# Quantum Computation Quantum phase estimation and the quantum Fourier transform

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# Encoding information in phases

In several quantum algorithms information is encoded in the relative phases of a quantum state.

The effect of Hadamard (once again)

$$\begin{aligned} H|x\rangle &= \frac{1}{\sqrt{2}}(|0\rangle + (-1)^{x}|1\rangle) &= \frac{1}{\sqrt{2}}\sum_{y\in 2}(-1)^{xy}|y\rangle \\ H^{\otimes n}|x\rangle &= \frac{1}{\sqrt{2^{n}}}\sum_{y\in 2^{n}}(-1)^{x\cdot y}|y\rangle \end{aligned}$$

is to encode information about the value of x into the phases  $(-1)^{x \cdot y}$  of basis states  $|y\rangle$ .

# Encoding information in phases

Of course, as a reversible gate, the Hadamard gate also decodes information from phases:

$$H^{\otimes n} \frac{1}{\sqrt{2^n}} \sum_{y \in 2^n} (-1)^{x \cdot y} |y\rangle = H^{\otimes n} (H^{\otimes n} |x\rangle)$$

$$= (H^{\otimes n} H^{\otimes n}) |x\rangle$$

$$= I|x\rangle$$

$$= |x\rangle$$

# Encoding information in phases

In general, phases are complex numbers

 $e^{2\pi i w}$ 

for any real  $w \in [0, 1[$ .

Of course,  $H^{\otimes n}$  cannot encode/decode information over such generic phases. The general situation can be described as follows:

## The phase estimation problem

Determine a good estimation of the phase parameter  $\boldsymbol{w}$  given a general quantum state

$$\frac{1}{\sqrt{2^n}} \sum_{y \in 2^n} e^{2\pi i w y} |y\rangle$$

# An algorithm for phase estimation

#### Notation

$$\mathbf{w} = 0.x_1x_2\cdots$$

is written in base 2 (i.e.  $w = x_1 2^{-1} + x_2 2^{-2} + \cdots$ ); thus

$$2^k w = x_1 x_2 \cdots x_k \cdot x_{k+1} x_{k+2} \cdots$$

and

$$\begin{array}{ll} e^{2\pi i(2^k w)} &=& e^{2\pi i(x_1 x_2 \cdots x_k \cdot x_{k+1} x_{k+2} \cdots)} \\ &=& e^{2\pi i(x_1 x_2 \cdots x_k)} e^{2\pi i(0 \cdot x_{k+1} x_{k+2} \cdots)} \\ &=& e^{2\pi i(0 \cdot x_{k+1} x_{k+2} \cdots)} \end{array}$$

because  $e^{2\pi iz} = 1$  for any integer z.

# Case A: 1-qubit state and $w = 0.x_1$

$$\begin{split} \frac{1}{\sqrt{2}} \sum_{y \in 2} e^{2\pi i (\mathbf{0}.\mathbf{x}_1)y} |y\rangle &= \frac{1}{\sqrt{2}} \sum_{y \in 2} e^{2\pi i (\frac{\mathbf{x}_1}{2})y} |y\rangle \\ &= \frac{1}{\sqrt{2}} \sum_{y \in 2} e^{\pi i (\mathbf{x}_1 y)} |y\rangle \\ &= \frac{1}{\sqrt{2}} \sum_{y \in 2} (-1)^{\mathbf{x}_1 y} |y\rangle \\ &= \frac{1}{\sqrt{2}} (|0\rangle + (-1)^{\mathbf{x}_1} |1\rangle) \end{split}$$

Clearly H will decode and retrieve  $x_1$  because

$$H\left(\frac{1}{\sqrt{2}}(|0\rangle + (-1)^{\mathsf{x}_1}|1\rangle)\right) = |\mathsf{x}_1\rangle$$

# Case B: 2-qubit state and $w = 0.x_1x_2$

#### Observe that

$$\frac{1}{\sqrt{2^2}} \sum_{y \in 2^2} e^{2\pi i (\mathbf{0}.\mathbf{x}_1 \mathbf{x}_2) y} |y\rangle \ = \ \left(\frac{|0\rangle + e^{2\pi i (\mathbf{0}.\mathbf{x}_2)} |1\rangle}{\sqrt{2}}\right) \otimes \left(\frac{|0\rangle + e^{2\pi i (\mathbf{0}.\mathbf{x}_1 \mathbf{x}_2)} |1\rangle}{\sqrt{2}}\right)$$

which means that  $x_2$ , but not  $x_1$ , can be retrieved from the first qubit through an application of H.

