Computational Physics

Assignment 5

Cheng-Ching Lin

1. Poisson Equation in 3D

(a). In Cartesian Coordinate

$$\begin{split} \nabla^2 \varphi &= \, -\frac{\rho}{\varepsilon_0} \\ \nabla^2 \phi \, + \frac{\rho}{\varepsilon_0} &= \, 0 \\ \frac{\varphi_{i+1,j,k}^n + \varphi_{i-1,j,k}^n - 2 \varphi_{i,j,k}^n}{h^2} \, + \, \frac{\varphi_{i,j+1,k}^n + \varphi_{i,j-1,k}^n - 2 \varphi_{i,j,k}^n}{h^2} \, + \frac{\varphi_{i,j,k+1}^n + \varphi_{i,j,k-1}^n - 2 \varphi_{i,j,k}^n}{h^2} \, + \frac{\rho_{i,j,k}}{\varepsilon_0} &= \, 0 \end{split}$$

For

$$x_{i} = h(i - 1) + x_{0}$$

$$y_{j} = h(j - 1) + y_{0}$$

$$z_{k} = h(k - 1) + z_{0}$$

$$t_{n} = \tau(n - 1) + t_{0}$$

We make time vary into the Poisson Equation. It will become some kinds of diffusion equation.

$$\begin{split} \nabla^2 \phi + \frac{\rho}{\varepsilon_0} &= \frac{\partial \phi}{\partial t} \\ \frac{ \phi_{i+1,j,k}^n + \phi_{i-1,j,k}^n - 2 \phi_{i,j,k}^n}{h^2} + \frac{\phi_{i,j+1,k}^n + \phi_{i,j-1,k}^n - 2 \phi_{i,j,k}^n}{h^2} + \frac{\phi_{i,j,k+1}^n + \phi_{i,j,k-1}^n - 2 \phi_{i,j,k}^n}{h^2} + \frac{\rho_{i,j,k}}{\varepsilon_0} \\ &= \frac{\phi_{i,j,k}^{n+1} - \phi_{i,j,k}^n}{\tau} \\ \frac{\tau}{h^2} \left(\phi_{i+1,j,k}^n + \phi_{i-1,j,k}^n + \phi_{i,j+1,k}^n + \phi_{i,j-1,k}^n + \phi_{i,j,k+1}^n + \phi_{i,j,k-1}^n - 6 \phi_{i,j,k}^n \right) + \phi_{i,j,k}^n + \frac{\tau \rho_{i,j,k}}{\varepsilon_0} = \phi_{i,j,k}^{n+1} \end{split}$$

For the stability of this equation and also to make this equation simpler, we choose $\frac{\tau}{h^2} = \frac{1}{6}$.

$$\frac{\left(\phi_{i+1,j,k}^n + \phi_{i-1,j,k}^n + \phi_{i,j+1,k}^n + \phi_{i,j-1,k}^n + \phi_{i,j,k+1}^n + \phi_{i,j,k-1}^n\right)}{6} + \frac{h^2 \rho_{i,j,k}}{6\varepsilon_0} = \phi_{i,j,k}^{n+1}$$

(b). In cylinder coordinate

In cylinder coordinate, the Laplace Operator will become this terms:

$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial z^2} = \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{\partial^2 \phi}{\partial r^2} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \psi^2} + \frac{\partial^2 \phi}{\partial z^2}$$

For

$$\begin{aligned} \mathbf{r}_{\mathrm{i}} &= h(i-1) + r_{0} \\ \psi_{\mathrm{j}} &= h(j-1) + \psi_{0} \\ \mathbf{z}_{\mathrm{k}} &= h(k-1) + z_{0} \\ \mathbf{t}_{\mathrm{n}} &= \tau(n-1) + t_{0} \\ 0 &\leq \mathrm{h} \leq 2\pi \end{aligned}$$

$$\begin{split} \frac{1}{r_i} \frac{\varphi_{i+1,j,k}^n - \varphi_{i-1,j,k}^n}{2h} + \frac{\varphi_{i+1,j,k}^n + \varphi_{i-1,j,k}^n - 2\varphi_{i,j,k}^n}{h^2} + \frac{1}{r_i^2} \frac{\varphi_{i,j+1,k}^n + \varphi_{i,j-1,k}^n - 2\varphi_{i,j,k}^n}{h^2} \\ + \frac{\varphi_{i,j,k+1}^n + \varphi_{i,j,k-1}^n - 2\varphi_{i,j,k}^n}{h^2} + \frac{\rho_{i,j,k}}{\varepsilon_0} &= \frac{\varphi_{i,j,k}^{n+1} - \varphi_{i,j,k}^n}{\tau} \\ \frac{\tau}{h^2 r_i^2} \left(\frac{h}{2} r_i \varphi_{i+1,j,k}^n - \frac{h}{2} r_i \varphi_{i-1,j,k}^n + r_i^2 \varphi_{i+1,j,k}^n + r_i^2 \varphi_{i-1,j,k}^n - 2r_i^2 \varphi_{i,j,k}^n + \varphi_{i,j+1,k}^n + \varphi_{i,j-1,k}^n - 2\varphi_{i,j,k}^n \right) \\ + r_i^2 \varphi_{i,j,k+1}^n + r_i^2 \varphi_{i,j,k-1}^n - 2r_i^2 \varphi_{i,j,k}^n \right) + \varphi_{i,j,k}^n + \frac{\tau \rho_{i,j,k}}{\varepsilon_0} &= \varphi_{i,j,k}^{n+1} \\ \frac{\tau}{h^2 r_i^2} \left(\frac{h}{2} r_i \varphi_{i+1,j,k}^n - \frac{h}{2} r_i \varphi_{i-1,j,k}^n + r_i^2 \varphi_{i-1,j,k}^n + r_i^2 \varphi_{i-1,j,k}^n + \varphi_{i,j+1,k}^n + \varphi_{i,j-1,k}^n + r_i^2 \varphi_{i,j,k-1}^n + r_i^2 \varphi_{i,j,k+1}^n \right) \\ + \frac{\tau}{h^2 r_i^2} \left(-2r_i^2 - 2 - 2r_i^2 \right) \varphi_{i,j,k}^n + \varphi_{i,j,k}^n + \frac{\tau \rho_{i,j,k}}{\varepsilon_0} = \varphi_{i,j,k}^{n+1} \end{split}$$

