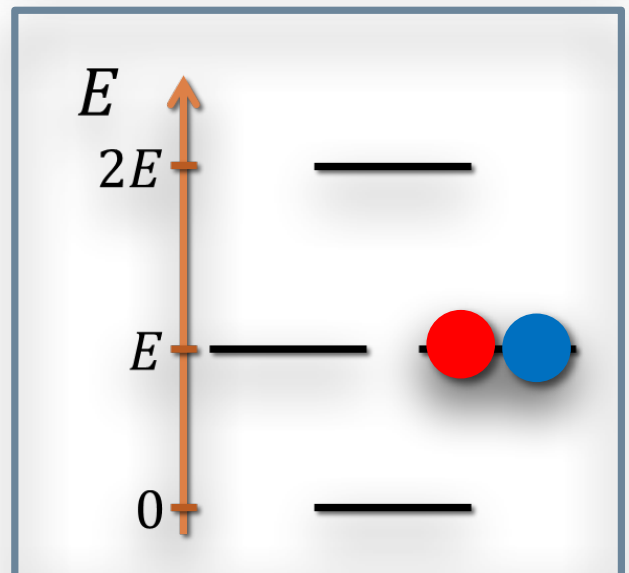
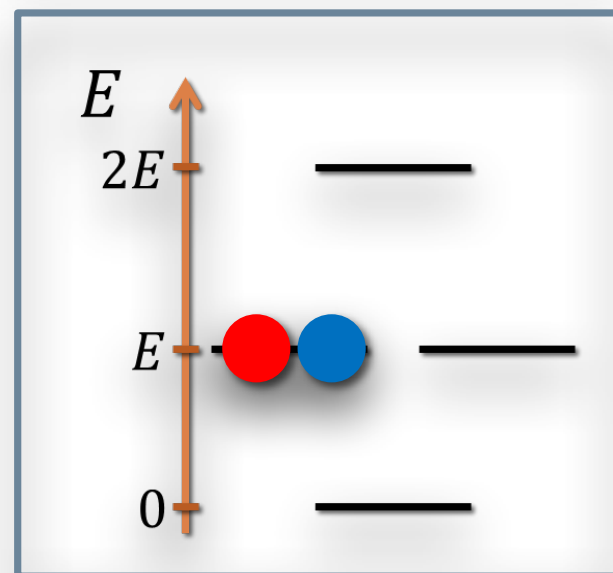
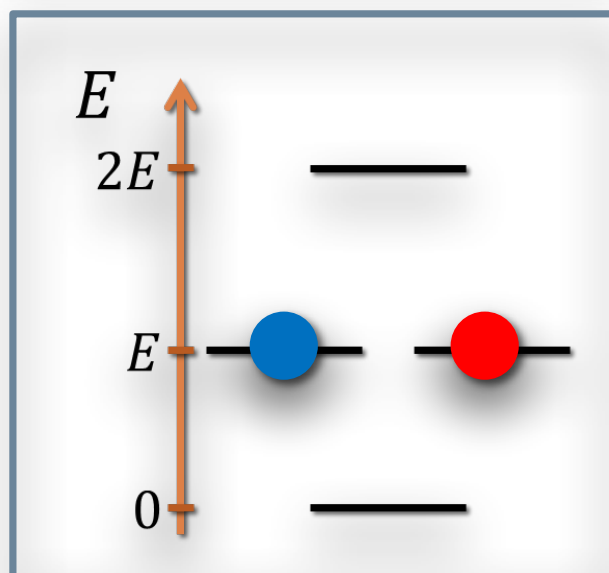
