retselis set2

April 18, 2021

Problem Set #2

Computational Electromagnetism and Applications ($\Upsilon\Phi\Upsilon203$)

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1 Exercise 2.1

Plot the fraction of the numerical to the analytical phase velocity (see slide 4-23 of S. Gedney) as a function of the angle of incidence θ for (a) $\Delta x = \Delta y = \lambda/10$ and (b) $\Delta x = \lambda/10$, $\Delta y = \lambda/20$. For both cases, choose the time step to be equal to 90% of the stability limit.

1.1 Solution

The stability limit for the two-dimensional explicit equation is given by the following formula:

$$\Delta t < \frac{1}{c\sqrt{\frac{1}{\Delta x^2} + \frac{1}{\Delta y^2}}}$$

which can be further simplified if $\Delta x = \Delta y = \Delta$:

$$\Delta t < \frac{\Delta}{c\sqrt{2}}$$

and from these two equations we can calculate the desired time stop to be some percentage of the stability limit.

From the dispersion relationship in discrete space:

$$\frac{1}{c^2 \Delta t^2} \sin^2 \left(\frac{\omega \Delta t}{2} \right) - \frac{1}{\Delta x^2} \sin^2 \left(\frac{\tilde{k}_x \Delta x}{2} \right) - \frac{1}{\Delta y^2} \sin^2 \left(\frac{\tilde{k}_y \Delta y}{2} \right) = 0$$

we can substitute $\tilde{k}_x = \tilde{k}cos(\theta)$ and $\tilde{k}_y = \tilde{k}sin(\theta)$, leading to:

$$\frac{1}{c^2 \Delta t^2} \sin^2 \left(\frac{\omega \Delta t}{2} \right) = \frac{1}{\Delta x^2} \sin^2 \left(\frac{\tilde{k} \cos(\theta) \Delta x}{2} \right) + \frac{1}{\Delta y^2} \sin^2 \left(\frac{\tilde{k} \sin(\theta) \Delta y}{2} \right)$$

 \tilde{k} can be computed using Newton-Raphson's method $\tilde{k}_{i+1} = \tilde{k}_i - f\left(\tilde{k}_i\right)/f'\left(\tilde{k}_i\right)$ which in this case can be rewritten as:

$$\tilde{k}_{i+1} = \tilde{k}_i - \frac{\frac{1}{\Delta x^2} \sin^2\left(\frac{\tilde{k}_i \cos(\theta) \Delta x}{2}\right) + \frac{1}{\Delta y^2} \sin^2\left(\frac{\tilde{k}_i \sin(\theta) \Delta y}{2}\right) - \frac{1}{c^2 \Delta t^2} \sin^2\left(\frac{\omega \Delta t}{2}\right)}{\frac{\cos(\theta)}{2\Delta x} \sin\left(\tilde{k}_i \cos(\theta) \Delta x\right) + \frac{\sin(\theta)}{2\Delta y} \sin\left(\tilde{k}_i \sin(\theta) \Delta y\right)}$$

We note that c values in the equation above are simplified (see the equation for the time step Δt . The numerical phase velocity is given by the following formula:

$$v_p = \frac{\omega}{\tilde{k}}$$

which can be rewritten as:

$$\frac{v_p}{c} = \frac{2\pi}{\lambda \tilde{k}}$$

We can then compute \tilde{k} for different values of the angle of incidence θ and plot those values to create the desired numerical phase velocity diagram. This procedure has been performed for both of the desired cases in the script which can be found below:

```
[1]: # Import script required packages
     import numpy as np
     import matplotlib.pyplot as plt
     # Make jupyter export images as .pdf files for higher quality
     from IPython.display import set_matplotlib_formats
     set_matplotlib_formats('png', 'pdf')
     def percent_stability_criterion(dx, dy, percent):
         # Computes a percent of the stability limit for the explicit 2D Formulation
         # Input: dx, dy, percent
         # Output: dt
         if dx == dy:
             dt = percent * dx / np.sqrt(2)
             dt = percent * 1 / (np.sqrt((1 / pow(dx, 2)) + (1 / pow(dy, 2))))
         return dt
     def k_hat_newton(theta, dx, dy, dt, lambd):
         # Computes k_hat using the Newton-Raphson method
         # Input: theta [deg], dx, dy, dt, l
         # Output: k hat
         theta = np.deg2rad(theta)
         # Define f and f dot
```

