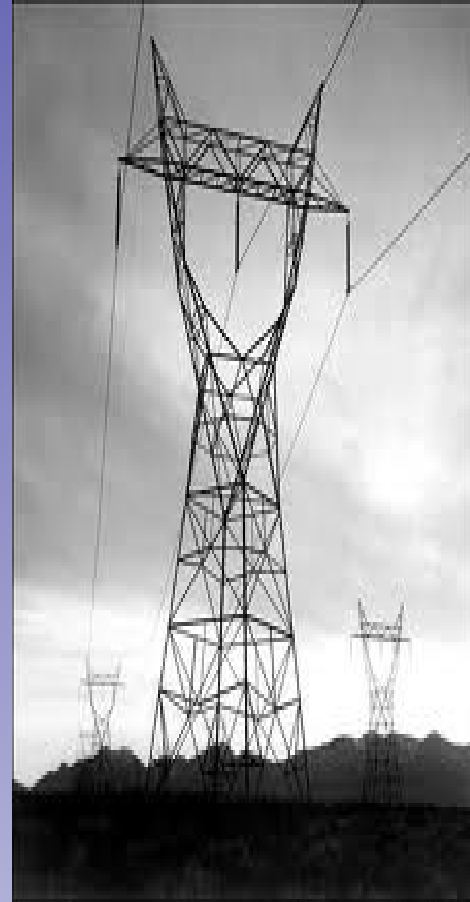


Conductors



Conductors

Outline

- I. Conductors and their electrical properties
- II. Examples using (i) spherical shell and (ii) a cavity in a spherical shell
- III. Force on a conductor with example
- IV. Applications

Conductors

Learning Objectives

- I. To learn about conductors.
- II. To learn to show the electrical properties of conductors.
- III. To learn about the force on a charged conductor.

Conductors

Learning Outcomes

- I. To be able to analytically show the electrical properties of conductors.
- II. To be able to calculate the electric field and potential in the presence of conductors
- III. To be able to calculate force on a charged conductor.

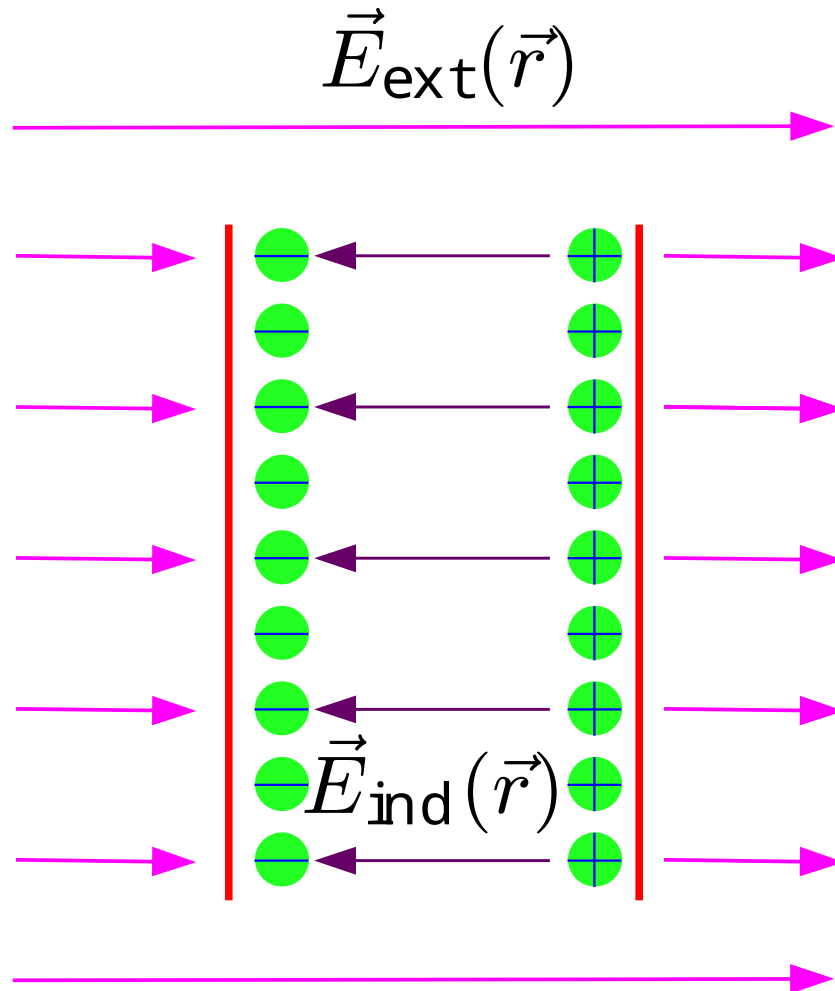
Conductors

The kinds of material based on their conducting properties are:

Material	Resistivity ($\Omega\text{-m}$)
Silver	1.6×10^{-8}
Copper	1.7×10^{-8}
Gold	2.4×10^{-8}
Iron	1.0×10^{-7}
Sea water	0.2
Polyethylene	2.0×10^{11}
Glass	$\sim 10^{12}$
Fused quartz	7.5×10^{17}

Conductors

Consider a rectangular block of conductor in the presence of an external electric field as shown in the figure.



Typical free electron response time $\sim 10^{-15}$ s to 10^{-12} s.

Conductors

Some of the properties of a conductor are:

1. $\vec{E}_{\text{net}}(\vec{r}) \equiv 0$ inside a conductor.
2. The volume free charge density, $\rho_{\text{inside}}^{\text{free}}(\vec{r}) = 0$ inside a conductor.
3. Any induced charges on a conductor can only reside on surface or surfaces of the conductor – as surface charge distribution, σ_{free} .
4. The entire volume and the surface of a conductor is an equipotential.
5. Just outside of the surface of a conductor, $\vec{E}_{\text{outside}}(\vec{r})$ is perpendicular to the surface.

Conductors

1. $\vec{E}_{\text{net}}(\vec{r}) \equiv 0$ inside a conductor:

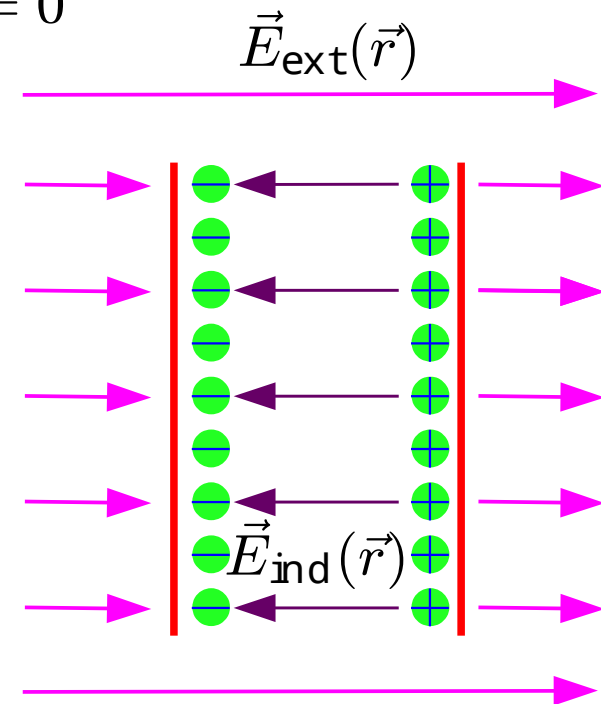
Consider a rectangular block of conductor in the presence of an external electric field $\vec{E}_{\text{ext}}(\vec{r})$.

In response to the external electric field, the free electrons redistribute themselves such that the net electric field inside the conductor is zero,

$$\vec{E}_{\text{net inside}}(\vec{r}) = \vec{E}_{\text{ext}}(\vec{r}) + \vec{E}_{\text{induced inside}}(\vec{r}) = 0$$

$$\Rightarrow \vec{E}_{\text{induced inside}}(\vec{r}) = -\vec{E}_{\text{ext}}(\vec{r})$$

So that $\vec{E}_{\text{net}}(\vec{r}) \equiv 0$ inside a conductor.



