Thermo Physics Review

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1 3D-Ideal Gas Law

$$PV = NkT \tag{1}$$

$$Nk = nR (2)$$

1.1 Equipartition Theorem

Only true if f is quadratic. E.g. If kinetic energy and potential energy are given by quadratic functions

$$y = nf \frac{1}{2}kT \tag{3}$$

$$f = 3$$
 Monatomic (4)

$$f = 5$$
 Diatomic (5)

2 1-st Law

$$U = Q + W \tag{6}$$

Values are (+) if work or heat is added or done on the system. Values are (-) if work or heat is removed from or done by the system.

$$dW = -PdV (7)$$

$$W = -\int_{V_{\perp}}^{V_f} P dV \tag{8}$$

2.1 Cycles

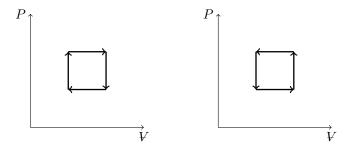


Figure 1: Left is Heat engine cycle. Right is Refrigerator cycle.

Efficiency is given by benefit divided by cost. Maximum efficiency is only found in the Caront Cycle.

$$Eff_{max} = 1 - \frac{T_C}{T_H} \tag{9}$$

3 Gas Expansion

Isothermal Temperature is constant T = const.

Abiabatic No heat is transferred $Q=0, PV^{\frac{2+f}{f}}=\mathrm{const}, \Delta S=0.$

$$C_{V} = \frac{\partial U}{\partial T} \quad C_{P} = \frac{\partial U}{\partial T} + P \left(\frac{\partial V}{\partial T}\right)_{P}$$
 (10)

$$C_V = \frac{1}{2}fNK$$
 $C_P = Nk\left(\frac{f+2}{2}\right)$ Only for ideal gasses (11)

4 Multiplicity

Einstien Solid

$$\Omega = \frac{(N+q-1)!}{q!(N-1)!} \approx \left(\frac{eq}{N}\right)^{N}$$

Two Level System

$$\Omega = \frac{N!}{N_{\uparrow}! N_{\downarrow}!} = \frac{N!}{N_{\uparrow}! (N - N_{\uparrow})!}$$

Ideal Gas

$$\Omega = \frac{V^N}{N!h^{3N}} \left[\frac{2\pi^{\frac{3N}{2}}}{\left(\frac{3N}{2} - 1\right)!} \cdot (2mU)^{\frac{3N-1}{2}} \right]$$

5 Entropy and Heat

Heat Capacities

$$C_V \equiv \left(\frac{\partial U}{\partial T}\right)_{N,V} \tag{12}$$

(13)

For an Einstein solid with $q\gg N$

$$C_V = \frac{\partial}{\partial T} (NkT) = Nk \tag{14}$$

For a monatomic ideal gas

$$C_V = \frac{\partial}{\partial T} \left(\frac{3}{2} NkT \right) = \frac{3}{2} Nk \tag{15}$$

Measuring Entropies

$$dS = \frac{dU}{T} = \frac{Q}{T} \tag{16}$$

$$dS = \frac{C_V dT}{T} \tag{17}$$

$$\Delta S = S_f - S_i = \int_{T_i}^{T_f} \frac{C_V}{T} dT \tag{18}$$

$$C_V \to 0 \quad \text{as} \quad T \to 0$$
 (19)

6 Thermo-Dynamic Potentials

$$S = S(N, V, U) \tag{20}$$

$$dS = \frac{\partial S}{\partial N}dN + \frac{\partial S}{\partial V}dV + \frac{\partial S}{\partial U}dU \tag{21}$$

$$= -\frac{\mu}{T}dN + \frac{P}{T}dV + \frac{1}{T}dU \tag{22}$$

$$dU = TdS + \mu dN \tag{23}$$

$$\frac{\partial U}{\partial S} = T \quad \frac{\partial U}{\partial N} = \mu \quad \frac{\partial U}{\partial V} = -P \tag{24}$$

Enthalpy $H \equiv U + PV$ T = 0

Helmholts Free Energy $F \equiv U - TS$ P = 0

Gibbs Gree Energy $G \equiv U - TS + PV$

7 Boltzmann Statistics

7.1 Partition Function

$$Z \equiv \sum_{s} e^{-\frac{E(s)}{kT}}$$
 (25)

Where s are the possible states of a single particle. Z cannot be measured but it contains all the information relevant to a system.

7.2 Boltzmann Distribution

$$P(s) = \frac{e^{-\frac{E(s)}{kT}}}{Z} \tag{26}$$

This gives the probability of a particle being in any given state s. Using this and the partition function, anything about a system can be derived.