

# Quantum Computation

## Revisiting the quantum Fourier transform

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## Recap

The previous lecture discussed an algorithm to extract the phase factor  $w \in [0, 1[$  from a generic  $n$ -qubit quantum state. Writing  $w$  as  $\frac{x}{2^n}$ , for  $x$  an integer representable in  $n$  qubits, the estimation process was described by

$$\frac{1}{\sqrt{2^n}} \sum_{y \in 2^n} e^{2\pi i (\frac{x}{2^n})y} |y\rangle \rightsquigarrow |x\rangle$$

Its inverse is **QFT**, the **quantum Fourier transform**, a most useful routine in Quantum Computation.

Let us revisit its construction in a systematic way.

# The quantum Fourier transform

Essentially, the QFT performs a **change-of-basis** operation which encodes information of computational basis states in **local phases**.

For 1 qubit state this is exactly what the Hadamard gate accomplishes:

$$H|0\rangle = \frac{1}{\sqrt{2}}(|0\rangle + 1|1\rangle) \qquad H|1\rangle = \frac{1}{\sqrt{2}}(|0\rangle + (-1)|1\rangle)$$

## QFT: 1 qubit

Thus,  $QFT_1 = H$ :

$$QFT_1 |0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + 1 |1\rangle) \qquad QFT_1 |1\rangle = \frac{1}{\sqrt{2}} (|0\rangle + (-1) |1\rangle)$$

Operation  $H^{-1}$  allows to extract information encoded in local phases

↓  
=  $H$

## Exercise

Let  $\omega_1 = e^{i2\pi\frac{1}{2}}$ . Show that  $QFT_1 |x\rangle = \frac{1}{\sqrt{2}} (|0\rangle + \omega_1^{1 \cdot x} |1\rangle)$

↓  
angle of  $\pi$  radians

## QFT: 2 qubits

Let  $\omega_2 = e^{i2\pi \frac{1}{4}}$

$$QFT_2 |00\rangle = \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{2 \cdot 0} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{1 \cdot 0} |1\rangle)$$

$$QFT_2 |01\rangle = \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{2 \cdot 1} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{1 \cdot 1} |1\rangle)$$

$$QFT_2 |10\rangle = \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{2 \cdot 2} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{1 \cdot 2} |1\rangle)$$

$$QFT_2 |11\rangle = \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{2 \cdot 3} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{1 \cdot 3} |1\rangle)$$

### Exercise

Compute the phase coefficients in the expressions above and use Bloch sphere to study  $QFT_2 |x\rangle$ .

## QFT: 2 qubits

### Hint

$$\begin{array}{ll}
 \omega_2^{2.0} &= 1 \\
 \omega_2^{2.1} &= -1 \\
 \omega_2^{2.2} &= 1 \\
 \omega_2^{2.3} &= -1
 \end{array}
 \qquad
 \begin{array}{ll}
 \omega_2^{1.0} &= 1 \\
 \omega_2^{1.1} &= e^{i\frac{\pi}{2}} \\
 \omega_2^{1.2} &= -1 \\
 \omega_2^{1.3} &= e^{i\frac{3}{2}\pi}
 \end{array}$$

Note that

- previously, information on  $|x\rangle$  previously encoded by vectors pointing to the poles; becomes encoded by vectors in the **xz-plane**
- for every  **$\omega_2$ -rotation** on the second qubit there are **two** such rotations on the first qubit

## QFT: 2 qubits

In order to derive a circuit for  $QFT_2$ , compute

$$\begin{aligned}
 QFT_2 |x\rangle &= \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{2 \cdot x} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{1 \cdot x} |1\rangle) \\
 &= \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{2(2x_1+x_2)} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{2x_1+x_2} |1\rangle) \\
 &= \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{4x_1+2x_2} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{2x_1+x_2} |1\rangle) \\
 &= \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{4x_1} \omega_2^{2x_2} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{2x_1} \omega_2^{x_2} |1\rangle) \\
 &= \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{2x_2} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + \omega_2^{2x_1} \omega_2^{x_2} |1\rangle) \\
 &= \underbrace{\frac{1}{\sqrt{2}} (|0\rangle + (-1)^{x_2} |1\rangle)}_{H|x_2\rangle} \otimes \underbrace{\frac{1}{\sqrt{2}} (|0\rangle + (-1)^{x_1} \omega_2^{x_2} |1\rangle)}_{\text{some controlled rot. on } H|x_1\rangle}
 \end{aligned}$$

## QFT: 2 qubits

Define

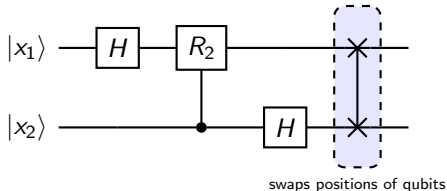
$$R_2 |0\rangle = |0\rangle; \quad \text{and} \quad R_2 |1\rangle = \omega_2 |1\rangle$$