Conductors

2. The volume free charge density, $\rho_{\text{inside}}^{\text{free}}(\vec{r}) = 0$ inside a conductor.

Since $E_{\text{inside}}(\vec{r})$ is zero inside the conductor, we have from Gauss' law,

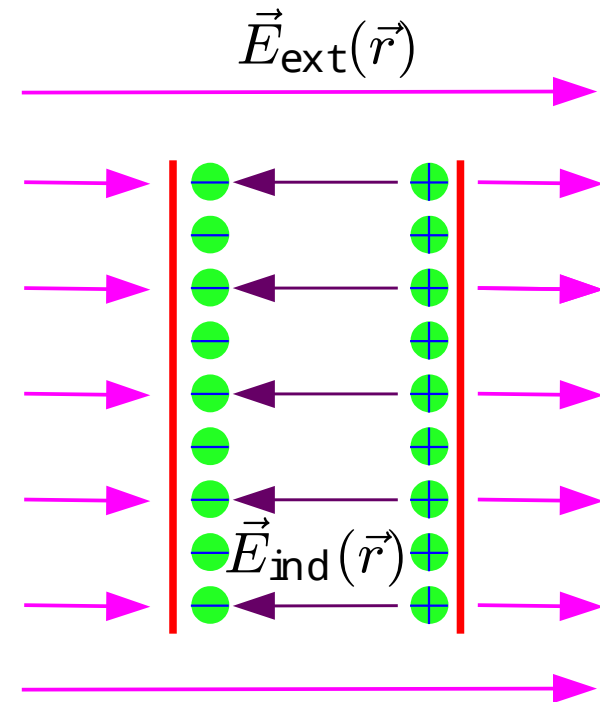
$$\vec{\nabla} \cdot \vec{E}_{\text{inside}}(\vec{r}) = \frac{\rho_{\text{inside}}^{\text{free}}(\vec{r})}{\epsilon_0}$$

Since,

$$\vec{E}_{\text{net inside}}(\vec{r}) = \vec{E}_{\text{ext}}(\vec{r}) + \vec{E}_{\text{induced inside}}(\vec{r}) = 0$$

So that

$$\begin{aligned}\vec{\nabla} \cdot \vec{E}_{\text{inside}}(\vec{r}) &= 0 \\ \Rightarrow \rho_{\text{inside}}^{\text{free}}(\vec{r}) &= 0\end{aligned}$$



The volume free charge density inside a conductor is zero.

Conductors

3. Any induced charges on a conductor can only reside on surface or surfaces of the conductor – as surface charge distribution, σ_{free} .

We have said that in the presence of an external electric field,

$$\rho_{\text{inside}}^{\text{free}}(\vec{r}) = 0$$

Therefore, any induced free charge must reside on the surface of the conductor as $\sigma_{\text{free}}(\vec{r})$.

Conductors

4. The entire volume and the surface of a conductor is an equipotential.

Consider two points a and b on the surface of a conductor. The potential difference between the two points is given by

$$\Delta V_{ab} = V(b) - V(a) = - \int_a^b \vec{E}(\vec{r}) \cdot d\vec{l} = 0$$

$$\Rightarrow V_a = V_b$$

If the integral is not zero then the charges would move. Therefore, the entire volume and the surface of a conductor is an equipotential.

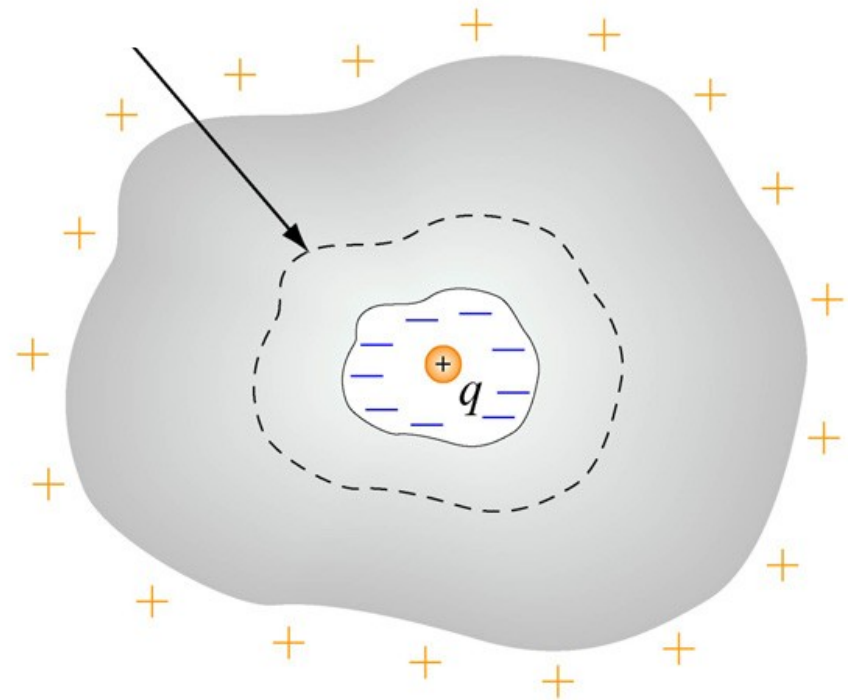
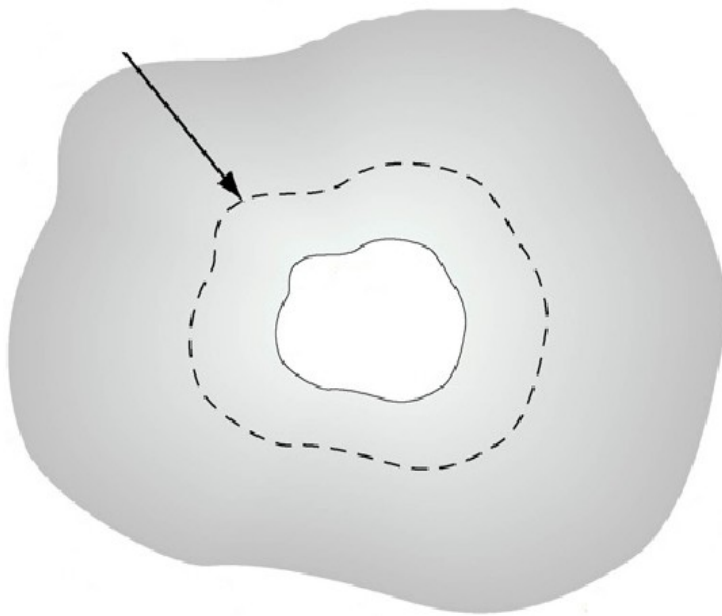
Conductors

5. Just outside of the surface of a conductor, $\vec{E}_{\text{outside}}(\vec{r})$ is perpendicular to the surface.

If $\vec{E}_{\text{tangential}}(\vec{r}) \neq 0$, then charges would move, which is inconsistent with electrostatics. Hence, just outside of the surface of a conductor, the electric field is perpendicular to the surface.