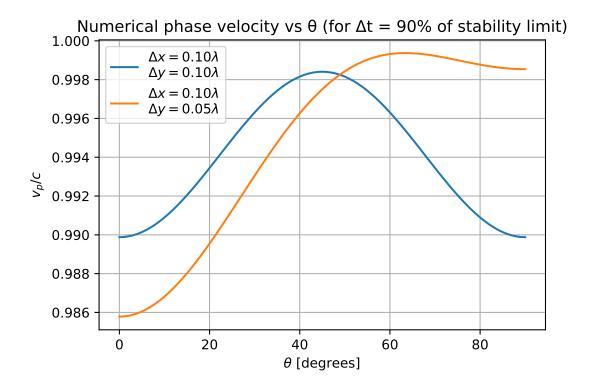
```
f = lambda khat: ((1 / pow(dx, 2)) * pow(np.sin(khat * np.cos(theta) * dx / u)
 \rightarrow 2), 2)) + (
                (1 / pow(dy, 2)) * pow(np.sin(khat * np.sin(theta) * dy / 2),
\rightarrow 2)) - (
                                  (1 / pow(dt, 2)) * pow(np.sin(np.pi * dt / _ _
\rightarrowlambd), 2))
    fdot = lambda khat: (np.cos(theta) / (2 * dx)) * np.sin(khat * np.
\rightarrowcos(theta) * dx) + (
                np.sin(theta) / (2 * dy)) * np.sin(khat * np.sin(theta) * dy)
    # Stopping criteria
    n = 10 # Number of significant digits to be computed
    max_repetitions = 1000
    es = 0.5 * pow(10, (2 - n)) # Scarborough Criterion
    ea = 100
    k_prev = 2 * np.pi / lambd
    repetitions = 0
    # Main Newton-Raphson loop
    while ea > es:
        repetitions = repetitions + 1
        k_next = k_prev - (f(k_prev) / fdot(k_prev))
        ea = np.fabs((k_next - k_prev) * 100 / k_next)
        k_prev = k_next
        if repetitions > max_repetitions:
            print('Max repetitions reached without achieving desired accuracy⊔

    for E!')
            break
    k_hat = k_next
    return k_hat
def phase_velocity_vs_theta_diagram(dx, dy, dt, lambd, min_theta, max_theta):
    \# Creates a numerical phase velocity diagram vs theta for theta interval
→ [min theta, max theta]
    # Input: dx, dy, dt, lambd, min_theta, max_theta
    # Output: None (can plot figure using plt.show())
    step_size = 1
    total_size = int(((max_theta - min_theta) / step_size) + 1)
    vpc_values = np.zeros(total_size)
    # Compute k and v_p/c values
    for i in range(0, total_size):
        theta = i * step size
        k = k_hat_newton(theta, dx, dy, dt, lambd)
        vpc_values[i] = 2 * np.pi / (lambd * k)
    # Plot output
    plt.plot(vpc_values, label="$\Delta x = %.2f \lambda$\n$\Delta y = %.2f_\
 \rightarrow \lambda$" % (dx, dy))
```

```
plt.title("Numerical phase velocity vs (for \Delta t = 90\% of stability limit)")
    plt.ylabel(r"$v_{p}/c$")
    plt.xlabel(r"$\theta$ [degrees]")
    plt.legend()
    return None
theta_start = 0
theta_end = 90
# Case A: Delta_x = Delta_y = lambda/10, Delta_t = 0.9 C.L.
Delta_x = lmbd / 10
Delta_y = lmbd / 10
Delta_t = percent_stability_criterion(Delta_x, Delta_y, 0.9)
phase\_velocity\_vs\_theta\_diagram(Delta\_x,\ Delta\_y,\ Delta\_t,\ lmbd,\ theta\_start, \\ \\ \sqcup

→theta_end)

# Case B: Delta_x = lambda/10, Delta_y = lambda/20, Delta_t = 0.9 C.L.
lmbd = 1
Delta x = lmbd / 10
Delta_y = lmbd / 20
Delta_t = percent_stability_criterion(Delta_x, Delta_y, 0.9)
phase_velocity_vs_theta_diagram(Delta_x, Delta_y, Delta_t, lmbd, theta_start,_u
→theta_end)
# Show resulting plot
plt.grid()
plt.show()
```