Intuitively,  $R_2$  rotates a vector in the  $xz$ -plane  $\frac{\pi}{2}$  radians

It yields a **controlled**- $R_2$  operation by  $|x\rangle |0\rangle \mapsto |x\rangle |0\rangle$  and  $|x\rangle |1\rangle \mapsto R_2 |x\rangle |1\rangle$ . or, equivalently,

$$|0\rangle |x_2\rangle \mapsto |0\rangle |x_2\rangle \quad |1\rangle |x_2\rangle \mapsto \omega_2^{x_2} |1\rangle |x_2\rangle$$

Putting all pieces together to derive the QFT circuit for 2 qubits:





## QFT: 3 qubits

$$QFT_3 |x\rangle = \frac{1}{\sqrt{2}} (|0\rangle + \omega_3^{4 \cdot x} |1\rangle) \otimes (|0\rangle + \omega_3^{2 \cdot x} |1\rangle) \otimes (|0\rangle + \omega_3^{1 \cdot x} |1\rangle)$$

for  $\omega_n = e^{i2\pi \cdot \frac{1}{2^n}}$ .

**N.B.** In the sequel the normalisation factor  $\frac{1}{\sqrt{2}}$  will be dropped in each state to make notation easier on the eyes

## QFT: 3 qubits

In order to derive a circuit for  $QFT_3$ , we observe

$$\omega_n^2 = \omega_{n-1} \quad \text{and thus} \quad \omega_n^{2^{n-1}} = e^{i\pi} = -1$$

and recall that a binary number  $x_1 \dots x_n$  represents the natural number  $2^{n-1} \cdot x_1 + \dots + 2^0 \cdot x_n$ .

Thus, compute  $QFT_3$  as follows:

# QFT: 3 Qubits

$$\begin{aligned}
& QFT_3 |x\rangle \\
&= (|0\rangle + \omega_3^{4 \cdot x} |1\rangle) \otimes (|0\rangle + \omega_3^{2 \cdot x} |1\rangle) \otimes (|0\rangle + \omega_3^{1 \cdot x} |1\rangle) \\
&= (|0\rangle + (-1)^x |1\rangle) \otimes (|0\rangle + \omega_3^{2 \cdot x} |1\rangle) \otimes (|0\rangle + \omega_3^{1 \cdot x} |1\rangle) \\
&= (|0\rangle + (-1)^{x_3} |1\rangle) \otimes (|0\rangle + \omega_3^{2 \cdot x} |1\rangle) \otimes (|0\rangle + \omega_3^{1 \cdot x} |1\rangle) \\
&= H|x_3\rangle \otimes (|0\rangle + \omega_3^{2 \cdot (4x_1 + 2x_2 + x_3)} |1\rangle) \otimes (|0\rangle + \omega_3^{1 \cdot x} |1\rangle) \\
&= H|x_3\rangle \otimes (|0\rangle + \omega_3^{2 \cdot (4x_1 + 2x_2)} \omega_3^{2 \cdot x_3} |1\rangle) \otimes (|0\rangle + \omega_3^{1 \cdot x} |1\rangle) \\
&= H|x_3\rangle \otimes (|0\rangle + \omega_2^{2 \cdot (2x_1 + x_2)} \omega_2^{x_3} |1\rangle) \otimes (|0\rangle + \omega_3^{4x_1 + 2x_2 + x_3} |1\rangle) \\
&= H|x_3\rangle \otimes (|0\rangle + \omega_2^{2 \cdot (2x_1 + x_2)} \omega_2^{x_3} |1\rangle) \otimes (|0\rangle + \omega_3^{4x_1 + 2x_2} \omega_3^{x_3} |1\rangle) \\
&= H|x_3\rangle \otimes (|0\rangle + \omega_2^{2 \cdot (2x_1 + x_2)} \omega_2^{x_3} |1\rangle) \otimes (|0\rangle + \omega_3^{2 \cdot (2x_1 + x_2)} \omega_3^{x_3} |1\rangle) \\
&= H|x_3\rangle \otimes \underbrace{(|0\rangle + \omega_2^{2 \cdot (2x_1 + x_2)} \omega_2^{x_3} |1\rangle) \otimes (|0\rangle + \omega_2^{2x_1 + x_2} \omega_3^{x_3} |1\rangle)}_{\text{some controlled-rotations on } QFT_2|x_1x_2\rangle}
\end{aligned}$$

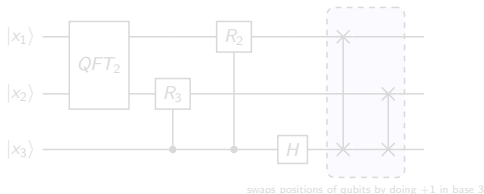
## QFT: 3 qubits

Take  $R_3 |0\rangle = |0\rangle$  and  $R_3 |1\rangle = \omega_3 |1\rangle$ . Intuitively,  $R_3$  rotates a vector in the  $xz$ -plane 'one 2<sup>3</sup>-th of the unit circle'.

