The enigmatic wavefunction



Google Doodle, December 11, 2017



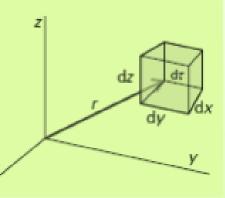
Born Interpretation

Classical wave equation:

$$\Psi(x,t)$$
 = Amplitude and $|\Psi(x,t)|^2$ = Intensity

Quantum mechanical system:

- The state is completely specified by a wavefunction $\Psi(x,t)$, which can be complex
- All possible information can be derived from $\Psi(x,t)$
- Intensity is equivalent to Probability.
- $|\Psi(x,t)|^2 = \rho(x)$, probability density.
- $|\Psi(x,t)|^2 dx = P(x)$, probability for finding the particle between x and x + dx
- In 3 dimensions, $P(x, y, z) = |\Psi(x,y,z,t)|^2 d\tau$



Laws of Quantum Mechanics

The average value of the observable corresponding to operator is

$$\langle a \rangle = \int \Psi * \hat{A} \Psi d\nu$$

Classical correspondence: Average values for a distribution function P(x):

$$\langle x \rangle = \int_{-\infty}^{\infty} x P(x) \cdot dx$$
 and $\langle x^2 \rangle = \int_{-\infty}^{\infty} x^2 P(x) \cdot dx$

Quantum mechanical analogue:

$$\left\langle a\right\rangle = \int_{-\infty}^{+\infty} \widehat{A}.P(x) \, dx = \int_{-\infty}^{+\infty} \widehat{A}. \left|\Psi\right|^2 \, dx \approx \int_{all \ space} \Psi^* \widehat{A} \Psi \, dx = \left\langle \Psi \middle| \widehat{A} \middle| \Psi \right\rangle$$

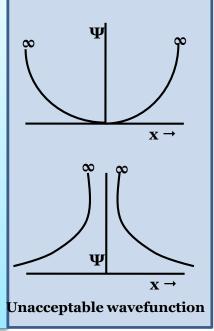
Normalization of Wavefunction

Since $\Psi^*\Psi d\tau$ is the probability, the total probability of finding the particle somewhere in space has to be unity

$$\iiint_{all \ space} \Psi^*(x, y, z).\Psi(x, y, z) dx dy dz$$

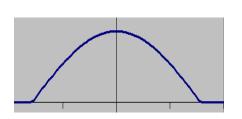
$$= \int_{all \ space} \Psi^* \Psi d\tau = \langle \Psi | \Psi \rangle = 1$$

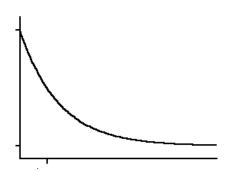
Divergent functions i.e. $\rightarrow \infty$: Ψ can not be normalized, and therefore is NOT an acceptable wave function. However, a constant value $C \neq 1$ is perfectly acceptable.

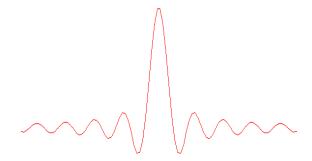


 Ψ must vanish at $\pm \infty$, or more appropriately at the boundaries and Ψ must be finite

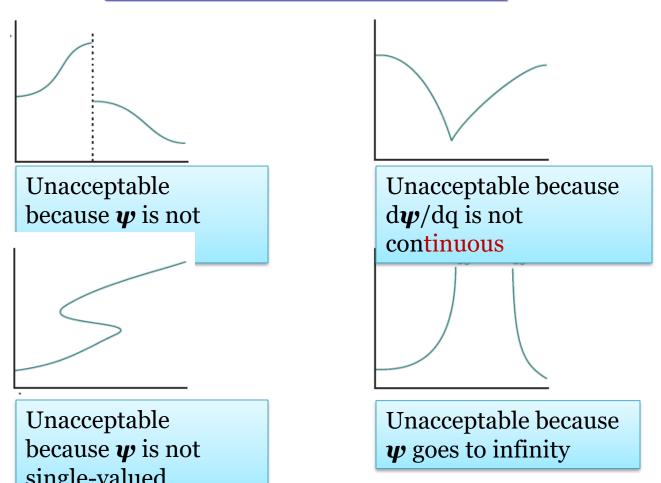
Acceptable wavefunctions







Restrictions on wavefunction



Restrictions on wavefunction

 $oldsymbol{\psi}$ must be a solution of the Schrodinger equation

 ψ must be normalizable: ψ must be finite and \rightarrow 0 at boundaries/ $\pm \infty$

 Ψ must be a continuous function of x,y,z

d**Ψ**/dq must be must be continuous in q

 Ψ must be single-valued

 Ψ must be quadratically-intergrable (square of the wavefunction should be integrable)

Origin of quantization

Quantum Mechanics

Examples of Exactly Solvable Systems

- 1. Free Particle
- 2. Particle in a Square-Well Potential
- 3. Hydrogen Atom

Time-independent Schrodinger equation

$$\widehat{H}\psi = E\psi$$

$$\left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right] \psi(x) = E \cdot \psi(x)$$

For a free particle V(x)=0There are no external forces acting

$$-\frac{\hbar^2}{2m}\frac{\partial^2}{\partial x^2}\psi(x) = E \cdot \psi(x)$$

$$-\frac{\hbar^2}{2m}\frac{d^2}{dx^2}\psi(x) = E \cdot \psi(x)$$

Second-order linear differential equation

Let us assume

$$\psi(x) = A\sin kx + B\cos kx$$

Trial Solution

$$y(x) = A\sin kx + B\cos kx$$

$$\frac{d}{dx}y(x) = \frac{d}{dx}(A\sin kx + B\cos kx) = k(A\cos kx - B\sin kx)$$

$$\frac{d^2}{dx^2}y(x) = -k^2(A\sin kx + B\cos kx) = -k^2y(x)$$

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$$\frac{\hbar^2}{2m}k^2\psi(x) = E \cdot \psi(x) \quad \Rightarrow \quad E = \frac{\hbar^2k^2}{2m} \quad \Rightarrow \quad k = \pm \frac{\sqrt{2mE}}{\hbar}$$

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de Broglie wave

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$$E = \frac{\hbar^2 k^2}{2m}$$
 There are no restrictions on k
 E can have any value
Energies of free particles are

continuous
$$\frac{\sqrt{2mE}}{\psi(x)} = A\sin\frac{\sqrt{2mE}}{\hbar}x + B\cos\frac{\sqrt{2mE}}{\hbar}x$$

Particle in 1-D Square-Well Potential

