



1. Angular Momentum Solutions

Practice set-1

1. The Hamiltonian of an electron in a constant magnetic field \vec{B} is given by $H = \mu \vec{\sigma} \cdot \vec{B}$, where μ is a positive constant and $\vec{\sigma} = (\sigma_1, \sigma_2, \sigma_3)$ denotes the Pauli matrices. Let $\omega = \mu B / \hbar$ and I be the 2×2 unit matrix. Then the operator $e^{iHt/\hbar}$ simplifies to

[NET JUNE 2011]

- A. $I \cos \frac{\omega t}{2} + \frac{i \vec{\sigma} \cdot \vec{B}}{B} \sin \frac{\omega t}{2}$ B. $I \cos \omega t + \frac{i \vec{\sigma} \cdot \vec{B}}{B} \sin \omega t$
 C. $I \sin \omega t + \frac{i \vec{\sigma} \cdot \vec{B}}{B} \cos \omega t$ D. $I \sin 2\omega t + \frac{i \vec{\sigma} \cdot \vec{B}}{B} \cos 2\omega t$

Solution: $H = \mu \vec{\sigma} \cdot \vec{B}$ where $\vec{\sigma} = (\sigma_1, \sigma_2, \sigma_3)$ are pauli spin matrices and \vec{B} are constant magnetic field. $\vec{\sigma} = (\sigma_1 \hat{i}, \sigma_2 \hat{j}, \sigma_3 \hat{k})$, $\vec{B} = (B_x \hat{i} + B_y \hat{j} + B_z \hat{k})$ and Hamiltonian $H = \mu \vec{\sigma} \cdot \vec{B}$ in matrices form is given by

$$H = \mu \begin{pmatrix} B_z & B_x - iB_y \\ B_x + iB_y & -B_z \end{pmatrix}.$$

Eigenvalue of given matrices are given by $+\mu B$ and $-\mu B$. H matrices are not diagonals so $e^{iHt/\hbar}$ is equivalent to

$$S^{-1} \begin{pmatrix} e^{\frac{i\mu B t}{\hbar}} & 0 \\ 0 & e^{-\frac{i\mu B t}{\hbar}} \end{pmatrix} S$$

where S is unitary matrices

$$\text{and } S^{-1} = S = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{pmatrix}.$$

$$S^{-1} \begin{pmatrix} e^{\frac{i\mu B t}{\hbar}} & 0 \\ 0 & e^{-\frac{i\mu B t}{\hbar}} \end{pmatrix} S = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} e^{\frac{i\mu B t}{\hbar}} & 0 \\ 0 & e^{-\frac{i\mu B t}{\hbar}} \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{pmatrix},$$

where $\omega = \mu B / \hbar$.

$$e^{iHt/\hbar} = \begin{pmatrix} \cos \omega t & i \sin \omega t \\ i \sin \omega t & \cos \omega t \end{pmatrix},$$

which is equivalent to $I \cos \omega t + i \sigma_x \sin \omega t$ can be written as $I \cos \omega t + \frac{i \vec{\sigma} \cdot \vec{B}}{B} \sin \omega t$, where $\sigma_x = \frac{i \vec{\sigma} \cdot \vec{B}}{B}$
The correct option is (b)

2. In a system consisting of two spin $\frac{1}{2}$ particles labeled 1 and 2, let $\vec{S}^{(1)} = \frac{\hbar}{2} \vec{\sigma}^{(1)}$ and $\vec{S}^{(2)} = \frac{\hbar}{2} \vec{\sigma}^{(2)}$ denote the corresponding spin operators. Here $\vec{\sigma} \equiv (\sigma_x, \sigma_y, \sigma_z)$ and $\sigma_x, \sigma_y, \sigma_z$ are the three Pauli matrices.

In the standard basis the matrices for the operators $S_x^{(1)} S_y^{(2)}$ and $S_y^{(1)} S_x^{(2)}$ are respectively,

[NET JUNE 2011]

- A. $\frac{\hbar^2}{4} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \frac{\hbar^2}{4} \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$
 B. $\frac{\hbar^2}{4} \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix}, \frac{\hbar^2}{4} \begin{pmatrix} -i & 0 \\ 0 & i \end{pmatrix}$
 C. $\frac{\hbar^2}{4} \begin{pmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \end{pmatrix}, \frac{\hbar^2}{4} \begin{pmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & -i & 0 \\ 0 & i & 0 & 0 \\ i & 0 & 0 & 0 \end{pmatrix}$
 D. $\frac{\hbar^2}{4} \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \end{pmatrix}, \frac{\hbar^2}{4} \begin{pmatrix} 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$

