# **Quantum Computation**Revisiting the quantum Fourier transform

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Recalling the basic idea

The previous lecture discussed an algorithm to extract the phase factor  $w \in [0,1[$  from a generic *n*-qubit quantum state. Writing w as  $\frac{x}{2^n}$ , for x an integer representable in n qubits, the estimation process was described bγ

$$\frac{1}{\sqrt{2^n}} \sum_{y \in 2^n} e^{2\pi i (\frac{x}{2^n}) y} |y\rangle \quad \rightsquigarrow \quad |x\rangle$$

Its inverse is QFT, the quantum Fourier transform, a most useful routine in Quantum Computation.

Let us revisit its construction in a systematic way.

Recalling the basic idea

## The quantum Fourier transform

Essentially, the QFT performs a change-of-basis operation which encodes information of computational basis states in local phases.

For 1 qubit state this is exactly what the Hadamard gate accomplishes:

$$H\ket{0} = \frac{1}{\sqrt{2}} \left( \ket{0} + \frac{1}{1} \ket{1} \right)$$
  $H\ket{1} = \frac{1}{\sqrt{2}} \left( \ket{0} + \frac{(-1)}{1} \ket{1} \right)$ 

Thus,  $QFT_1 = H$ :

Recalling the basic idea

$$\textit{QFT}_1 \left| 0 \right\rangle = \tfrac{1}{\sqrt{2}} \big( \left| 0 \right\rangle + \tfrac{1}{1} \left| 1 \right\rangle \big) \qquad \qquad \textit{QFT}_1 \left| 1 \right\rangle = \tfrac{1}{\sqrt{2}} \big( \left| 0 \right\rangle + \tfrac{(-1)}{1} \left| 1 \right\rangle \big)$$

Operation  $H^{-1}$  allows to extract information encoded in local phases



#### Exercise

Let  $\omega_1=e^{i2\pi\frac{1}{2}}.$  Show that  $QFT_1\ket{x}=\frac{1}{\sqrt{2}}\Big(\ket{0}+\omega_1^{1.\mathsf{x}}\ket{1}\Big)$ 



angle of  $\pi$  radians

Note that  $\omega_1$  represents a rotation of  $\pi$  radians, diving the unit circle into two slices...

Actually, the two  $2^{th}$ -roots of the identity are

$$\omega_1^0=1$$
 and  $\omega_1^1=e^{rac{i2\pi}{2}}=e^{i\pi}=-1$ 

Also note that

Recalling the basic idea

$$\omega_1^{1.x} = e^{\frac{i2\pi x}{2}} = e^{i2\pi \frac{x}{2}} = e^{i2\pi(0.x)}$$

as used in the previous lecture.

Let 
$$\omega_2 = e^{i2\pi \frac{1}{4}}$$

$$\begin{split} QFT_2 \left| 00 \right\rangle &= \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{2 \cdot 0} \left| 1 \right\rangle \right) \otimes \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{1 \cdot 0} \left| 1 \right\rangle \right) \\ QFT_2 \left| 01 \right\rangle &= \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{2 \cdot 1} \left| 1 \right\rangle \right) \otimes \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{1 \cdot 1} \left| 1 \right\rangle \right) \\ QFT_2 \left| 10 \right\rangle &= \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{2 \cdot 2} \left| 1 \right\rangle \right) \otimes \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{1 \cdot 2} \left| 1 \right\rangle \right) \\ QFT_2 \left| 11 \right\rangle &= \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{2 \cdot 3} \left| 1 \right\rangle \right) \otimes \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{1 \cdot 3} \left| 1 \right\rangle \right) \end{split}$$

In general

$$\mathit{QFT}_2\ket{\mathbf{x}} = rac{1}{\sqrt{2}}ig(\ket{0} + \omega_2^{2\cdot\mathbf{x}}\ket{1}ig)\otimes rac{1}{\sqrt{2}}ig(\ket{0} + \omega_2^{1\cdot\mathbf{x}}\ket{1}ig)$$