## The phase rotator

$$R_2 = \begin{bmatrix} 1 & 0 \\ 0 & e^{\frac{2\pi i}{4}} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & e^{2\pi i(0.01)} \end{bmatrix}$$

where 0.01 is in base 2 (thus, equal to  $2^{-2}$ ).

# Case B: 2-qubit state and $w = 0.x_1x_2$

Taking  $x_2 = 1$  and applying the inverse of the phase rotator to the second qubit, yields

$$R_{2}^{-1} \begin{pmatrix} \frac{|0\rangle + e^{2\pi i(0.x_{1}1)}|1\rangle}{\sqrt{2}} \end{pmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & e^{-2\pi i(0.01)} \end{bmatrix} \begin{pmatrix} \frac{|0\rangle + e^{2\pi i(0.x_{1}1)}|1\rangle}{\sqrt{2}} \end{pmatrix}$$
$$= \frac{|0\rangle + e^{2\pi i(0.x_{1}1 - 0.01)}|1\rangle}{\sqrt{2}}$$
$$= \frac{|0\rangle + e^{2\pi i(0.x_{1}1)|1\rangle}}{\sqrt{2}}$$

## Concluding

- $x_1$  can now be determined by an application of H, as before.
- Moreover, the decision to apply R before the application of H depends on x<sub>2</sub> being 1 or 0, respectively.
- Thus, to find  $w = 0.x_1x_2$  it is enough to apply a controlled version of R, precisely controlled by the state of the first qubit.



# Case B: 2-qubit state and $w = 0.x_1x_2$

#### The circuit

$$\begin{array}{c|c} \frac{1}{\sqrt{2}} \left( |0\rangle + e^{2\pi i (0.x_2)} |1\rangle \right) & \hline & H \\ \hline \\ \frac{1}{\sqrt{2}} \left( |0\rangle + e^{2\pi i (0.x_1x_2)} |1\rangle \right) & \hline & R_2 \\ \hline \\ |x_2\rangle \left( \frac{|0\rangle + e^{2\pi i (0.x_1x_2)} |1\rangle}{\sqrt{2}} \right) & |x_2\rangle \left( \frac{|0\rangle + e^{2\pi i (0.x_1)} |1\rangle}{\sqrt{2}} \right) \end{array}$$

#### The state is now

$$\begin{split} &\frac{1}{\sqrt{2^3}} \sum_{y \in 2^3} e^{2\pi i (\mathbf{0}.\mathbf{x}_1 \mathbf{x}_2 \mathbf{x}_3) y} |y\rangle &= \\ &= \left( \frac{|0\rangle + e^{2\pi i (\mathbf{0}.\mathbf{x}_3)} |1\rangle}{\sqrt{2}} \right) \otimes \left( \frac{|0\rangle + e^{2\pi i (\mathbf{0}.\mathbf{x}_2 \mathbf{x}_3)} |1\rangle}{\sqrt{2}} \right) \otimes \left( \frac{|0\rangle + e^{2\pi i (\mathbf{0}.\mathbf{x}_1 \mathbf{x}_2 \mathbf{x}_3)} |1\rangle}{\sqrt{2}} \right) \end{split}$$

In this case the third qubit has to conditionally rotate both  $x_2$  and  $x_3$ , leading to the following circuit

$$\frac{1}{\sqrt{2}} \left( |0\rangle + e^{2\pi i (0.x_3)} |1\rangle \right) - H - |x_3\rangle$$

$$\frac{1}{\sqrt{2}} \left( |0\rangle + e^{2\pi i (0.x_2x_3)} |1\rangle \right) - |x_2\rangle$$

$$\frac{1}{\sqrt{2}} \left( |0\rangle + e^{2\pi i (0.x_1x_2x_3)} |1\rangle \right) - |x_1\rangle$$

# Going generic

Gate  $R_3$  in the circuit is an instance of a 1-qubit phase rotator

$$R_k = \begin{bmatrix} 1 & 0 \\ 0 & e^{\frac{2\pi i}{2^k}} \end{bmatrix}$$

whose inverse acts as

$$R_k^{-1}|0\rangle = |0\rangle$$
  

$$R_k^{-1}|1\rangle = e^{-2\pi i (0.0 \cdots 1)}|1\rangle$$

with 1 in  $0.0 \cdots 1$  appearing in position k.

