

An Improved ANS Transaction Template

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INTRODUCTION

The Department of Homeland Security (DHS) continues to fund research (through the Domestic Nuclear Detection Office (DNDO)) for replacement detectors for nuclear material. These detectors are designed to replace the ^3He detectors currently installed as radiation portal monitors. DHS / DNDO (along with PNNL) has determined a set of objectives that replacement technologies should meet (summerized in Table I). These detec-

TABLE I: Replacement Detector Requirements

Parameter	Specification
Absolute neutron detection efficiency	2.5 cps/ng of ^{252}Cf
Intrinsic gamma-neutron detection efficiency	$\epsilon_{int, \gamma n} \leq 10^{-6}$
Gamma absolute rejection ratio for neutrons (GARRn)	$0.9 \leq \text{GARRn} \leq 1.1$ at 10 mR/h exposure
Cost	\$ 30,000 per system

tors must be able to effectively discriminate between gammas (which can occur in medical isotopes) and neutrons (indictive of special nuclear material).

A possible replacment technology being investigated is ^6Li (thermal cross section of 940 barns) embeded in a scinillating polymer matrix. Upon the absorbiton of a neutron ^6Li fissions into two fragments; an alpha particle of energy 2.05 Mev and a triton of energy 2.73 MeV. This large amount of kinetic energy liberated (Q-value 4.78 MeV) makes ^6Li an attractive alternative compared to other reactions such as the ^3He reaction which has a Q-value of 0.756 MeV or ^{10}B (Q-value of 2.31 MeV). In addition, the the alpha and triton produced are charged ions and relatively easy to stop and detect.

THEORY

Neutron detectors often utilize a material doped with an isotope of large thermal cross section for absorption such as ^3He , ^6Li or ^{10}B . When these materials absorb a neutron the nucleus of the isotope becomes unstable and fissions into reaction products. These reaction products (having an initial kinetic energy equal to the Q-value of the neutron absorption reaction) travel through the material, transferring their kinetic energy to the material. Photon interactions in the detector occur when a photon scatters off a single electron in a Compton scatter-

ing event and transfers a portion of its energy to the electron predominately through Compton scattering. This Compton electron then produces a cascade of secondary electrons in the material, which, depending upon the energy, may or may not deposit a majority of its energy in the detector. The difference in the transfer of kinetic energy from charged particle to electrons and from photon interactions (Compton scattering) to electrons introduces an opportunity to exploit the difference in energy deposition in order to maximize the discrimination between neutron and photon interactions in a detector.

DESCRIPTION OF WORK

The energy deposition was investigated using simulation and then benchmarking the simulation against the measured performance of the detector.

Detector Simulation

The detector geometry was simulated as a single layer of neutron absorbing material mounted atop of a non-scintillating material (PMMA). The initial events for runs where chosen by setting up a particle gun for thermal (0.025 eV) neutrons upon the detector and for both gammas resulting from a ^{60}Co decay.

The user of the GEANT4 toolkit is responsible for selecting the proper physics processes to model in the `PhysicsList`. Thus, extensive use of `G4ModularPhysicsList` was employed to handle the assigning of the physics processes to each particle in the correct order. The physics lists chosen for this simulation are listed below:

- **G4EmStandardPhysics** The electromagnetic physics defines the electrons, muons, and taus along with their corresponding neutrinos. For electrons, the primary concern of this simulation, multiple scattering, electron ionization, and electron bremsstrahlung processes were assigned. In addition the positron is defined and the multiple scattering process, electron ionization process, electron bremsstrahlung process and positron annihilation is assigned [1].
- **G4EmLivermorePhysics** The Livermore physics process extend the `EMStandardPhysics` down to low (250 eV) energies. Even lower energies can be reached by including `G4DNAPhysics`. The physics processes extended with `G4EmLivermorePhysics` are the photo-electric effect, Compton scattering, Rayleigh scattering, gamma conversion, Ionisation and Bremsstrahlung[1].