Conductors

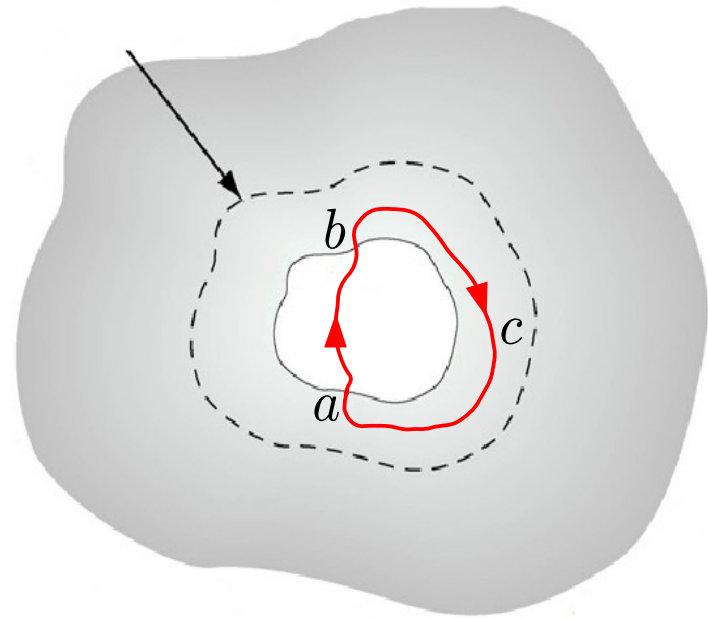
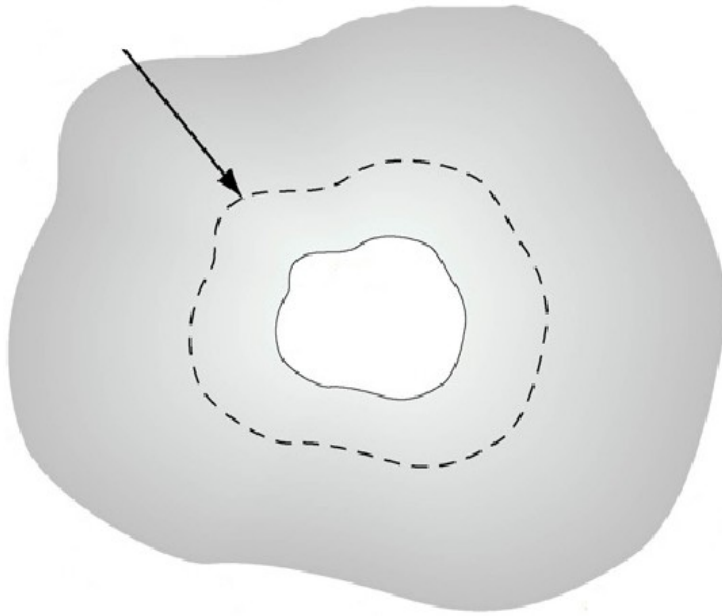
- (i) The potential $V(\vec{r})$ inside the cavity of a conductor is constant if there is no charge inside the cavity. The corresponding electric field $\vec{E}(\vec{r})$ inside the cavity is zero.
- (ii) A charge q placed inside the cavity induces a surface charge on the inner wall of the cavity and $q_{\text{ind}} = -q$.



Conductors

Consider a conductor with a cavity. Find $\vec{E}(\vec{r})$ and $V(\vec{r})$ inside the cavity.

Inside the conductor, $\vec{E}(\vec{r}) = 0$ and $V(\vec{r}) = \Phi$, a constant



In order to show that $\vec{E}(\vec{r}) = 0$ and $V(\vec{r}) = \Phi$ inside the cavity, assume $\vec{E}(\vec{r})$ to be non-zero and consider the line integral along the contour $a \rightarrow b \rightarrow c \rightarrow a$ such that the line element along $a \rightarrow b$ follows $\vec{E}(\vec{r})$.

Conductors

That is

$$\oint_C \vec{E}(\vec{r}) \cdot d\vec{l} = \int_{a \rightarrow b} \vec{E}(\vec{r}) \cdot d\vec{l} + \int_{b \rightarrow c \rightarrow a} \vec{E}(\vec{r}) \cdot d\vec{l}$$

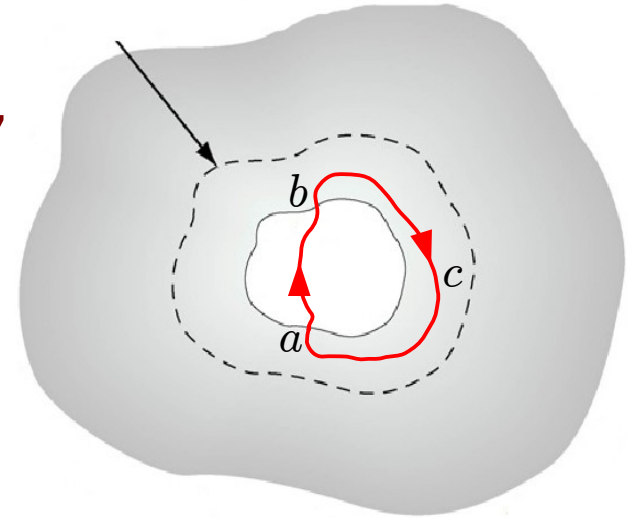
Since $\vec{E}(\vec{r})$ is zero inside the conductor,

$$\int_{b \rightarrow c \rightarrow a} \vec{E}(\vec{r}) \cdot d\vec{l} = 0$$

and as $\vec{E}(\vec{r})$ is along $d\vec{l}$, the integral

$$\int_{a \rightarrow b} \vec{E}(\vec{r}) \cdot d\vec{l} = \int_{a \rightarrow b} E dl$$

can be zero only if $\vec{E}(\vec{r})$ is zero inside the cavity. In addition, since the inner wall of the cavity is at the potential $V(\vec{r}) = \Phi$, we must have the same potential inside the cavity.



Conductors

A charge q placed inside the cavity induces a surface charge on the inner wall of the cavity and $q_{\text{ind}} = -q$.

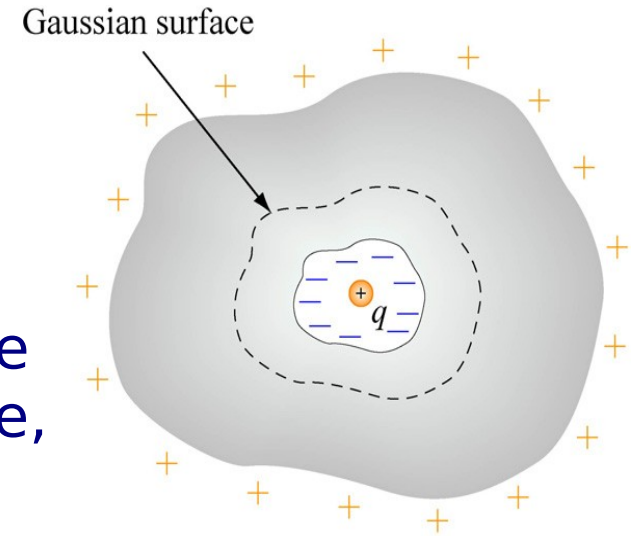
Consider a Gaussian surface inside the conductor as shown in the figure. Using Gauss' law and the fact that $\vec{E}(\vec{r}) = 0$ inside the conductor, we have

$$\oint_S \vec{E} \cdot d\vec{S} = \frac{1}{\epsilon_0} \int_V \rho dV = 0$$

So that the total charge enclosed inside the volume must be zero, and therefore,

$$q_{\text{ind}} = -q$$

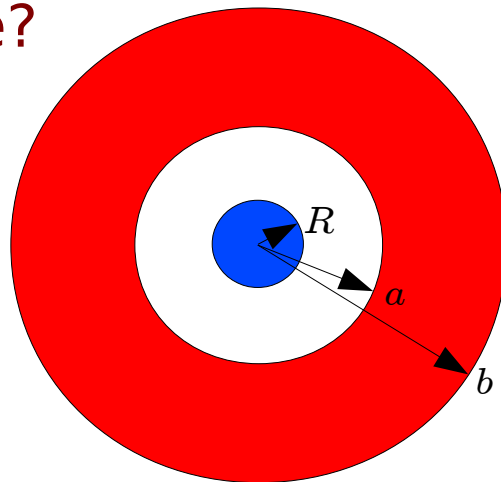
on the inner wall of the cavity. Consequently, the outer surface of the conductor gets a positive charge q .



Conductors

Example 2.35: A metal sphere of radius R , carrying charge q , is surrounded by a thick concentric metal shell (inner radius a , outer radius b). The shell carries no net charge.

- (i) Find the surface charge density σ at R , at a , and at b .
- (ii) Find the potential at the center of the sphere, using infinity as reference.
- (iii) Now the outer surface is touched to a grounding wire which lowers its potential to zero. How do your answers to (i) and (ii) change?