Solution: $S_x^{(1)} S_y^{(2)} = \frac{\hbar^2}{4} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \otimes \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \Rightarrow \frac{\hbar^2}{4} \begin{pmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \end{pmatrix}$
 $S_y^{(1)} S_x^{(2)} = \frac{\hbar^2}{4} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \otimes \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \Rightarrow \frac{\hbar^2}{4} \begin{pmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & -i & 0 \\ 0 & i & 0 & 0 \\ i & 0 & 0 & 0 \end{pmatrix}$

The correct option is c

3. These two operators of above QUESTION satisfy the relation

[NET JUNE 2011]

- A. $\{S_x^{(1)} S_y^{(2)}, S_y^{(1)} S_x^{(2)}\} = S_z^{(1)} S_z^{(2)}$
 B. $\{S_x^{(1)} S_y^{(2)}, S_y^{(1)} S_x^{(2)}\} = 0$
 C. $[S_x^{(1)} S_y^{(2)}, S_y^{(1)} S_x^{(2)}] = i S_z^{(1)} S_z^{(2)}$
 D. $[S_x^{(1)} S_y^{(2)}, S_y^{(1)} S_x^{(2)}] = 0$

Solution: We have matrix $S_x^{(1)} S_y^{(2)}$ and $S_y^{(1)} S_x^{(2)}$ from question 6(A) so commutation is given by $[S_x^{(1)} S_y^{(2)}, S_y^{(1)} S_x^{(2)}] = 0$.
The correct option is (d)

4. The component along an arbitrary direction \hat{n} , with direction cosines (n_x, n_y, n_z) , of the spin of a spin $-\frac{1}{2}$ particle is measured. The result is

[NET JUNE 2012]

- A. 0
 B. $\pm \frac{\hbar}{2} n_z$
 C. $\pm \frac{\hbar}{2} (n_x + n_y + n_z)$
 D. $\pm \frac{\hbar}{2}$

Solution: $S_x = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, S_y = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, S_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

$\vec{n} = n_x \hat{i} + n_y \hat{j} + n_z \hat{k}$ and $n_x^2 + n_y^2 + n_z^2 = 1, \vec{S} = S_x \hat{i} + S_y \hat{j} + S_z \hat{k}$

$\vec{n} \cdot \vec{S} = n_x \begin{pmatrix} 0 & \frac{\hbar}{2} \\ \frac{\hbar}{2} & 0 \end{pmatrix} + n_y \begin{pmatrix} 0 & -\frac{i\hbar}{2} \\ \frac{i\hbar}{2} & 0 \end{pmatrix} + n_z \begin{pmatrix} \frac{\hbar}{2} & 0 \\ 0 & -\frac{\hbar}{2} \end{pmatrix}$

$\vec{n} \cdot \vec{S} = \begin{pmatrix} n_z \frac{\hbar}{2} & \frac{\hbar}{2} (n_x - in_y) \\ \frac{\hbar}{2} (n_x + in_y) & -n_z \frac{\hbar}{2} \end{pmatrix}$

Let λ is eigen value of $\vec{n} \cdot \vec{S}$

$$\begin{vmatrix} n_z \frac{\hbar}{2} - \lambda & \frac{\hbar}{2} (n_x - in_y) \\ \frac{\hbar}{2} (n_x + in_y) & -n_z \frac{\hbar}{2} - \lambda \end{vmatrix} = 0$$

$\Rightarrow -\left(\frac{n_z \hbar}{2} - \lambda\right)\left(\frac{n_z \hbar}{2} + \lambda\right) - \frac{\hbar^2}{4} (n_x^2 + n_y^2) = 0 \Rightarrow -\left(\frac{n_z^2 \hbar^2}{4} - \lambda^2\right) - \frac{\hbar^2}{4} (n_x^2 + n_y^2) = 0$

$\Rightarrow -\frac{\hbar^2}{4} (n_x^2 + n_y^2 + n_z^2) + \lambda^2 = 0 \Rightarrow \lambda = \pm \frac{\hbar}{2}$

The correct option is (d)

5. In a basis in which the z -component S_z of the spin is diagonal, an electron is in a spin state $\psi = \begin{pmatrix} (1+i)/\sqrt{6} \\ \sqrt{2/3} \end{pmatrix}$. The probabilities that a measurement of S_2 will yield the values $\hbar/2$ and $-\hbar/2$ are, respectively,