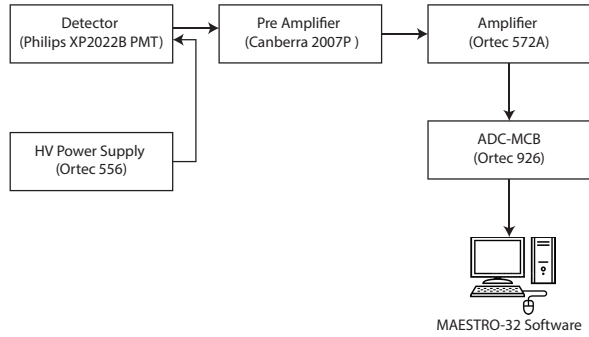


Fig. 1: Pulse height measurement electronics

- **HadronPhysicsQGSP_BERT_HP** Hadronic physics are included to model the nuclear interactions. The chosen list is a Quark Gluon String Model for energies in the 5-25 GeV range, with a Bertini cascade model until 20 MeV. Once a hadron has an energy of 20 MeV the high precision cross section driven models are applied[2].
- **G4IonPhysics** Finally, to handle the transport of the charged ions resulting from an ${}^6\text{Li}(n, \alpha){}^3\text{H}$ interaction the **G4IonPhysics** list was used.

Fabricated Detector Measurements

The pulse height of a radiation event is proportional to the energy deposition of the event[3].

Samples are mounted to a Philips XP2022B 10 stage PMT with silicone based optical grease (Saint Gobain BC-630), which is attached to a Canberra 2007P base (also serving as a preamplifier). The PMT's voltage is supplied by an Ortec 556 high voltage power supply, with the power being supplied to the Canberra 2007P pre amplifier base by the Ortec 571 amplifier. The output signal of the Canberra 2007P base feeds into an Ortec 572A amplifier for pulse shaping and amplification. The amplified signal is then inputted to an Ortec 926 MCB-ADC, and can then be read using the MAESTRO-32 software. A schematic of the setup is shown in Figure 1. A 100 μCi ${}^{60}\text{Co}$ source was used to measure the gamma energy deposition of the films, while a moderated ${}^{252}\text{Cf}$ source was used for the neutrons. The ${}^{252}\text{Cf}$ irradiator contains two detector wells; a lead well and a cadmium well. The lead well measures the response of neutron of all energies, while the cadmium well measures the responses of fast neutrons. The two responses are then subtracted, as shown in Figure 2.

RESULTS

CONCLUSIONS

The included ANS style file and this clear example file are a panacea for the hours of headache that invariably results from formatting a document in Microsoft Word.