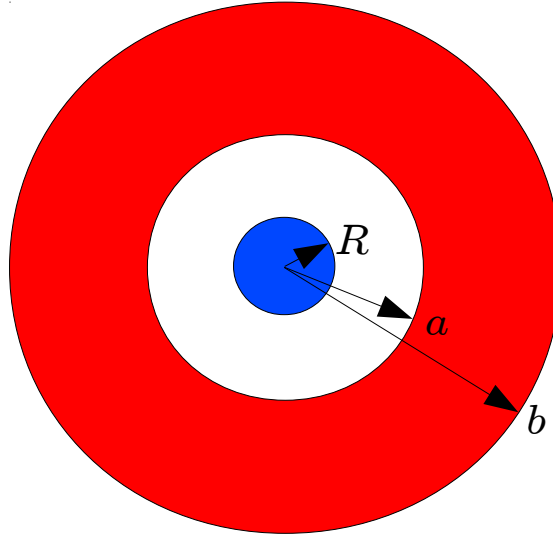
Conductors

(i) The surface charge densities at the various surfaces are:

On the surface of the sphere of radius R , $\sigma_R = \frac{q}{4\pi R^2}$

On the inner surface of the spherical shell $\sigma_a = -\frac{q}{4\pi a^2}$

On the outer surface of the spherical shell $\sigma_b = \frac{q}{4\pi b^2}$



Conductors

(ii) The potential at the center of the sphere is given by

$$V_0 = - \int_{\infty}^0 \vec{E}(\vec{r}) \cdot d\vec{l} = - \frac{q}{4\pi\epsilon_0} \int_{\infty}^b \frac{1}{r^2} dr - \frac{q}{4\pi\epsilon_0} \int_a^R \frac{1}{r^2} dr$$

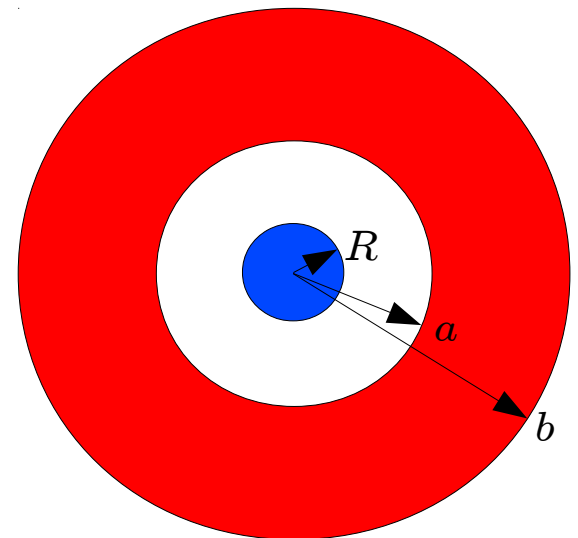
or, $V_0 = \frac{q}{4\pi\epsilon_0} \left[\frac{1}{b} - \frac{1}{a} + \frac{1}{R} \right]$

(iii) When the outer surface of the shell is grounded then the charge density on the outside shell becomes zero.

Now, the potential at the center of the sphere becomes

$$V_0 = - \int_b^0 \vec{E}(\vec{r}) \cdot d\vec{l} = - \frac{q}{4\pi\epsilon_0} \int_a^R \frac{1}{r^2} dr$$

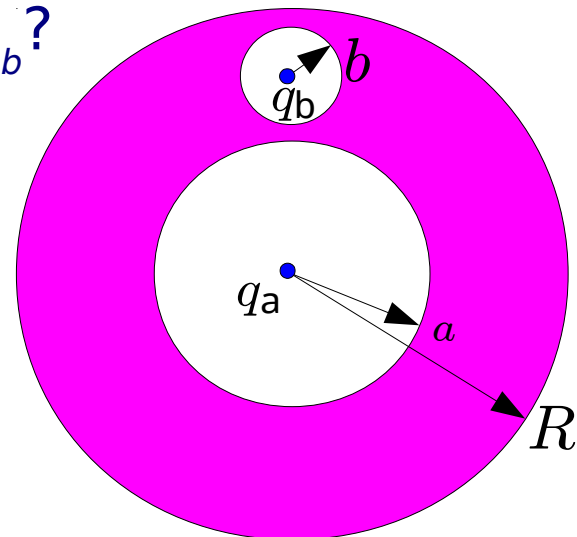
or, $V_0 = \frac{q}{4\pi\epsilon_0} \left[\frac{1}{R} - \frac{1}{a} \right]$



Conductors

Example 2.36: Two spherical cavities, of radii a and b are hollowed out from the interior of a (neutral) conducting sphere of radius R . At the centre of each cavity a point charge is placed – call these charges q_a and q_b .

- (i) Find the surface charge density σ at R , at a , and at b .
- (ii) What is the field outside the conductor?
- (iii) What is the field within each cavity?
- (iv) What is the force on q_a and q_b ?



Conductors

(i) The surface charge densities at the various surfaces are:

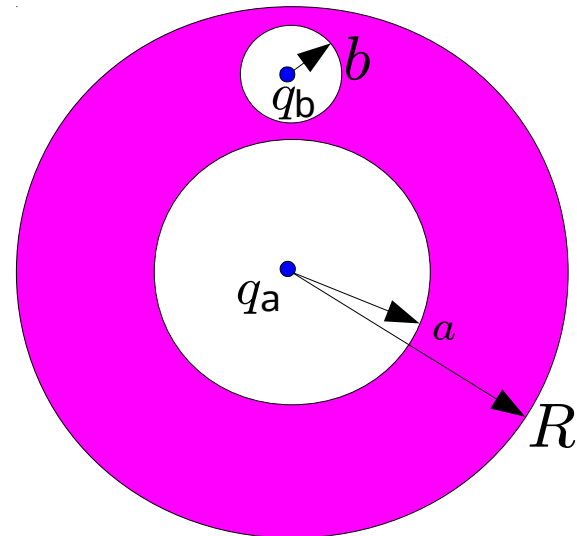
The inner surface of the cavity of radius a , $\sigma_a = -\frac{q_a}{4\pi a^2}$

The inner surface of the cavity of radius b , $\sigma_b = -\frac{q_b}{4\pi b^2}$

The surface of the sphere of radius R , $\sigma_R = \frac{q_a + q_b}{4\pi R^2}$

(ii) The electric field outside the conductor is:

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q_a + q_b}{r^2} \hat{r}, \quad r > R$$



Conductors

(iii) The electric field inside cavity a ,

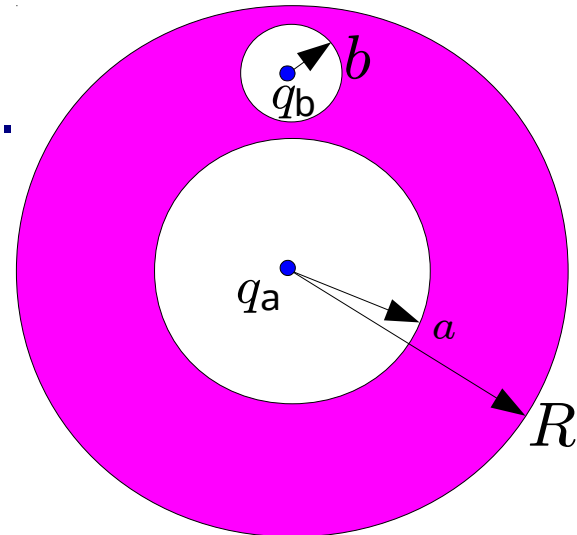
$$\vec{E}_a = \frac{1}{4\pi\epsilon_0} \frac{q_a}{r_a^2} \hat{r}_a, \quad r_a < a$$

The electric field inside cavity b ,

$$\vec{E}_b = \frac{1}{4\pi\epsilon_0} \frac{q_b}{r_b^2} \hat{r}_b, \quad r_b < b$$