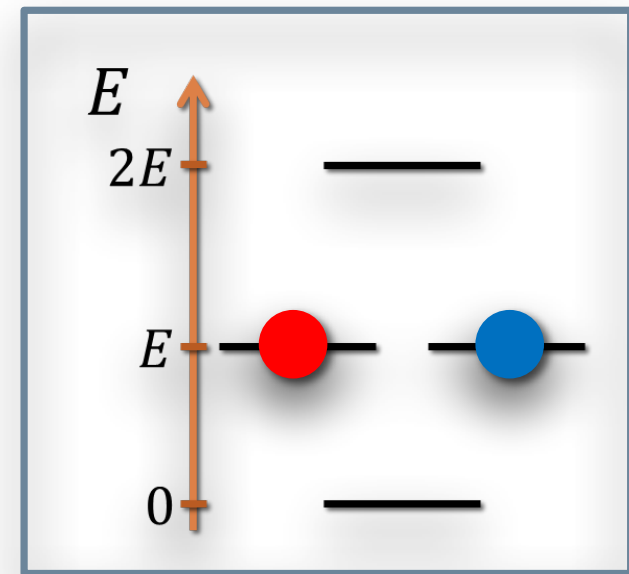
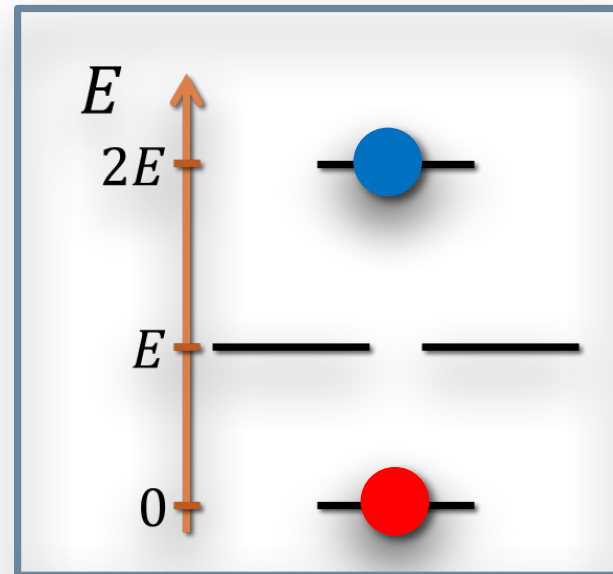
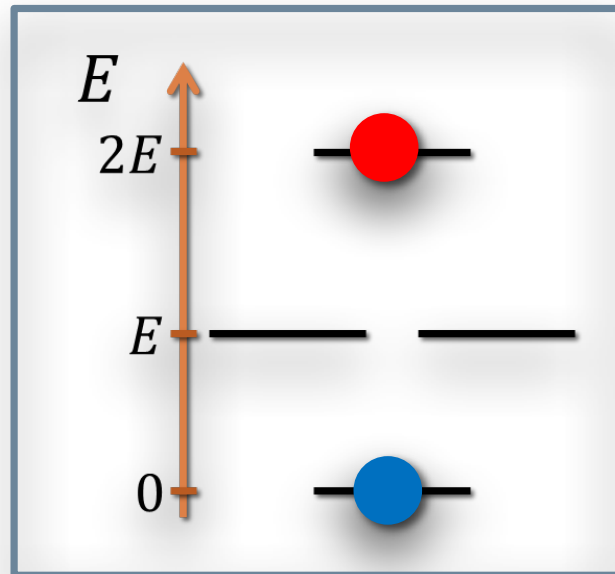