#### Exercise

Show that, for  $\mathbf{x} = |x_1x_2\rangle$ ,  $QFT_2|x\rangle$  can be written as

$$extit{QFT}_2 \ket{x} = rac{1}{\sqrt{2}} ig( \ket{0} + e^{i2\pi(0.x1)} \ket{1} ig) \otimes rac{1}{\sqrt{2}} ig( \ket{0} + e^{i2\pi(0.x_1x_2)} \ket{1} ig)$$

#### Exercise

Compute the phase coeficients in the expressions above and use Bloch sphere to study  $QFT_2|x\rangle$ .

#### Hint

$$\begin{array}{lllll} \omega_2^{2.0} & = & 1 & & \omega_2^{1.0} & = & 1 \\ \omega_2^{2.1} & = & -1 & & \omega_2^{1.1} & = & e^{i\frac{\pi}{2}} \\ \omega_2^{2.2} & = & 1 & & \omega_2^{1.2} & = & -1 \\ \omega_2^{2.3} & = & -1 & & \omega_2^{1.3} & = & e^{i\frac{3}{2}\pi} \end{array}$$

#### Note that

- for every  $\omega_2$ -rotation on the second qubit there are *two* such rotations on the first qubit
- $\omega_2^2 = \omega_1$ , or, in general,  $\omega_n^2 = \omega_{n-1}$

In order to derive a circuit for  $QFT_2$ , compute

$$\begin{split} QFT_2 \left| x \right\rangle &= \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{2 \cdot x} \left| 1 \right\rangle \right) \otimes \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{1 \cdot x} \left| 1 \right\rangle \right) \\ &= \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{2(2x_1 + x_2)} \left| 1 \right\rangle \right) \otimes \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{2x_1 + x_2} \left| 1 \right\rangle \right) \\ &= \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{4x_1 + 2x_2} \left| 1 \right\rangle \right) \otimes \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{2x_1 + x_2} \left| 1 \right\rangle \right) \\ &= \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{4x_1} \omega_2^{2x_2} \left| 1 \right\rangle \right) \otimes \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{2x_1} \omega_2^{x_2} \left| 1 \right\rangle \right) \\ &= \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{2x_2} \left| 1 \right\rangle \right) \otimes \frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \omega_2^{2x_1} \omega_2^{x_2} \left| 1 \right\rangle \right) \\ &= \underbrace{\frac{1}{\sqrt{2}} \left( \left| 0 \right\rangle + \left( -1 \right)^{x_2} \left| 1 \right\rangle \right)}_{\text{Some controlled rot. on } H \left| x1 \right\rangle}_{\text{some controlled rot. on } H \left| x1 \right\rangle} \end{split}$$

Define

$$R_2 \ket{0} = \ket{0}$$
 and  $R_2 \ket{1} = \omega_2 \ket{1}$ 

which rotates a vector in the xz-plane  $\frac{\pi}{2}$  radians

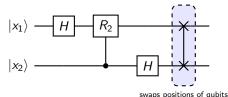
It yields a controlled- $R_2$  operation

$$|x\rangle |0\rangle \mapsto |x\rangle |0\rangle \qquad |x\rangle |1\rangle \mapsto R_2 |x\rangle |1\rangle$$

or, equivalently,

$$|0\rangle |y\rangle \mapsto |0\rangle |y\rangle \qquad |1\rangle |y\rangle \mapsto \omega_2^{\mathsf{y}} |1\rangle |y\rangle$$

Putting all pieces together to derive the QFT circuit for 2 qubits:



QFT on 3 qubits

$$\mathit{QFT}_3 \left| \mathbf{x} \right\rangle = \tfrac{1}{\sqrt{2}} \left( \left. \left| 0 \right\rangle + \omega_3^{4 \cdot \mathbf{x}} \left| 1 \right\rangle \right. \right) \otimes \left( \left. \left| 0 \right\rangle + \omega_3^{2 \cdot \mathbf{x}} \left| 1 \right\rangle \right. \right) \otimes \left( \left. \left| 0 \right\rangle + \omega_3^{1 \cdot \mathbf{x}} \left| 1 \right\rangle \right. \right)$$

for  $\omega_n = e^{i2\pi \cdot \frac{1}{2^n}}$ .