# Going generic

The output state of the circuit is

$$|x_3x_2x_1\rangle$$

Thus, relabelling the qubits in reverse order, this provides an efficient circuit to estimate the phase (actually, to give a totally accurate estimation ...), by computing

$$\frac{1}{\sqrt{2^n}} \sum_{y \in 2^n} e^{2\pi i (\frac{x}{2^n}) y} |y\rangle \quad \rightsquigarrow \quad |x\rangle$$

# Inverting ...

The inverse of the phase estimation transformation computes

$$|x\rangle \quad \leadsto \quad \frac{1}{\sqrt{2^n}} \sum_{y \in 2^n} e^{2\pi i (\frac{x}{2^n})y} |y\rangle$$

which is obtained by taking the inverses of each gate and building the circuit in reverse order.

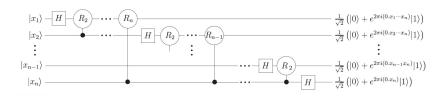
The result is formally identical to the discrete Fourier transform.

# The quantum Fourier transform

QFT on basis states  $|0\rangle, |1\rangle \cdots |k-1\rangle$ 

$$QFT_k(|x\rangle) = \frac{1}{\sqrt{k}} \sum_{y=0}^{k-1} e^{2\pi i (\frac{x}{k})y} |y\rangle$$

#### The circuit



## Complexity (number of gates)

- one H plus n-1 conditional rotations on the first qubit
- one H plus n-2 conditional rotations on the second qubit
- ..

$$n + (n-1) + (n-2) + \cdots + 1 = \frac{n(n-1)}{2}$$

• plus  $\frac{n}{2}$  swaps (each implemented by 3 CNOT gates)

Thus

$$\frac{n(n-1)}{2} + 3x\frac{n}{2} = \frac{n^2 + 2n}{2} \approx \mathcal{O}(n^2)$$

# The quantum Fourier transform

## Complexity (number of gates)

$$\frac{n(n-1)}{2} + 3x\frac{n}{2} = \frac{n^2 + 2n}{2} \approx \mathcal{O}(n^2)$$

which compares to the classical case for the Fast FT:  $O(n2^n)$ 

The result is impressive: the quantum version requires exponentially less operations to compute the Fourier transform than the (best) classical one.

- However, typical uses (e.g. in speech recognition) are limited by the impossibility of directly measuring the Fourier transformed amplitudes of the original state.
- This requires a subtler use of QFT in practice: the phase estimation procedure, underlying many quantum algorithms, is one of them.

- The circuit for  $QFT_k$  computes the QFT for k a power of 2, i.e  $k=2^n$
- The phase estimation algorithm works only when the phase is of the form  $w = 0.x_1x_2\cdots x_n$ , i.e.  $\frac{x}{2n}$  for some integer x

However, it can be shown that, for an arbitrary w, the algorithm will compute x such that  $\frac{x}{2n}$  is closest to w with high probability.

## The question

What is the error emerging when w is not an integer multiple of  $\frac{1}{2n}$ ?

 $QFT^{-1}$  computes some superposition

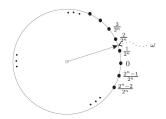
$$\sum_{x} \alpha_{x}(\mathbf{w})|x\rangle$$

which represents the values of x that once measured gives a good estimate of w, outputing x with probability  $|\alpha_x(w)|^2$ .

This output x corresponds to an estimate

$$\tilde{w} = \frac{x}{2^n}$$

Consider w an integer not multiple of  $\frac{1}{2^n}$ , and let  $\hat{w}$  be the nearest integer multiple of  $\frac{1}{2^n}$  to w, i.e.  $\hat{w} = \frac{\hat{x}}{2^n}$  is the closest number of this form to w.



## Theorem

The phase estimation algorithm returns  $\hat{\chi}$  with probability at least  $\frac{4}{\pi^2}$ , i.e. the algorithm outputs an estimate  $\hat{\chi}$  with the given probability such that

$$\left|\frac{\hat{x}}{2^n} - \mathbf{w}\right| \leq \frac{1}{2^{n+1}}$$

#### **Theorem**

If 
$$\frac{x}{2^n} \le \mathbf{w} \le \frac{x+1}{2^n}$$

The phase estimation algorithm returns either x or x+1 with probability at least  $\frac{8}{\pi^2}$  i.e. the algorithm outputs an estimate  $\hat{x}$  with the given probability such that

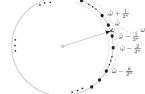
$$\left|\frac{\hat{\chi}}{2^n} - \mathbf{w}\right| = \frac{1}{2^n}$$

# The reverse question

How many qubits are required to get w accurate to n bits, with a probability p below a certain level?