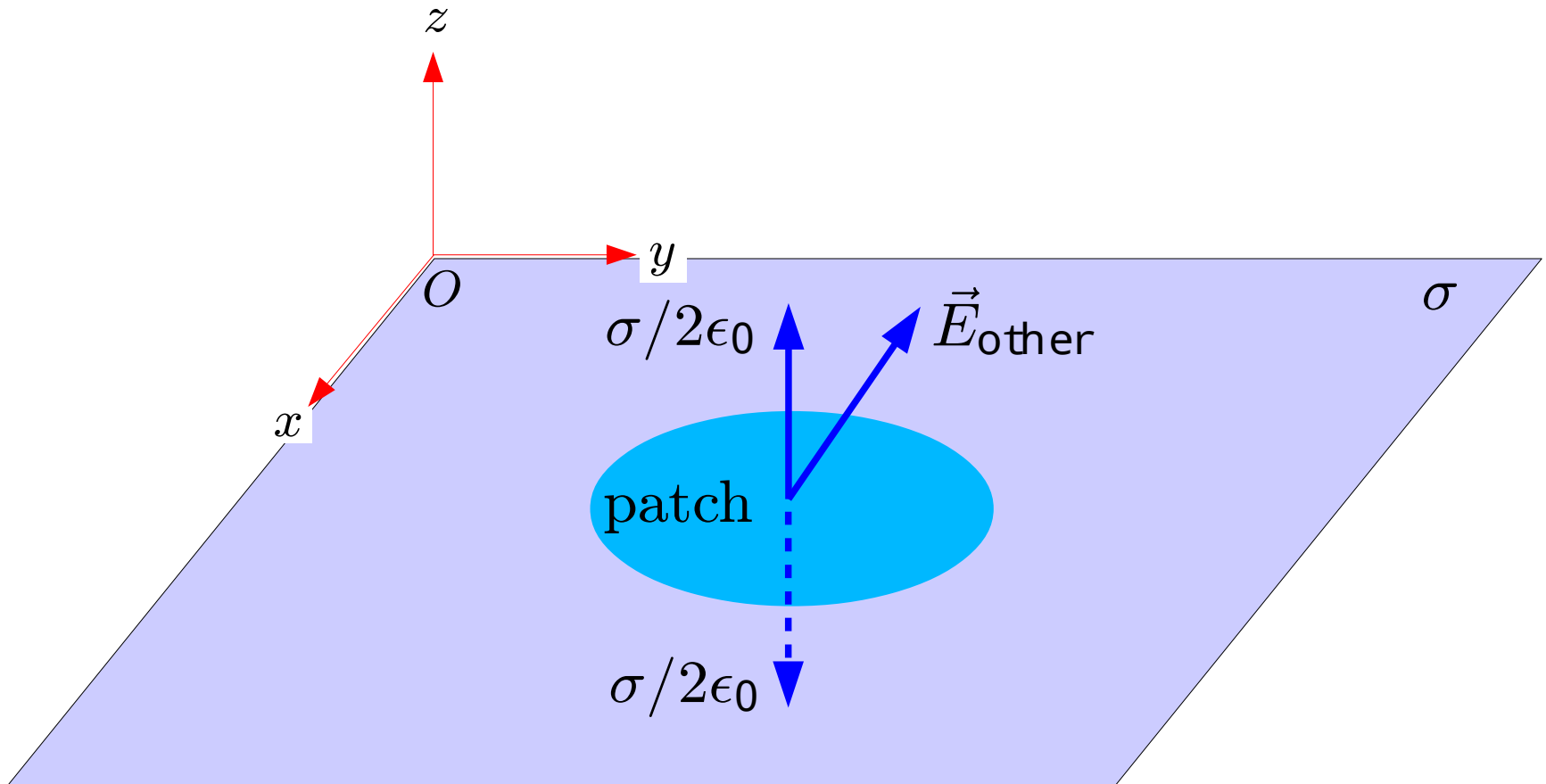
The vectors \vec{r}_a and \vec{r}_b are measured from the respective centers of the cavities.

(iv) The force on q_a and q_b is zero.



Conductors : Force on a Conductor

Consider a small patch of area A on the conducting surface with a surface charge density σ . What is the electrostatic force acting on this patch?



Conductors : Force on a Conductor

The electric fields, above and below the surface, due to the patch are

$$\vec{E}_{\text{above}}^{\text{patch}} = \frac{\sigma}{2\epsilon_0} \hat{n} \quad \text{and} \quad \vec{E}_{\text{below}}^{\text{patch}} = -\frac{\sigma}{2\epsilon_0} \hat{n}$$

Let the electric field at the patch due to the rest of the surface be \vec{E}_{other} . Then,

$$\vec{E}_{\text{above}} = \vec{E}_{\text{other}} + \vec{E}_{\text{above}}^{\text{patch}}$$

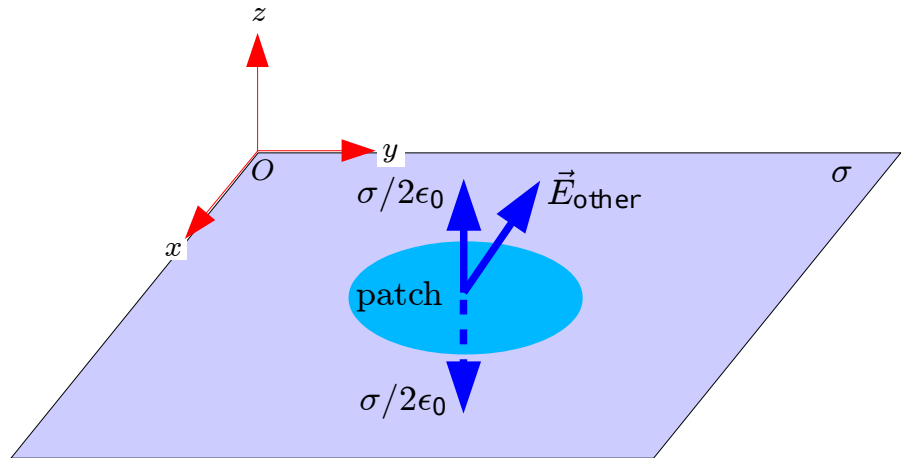
So that,

$$\vec{E}_{\text{above}} = \vec{E}_{\text{other}} + \frac{\sigma}{2\epsilon_0} \hat{n}$$

and

$$\vec{E}_{\text{below}} = \vec{E}_{\text{other}} - \frac{\sigma}{2\epsilon_0} \hat{n}$$

Hence,
$$\vec{E}_{\text{other}} = \frac{1}{2} \left(\vec{E}_{\text{above}} + \vec{E}_{\text{below}} \right)$$



That is, the electric field at the patch due to the rest of the charges is the average of \vec{E}_{above} and \vec{E}_{below} .

Conductors : Force on a Conductor

To apply the previous result, we note that for the conducting surface with charge density σ , we have

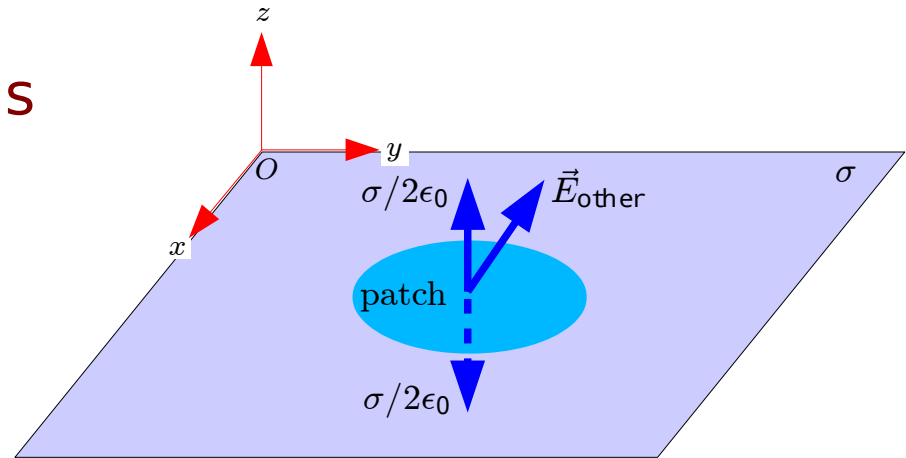
$$\vec{E}_{\text{above}} = \frac{\sigma}{\epsilon_0} \hat{n}, \quad \vec{E}_{\text{below}} = 0$$

Therefore, the average field is

$$\vec{E}_{\text{other}} = \frac{1}{2} \left(\frac{\sigma}{\epsilon_0} \hat{n} + \vec{0} \right)$$