State Index (i)	State Energy	State degeneracy (g_i)
1	$2E$	1
2	E	2
3	0	1

Statistical Physics

Different arrangements giving energy E

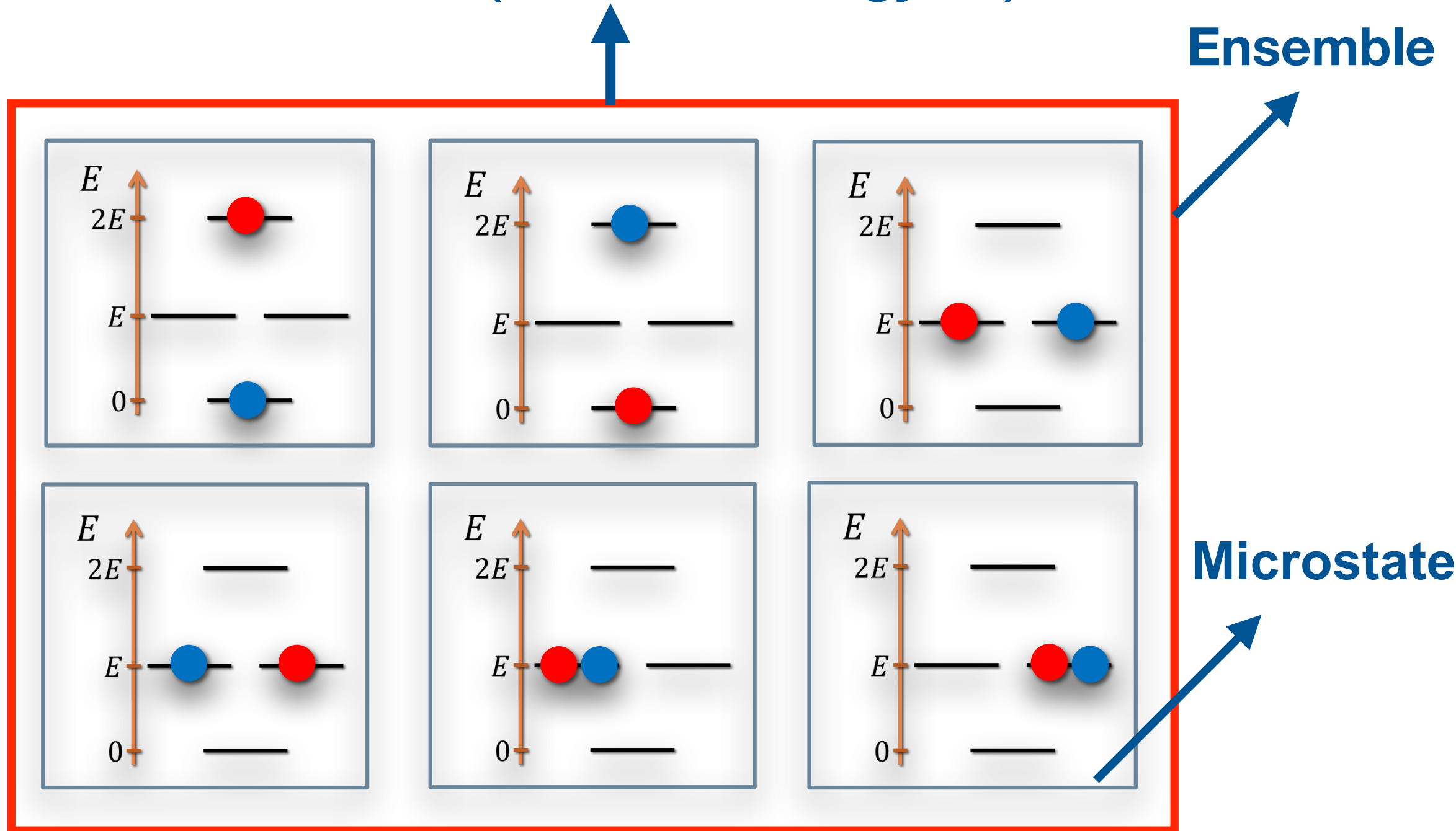
What are the different ways in which we get the total energy to be $2E$?