It yields a **controlled- $R_3$**  operation defined by  $|x\rangle |0\rangle \mapsto |x\rangle |0\rangle$  and  $|x\rangle |1\rangle \mapsto R_3 |x\rangle |1\rangle$ . Equivalently

$$|0\rangle |y\rangle \mapsto |0\rangle |y\rangle \text{ and } |1\rangle |y\rangle \mapsto \omega_3^y |1\rangle |y\rangle$$

Putting all pieces together we derive the QFT circuit for 3 qubits



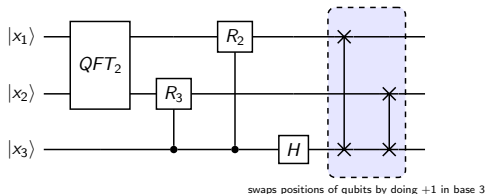
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Putting all pieces together we derive the QFT circuit for 3 qubits



## QFT: $n$ qubits

Calculation easily extends to  $QFT_n$  (*in lieu* of  $QFT_3$ ) :

Let  $\omega_n = e^{i2\pi \cdot \frac{1}{2^n}}$  (division of the **unit circle** in  $2^n$  slices)

$$QFT_n |x\rangle = \frac{1}{\sqrt{2}} (|0\rangle + \omega_n^{2^{n-1} \cdot x} |1\rangle) \otimes \cdots \otimes (|0\rangle + \omega_n^{2^0 \cdot x} |1\rangle)$$

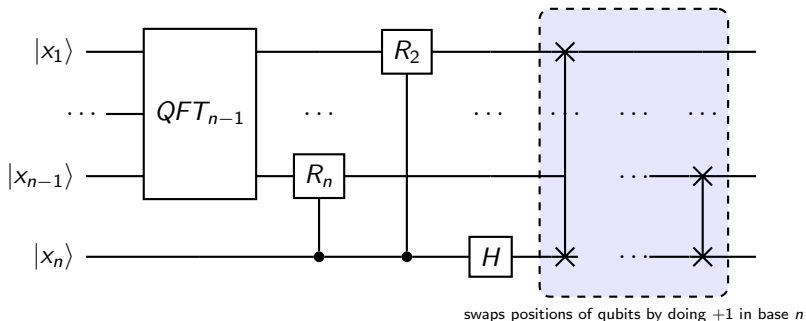
Take  $R_n |0\rangle = |0\rangle$  and  $R_n |1\rangle = \omega_n |1\rangle$ . Intuitively,  $R_n$  rotates a vector in the  $xz$ -plane '**one  $2^n$ -th** of the unit circle'

It yields a **controlled- $R_n$**  operation defined by  $|x\rangle |0\rangle \mapsto |x\rangle |0\rangle$  and  $|x\rangle |1\rangle \mapsto R_n |x\rangle |1\rangle$ . Equivalently

$$|0\rangle |y\rangle \mapsto |0\rangle |y\rangle \quad \text{and} \quad |1\rangle |y\rangle \mapsto \omega_n^y |1\rangle |y\rangle$$

## QFT: $n$ qubits

This suggests a recursive definition for the general  $QFT$  circuit:



## An equivalent formulation of QFT

Although we have been working with

$$QFT_n |x\rangle = \frac{1}{\sqrt{2}}(|0\rangle + \omega_n^{2^{n-1} \cdot x} |1\rangle) \otimes \cdots \otimes \frac{1}{\sqrt{2}}(|0\rangle + \omega_n^{1 \cdot x} |1\rangle)$$

we are already familiar with an equivalent, useful definition

$$QFT_n |x\rangle = \frac{1}{\sqrt{2^n}} \sum_{k=0}^{2^n-1} \omega_n^{x \cdot k} |k\rangle$$

Examples with  $n = 1$  and  $n = 2$

$$QFT_1 |x\rangle = \frac{1}{\sqrt{2}}(|0\rangle + \omega_1^x |1\rangle)$$

$$QFT_2 |x\rangle = \frac{1}{\sqrt{2^2}}(|00\rangle + \omega_2^x |01\rangle + \omega_2^{2 \cdot x} |10\rangle + \omega_2^{3 \cdot x} |11\rangle)$$