

# Quantum Computation

## (Lecture 3)

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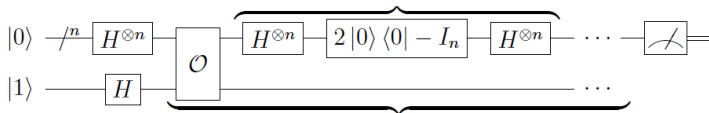
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# Grover's algorithm

Recall Grover's algorithm:

- Prepare the initial state:  $|0\rangle^{\otimes n}|1\rangle$
- Apply  $H^{\otimes n} \otimes H$  to yield  $\frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle|-\rangle$
- Apply the Grover iterator  $G$  to  $\frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle|-\rangle$  a **suitable number of times** to obtain state  $|a\rangle|-\rangle$  with high probability
- Measure the first  $n$  qubits to retrieve  $|a\rangle$



## A geometric perspective on $G$

**Initial state:**  $|\psi\rangle = \frac{1}{\sqrt{N}}|a\rangle + \sqrt{\frac{N-1}{N}}|r\rangle$

The repeated application of  $G$  leaves the system in the 2-dimensional subspace of the original  $N$ -dimensional space, spanned by  $|a\rangle$  and  $|r\rangle$ .

Another basis is given by  $|\psi\rangle$  and the state **orthogonal** to  $|\psi\rangle$ :

$$|\bar{\psi}\rangle = \sqrt{\frac{N-1}{N}}|a\rangle - \frac{1}{\sqrt{N}}|r\rangle$$

Define an angle  $\theta$  st  $\sin \theta = \frac{1}{\sqrt{N}}$  (and, of course,  $\cos \theta = \sqrt{\frac{N-1}{N}}$ ), and express both bases as

$$\begin{aligned} |\psi\rangle &= \sin \theta |a\rangle + \cos \theta |r\rangle & |\bar{\psi}\rangle &= \cos \theta |a\rangle - \sin \theta |r\rangle \\ |a\rangle &= \sin \theta |\psi\rangle + \cos \theta |\bar{\psi}\rangle & |r\rangle &= \cos \theta |\psi\rangle - \sin \theta |\bar{\psi}\rangle \end{aligned}$$

## A geometric perspective on $G$

$G$  has two components:

- $V$  which applies a phase shift to  $|a\rangle$ : reflection over  $|r\rangle$ .
- $W$  which applies a phase shift to all vectors in the subspace orthogonal to  $|\psi\rangle$ : reflection over  $|\psi\rangle$ .

Let's express the action of  $V$  in the basis  $|\psi\rangle, |\bar{\psi}\rangle$  to perform afterwards the second reflection:

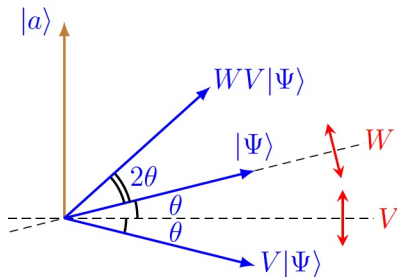
$$\begin{aligned} V|\psi\rangle &= -\sin\theta|a\rangle + \cos\theta|r\rangle \\ &= -\sin\theta(\sin\theta|\psi\rangle + \cos\theta|\bar{\psi}\rangle) + \cos\theta(\cos\theta|\psi\rangle - \sin\theta|\bar{\psi}\rangle) \\ &= -\sin^2\theta|\psi\rangle - \sin\theta\cos\theta|\bar{\psi}\rangle + \cos^2\theta|\psi\rangle - \cos\theta\sin\theta|\bar{\psi}\rangle \\ &= (-\sin^2\theta + \cos^2\theta)|\psi\rangle - 2\sin\theta\cos\theta|\bar{\psi}\rangle \\ &= \cos 2\theta|\psi\rangle - \sin 2\theta|\bar{\psi}\rangle \end{aligned}$$