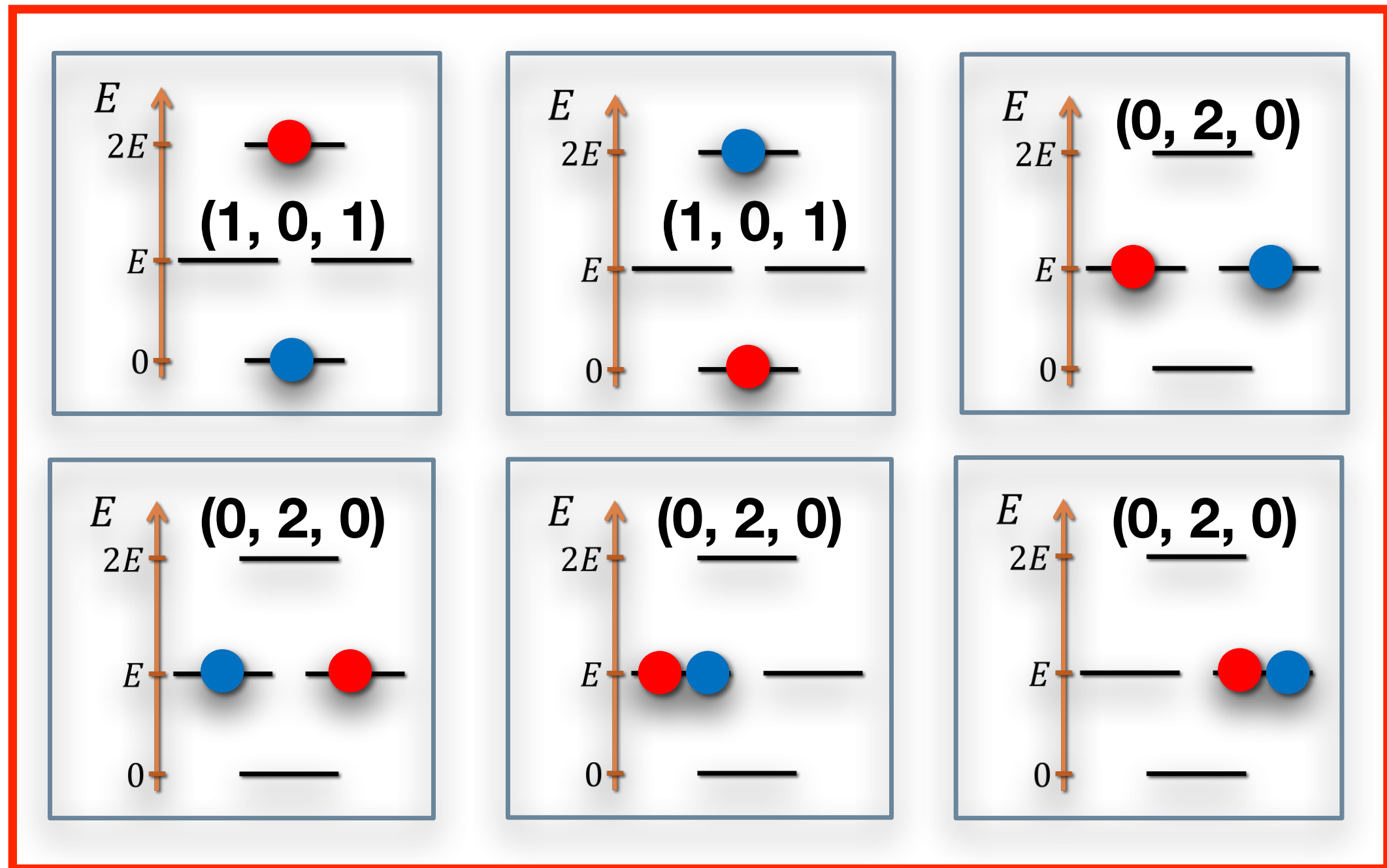
Statistical Physics

Microstates, ensemble, and the macrostate

Macrostate (with total energy $2E$)