#### N.B.

In the sequel the normalisation factor  $\frac{1}{\sqrt{2}}$  will be dropped in each state to increase readability

QFT on 3 qubits

Recalling that a binary number  $x_1 \dots x_n$  represents the natural number

$$2^{n-1} \cdot x_1 + \cdots + 2^0 \cdot x_n$$

and that

$$\omega_n^2 = \omega_{n-1}$$
 and thus  $\omega_n^{2^{n-1}} = e^{i\pi} = -1$ 

 $QFT_3$ 

## QFT: 3 Qubits

$$\begin{aligned} &QFT_{3} \left| x \right\rangle \\ &= \left( \left| 0 \right\rangle + \omega_{3}^{4 \cdot x} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot x} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{1 \cdot x} \left| 1 \right\rangle \right) \\ &= \left( \left| 0 \right\rangle + \left( -1 \right)^{x} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot x} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{1 \cdot x} \left| 1 \right\rangle \right) \\ &= \left( \left| 0 \right\rangle + \left( -1 \right)^{x_{3}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot x} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{1 \cdot x} \left| 1 \right\rangle \right) \\ &= \left( \left| 0 \right\rangle + \left( -1 \right)^{x_{3}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot x} \left| 1 \right\rangle \right) \\ &= H \left| x_{3} \right\rangle \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot \left( 4x_{1} + 2x_{2} + x_{3} \right)} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{1 \cdot x} \left| 1 \right\rangle \right) \\ &= H \left| x_{3} \right\rangle \otimes \left( \left| 0 \right\rangle + \omega_{2}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{2}^{x_{3}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{4x_{1} + 2x_{2} + x_{3}} \left| 1 \right\rangle \right) \\ &= H \left| x_{3} \right\rangle \otimes \left( \left| 0 \right\rangle + \omega_{2}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{2}^{x_{3}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{4x_{1} + 2x_{2}} \omega_{3}^{x_{3}} \left| 1 \right\rangle \right) \\ &= H \left| x_{3} \right\rangle \otimes \left( \left| 0 \right\rangle + \omega_{2}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{2}^{x_{3}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{3}^{x_{3}} \left| 1 \right\rangle \right) \\ &= H \left| x_{3} \right\rangle \otimes \left( \left| 0 \right\rangle + \omega_{2}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{2}^{x_{3}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{3}^{x_{3}} \left| 1 \right\rangle \right) \\ &= H \left| x_{3} \right\rangle \otimes \left( \left| 0 \right\rangle + \omega_{2}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{2}^{x_{3}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{3}^{x_{3}} \left| 1 \right\rangle \right) \\ &= H \left| x_{3} \right\rangle \otimes \left( \left| 0 \right\rangle + \omega_{2}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{2}^{x_{3}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{3}^{x_{3}} \left| 1 \right\rangle \right) \\ &= H \left| x_{3} \right\rangle \otimes \left( \left| 0 \right\rangle + \omega_{2}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{2}^{x_{3}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{3}^{x_{3}} \left| 1 \right\rangle \right) \\ &= H \left| x_{3} \right\rangle \otimes \left( \left| 0 \right\rangle + \omega_{2}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{2}^{x_{3}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{3}^{x_{3}} \left| 1 \right\rangle \right) \\ &= H \left| x_{3} \right\rangle \otimes \left( \left| 0 \right\rangle + \omega_{2}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{2}^{x_{3}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + \omega_{3}^{2 \cdot \left( 2x_{1} + x_{2} \right)} \omega_{3}$$