[NET JUNE 2013]

A. 1/2 and 1/2

B. 2/3 and 1/3

C. 1/4 and 3/4

D. 1/3 and 2/3

Solution: Eigen state of S_z is $|\phi_1\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ and $|\phi_2\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ corresponds to Eigen value $\frac{\hbar}{2}$ and $-\frac{\hbar}{2}$ respectively.

$$P\left(\frac{\hbar}{2}\right) = \frac{|\langle \phi_1 | \psi \rangle|^2}{\langle \psi | \psi \rangle} = \frac{|1+i|^2}{\sqrt{6}} = \frac{2}{6} = \frac{1}{3}, \quad P\left(-\frac{\hbar}{2}\right) = \frac{|\langle \phi_2 | \psi \rangle|^2}{\langle \psi | \psi \rangle} = \frac{2}{3}$$

The correct option is (d)

6. A spin $-\frac{1}{2}$ particle is in the state $\chi = \frac{1}{\sqrt{11}} \begin{pmatrix} 1+i \\ 3 \end{pmatrix}$ in the eigenbasis of S^2 and S_2 . If we measure S_z , the probabilities of getting $+\frac{\hbar}{2}$ and $-\frac{\hbar}{2}$, respectively are

[NET DEC 2013]

A. $\frac{1}{2}$ and $\frac{1}{2}$

B. $\frac{2}{11}$ and $\frac{9}{11}$

C. 0 and 1

D. $\frac{1}{11}$ and $\frac{3}{11}$

Solution: $P\left(\frac{\hbar}{2}\right) = \left| \frac{1}{\sqrt{11}} (10) \begin{pmatrix} 1+i \\ 3 \end{pmatrix} \right|^2 = \frac{1}{11} \times 2 = \frac{2}{11} \quad \because \langle \psi | \psi \rangle = 1$

$P\left(-\frac{\hbar}{2}\right) = \left| \frac{1}{\sqrt{11}} (01) \begin{pmatrix} 1+i \\ 3 \end{pmatrix} \right|^2 = \frac{9}{11}$

i.e. probability of S_z getting $\left(\frac{\hbar}{2}\right)$ and $\left(-\frac{\hbar}{2}\right)$

The correct option is b

7. Let $\vec{\sigma} = (\sigma_1, \sigma_2, \sigma_3)$, where $\sigma_1, \sigma_2, \sigma_3$ are the Pauli matrices. If \vec{a} and \vec{b} are two arbitrary constant vectors in three dimensions, the commutator $[\vec{a} \cdot \vec{\sigma}, \vec{b} \cdot \vec{\sigma}]$ is equal to (in the following I is the identity matrix)

[NET DEC 2014]

A. $(\vec{a} \cdot \vec{b})(\sigma_1 + \sigma_2 + \sigma_3)$

B. $2i(\vec{a} \times \vec{b}) \cdot \vec{\sigma}$

C. $(\vec{a} \cdot \vec{b})I$

D. $|\vec{a}||\vec{b}|I$

Solution: $\vec{a} = a_1\hat{i} + a_2\hat{j} + a_3\hat{k}, \vec{b} = b_1\hat{i} + b_2\hat{j} + b_3\hat{k}, \sigma = \sigma_x\hat{i} + \sigma_y\hat{j} + \sigma_z\hat{k}$

$$\begin{aligned} [\vec{a} \cdot \vec{\sigma}, \vec{b} \cdot \vec{\sigma}] &= [a_1\sigma_x + a_2\sigma_y + a_3\sigma_z, b_1\sigma_x + b_2\sigma_y + b_3\sigma_z] \\ [\vec{a} \cdot \vec{\sigma}, \vec{b} \cdot \vec{\sigma}] &= a_1b_1[\sigma_x, \sigma_x] + a_1b_2[\sigma_x, \sigma_y] + a_1b_3[\sigma_x, \sigma_z] + a_2b_1[\sigma_y, \sigma_x] + a_2b_2[\sigma_y, \sigma_y] \\ &\quad + a_2b_3[\sigma_y, \sigma_z] + a_3b_1[\sigma_z, \sigma_x] + a_3b_2[\sigma_z, \sigma_y] + a_3b_3[\sigma_z, \sigma_z] \\ &= a_1b_1 \cdot 0 + a_1b_2 \cdot 2i\sigma_z - 2ia_1b_3\sigma_y - a_2b_1 \cdot 2i\sigma_z + 0 + a_2b_3 \cdot 2i\sigma_x + a_3b_1 \cdot 2i\sigma_y - a_3b_2 \cdot 2i\sigma_x + 0 \\ &\Rightarrow [\vec{a} \cdot \vec{\sigma}, \vec{b} \cdot \vec{\sigma}] = 2i(\vec{a} \times \vec{b}) \cdot \vec{\sigma} \end{aligned}$$