2 Exercise 2.2

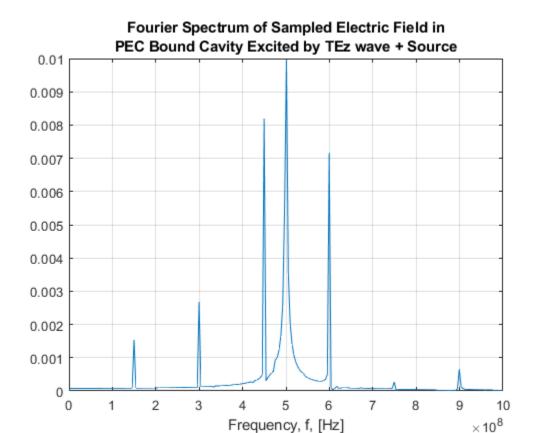
For the attached program fdtd2D_TEz.m make the necessary changes (there are some parts for which you have to add some lines of code) and compute using a Fourier transformation the resonant frequencies for an empty cavity (see slides 4-39 of S. Gedney). What happens to these frequencies if a conducting cylinder is added in this cavity, by making the relevant changes to the code? (starting from line 100)

2.1 Solution

```
Grid parameters
ie=100;
               %number of grid cells in x-direction
               %number of grid cells in y-direction
je=100;
ib=ie+1;
jb=je+1;
is=40:
               %location of source
js=je/2;
               %location of source
ir=15;
               %location of field recording
jr=15;
               %location of field recording
dx=0.01;
               %space increment of square lattice
dy=0.01;
nmax=10000;
               %total number of time steps
Material parameters
   This is used for adding a lossy material in the cavity
media=2;
%The first element corresponds to media 1 and the second to media 2:
eps=[1.0 1.0]; %relative electric permittivity
sig=[0.0 1.0e+7];
             %electric conductivity (S/m)
mur=[1.0 1.0];
             %relative permeability
sim=[0.0 0.0];
             % equivalent magnetic loss (\Omega/m)
Wave excitation
rtau=500.0e-9;
tau=rtau/dt;
delay=3*tau;
source=zeros(1,nmax);
for n=1:7.0*tau
 source(n)=sin(omega*(n-delay)*dt)*exp(-((n-delay)^2/tau^2));
```

```
end
Field arrays
              ***************
ex=zeros(ie,jb);
              %fields in main grid
ey=zeros(ib, je);
hz=zeros(ie, je);
erec=zeros(1,nmax);
Updating coefficients
for i=1:media
 eaf =dt*sig(i)/(2.0*epsz*eps(i));
 ca(i)=(1.0-eaf)/(1.0+eaf);
 cb(i)=dt/epsz/eps(i)/dx/(1.0+eaf);
 haf =dt*sim(i)/(2.0*muz*mur(i));
 da(i)=(1.0-haf)/(1.0+haf);
 db(i)=dt/muz/mur(i)/dx/(1.0+haf);
end
Geometry specification (main grid)
Initialize entire main grid to free space
caex(1:ie,1:jb)=ca(1);
cbex(1:ie,1:jb)=cb(1);
caey(1:ib,1:je)=ca(1);
cbey(1:ib,1:je)=cb(1);
dahz(1:ie,1:je)=da(1);
dbhz(1:ie,1:je)=db(1);
BEGIN TIME-STEPPING LOOP
for n=1:nmax
   Update electric fields (EX and EY) in main grid
```

```
ex(:,2:je)=caex(:,2:je).*ex(:,2:je)+...
       cbex(:,2:je).*(hz(:,2:je)-hz(:,1:je-1));
ey(2:ie,:)=caey(2:ie,:).*ey(2:ie,:)+...
       cbey(2:ie,:).*(hz(1:ie-1,:)-hz(2:ie,:));
ey(is,2:je) = ey(is,2:je) + source(n)/dy;
field_recording(n) = ey(ir,jr);
Update magnetic fields (HZ) in main grid
hz(1:ie,1:je)=dahz(1:ie,1:je).*hz(1:ie,1:je)+...
         dbhz(1:ie,1:je).*(ex(1:ie,2:jb)-ex(1:ie,1:je)+...
                      ey(1:ie,1:je)-ey(2:ib,1:je));
END TIME-STEPPING LOOP
end
Y = fft(field recording);
Fs = 1/dt;
L = nmax;
P2 = abs(Y/L);
P1 = P2(1:L/2+1);
P1(2:end-1) = 2*P1(2:end-1);
f = Fs*(0:(L/2))/L;
%plot inline --format=svg
figure;
plot(f, P1);
xlim([0,1*10e8]);
xlabel('Frequency, f, [Hz]');
title({'Fourier Spectrum of Sampled Electric Field in',...
  'PEC Bound Cavity Excited by TEz wave + Source'});
grid on;
```



From the figure above, we can determine the frequencies for each cavity mode in order to check the FDTD approximation, as seen in the table below:

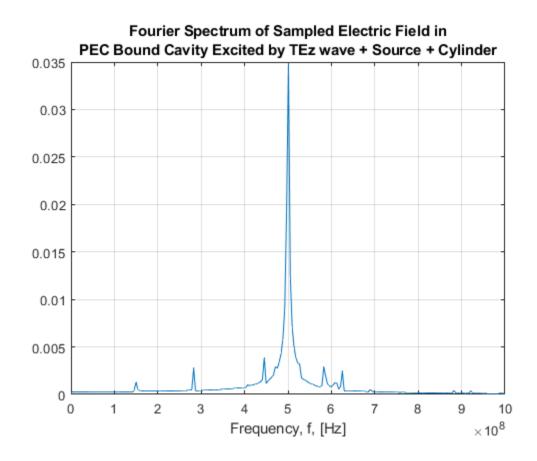
Cavity Mode	Approximation (FDTD)	Exact
TE_{10}	149.9 MHz	150.0 MHz
TE_{20}	$299.8~\mathrm{MHz}$	$300.0~\mathrm{MHz}$
TE_{30}	$449.7~\mathrm{MHz}$	$450.0~\mathrm{MHz}$
TE_{40}	$599.6~\mathrm{MHz}$	$600.0~\mathrm{MHz}$
TE_{50}	$749.4~\mathrm{MHz}$	$750.0~\mathrm{MHz}$
TE_{60}	$899.3~\mathrm{MHz}$	900.0 MHz

We can also observe an extra peak at 501.1 MHz, which corresponds to the frequency of the source which we included in the program (500 MHz). In any case, we note that the resulting error is below 1%. Let us now include the conducting cylinder:

```
[2]: ex=zeros(ie,jb); %fields in main grid ey=zeros(ib,je);
```

```
hz=zeros(ie,je);
erec=zeros(1,nmax);
Updating coefficients
for i=1:media
 eaf =dt*sig(i)/(2.0*epsz*eps(i));
 ca(i)=(1.0-eaf)/(1.0+eaf);
 cb(i)=dt/epsz/eps(i)/dx/(1.0+eaf);
 haf =dt*sim(i)/(2.0*muz*mur(i));
 da(i)=(1.0-haf)/(1.0+haf);
 db(i)=dt/muz/mur(i)/dx/(1.0+haf);
end
Geometry specification (main grid)
Initialize entire main grid to free space
caex(1:ie,1:jb)=ca(1);
cbex(1:ie,1:jb)=cb(1);
caey(1:ib, 1:je)=ca(1);
cbey(1:ib,1:je)=cb(1);
dahz(1:ie,1:je)=da(1);
dbhz(1:ie,1:je)=db(1);
   Add cylinder made of medium 2
          % diameter of cylinder
diam=20;
for i=1:ie
for j=1:je
 dist2=(i+0.5-icenter)^2 + (j-jcenter)^2;
 if dist2 <= rad^2</pre>
   caex(i,j)=ca(2);
   cbex(i,j)=cb(2);
 end
 dist2=(i-icenter)^2 + (j+0.5-jcenter)^2;
 if dist2 <= rad^2</pre>
```

```
caey(i,j)=ca(2);
  cbey(i,j)=cb(2);
 end
end
end
BEGIN TIME-STEPPING LOOP
for n=1:nmax
Update electric fields (EX and EY) in main grid
ex(:,2:je)=caex(:,2:je).*ex(:,2:je)+...
      cbex(:,2:je).*(hz(:,2:je)-hz(:,1:je-1));
ey(2:ie,:)=caey(2:ie,:).*ey(2:ie,:)+...
      cbey(2:ie,:).*(hz(1:ie-1,:)-hz(2:ie,:));
\%\% put one line here for the source, which extends in j = 2:je;
\%ex(is,2:je) = ex(is,2:je) + source(n)/dx;
ey(is,2:je) = ey(is,2:je) + source(n)/dx;
%%% put one line here to record the field
field_recording(n) = ey(ir,jr);
Update magnetic fields (HZ) in main grid
hz(1:ie,1:je)=dahz(1:ie,1:je).*hz(1:ie,1:je)+...
        dbhz(1:ie,1:je).*(ex(1:ie,2:jb)-ex(1:ie,1:je)+...
                  ey(1:ie,1:je)-ey(2:ib,1:je));
END TIME-STEPPING LOOP
end
Y = fft(field_recording);
Fs = 1/dt;
L = nmax;
```



We note that addition of the cylinder results to the resonant frequencies appearing at different and lower frequencies. We can also observe that the frequency of the source dominates the spectrum, having the highest amplitude among the different frequencies.