If we do what we just do in (a) and make it simpler, then we need to solve the equation.

$$\frac{\tau}{h^2 r_i^2} (4r_i^2 + 2) = 1$$

And it's very complicated, so it shows that cylinder coordinate doesn't prefer to use in numerical solution.

(c). In spherical coordinate

In cylinder coordinate, the Laplace Operator will become this terms:

$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial z^2} = \frac{2}{r} \frac{\partial \phi}{\partial r} + \frac{\partial^2 \phi}{\partial r^2} + \frac{\cot \theta}{r^2} \frac{\partial \phi}{\partial \theta} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} + \frac{1}{r^2 sin\theta} \frac{\partial^2 \phi}{\partial \psi^2}$$

For

$$r_{i} = h(i-1) + r_{0}$$

$$\theta_{j} = h(j-1) + \theta_{0}$$

$$\psi_{k} = h(k-1) + \psi_{0}$$

$$t_{n} = \tau(n-1) + t_{0}$$

$$0 \le h \le 2\pi$$

$$\begin{split} \frac{2}{r_{i}} \frac{\Phi_{i+1,j,k}^{n} - \Phi_{i-1,j,k}^{n}}{2h} + \frac{\Phi_{i+1,j,k}^{n} + \Phi_{i-1,j,k}^{n} - 2\Phi_{i,j,k}^{n}}{h^{2}} + \frac{\cot\theta_{j}}{r_{i}^{2}} \frac{\Phi_{i,j+1,k}^{n} - \Phi_{i,j-1,k}^{n}}{2h} \\ + \frac{1}{r_{i}^{2}} \frac{\Phi_{i,j+1,k}^{n} + \Phi_{i,j-1,k}^{n} - 2\Phi_{i,j,k}^{n}}{h^{2}} + \frac{1}{r_{i}^{2} \sin\theta_{j}} \frac{\Phi_{i,j,k+1}^{n} + \Phi_{i,j,k-1}^{n} - 2\Phi_{i,j,k}^{n}}{h^{2}} + \frac{\Phi_{i,j,k}^{n}}{\tau} + \frac{\Phi_{i,j,k}^{n}}{\tau} \\ = \frac{\Phi_{i,j,k}^{n+1}}{\tau} \\ \frac{\tau \cos\theta_{j}}{r_{i}^{2}h^{2} \sin\theta_{j}} \left(r_{i}h tan\theta_{j} \Phi_{i+1,j,k}^{n} - r_{i}h tan\theta_{j} \Phi_{i-1,j,k}^{n} + r_{i}^{2} tan\theta_{j} \Phi_{r+1,i,j}^{n} + r_{i}^{2} tan\theta_{j} \Phi_{i-1,j,k}^{n} + \frac{h}{2} \Phi_{i,j+1,k}^{n} \\ - \frac{h}{2} \Phi_{i,j-1,k}^{n} + tan\theta_{j} \Phi_{i,j-1,k}^{n} + tan\theta_{j} \Phi_{i,j+1,k}^{n} + \frac{h^{2}}{\cos\theta_{j}} \Phi_{i,j,k-1}^{n} + \frac{h^{2}}{\cos\theta_{j}} \Phi_{i,j,k+1}^{n} \right) \\ + \frac{\tau \cos\theta_{j}}{r_{i}^{2}h^{2} \sin\theta_{i}} \left(-2r_{i}^{2} tan\theta_{j} - 2 tan\theta_{j} - 2 tan\theta_{j} - \frac{2h^{2}}{\cos\theta_{i}} \right) \Phi_{i,j,k}^{n} + \Phi_{i,j,k}^{n} + \frac{\tau \rho_{i,j,k}}{\varepsilon_{0}} = \Phi_{i,j,k}^{n+1} \end{split}$$

So,...

$$\frac{2\tau\cos\theta_{j}}{r_{i}^{2}h^{2}\sin\theta_{j}}\left(-2r_{i}^{2}\tan\theta_{j}-2\tan\theta_{j}-\frac{2h^{2}}{\cos\theta_{j}}\right)=-1$$

We can see that the equation is so hard to solve it.