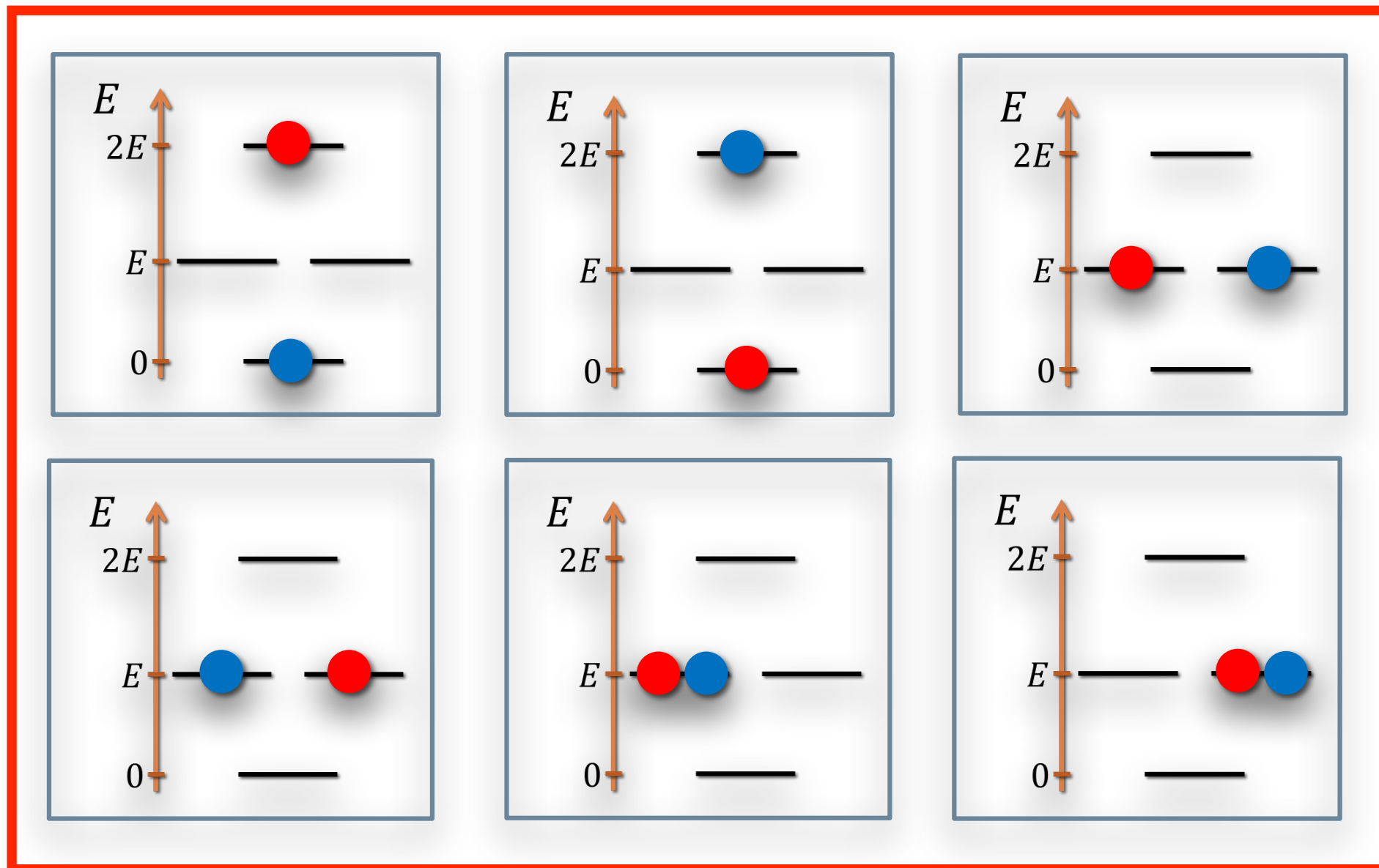


Configuration



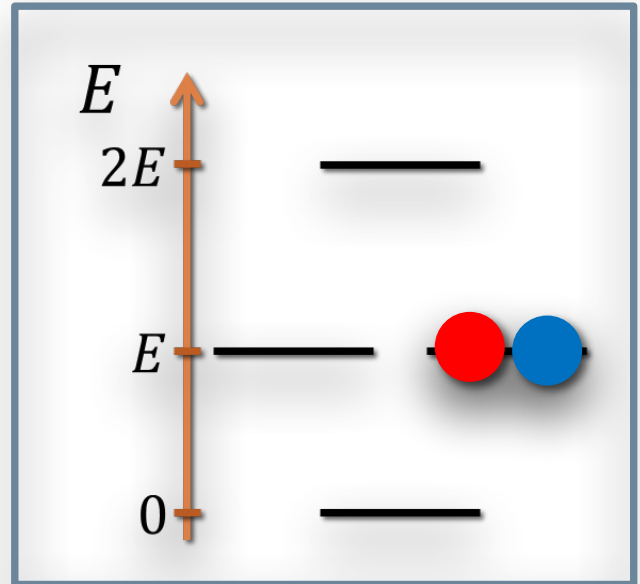
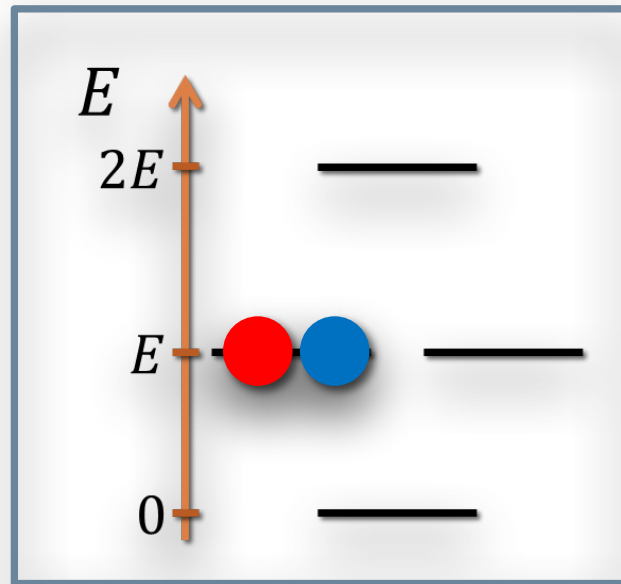
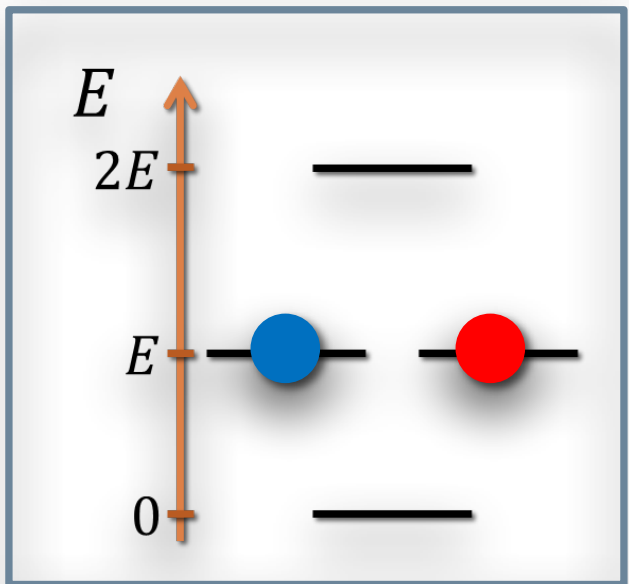
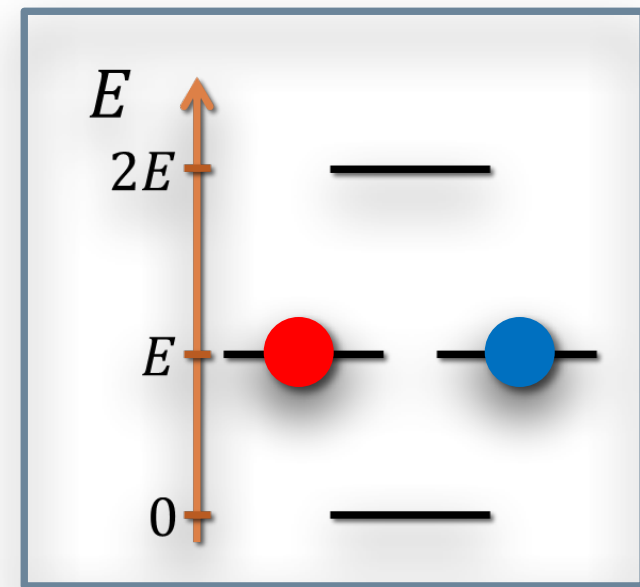