The correct option is (b)

8. If L_i are the components of the angular momentum operator \vec{L} , then the operator $\sum_{i=1,2,3} [\vec{L}, L_i]$ equals [NET JUNE 2015]

A. \vec{L}

B. $2\vec{L}$

C. $3\vec{L}$

D. $-\vec{L}$

Solution: Let $\vec{L} = L_x\hat{i} + L_y\hat{j} + L_z\hat{k}$

$$x = 1, y = 2, z = 3$$

$$[\vec{L}, L_x] = [L_y, L_x]j + [L_z, L_x]k = -i\hbar L_z\hat{j} + L_y\hat{k}i\hbar$$

$$[[\vec{L}, L_x], L_x] = i\hbar[-L_z, L_x]\hat{j} + [L_y, L_x]i\hbar - i\hbar \cdot i\hbar L_y\hat{j} - (i\hbar)L_z(i\hbar)L_z(i\hbar) \cdot \hat{k} = \hbar^2[L_y\hat{j} + L_z\hat{k}]$$

$$\text{similarly, } [[\vec{L}, L_y], L_y] = \hbar^2[L_x\hat{i} + L_z\hat{k}]$$

$$[[\vec{L}, L_z], L_z] = \hbar^2[L_x\hat{i} + L_y\hat{j}]$$

$$\sum_{i=1,2,3} [[L, L_i]L_i] = 2\hbar^2[L_x\hat{i} + L_y\hat{j} + L_z\hat{k}] = 2\vec{L} \quad \text{put } \hbar = 1$$

The correct option is (b)

9. The Hamiltonian for a spin- $\frac{1}{2}$ particle at rest is given by $H = E_0(\sigma_z + \alpha\sigma_x)$, where σ_x and σ_z are Pauli spin matrices and E_0 and α are constants. The eigenvalues of this Hamiltonian are [NET DEC 2015]

A. $\pm E_0\sqrt{1+\alpha^2}$

B. $\pm E_0\sqrt{1-\alpha^2}$

C. E_0 (doubly degenerate)

D. $E_0(1 \pm \frac{1}{2}\alpha^2)$

Solution: $H = E_0(\sigma_z + \alpha\sigma_x) = E_0 \left[\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \alpha \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right] \Rightarrow H = E_0 \begin{pmatrix} 1 & \alpha \\ \alpha & -1 \end{pmatrix}$ if λ is eigen value, then

$$H - \lambda I = 0 \Rightarrow E_0 \begin{pmatrix} (1-\lambda) & \alpha \\ \alpha & -(1+\lambda) \end{pmatrix} = 0, \quad \lambda = \pm E_0\sqrt{1+\alpha^2}$$

The correct option is (a)

10. If $\hat{L}_x, \hat{L}_y, \hat{L}_z$ are the components of the angular momentum operator in three dimensions the commutator $[\hat{L}_x, \hat{L}_x \hat{L}_y \hat{L}_z]$ may be simplified to

[NET JUNE 2016]

- A. $i\hbar L_x (\hat{L}_z^2 - \hat{L}_y^2)$ B. $i\hbar \hat{L}_z \hat{L}_y \hat{L}_x$
 C. $i\hbar L_x (2\hat{L}_z^2 - \hat{L}_y^2)$ D. 0

Solution:

$$\begin{aligned} : [L_x, L_x L_y L_z] &= L_x [L_x, L_y L_z] + [L_x, L_x] L_y L_z \\ &= L_x [L_x, L_y] L_z + L_x L_y [L_x, L_z] + 0 = L_x [i\hbar L_z] L_z + L_x L_y (-i\hbar L_y) \\ &= i\hbar L_x L_z^2 - i\hbar L_x L_y^2 = i\hbar L_x (L_z^2 - L_y^2) \end{aligned}$$

The correct option is (a)

11. The Hamiltonian of a spin $\frac{1}{2}$ particle in a magnetic field \vec{B} is given by $H = -\mu \cdot \vec{B} \cdot \vec{\sigma}$, where μ is a real constant and $\vec{\sigma} = (\sigma_x, \sigma_y, \sigma_z)$ are the Pauli spin matrices. If $\vec{B} = (B_0, B_0, 0)$ and the spin state at time $t = 0$ is an eigenstate of σ_x , then of the expectation values $\langle \sigma_x \rangle$, $\langle \sigma_y \rangle$ and $\langle \sigma_z \rangle$