The force per unit area \vec{f} is

$$\vec{f} = \sigma \vec{E}_{\text{other}} = \frac{\sigma^2}{2\epsilon_0} \hat{n}$$

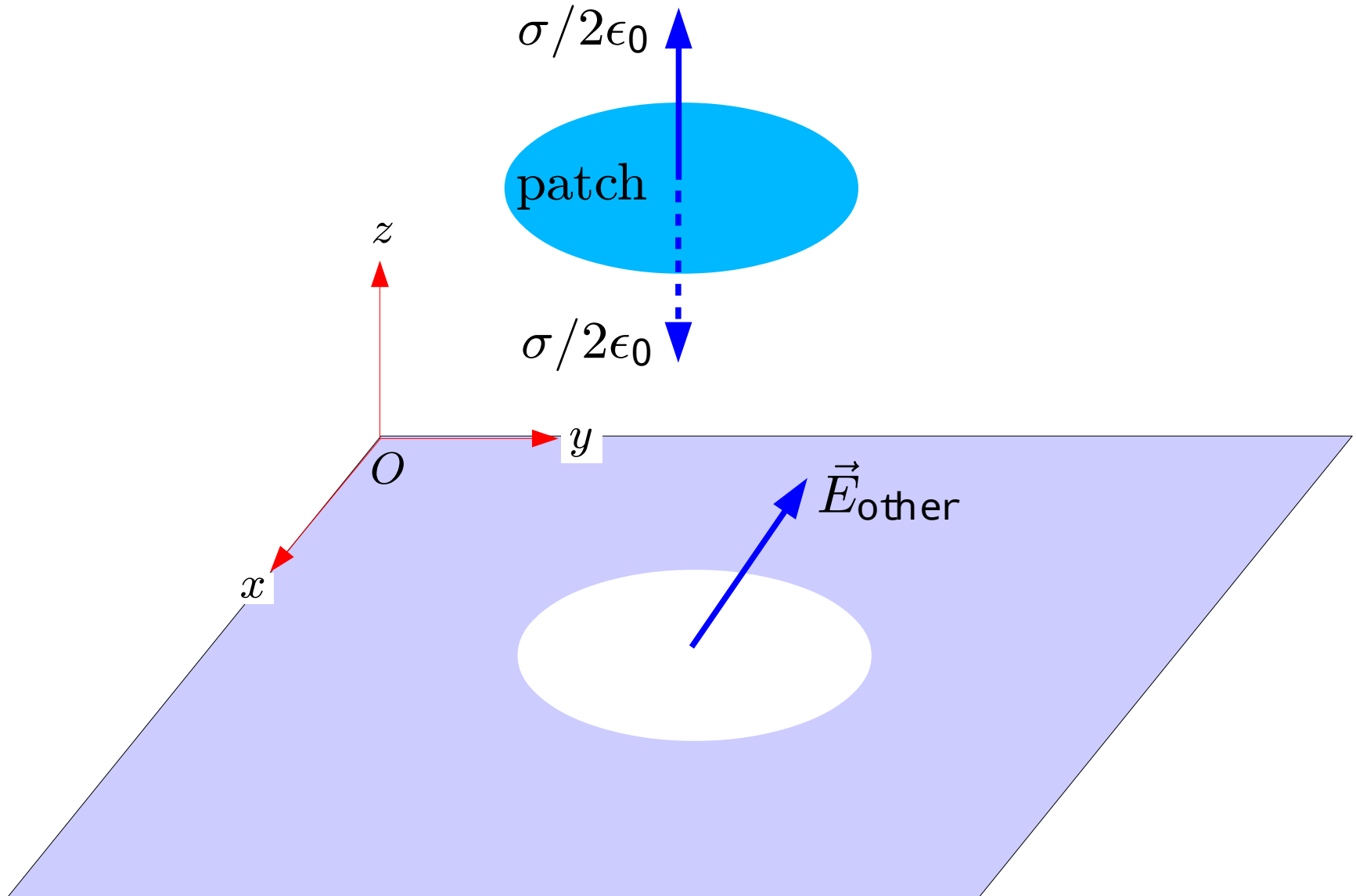


This is known as the electrostatic pressure. In terms of the electric field just above the conductor, the electrostatic pressure P is,