Therefore, the Cartesian coordinate always prefers to use in numerical solutions.

2. A charge in a box

First, I create the function for calculate the electric potential.

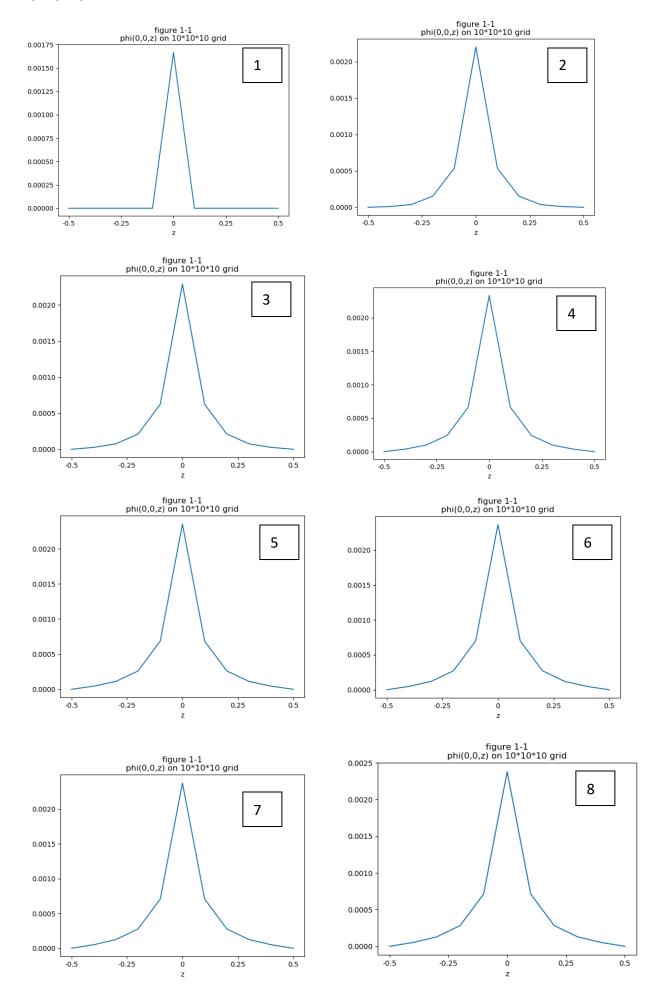
Code:

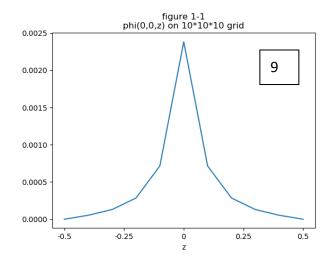
(1). Calculate the electric potential in 10*10*10 grid and 100*100*100 grid, and also plot the electric potential on (0,0,z)

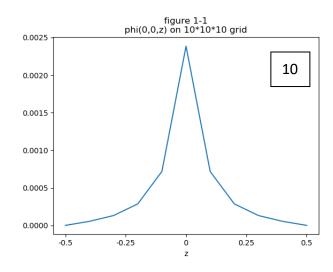
```
plt.show()
```

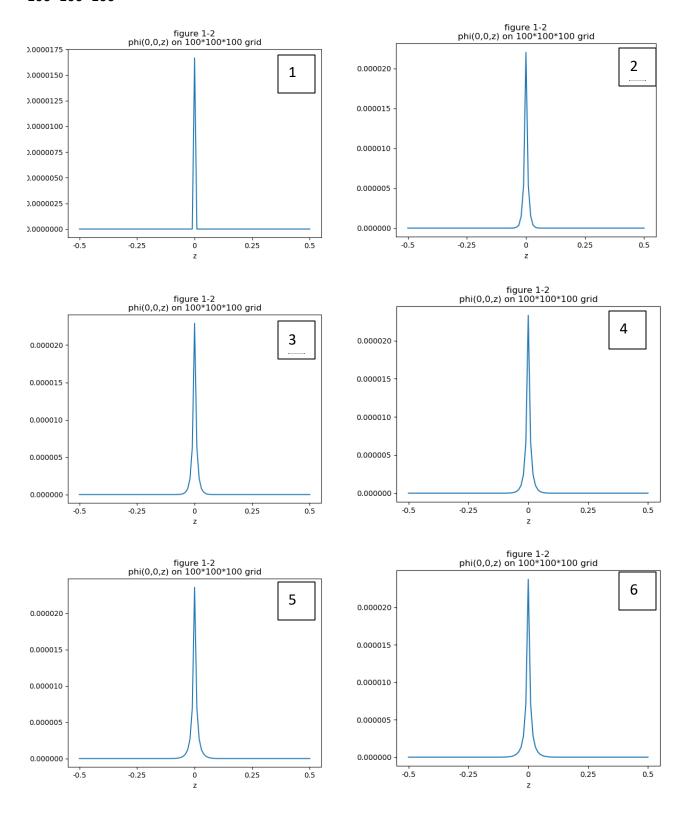
and the result:

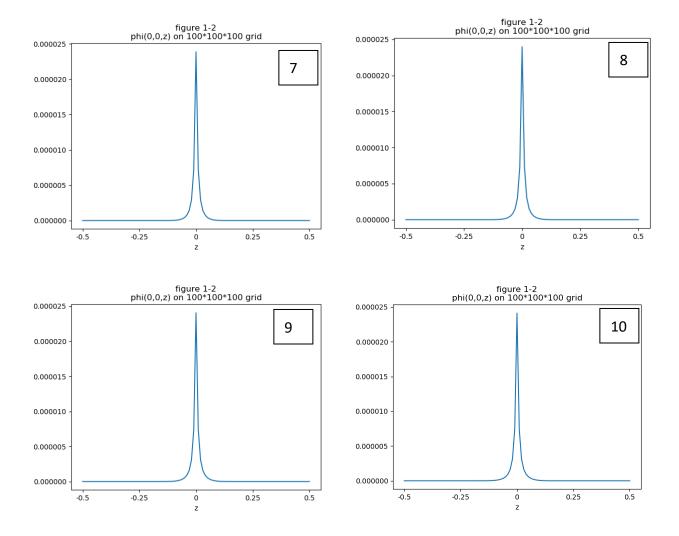
10*10*10







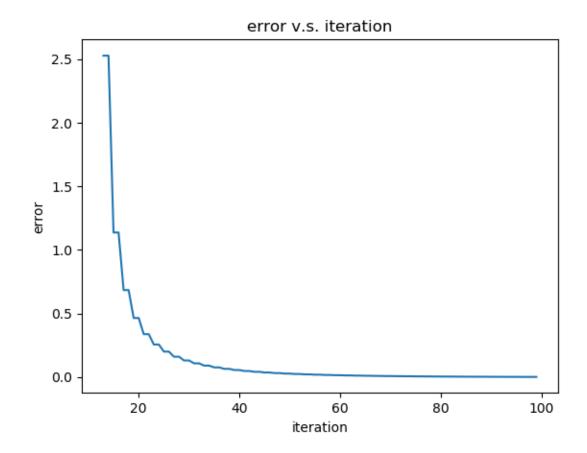


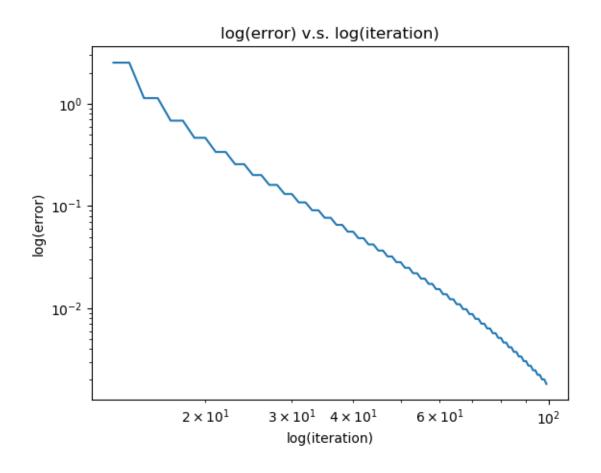


Because the equation is some kinds of distribution equation, we can see some kind of distribution in this figure. We can see that in the same time steps, the second figure is more sharper than first one, that is because the space resolutions is more higher, we can see that there is a peak (a charge) in the center.

(2). The error v.s. iterations

```
k = 0.5 - kdx * h
         i = 0.5 - i dx * h
              deltamax = deltatemp
```





(3). Calculate the electric potential in 10*10*10 grid and 100*100*100 grid, and also plot the electric potential on (0,0,z).

This time, the boundary has electric potential, and there is no charge in origin:

$$\frac{q}{\sqrt{x^2 + y^2 + z^2}}$$

First, I need to create another function to calculate the electric potential.

```
t = n*tau
```