[NET JUNE 2018]

- A. only $\langle \sigma_x \rangle$ changes with time B. only $\langle \sigma_y \rangle$ changes with time
 C. only $\langle \sigma_z \rangle$ changes with time D. all three change with time

Solution: $\langle \sigma_x \rangle$, $\langle \sigma_y \rangle$ and $\langle \sigma_z \rangle$ will changes with time because Eigen state of σ_x ie $\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$ and $\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$ and can be written in basis of eigen state of $H = -\mu \cdot \vec{B} \cdot \vec{\sigma} = -B_0 \begin{pmatrix} 0 & 1-i \\ 1+i & 0 \end{pmatrix}$
 The correct option is (d)

Future
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Practice Set-2

1. For a spin- s particle, in the eigen basis of \vec{S}^2, S_x the expectation value $\langle sm | S_x^2 | sm \rangle$ is

[GATE 2010]

- A. $\frac{\hbar^2 \{s(s+1) - m^2\}}{2}$ B. $\hbar^2 \{s(s+1) - 2m^2\}$
 C. $\hbar^2 \{s(s+1) - m^2\}$ D. $\hbar^2 m^2$

Solution:

$$\begin{aligned} \langle sm | S_x^2 | sm \rangle &= \frac{1}{4} \langle sm | (S_+ + S_-)^2 | sm \rangle = \frac{1}{4} \langle sm | S_+^2 + S_-^2 + S_+ S_- + S_- S_+ | sm \rangle \\ &= \frac{1}{4} \langle sm | S_+ S_- + S_- S_+ | sm \rangle = \frac{\hbar^2}{2} [s(s+1) - m^2] \quad [\because S_+ S_- + S_- S_+ = 2(S^2 - S_z^2)] \end{aligned}$$

The correct option is (a)

2. If L_x, L_y and L_z are respectively the x, y and z components of angular momentum operator L . The commutator $[L_x L_y, L_z]$ is equal to

[GATE 2011]

- A. $i\hbar (L_x^2 + L_y^2)$ B. $2i\hbar L_z$
 C. $i\hbar (L_x^2 - L_y^2)$ D. 0

Solution: $[L_x L_y, L_z] = L_x [L_y L_z] + [L_x, L_z] L_y = i\hbar (L_x^2 - L_y^2)$
 The correct option is (c)

3. Which one of the following commutation relations is NOT CORRECT? Here, symbols have their usual meanings.

[GATE 2013]

- A. $[L^2, L_z] = 0$ B. $[L_x, L_y] = i\hbar L_z$
 C. $[L_z, L_+] = \hbar L_+$ D. $[L_z, L_-] = \hbar L_-$

Solution: The correct option is (d)

4. A spin-half particle is in a linear superposition $0.8|\uparrow\rangle + 0.6|\downarrow\rangle$ of its spin-up and spindown states. If $|\uparrow\rangle$ and $|\downarrow\rangle$ are the eigenstates of σ_z , then what is the expectation value up to one decimal place, of the operator $10\sigma_z + 5\sigma_x$? Here, symbols have their usual meanings.

[GATE 2013]

Solution:

$$\begin{aligned} \psi &= .8|\uparrow\rangle + .6|\downarrow\rangle = 0.8 \begin{pmatrix} 1 \\ 0 \end{pmatrix} + 0.6 \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 0.8 \\ 0.6 \end{pmatrix} \\ \text{Operator } A &= 10\sigma_z + 5\sigma_x = 10 \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + 5 \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \Rightarrow A = \begin{pmatrix} 10 & 5 \\ 5 & -10 \end{pmatrix} \\ \langle A \rangle &= \langle \psi | A | \psi \rangle = \begin{pmatrix} 0.8 & 0.6 \end{pmatrix} \begin{pmatrix} 10 & 5 \\ 5 & -10 \end{pmatrix} \begin{pmatrix} 0.8 \\ 0.6 \end{pmatrix} = (8.8 - 1.2) = 7.6 \end{aligned}$$

5. If \vec{L} is the orbital angular momentum and \vec{S} is the spin angular momentum, then $\vec{L} \cdot \vec{S}$ does not commute with

[GATE 2014]

A. S_z B. L^2 C. S^2 D. $(\vec{L} + \vec{S})^2$

Solution: The correct option is (d)