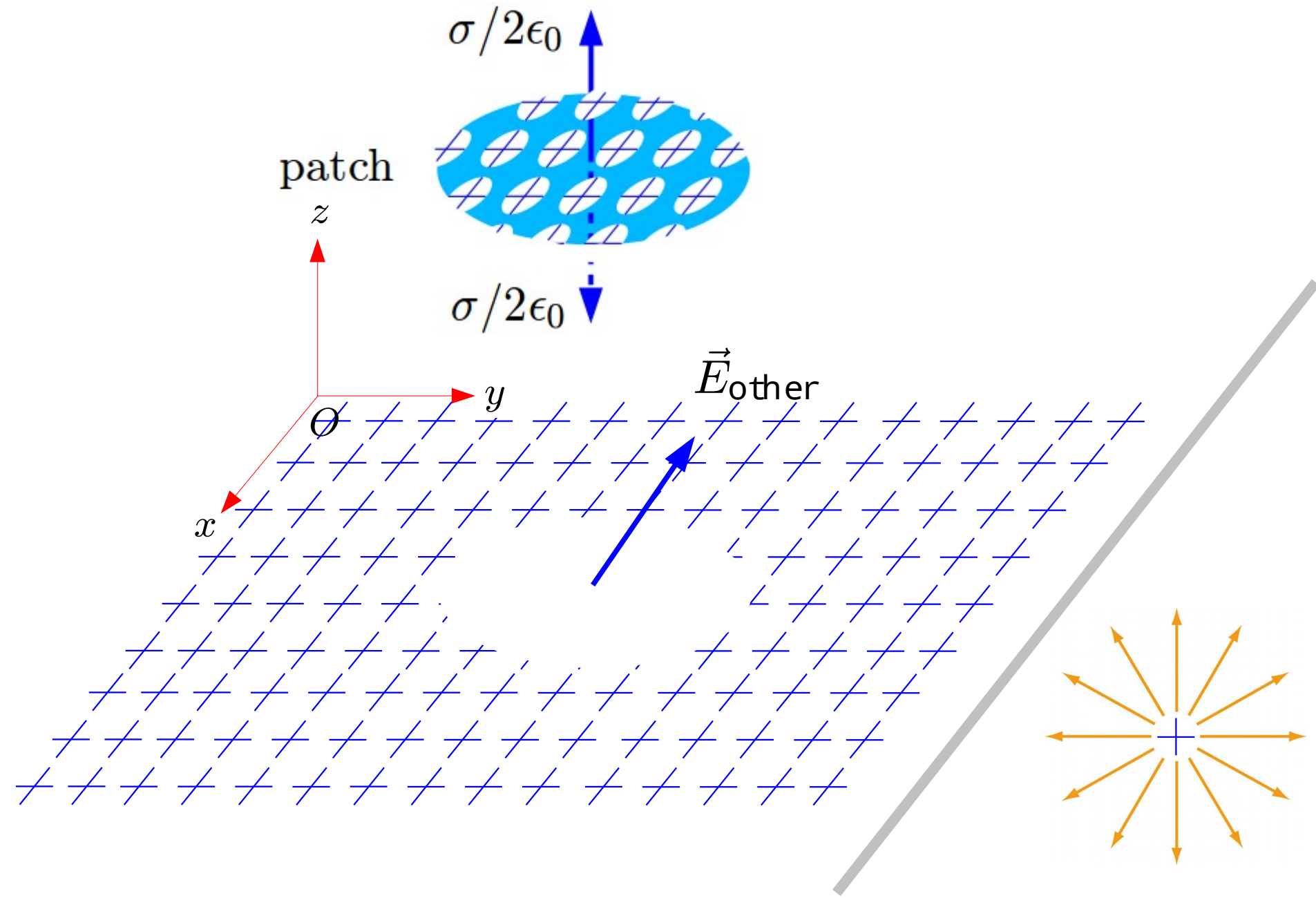
$$P = \frac{1}{2} \epsilon_0 E^2$$

Conductors : Force on a Conductor

More About Electrostatic Pressure



Conductors : Force on a Conductor



Conductors : Force on a Conductor

It can be shown that \vec{E}_{other} is equal to

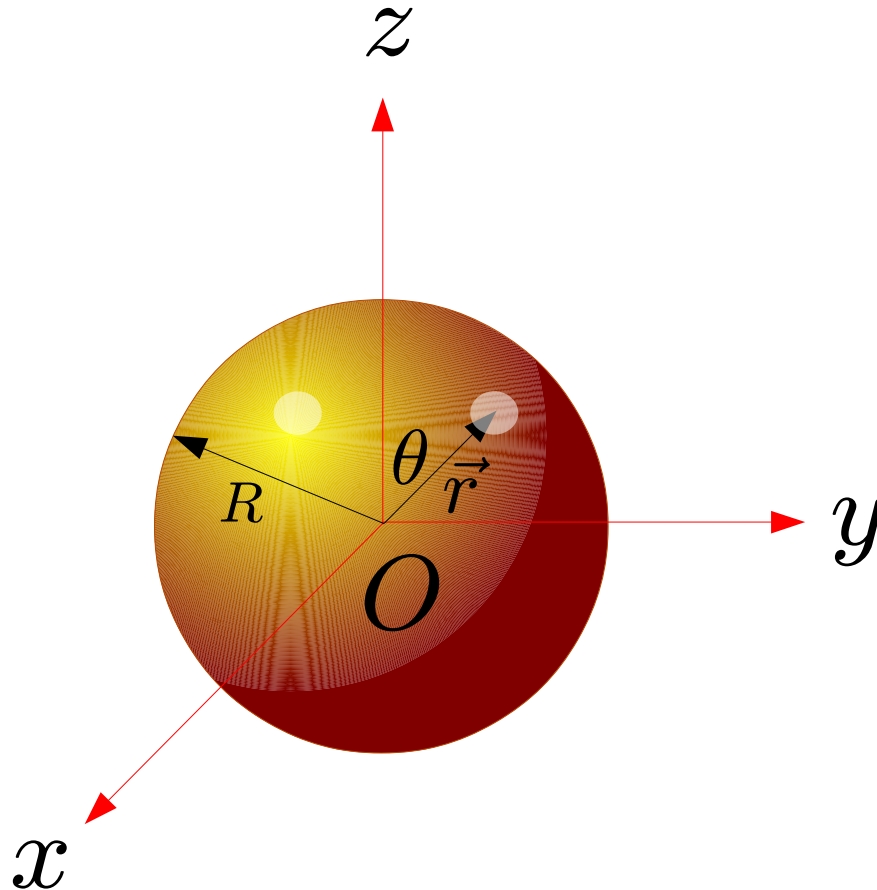
$$\vec{E}_{\text{other}} = \frac{\sigma}{2\epsilon_0} \hat{n}$$

From which we recover the previous result for the electrostatic pressure,

$$P = \frac{1}{2} \epsilon_0 E^2$$

Conductors : Force on a Conductor

Example 2.38: A metal sphere of radius R carries a total charge Q . What is the force of repulsion between the “northern” hemisphere and the “southern” hemisphere?



Conductors : Force on a Conductor

Example 2.38: A metal sphere of radius R carries a total charge Q . What is the force of repulsion between the “northern” hemisphere and the “southern” hemisphere?

The surface charge density on the sphere is

$$\sigma = \frac{Q}{4\pi R^2}$$

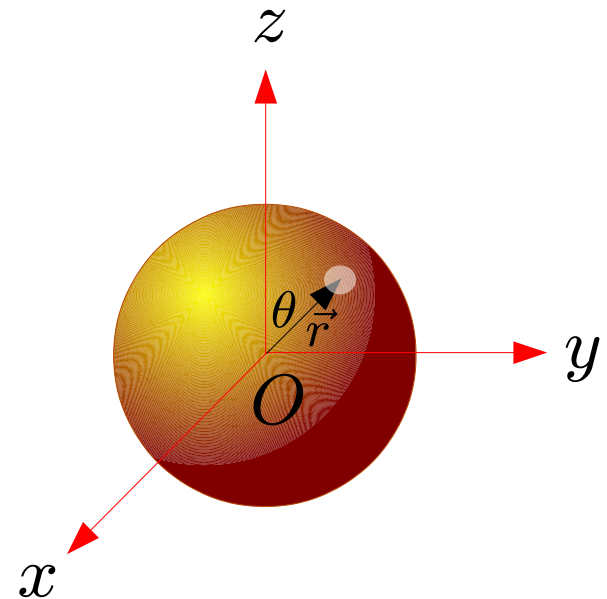
Consider an elemental patch of area da containing charge $dq = \sigma da$, as shown in the figure. The electric fields, outside and inside the sphere are

$$\vec{E}_{\text{inside}} = \vec{0}, \quad r < R$$

and $\vec{E}_{\text{outside}} = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \hat{r}, \quad r > R$

Therefore,

$$\vec{E}_{\text{avg}} = \vec{E}_{\text{other}} = \frac{1}{2} \frac{1}{4\pi\epsilon_0} \frac{Q}{R^2} \hat{r},$$



Conductors : Force on a Conductor

The z-component of the force per unit area on da is

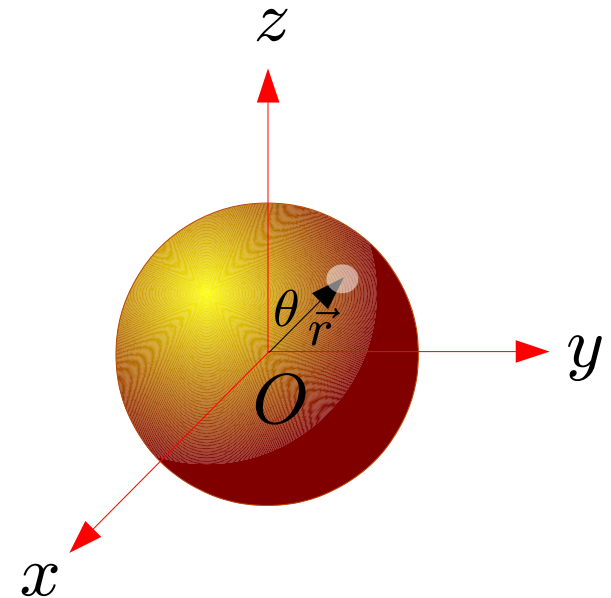
$$f_z = \sigma(\vec{E}_{\text{avg}})_z = \frac{Q}{4\pi R^2} \left(\frac{1}{2} \frac{1}{4\pi\epsilon_0} \frac{Q}{R^2} \hat{r} \cdot \hat{k} \right)$$

Thus, the total force on the “northern” hemisphere

$$\begin{aligned} F_z &= \int f_z da = \frac{Q}{4\pi R^2} \frac{1}{2} \frac{1}{4\pi\epsilon_0} \frac{Q}{R^2} \int_0^{\pi/2} (\cos \theta) R^2 \sin \theta d\theta d\phi \\ &= \left(\frac{Q}{4R} \right)^2 \frac{1}{\pi\epsilon_0} \int_0^{\pi/2} \cos \theta \sin \theta d\theta \\ &= \left(\frac{Q}{4R} \right)^2 \frac{1}{\pi\epsilon_0} \left(\frac{1}{2} \sin^2 \theta \right) \Big|_0^{\pi/2} \end{aligned}$$

The total force F_z is given by

$$F_z = \frac{Q^2}{32\pi\epsilon_0 R^2}$$



Conductors : Force on a Conductor

The surface charge density on the sphere is

$$\sigma = \frac{Q}{4\pi R^2}$$

Consider an elemental patch of area da containing charge $dq = \sigma da$. The electric field at the patch due to the rest of the charges is

$$\vec{E}_{\text{other}} = \frac{1}{2} \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \hat{r},$$