After this code, we need to calculate the electric potential in 10*10*10 grid and 100*100*100 grid. Code:

```
j = 1.5 - jdx*h
```

```
for idx in range(301):
    i = 1.5 - idx * h
        temp = phi_2(i, j, k, n, pp, h)
        pp[idx, jdx, kdx, ndx] = temp

plt.plot(pp[150, 150, :, ndx])

plt.xticks([0, 50, 100, 150, 200, 250, 300], [-1.5, -1.0, -0.5, 0, 0.5, 1.0, 1.5])

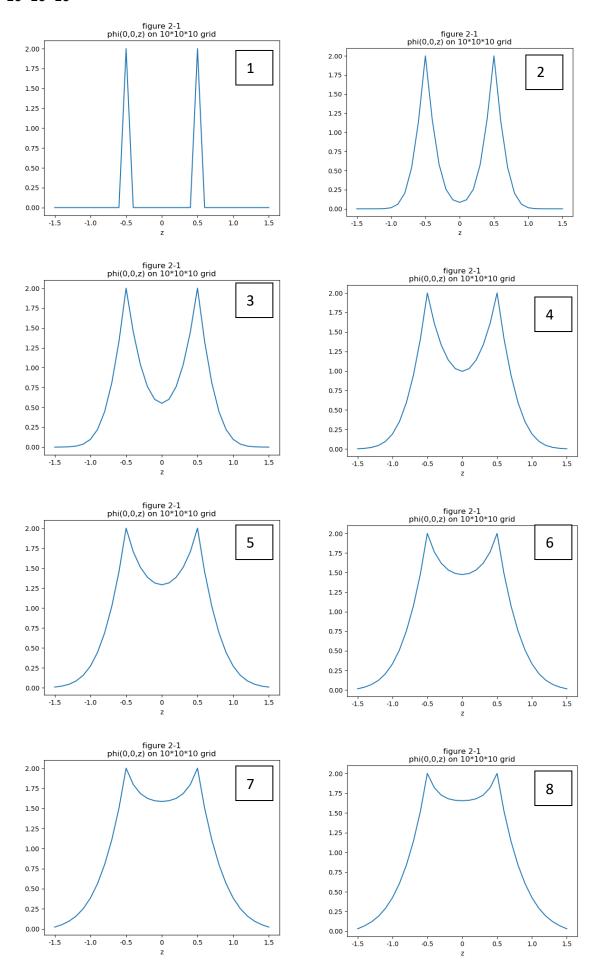
plt.title("figure 2-2\nphi(0,0,z) on 100*100*100 grid")

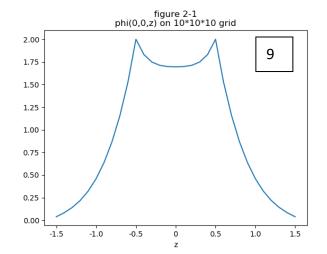
plt.xlabel("z")

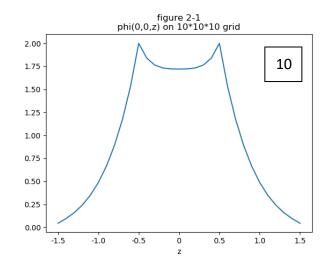
plt.show()
```

And the result:

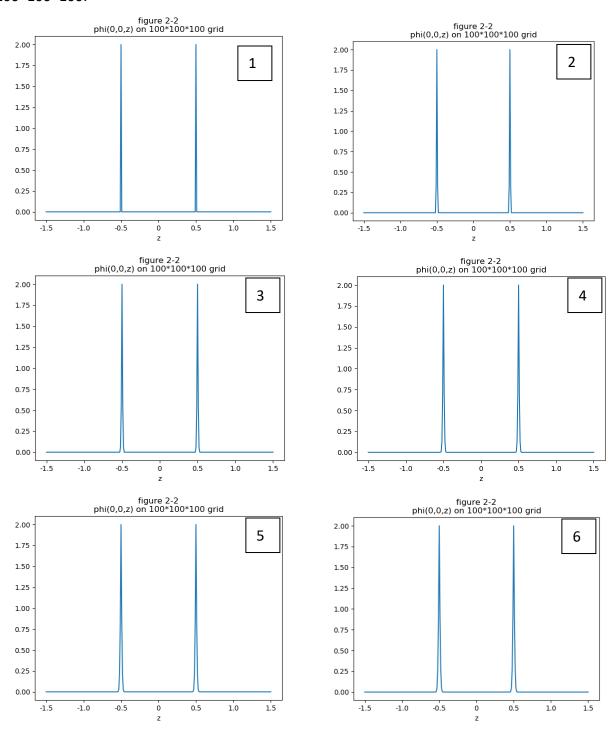
10*10*10

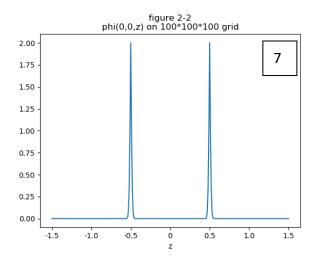


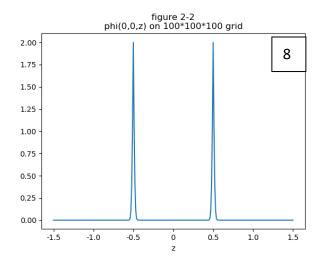


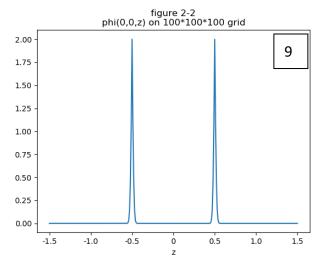


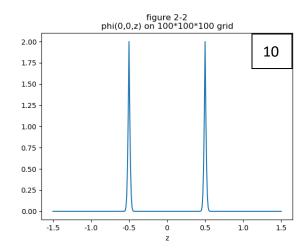
100*100*100:





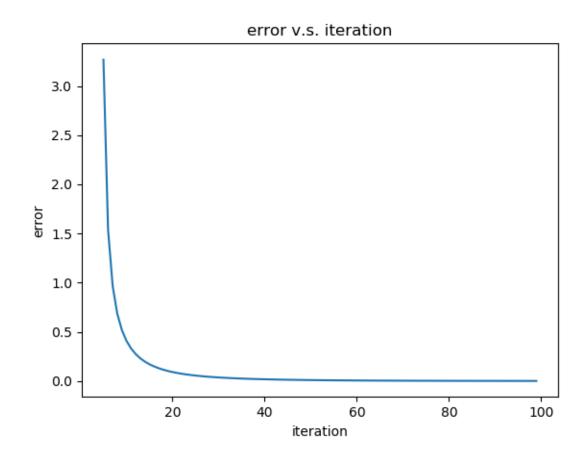


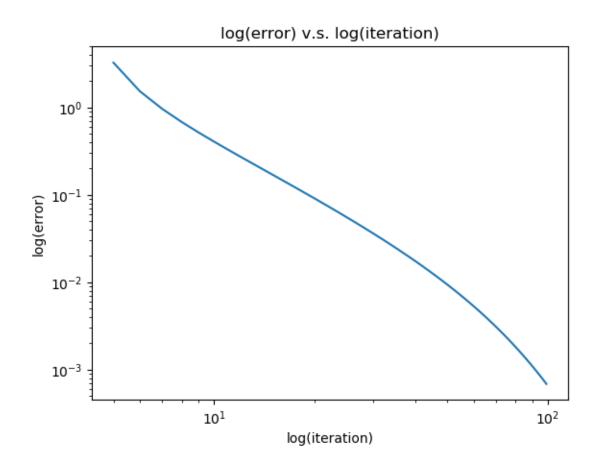




(4). The max difference v.s. iterations

```
i = 0.5 - idx*h
```





3. A box in a box

(1). Find E on xy-plane(z=0):

First, I create the function to calculate the phi.

```
elif i =0.5 and j \leq 0.5 and j \geq -0.5 and k \leq0.5 and k \geq -0.5:#inside box, I use poisson equation
 h^*h/6+(pp[xi+l,yj,zk,n-l]+pp[xi-l,yj,zk,n-l]+pp[xi,yj+l,zk,n-l]+pp[xi,yj-l,zk,n-l]+pp[xi,yj,zk+l,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi,zk,n-l]+pp[xi
 h^*h/6+(pp[xi+l,yj,zk,n-l]+pp[xi-l,yj,zk,n-l]+pp[xi,yj+l,zk,n-l]+pp[xi,yj-l,zk,n-l]+pp[xi,yj,zk+l,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yj,zk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk,n-l]+pp[xi,yk-l,zk
                                                                             ans =
h*h/6+(pp[xi+1,yj,zk,n-1]+pp[xi-1,yj,zk,n-1]+pp[xi,yj+1,zk,n-1]+pp[xi,yj-1,zk,n-1]+pp[xi,yj,zk+1,n-1]+pp[xi,yj,zk-1,
```

```
ans = 0
else: #I use the laplace equation
ans =

(pp[xi+1,yj,zk,n-1]+pp[xi-1,yj,zk,n-1]+pp[xi,yj+1,zk,n-1]+pp[xi,yj-1,zk,n-1]+pp[xi,yj,zk+1,n-1]+pp[xi,yj,zk-1,n-1])/6
return ans
```

And then, I start to find E. The E equation is:

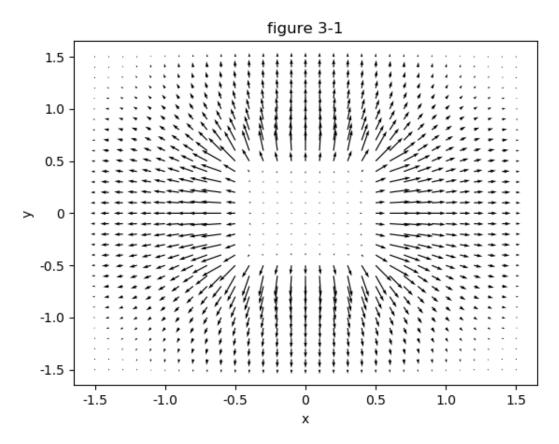
$$\vec{E} = -\nabla \phi$$

$$\vec{E} = -\left(\frac{\partial \phi}{\partial x}\hat{x} + \frac{\partial \phi}{\partial y}\hat{y} + \frac{\partial \phi}{\partial z}\hat{z}\right)$$

$$\vec{E} = -\left(\frac{\phi_{i+1} - \phi_{i-1}}{2h}\hat{x} + \frac{\phi_{j+1} - \phi_{j-1}}{2h}\hat{y} + \frac{\phi_{k+1} - \phi_{k-1}}{2h}\hat{z}\right)$$

```
k = 1.5 - kdx * h
```

And the result:



We can see that the space between two boxes, the electric field flows from inner to outer, and all symmetric at origin.