Actually, the crucial choice is the value of n (number of qubits used) to ensure the estimation is close enough.

For  $p=1-\frac{1}{2(k-1)}$ , the algorithm returns one of the 2k closest integer multiples of  $\frac{1}{2n}$ , i.e.



which means that  $|w - \hat{w}| \leq \frac{k}{2^n}$ .

# The reverse question

Thus, to estimate  $\hat{w}$  such that  $|w - \hat{w}| \leq \frac{1}{2^r}$  with probability at least

$$1-\frac{1}{2^m}$$

the maximum number of qubits required is

$$n = r + m + 1$$

• In practice a much smaller error is obtained: for example, with probability at least  $\frac{8}{\pi^2}$ , the error will be at most

$$\frac{1}{2^{r+m}}$$

## **Exercises**

Recall the definition of QFT on K basis states:

$$QFT_{K}(|x\rangle) = \frac{1}{\sqrt{K}} \sum_{y=0}^{K-1} e^{2\pi i (\frac{x}{K})y} |y\rangle$$

#### Exercise 1

Compute  $QFT_K(|00\cdots 0\rangle)$ .

#### Exercise 2

Verify the following equality, used in the slides but not proved.

$$\begin{aligned} QFT_{\mathcal{K}}(|x_1\cdots x_n\rangle) &= \\ &\left(\frac{|0\rangle + e^{2\pi i(0.x_n)}|1\rangle}{\sqrt{2}}\right) \otimes \left(\frac{|0\rangle + e^{2\pi i(0.x_nx_{n-1})}|1\rangle}{\sqrt{2}}\right) \cdots \otimes \cdots \left(\frac{|0\rangle + e^{2\pi i(0.x_1x_2\cdots x_n)}|1\rangle}{\sqrt{2}}\right) \end{aligned}$$

## **Exercises**

Hint to Exercise 2: The case of  $QFT_4$  applied to  $|x\rangle = |x_1x_2\rangle$ 

$$QFT_4(|x\rangle) = \frac{1}{2} \sum_{y=0}^{3} e^{2\pi i x y 2^{-2}} |y\rangle$$
$$= \frac{1}{2} \sum_{y_1, y_2=0}^{1} e^{2\pi i x (y_1 2^{-1} + y_2 2^{-2})} |y_1 y_2\rangle$$

because, for  $|y\rangle = |y_1y_2\rangle$ ,

$$\frac{y}{2^n} = \sum_{i=1}^n y_i 2^{-i}$$

## Exercises

# Hint to Exercise 2: The case of $QFT_4$ applied to $|x\rangle = |x_1x_2\rangle$

$$\begin{split} \cdots &= \frac{1}{2} \sum_{y_1, y_2 = 0}^{1} (e^{2\pi i x y_1 2^{-1}} | y_1 \rangle \otimes e^{2\pi i x y_2 2^{-2}} | y_2 \rangle) \\ &= \frac{1}{2} \sum_{y_1 = 0}^{1} (e^{2\pi i x y_1 2^{-1}} | y_1 \rangle \otimes \sum_{y_2 = 0}^{1} e^{2\pi i x y_2 2^{-2}} | y_2 \rangle) \\ &= \frac{(|0\rangle + e^{2\pi i x 2^{-1} | 1 \rangle})}{\sqrt{2}} \otimes \frac{(|0\rangle + e^{2\pi i x 2^{-2} | 1 \rangle})}{\sqrt{2}} \\ &= \frac{(|0\rangle + e^{2\pi i (x_1, x_2) | 1 \rangle})}{\sqrt{2}} \otimes \frac{(|0\rangle + e^{2\pi i (0, x_1 x_2) | 1 \rangle})}{\sqrt{2}} \\ &= \frac{(|0\rangle + e^{2\pi i (0, x_2) | 1 \rangle})}{\sqrt{2}} \otimes \frac{(|0\rangle + e^{2\pi i (0, x_1 x_2) | 1 \rangle})}{\sqrt{2}} \end{split}$$

because,  $e^{2\pi i(a.b)} = e^{2\pi ia}e^{2\pi i(0.b)} = e^{2\pi i(0.b)}$ 