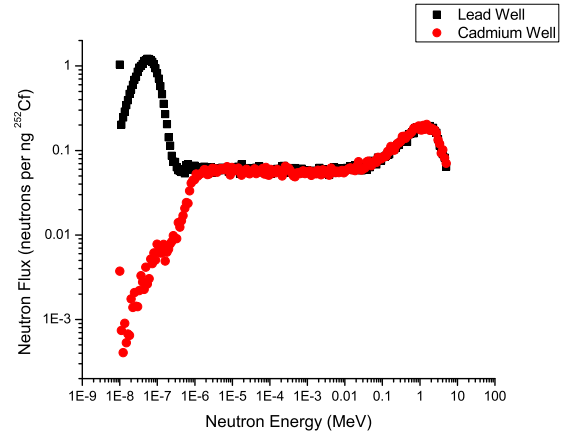


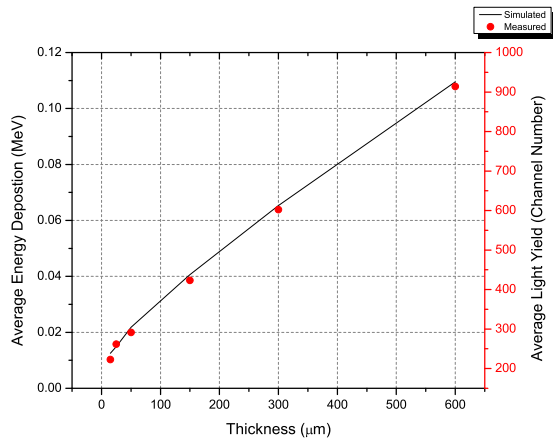
Fig. 2: Simulated Lead and Cadmium Well Spectra. The difference between the two spectra is the thermal neutron response and is what is measured.

ACKNOWLEDGMENTS

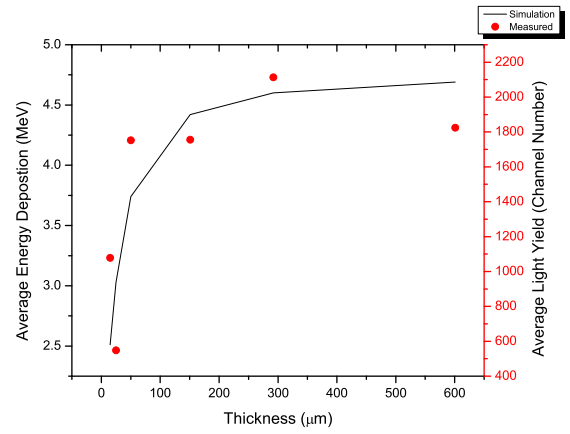
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REFERENCES

1. CERN, "Physics Lists EM constructors in Geant4 9.3," http://geant4.cern.ch/geant4/collaboration/working_groups/electromagnet (Feb. 2012).
2. CERN, "Reference Physics Lists," http://geant4.cern.ch/support/proc_mod_catalog/physics_lists/referencePL (Oct. 2008).
3. J. B. BIRKS, "Scintillations from Organic Crystals: Specific Fluorescence and Relative Response to Different Radiations," *Proceedings of the Physical Society. Section A*, **64**, 10, 874–877 (Oct. 1951).

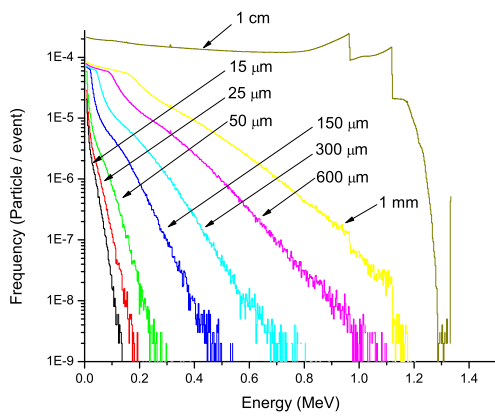


(a) Gamma (^{60}Co)

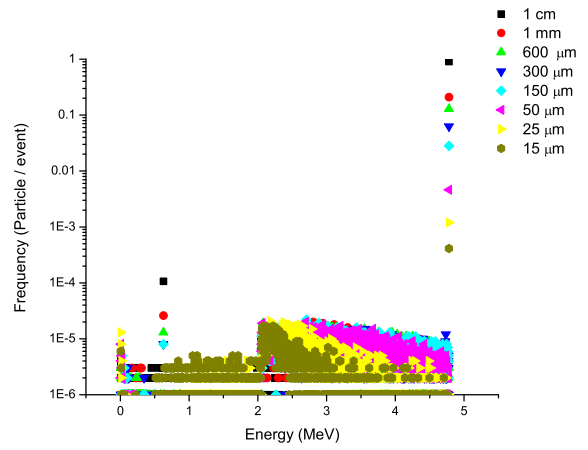


(b) Neutrons

Fig. 3: Average Energy Deposition and Measured Light Yield



(a) Gamma (^{60}Co)



(b) Neutrons

Fig. 4: Simulated Energy Deposition