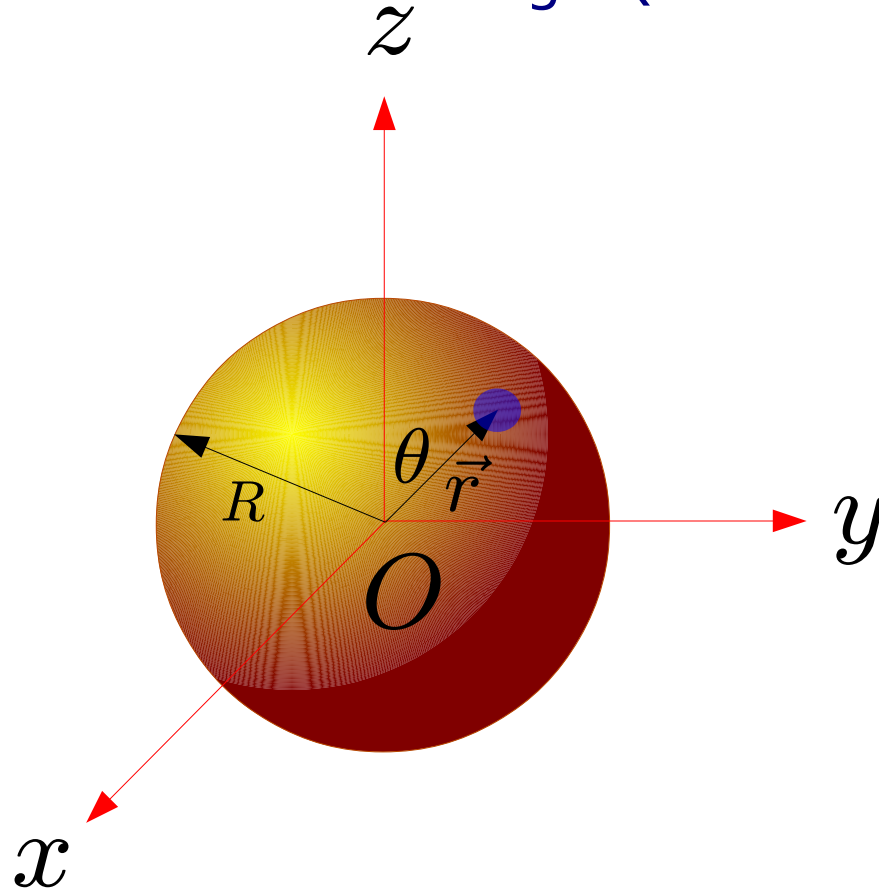
The z-component of the force per unit area on da is

$$f_z = \sigma (\vec{E}_{\text{other}})_z = \frac{Q}{4\pi R^2} \left(\frac{1}{2} \frac{1}{4\pi\epsilon_0} \frac{Q}{R^2} \hat{r} \cdot \hat{k} \right)$$

The rest of the calculation is similar to the previous approach.

Conductors : Force on a Conductor

Example 2.43: Find the net force that the southern hemisphere of a uniformly charged sphere exerts on the northern hemisphere. Express your answer in terms of the radius R and the total charge Q ?



Conductors : Force on a Conductor

Since the sphere contains a volume charge density, the electric field doesn't suffer any discontinuity, and therefore, the force (per unit volume) on a volume element $d\tau$ at \vec{r} containing charge $dq = \rho d\tau$ is given by

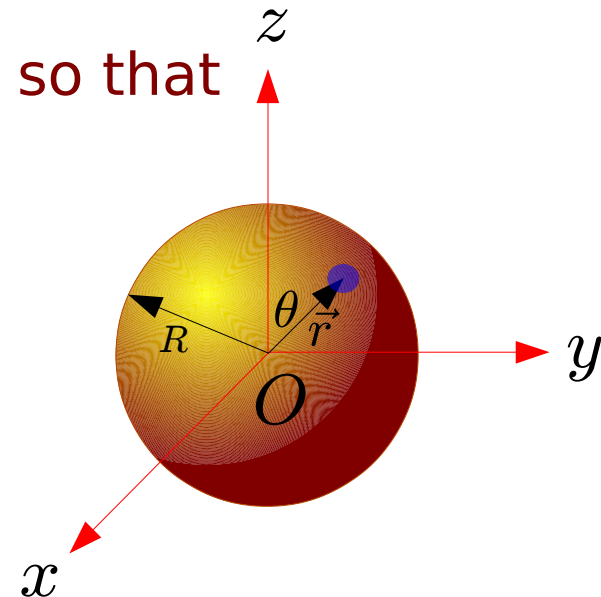
$$\vec{f} = \rho \vec{E}$$

The electric field inside the sphere is

$$\vec{E} = \frac{Qr}{4\pi\epsilon_0 R^3} \hat{r}, \quad r < R$$

Only z-component of the force survives, so that

$$f_z = \frac{3Q}{4\pi R^3} \left(\frac{Qr}{4\pi\epsilon_0 R^3} \hat{r} \cdot \hat{k} \right)$$



Conductors : Force on a Conductor

Thus, the total force on the “northern” hemisphere

$$F_z = \int f_z d\tau = \frac{3Q^2}{\epsilon_0(4\pi R^3)^2} \int_0^R r dr \int_0^{\pi/2} r^2 \cos \theta \sin \theta d\theta \int_0^{2\pi} d\phi$$

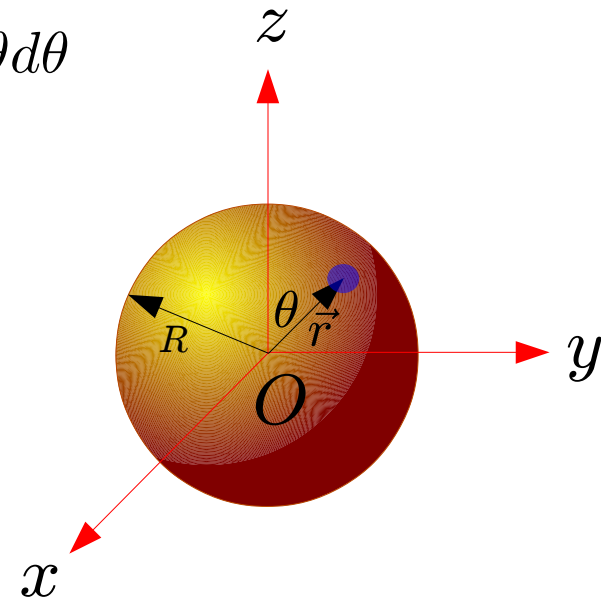
$$= \frac{3Q^2}{\epsilon_0(4\pi R^3)^2} 2\pi \int_0^R r^3 dr \int_0^{\pi/2} \cos \theta \sin \theta d\theta$$

$$= \frac{3Q^2}{\epsilon_0(4\pi R^3)^2} 2\pi \frac{R^4}{4} \left(\frac{1}{2} \sin^2 \theta \right) \Big|_0^{\pi/2}$$

$$= \frac{3Q^2}{\epsilon_0(4\pi R^3)^2} 2\pi \frac{R^4}{4} \frac{1}{2}$$

The total force F_z is given by

$$F_z = \frac{3Q^2}{64\pi\epsilon_0 R^2}$$



Conductors

Applications

- I. Use of copper in making electrical wires

Ordinary cables



Windmill generator



Superconducting cables



Conductors

Summary

1. $\vec{E}_{\text{net}}(\vec{r}) \equiv 0$ inside a conductor.
2. The volume free charge density, $\rho_{\text{inside}}^{\text{free}}(\vec{r}) = 0$ inside a conductor.
3. Any induced charges on a conductor can only reside on surface or surfaces of the conductor – as surface charge distribution, σ_{free} .
4. The entire volume and the surface of a conductor is an equipotential.
5. Just outside of the surface of a conductor, $\vec{E}_{\text{outside}}(\vec{r})$ is perpendicular to the surface.
6. The force per unit area is $\vec{f} = \sigma \vec{E}_{\text{other}} = \frac{\sigma^2}{2\epsilon_0} \hat{n}$
7. The electrostatic pressure is $P = \frac{1}{2} \epsilon_0 E^2$