(2).to find charge distribution:

The equation of charge distribution is:

$$\nabla \cdot \vec{\mathbf{E}} = \frac{\rho}{\varepsilon_0}$$

For $\,\rho=\frac{q}{v}$ is charge density, and we suppose $\epsilon_0=1.$ So, the equation will become this:

$$\begin{split} \frac{\partial \overrightarrow{E}}{\partial x} + \frac{\partial \overrightarrow{E}}{\partial y} + \frac{\partial \overrightarrow{E}}{\partial z} &= \frac{q}{V} \\ \frac{E_{i+1} - E_{i-1}}{2h} + \frac{E_{j+1} - E_{j-1}}{2h} + \frac{E_{k+1} - E_{k-1}}{2h} &= \frac{q}{V} \end{split}$$

```
ex[idx, jdx] = 0
ex[idx,jdx] = -(pp[idx+1,jdx,10,100])/(2*h)
```

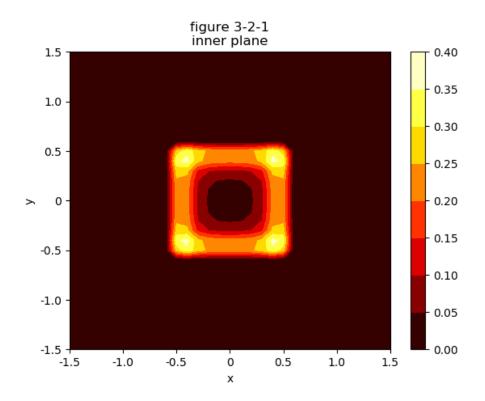
```
ey[idx, jdx] = -(-pp[idx, jdx-1, 10, 100])/(2*h)
ex[idx,jdx] = -(pp[idx+1,jdx,10,100]-pp[idx-1,jdx,10,100])/(2*h)
ey[idx,jdx] = -(pp[idx,jdx+1,10,100]-pp[idx,jdx-1,10,100])/(2*h)
eyo[idx,jdx] = -(pp[idx,jdx+1,1,99]-pp[idx,jdx-1,1,99])/(2*h*6)
```

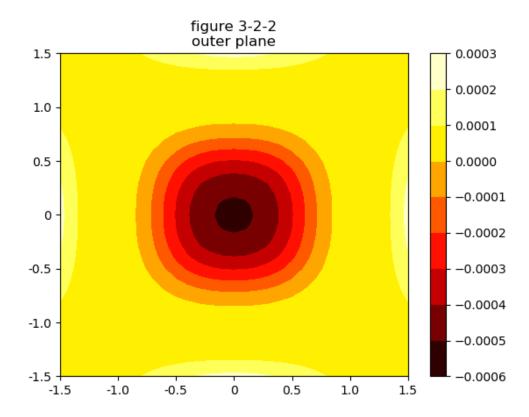
```
(eyo[idx, jdx + 1]) / (2 * h))
```

(eyo[idx, jdx + 1]) / (2 * h))

```
aa = plt.contourf(qo, cmap = plt.cm.hot)
plt.colorbar(aa, orientation="vertical")
plt.xticks([0, 5, 10, 15, 20, 25, 30], [-1.5, -1.0, -0.5, 0, 0.5, 1.0, 1.5])
plt.yticks([0, 5, 10, 15, 20, 25, 30], [-1.5, -1.0, -0.5, 0, 0.5, 1.0, 1.5])
plt.title("figure 3-2-2\nouter plane")
plt.show()
```

and the result:





We can see the figure on the inner plane, the charge will more on the edge of plane, especially on the corner. On the outer plane, we can see that the charge will be crowed on the center for the projection of inner box.

(3). Find E on xy-plane(z=0):

This time, we move the inner box to place at a distance of 0.3 from one of the sides of the outer box. First, I need to create a function to calculate phi, and I move the box to $(0.5^{\circ}-0.5, -0.2^{\circ}-1.2, 0.5^{\circ}-0.5)$. Code:

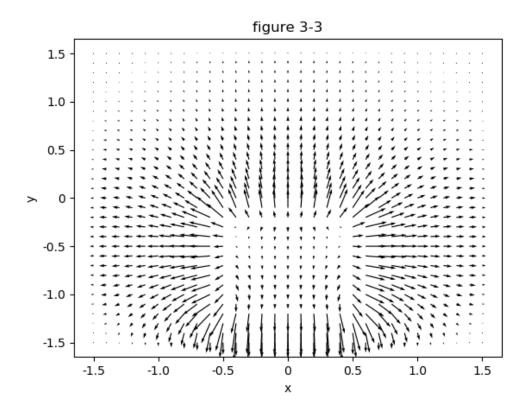
```
t = n*tau
```

and I start to find E.

```
n = ndx
        if jdx >= 30 or idx >= 30:
        ex[idx,jdx] = -(pp[idx+1,jdx,15,1000]-pp[idx-1,jdx,15,1000])/(2*h)
```

```
plt.ylabel("y")
plt.xlabel("x")
plt.show()
```

the result:



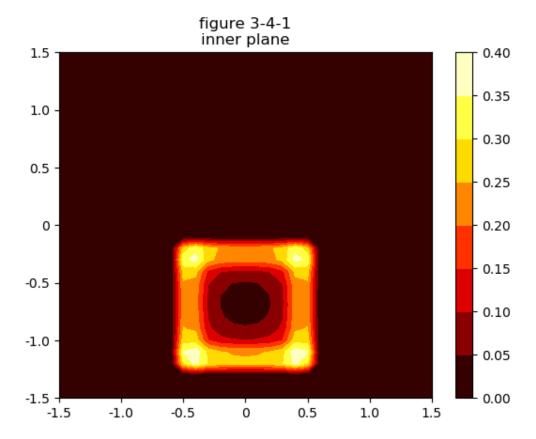
Because the inner box is not at the center anymore, we can see that the electric field is stronger when the plane is closer to outer box. Also, we can see inside the inner box, the electric field is not symmetry anymore.