## A geometric perspective on $G$

Then, the second reflection over  $|\psi\rangle$  yields the effect of the Grover iterator:

$$G|\psi\rangle = \cos 2\theta |\psi\rangle + \sin 2\theta |\bar{\psi}\rangle$$

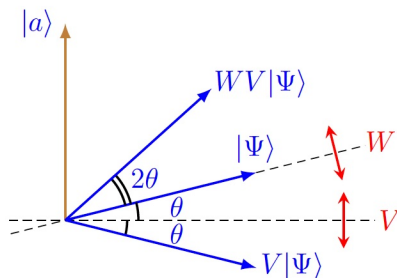
which boils down to a  $2\theta$  rotation:



## What's behind the scenes?

- The key is the selective shifting of the phase of one state of a quantum system, one that satisfies some condition, at each iteration.
- Performing a phase shift of  $\pi$  is equivalent to multiplying the amplitude of that state by  $-1$ : the amplitude for that state changes, but the probability of being in that state remains the same
- Subsequent transformations take advantage of that difference in amplitude to single out that state and increase the associated probability.
- This would **not be possible if the amplitudes were probabilities**, not holding extra information regarding the phase of the state in addition to the probability — it's a **quantum feature**.

# How many times should $G$ be applied?



From this picture, we may also conclude that the **angular distance to cover** towards an amplitude maximizing the probability of finding the correct solution is

$$\frac{\pi}{2} - \theta = \frac{\pi}{2} - \arcsin\left(\frac{1}{\sqrt{N}}\right)$$

## How many times should $G$ be applied?

Thus, the ideal number of iterations is

$$t = \left\lceil \frac{\frac{\pi}{2} - \arcsin \frac{1}{\sqrt{N}}}{2\theta} \right\rceil$$

A lower bound for  $\theta$  gives an upper bound for  $t$   
 — for  $N$  large  $\theta \approx \sin \theta = \frac{1}{\sqrt{N}}$ . Thus,

$$t = \frac{\frac{\pi\sqrt{N}-2}{2\sqrt{N}}}{\frac{2}{\sqrt{N}}} \approx \frac{\pi}{4}\sqrt{N}$$

So,  $G$  applied  $t$  times leaves the system within an angle  $\theta$  of  $|a\rangle$ . Then, a measurement in the computational basis yields the correct solution with probability

$$\|\langle a|G^t|\psi\rangle\| \geq \cos^2 \theta = 1 - \sin^2 \theta = \frac{N-1}{N}$$

which, for large  $N$ , is very close to 1.



## How many times should $G$ be applied?

For an **alternative computation**, recall

$$G|\psi\rangle = \cos 2\theta|\psi\rangle + \sin 2\theta|\bar{\psi}\rangle$$

By induction (prove it!), after  $k$  iterations,

$$\begin{aligned} G^k|\psi\rangle &= \cos(2k\theta)|\psi\rangle + \sin(2k\theta)|\bar{\psi}\rangle \\ &= \sin(2k+1)\theta|a\rangle + \cos(2k+1)\theta|r\rangle \end{aligned}$$

Thus, to maximize the probability of obtaining  $|a\rangle$ ,  $k$  is selected st

$$\sin((2k+1)\theta) \approx 1 \quad \text{i.e.} \quad (2k+1)\theta \approx \frac{\pi}{2}$$

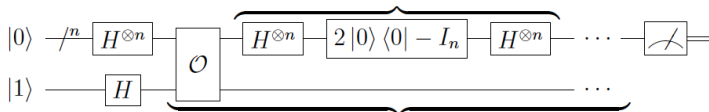
which leads to

$$k \approx \frac{\pi}{4\theta} - \frac{1}{2} \approx \frac{\pi}{4}\sqrt{N} \approx t$$

# Grover's algorithm ( $\mathcal{O}(\sqrt{N})$ )

Revisit our first slide:

- Prepare the initial state:  $|0\rangle^{\otimes n}|1\rangle$
- Apply  $H^{\otimes n} \otimes H$  to yield  $\frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle|-\rangle$
- Apply the Grover iterator  $G$  to  $\frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle|-\rangle$ ,  $t \approx \frac{\pi}{4}\sqrt{N}$  times, leading approximately to state  $|a\rangle|-\rangle$
- Measure the first  $n$  qubits to retrieve  $|a\rangle$



Execution time wrt (classical) exhaustive search:

from  $\mathcal{O}(N)$  to  $\mathcal{O}(\sqrt{N})$

## Multiple solutions

Assume there are  $M$  (out of  $2^n = N$ ) input strings evaluating to 0 by  $f$

$$|\psi\rangle = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle = \underbrace{\sqrt{\frac{M}{N}} |s\rangle}_{\text{solution}} + \underbrace{\sqrt{\frac{N-M}{N}} |r\rangle}_{\text{the rest}}$$

where

$$|s\rangle = \frac{1}{\sqrt{M}} \sum_{x \text{ solution}} |x\rangle \quad \text{and} \quad |r\rangle = \frac{1}{\sqrt{N-M}} \sum_{x \text{ no solution}} |x\rangle$$

## Multiple solutions

$$t = \left\lfloor \frac{\frac{\pi}{2} - \arcsin \sqrt{\frac{M}{N}}}{2\theta} \right\rfloor$$

which, for  $N$  large,  $M \ll N$  (thus  $\theta \approx \sin \theta$ ), yields

$$t \approx \frac{\pi}{4} \sqrt{\frac{N}{M}}$$

The probability to retrieve a correct solution is

$$\| \langle s | G^t | \psi \rangle \|^2 \geq \cos^2 \theta = 1 - \sin^2 \theta = \frac{N - M}{N}$$

which, for  $M = \frac{N}{2}$  yields  $\frac{1}{2}$ , but for  $M \ll N$ , is again close to 1.

## Multiple solutions

Computing the effect of  $G$ :  $2\theta$

$$\sin 2\theta = 2\sqrt{\frac{N-M}{N}} = 2\frac{\sqrt{M(N-M)}}{N}$$

$$2\theta = \arcsin\left(2\frac{\sqrt{M(N-M)}}{N}\right)$$

$M$ (out of 100)	$\arcsin \theta$
0	0
1	0.198
20	0.8
40	0.979
50	1
60	0.979
80	0.8
99	0.198
M	0

## Multiple solutions

Surprisingly, the rotation in each iteration decreases from  $M = \frac{N}{2}$  to  $N$ , and the number of iterations consequently increases, although one would expect to be easier to find a correct solution if their number increases!

**Solution: resort to draft paper!**

To double the number of elements in the search space, by adding  $N$  extra elements, none of which being a solution.

# The technique: Amplitude amplification

Grover's algorithm made use of

$$H^{\otimes n}|00 \dots 0\rangle$$

to prepare a **uniform** superposition of potential solutions.

In general, one may resort to any program  $K$  to map the solution space to any **superposition of guesses**, plus some extra qubits to be used as **draft paper**:

$$K|00 \dots 0\rangle = \sum_x \alpha_x |x\rangle |\text{draft}(x)\rangle$$

## The technique: Amplitude amplification

$$|\psi\rangle = \sum_{x \text{ solution}} \alpha_x |x\rangle |\text{draft}(x)\rangle + \sum_{x \text{ no solution}} \alpha_x |x\rangle |\text{draft}(x)\rangle$$

yielding the following probabilities:

$$p_s = \sum_{x \text{ solution}} \|\alpha_x\|^2 \quad \text{and} \quad p_{ns} = \sum_{x \text{ no solution}} \|\alpha_x\|^2 = 1 - p_s$$

Of course, amplification has no use if  $p_s \in \{0, 1\}$ .