Take  $R_3 |0\rangle = |0\rangle$  and  $R_3 |1\rangle = \omega_3 |1\rangle$ . Intuitively,  $R_3$  rotates a vector in the xz-plane 'one 2³-th of the unit circle'. It yields a controlled- $R_3$  operation defined by

$$|x\rangle |0\rangle \mapsto |x\rangle |0\rangle$$
 and  $|x\rangle |1\rangle \mapsto R_3 |x\rangle |1\rangle$ 

#### Equivalently

$$|0\rangle |y\rangle \mapsto |0\rangle |y\rangle$$
 and  $|1\rangle |y\rangle \mapsto \omega_3^{y} |1\rangle |y\rangle$ 

Putting all pieces together we derive the QFT circuit for 3 qubits



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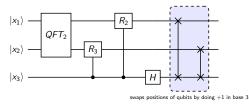
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$$|0\rangle |y\rangle \mapsto |0\rangle |y\rangle$$
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Putting all pieces together we derive the QFT circuit for 3 qubits



Calculation easily extends to  $QFT_n$  (in lieu of  $QFT_3$ ):

Let  $\omega_n = e^{i2\pi \cdot \frac{1}{2^n}}$  (division of the unit circle in  $2^n$  slices)

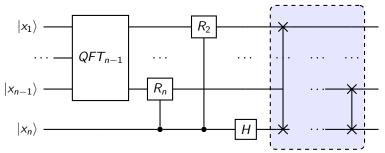
$$QFT_{n}|\mathbf{x}\rangle = \frac{1}{\sqrt{2}}(|0\rangle + \omega_{n}^{2^{n-1} \cdot \mathbf{x}}|1\rangle) \otimes \cdots \otimes (|0\rangle + \omega_{n}^{2^{0} \cdot \mathbf{x}}|1\rangle)$$

Take  $R_n|0\rangle = |0\rangle$  and  $R_n|1\rangle = \omega_n|1\rangle$ . Intuitively,  $R_n$  rotates a vector in the xz-plane 'one  $2^n$ -th of the unit circle'

It yields a controlled- $R_n$  operation defined by  $|x\rangle |0\rangle \mapsto |x\rangle |0\rangle$  and  $|x\rangle |1\rangle \mapsto R_n |x\rangle |1\rangle$ . Equivalently

$$|0\rangle |y\rangle \mapsto |0\rangle |y\rangle$$
 and  $|1\rangle |y\rangle \mapsto \omega_n^y |1\rangle |y\rangle$ 

This suggests a recursive definition for the general QFT circuit:



swaps positions of qubits by doing +1 in base n

## An equivalent formulation of QFT

Although we have been working with

$$QFT_n|x\rangle = \frac{1}{\sqrt{2}}(|0\rangle + \omega_n^{2^{n-1}\cdot x}|1\rangle) \otimes \cdots \otimes \frac{1}{\sqrt{2}}(|0\rangle + \omega_n^{1\cdot x}|1\rangle)$$

we are already familiar with an equivalent, useful definition

$$QFT_n |x\rangle = \frac{1}{\sqrt{2^n}} \sum_{k=0}^{2^n-1} \omega_n^{k \cdot x} |k\rangle$$

Examples with n=1 and n=2

$$\begin{split} & \textit{QFT}_1 \left| x \right\rangle = \frac{1}{\sqrt{2}} \big( \left| 0 \right\rangle + \omega_1^x \left| 1 \right\rangle \big) \\ & \textit{QFT}_2 \left| x \right\rangle = \frac{1}{\sqrt{2^2}} \big( \left| 00 \right\rangle + \omega_2^x \left| 01 \right\rangle + \omega_2^{2 \cdot x} \left| 10 \right\rangle + \omega_2^{3 \cdot x} \left| 11 \right\rangle \big) \end{split}$$