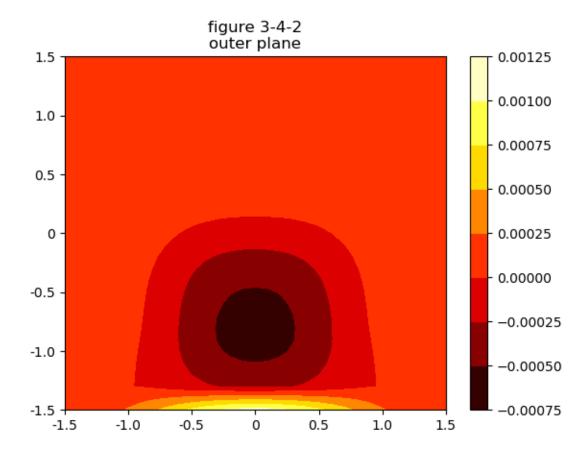
(4).charge distribution

```
def test_4():
    pp = np.zeros((31, 31, 31, 101))
    ex = np.zeros((31, 31))
    ey = np.zeros((31, 31))
    exo = np.zeros((31, 31))
    eyo = np.zeros((31, 31))
    qi = np.zeros((31,31))
    qo = np.zeros((31,31))
    #calculate electric potential
    for ndx in range(101):
        n = ndx
        for kdx in range(31):
        k = 1.5 - kdx * h
```

```
exo[idx,jdx] = -(pp[idx+1,jdx,1,99]-pp[idx-1,jdx,1,99])/(2*h*6)
eyo[idx,jdx] = -(pp[idx,jdx+1,1,99]-pp[idx,jdx-1,1,99])/(2*h*6)
elif jdx >= 30:
```

The result:





These are just the same as previous one, but we can see that the outer plane closer to inner box, the electric field is stronger than others are.

(5).

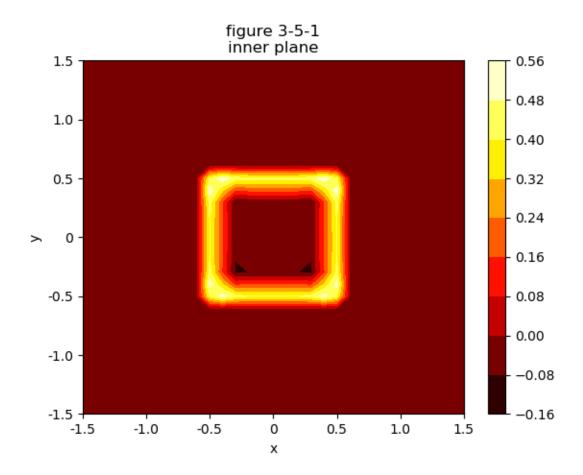
I want to see what will happen on the closest plane on inner and outer box, so I find charge distribution again.

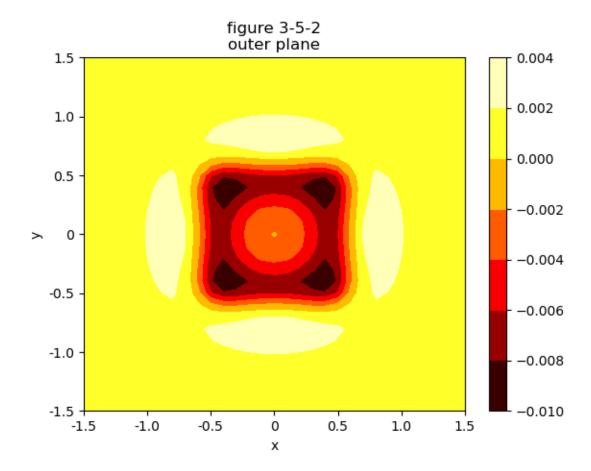
```
def test_5():
    pp = np.zeros((31, 31, 31, 1001))
    ex = np.zeros((31, 31))
    ey = np.zeros((31, 31))
    exo = np.zeros((31, 31))
    eyo = np.zeros((31, 31))
    qi = np.zeros((31, 31))
    qo = np.zeros((31, 31))
    for ndx in range(101):
        n = ndx
        for kdx in range(31):
        k = 1.5 - kdx * h
        for jdx in range(31):
        j = 1.5 - jdx * h
        for idx in range(31):
```

```
eyo[idx,jdx] = -(pp[1,jdx+1,idx,99]-pp[1,jdx-1,idx,99])/(2*h*6)
```

exo[idx,jdx] = -(pp[1,jdx,idx+1,99]-pp[1,jdx,idx-1,99])/(2*h*6)

The result:





The result is the same as (2). But the charge is more stronger than ever, and the outer plane is much clear to see the projection of inner box.