## The technique: Amplitude amplification

Otherwise ( $0 < p_s < 1$ ), the amplitudes of **solution** inputs should be amplified.

First, express

$$|\psi\rangle = \sqrt{p_s}|\psi_s\rangle + \sqrt{p_{ns}}|\psi_{ns}\rangle$$

for the **normalised** components

$$|\psi_s\rangle = \sum_{x \text{ solution}} \frac{\alpha_x}{\sqrt{p_s}} |x\rangle |\text{draft}(x)\rangle$$

$$|\psi_{ns}\rangle = \sum_{x \text{ solution}} \frac{\alpha_x}{\sqrt{p_{ns}}} |x\rangle |\text{draft}(x)\rangle$$

which rewrites to

$$|\psi\rangle = \sin \theta |\psi_s\rangle + \cos \theta |\psi_{ns}\rangle$$

for  $\theta \in [0, \frac{\pi}{2}]$  such that  $\sin^2 \theta = p_s$ .

# The technique: Amplitude amplification

A generic **search iterator** is built as

$$S = K P K^{-1} V = W_K V$$

where

$$W_K |\psi\rangle = |\psi\rangle$$

$$W_K |\phi\rangle = -|\phi\rangle \quad \text{for all states orthogonal to } |\psi\rangle$$

The sets  $\{|\psi_s\rangle, |\psi_{ns}\rangle\}$  and  $\{|\psi\rangle, |\overline{\psi}\rangle\}$  are bases for the relevant 2-dimensional subspace.

## The technique: Amplitude amplification

As expected, starting in  $|\psi\rangle$ , the oracle produces

$$-\sin\theta|\psi_s\rangle + \cos\theta|\psi_{ns}\rangle = \cos(2\theta)|\psi\rangle - \sin(2\theta)|\bar{\psi}\rangle$$

which, followed by the amplifier, yields

$$\cos(2\theta)|\psi\rangle + \sin(2\theta)|\bar{\psi}\rangle$$

i.e. the effect of iterator  $S$  is

$$S|\psi\rangle = \cos(2\theta)|\psi\rangle + \sin(2\theta)|\bar{\psi}\rangle$$

which can be expressed in the basis  $\{|\psi_s\rangle, |\psi_{ns}\rangle\}$  as

$$S|\psi\rangle = \sin(3\theta)|\psi_s\rangle + \cos(3\theta)|\psi_{ns}\rangle$$

## The technique: Amplitude amplification

The repeated application of  $S$  a total of  $k$  times rotates the initial state  $|\psi\rangle$  to

$$S^k|\psi\rangle = \sin((2k+1)\theta)|\psi_s\rangle + \cos((2k+1)\theta)|\psi_{ns}\rangle$$

For the correct number of iterations, this procedure reaches a state such that a measurement will return an element of the subspace spanned by  $|\psi_s\rangle$  with a probability close to 1.

## The technique: Amplitude amplification

As before, to get that high probability, the smallest value for  $k$  one can choose is such that

$$(2k + 1)\theta \approx \frac{\pi}{2}$$

For a **small**  $\theta$ , as

$$\sin \theta = \sqrt{p_s} \approx \theta$$

the magnitude of the right number of iterations is

$$\mathcal{O}\left(\sqrt{\frac{1}{\theta}}\right)$$

because

$$(2k + 1)\sqrt{p_s} = \theta \Leftrightarrow k = \frac{\pi}{4\sqrt{p_s}} - \frac{1}{2}$$

## To follow

The algorithm requires that one knows **in advance** how many times iterator  **$S$**  is to be applied:

- For  $K = H$  (uniform sampling the input) this boils down to know the number of solutions of the search problem.
- For a generic  $K$  this amounts to know the probability with which  $K$  guesses a solution to the problem, i.e.  $\sin(\theta)$ .

To see ...

- **blind** search
- **estimate the amplitude** with which  $K$  maps  $|00 \cdots 0\rangle$  to the subspace of solutions