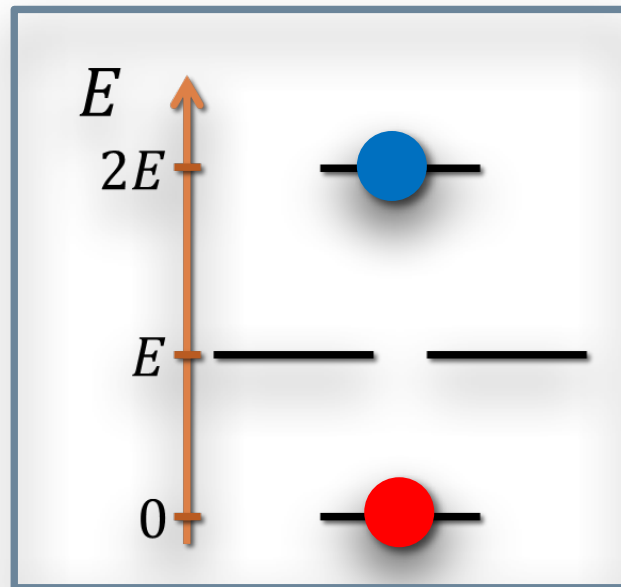
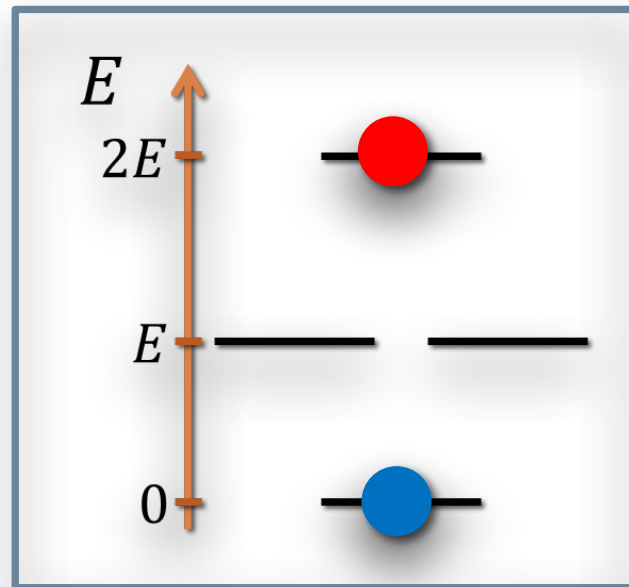
The *apriori* probability of occupancy of all **microstates** is equal.
So, the $(0,2,0)$ **configuration** is twice as probable as the $(1,0,1)$.