6. If L_+ and L_- are the angular momentum ladder operators then the expectation value of $(L_+L_- + L_-L_+)$ in the state $|l = 1, m = 1\rangle$ of an atom is \hbar^2

[GATE 2014]

Solution: $(L_+L_- + L_-L_+) = 2(L^2 - L_z^2) = 2(l(l+1) - m^2)\hbar^2 = 2\hbar^2$

7. The Pauli matrices for three spin $-\frac{1}{2}$ particles are $\vec{\sigma}_1, \vec{\sigma}_2$ and $\vec{\sigma}_3$, respectively. The dimension of the Hilbert space required to define an operator $\vec{O} = \vec{\sigma}_1 \cdot \vec{\sigma}_2 \times \vec{\sigma}_3$ is

[GATE 2015]

Solution: $\sigma_2 \times \sigma_3$ has dimension of 4 and $\sigma_1 \cdot \sigma_2 \times \sigma_3$ has dimension of $2 \times 4 = 8$

8. Let the Hamiltonian for two spin-1/2 particles of equal masses m , momenta \vec{p}_1 and \vec{p}_2 and positions \vec{r}_1 and \vec{r}_2 be $H = \frac{1}{2m}p_1^2 + \frac{1}{2m}p_2^2 + \frac{1}{2}m\omega^2(r_1^2 + r_2^2) + k\vec{\sigma}_1 \cdot \vec{\sigma}_2$, where $\vec{\sigma}_1$ and $\vec{\sigma}_2$ denote the corresponding Pauli matrices, $\hbar\omega = 0.1\text{eV}$ and $k = 0.2\text{eV}$. If the ground state has net spin zero, then the energy (in eV) is

[GATE 2015]

Solution: $H = \frac{1}{2m}p_1^2 + \frac{1}{2m}p_2^2 + \frac{1}{2}m\omega^2(r_1^2 + r_2^2) + k\vec{\sigma}_1 \cdot \vec{\sigma}_2$

$$\vec{\sigma} = \vec{\sigma}_1 + \vec{\sigma}_2 \Rightarrow \vec{\sigma}^2 = \sigma_1^2 + \sigma_2^2 + 2\vec{\sigma}_1 \cdot \vec{\sigma}_2 \Rightarrow 2\vec{\sigma}_1 \cdot \vec{\sigma}_2 = \vec{\sigma}^2 - \sigma_1^2 - \sigma_2^2$$

$$\Rightarrow 2\vec{\sigma}_1 \cdot \vec{\sigma}_2 = 0 - 3I - 3I = -6I \Rightarrow \vec{\sigma}_1 \cdot \vec{\sigma}_2 = -3$$

Now energy $E = 2 \times \frac{3}{2}\hbar\omega + k(-3) = 3 \times (0.1) + (0.2)(-3) = -0.3\text{eV}$

9. If \vec{s}_1 and \vec{s}_2 are the spin operators of the two electrons of a He atom, the value of $\langle \vec{s}_1 \cdot \vec{s}_2 \rangle$ for the ground state is

[GATE 2016]

A. $-\frac{3}{2}\hbar^2$ B. $-\frac{3}{4}\hbar^2$

C. 0

D. $\frac{1}{4}\hbar^2$

Solution: $\vec{s} = \vec{s}_1 + \vec{s}_2, s_1 = \frac{1}{2}, s_2 = \frac{1}{2}, s = 0, 1$

$$\langle \vec{s}_1 \cdot \vec{s}_2 \rangle = \frac{s(s+1)\hbar^2 - s_1(s_1+1)\hbar^2 - s_2(s_2+1)\hbar^2}{2}$$

For

$$s = 1, \langle \vec{s}_1 \cdot \vec{s}_2 \rangle = \frac{2\hbar^2 - \frac{3}{4}\hbar^2 - \frac{3}{4}\hbar^2}{2} = \frac{3}{4}\hbar^2$$

$$s = 0, \langle \vec{s}_1 \cdot \vec{s}_2 \rangle = \frac{0\hbar^2 - \frac{3}{4}\hbar^2 - \frac{3}{4}\hbar^2}{2} = -\frac{3}{4}\hbar^2$$

The correct option is (b)

10. σ_x, σ_y and σ_z are the Pauli matrices. The expression $2\sigma_x\sigma_y + \sigma_y\sigma_x$ is equal to

[GATE 2016]

A. $-3i\sigma_z$ B. $-i\sigma_z$ C. $i\sigma_z$ D. $3i\sigma_z$