Classical Particles

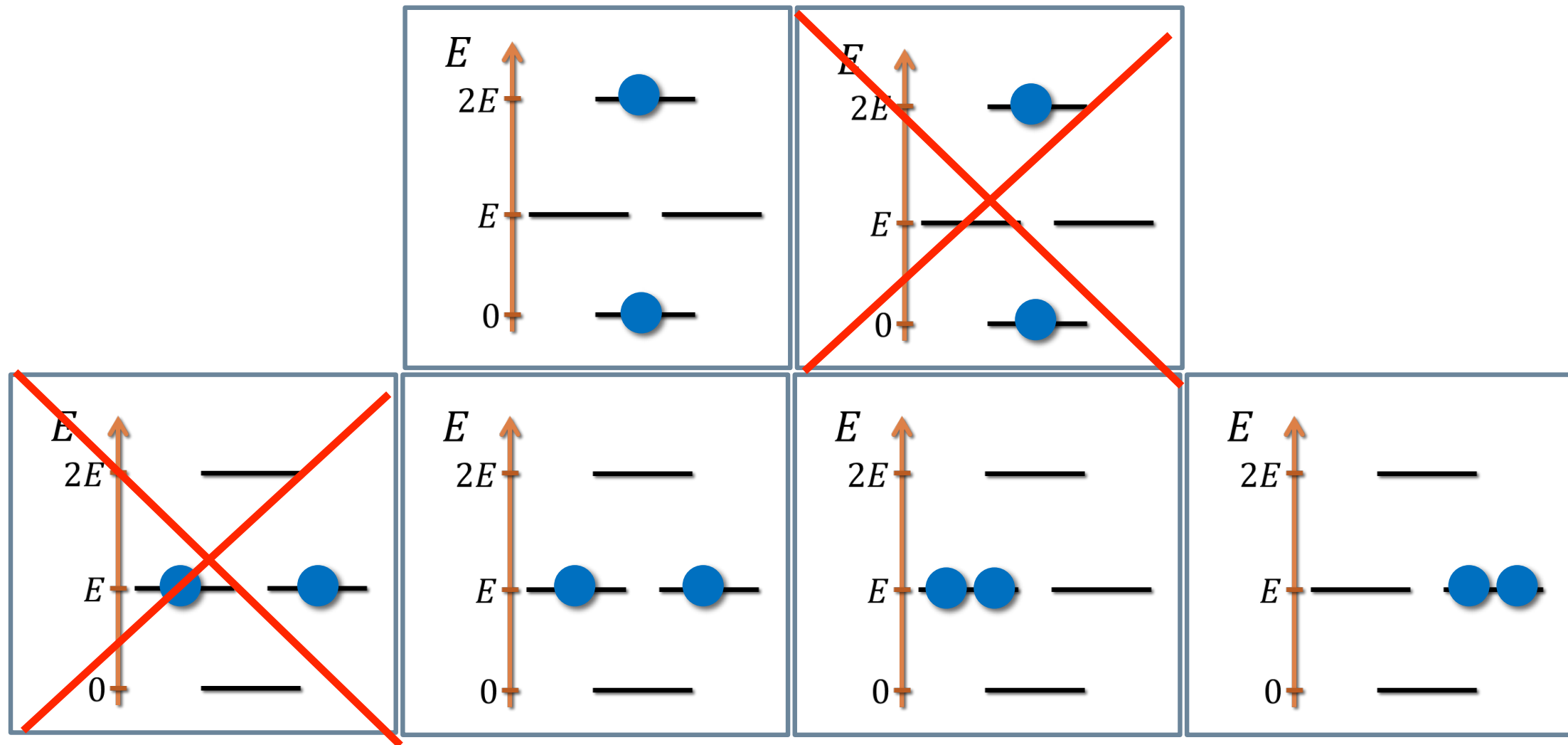


In this example, implicitly we assumed classical particles, which are distinguishable (red and blue) and which do not obey the Pauli's exclusion principle (microstate 5 and 6).

Bosons

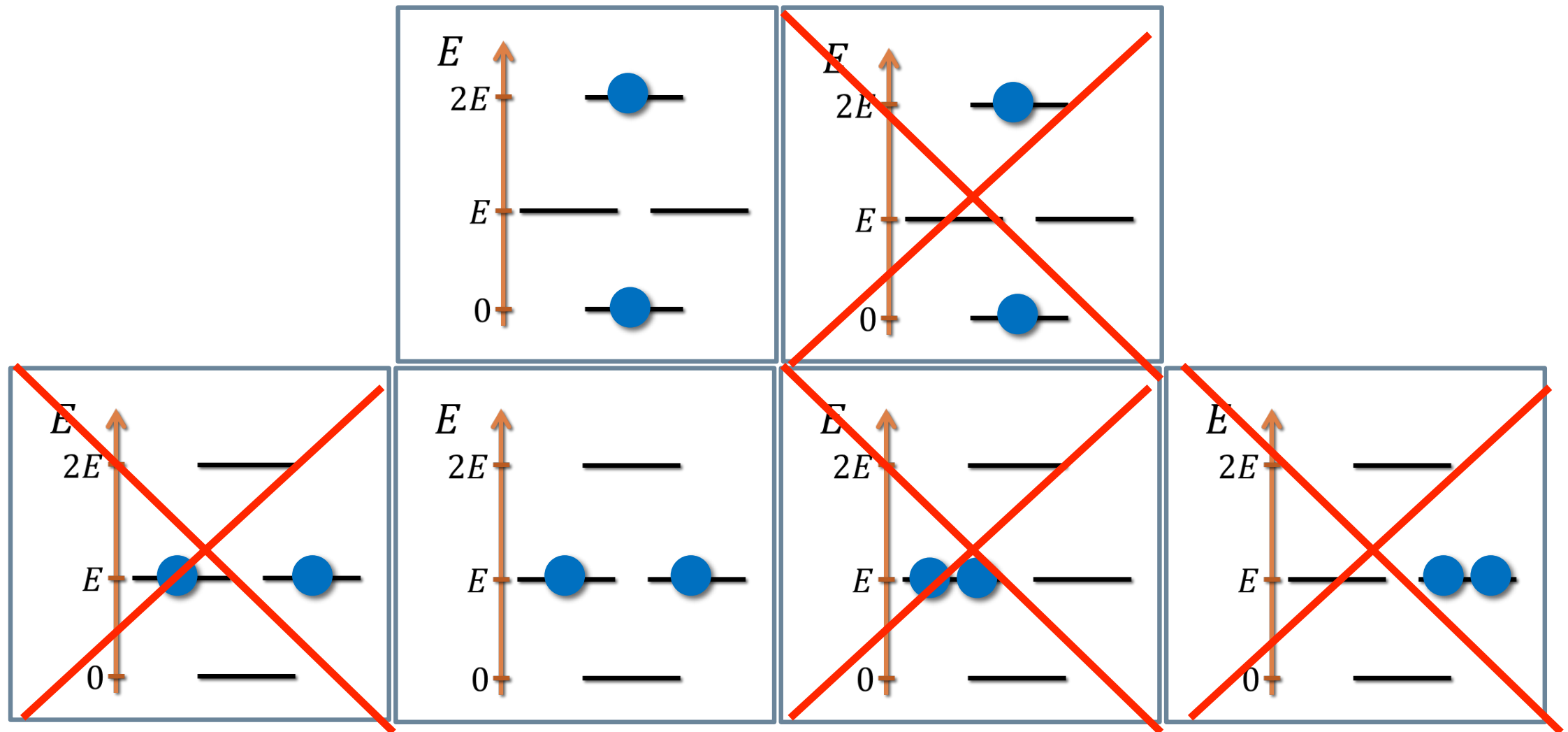


Bosons



The $(0,2,0)$ configuration is thrice as probable as the $(1,0,1)$.

Fermions



The $(0,2,0)$ configuration is equally probable as the $(1,0,1)$.

Statistical Physics

General problem (M-B, B-E, F-D)

Given a system with set of energy states E_i (0, E, 2E), each with degeneracy g_i (1, 2, 1)

with a fixed amount of energy, E (2E), and a fixed number of particles, N (2),

what is the most probable set of $N_i \equiv N(E_i)$ or most probable configuration?

In case of the classical particles, we found that the configuration $N_1 \equiv N(E_1) = N(2E) = 0$, $N_2 \equiv N(E_2) = N(E) = 2$ and $N_3 \equiv N(E_3) = N(0) = 0$ or (0,2,0) is the “more” probable one.

Statistical Physics

General problem (M-B, B-E, F-D)

$$\sum_{i=1}^{\infty} N_i = N$$

$$\sum_{i=1}^{\infty} E_i N_i = E$$

State Index (i)	State Energy (E_i)	State Degeneracy (g_i)	State Occupancy (N_i)
1	E_1	g_1	N_1
2	E_2	g_2	N_2
\vdots	\vdots	\vdots	\vdots
i	E_i	g_i	N_i
\vdots	\vdots	\vdots	\vdots

Here we are generalizing the problem for a larger system. We are looking for $\{N_i\} \equiv (N_1, N_2, \dots N_i \dots)$

Equilibrium Configuration

The configuration for which the number or **multiplicity** of microstates is maximum corresponds to the most probable configuration.

This configuration is called the **equilibrium configuration**.

If the total energy E of the system is provided to the system as heat, so that it reaches an equilibrium temperature T , then this configuration is the most probable one at **thermal equilibrium** at temperature T .

Equilibrium Configuration

So the general problem we are trying to solve is, given a system with

$$\sum_{i=1}^{\infty} N_i = N$$

$$\sum_{i=1}^{\infty} E_i N_i = E$$

What is the configuration $\{N_i\} \equiv (N_1, N_2, \dots N_i \dots)$ for which the multiplicity $Q(\{N_i\})$ is maximum?

For this, we have to first learn how to calculate $Q(\{N_i\})$.

Recommended Readings

Statistical Physics, Chapter 10

