

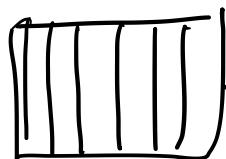
## Double slit

$$I = I_0 \underbrace{\frac{\sin^2 \beta}{\beta^2}}_{\text{diffraction}} \underbrace{\cos^2 \gamma}_{\text{interference}}$$

$$\beta = \frac{\pi b \sin \theta}{\lambda}$$

$$\gamma = \frac{\pi d \sin \theta}{\lambda}$$

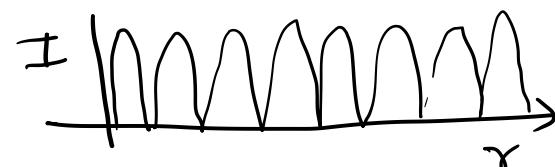
Effect of increasing # of slits.

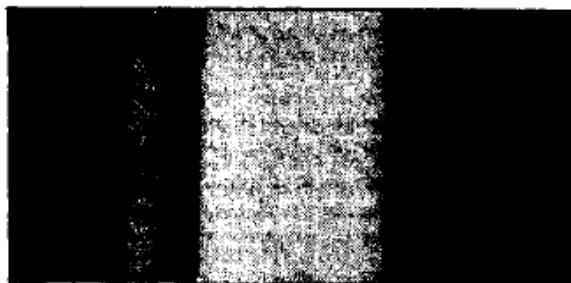
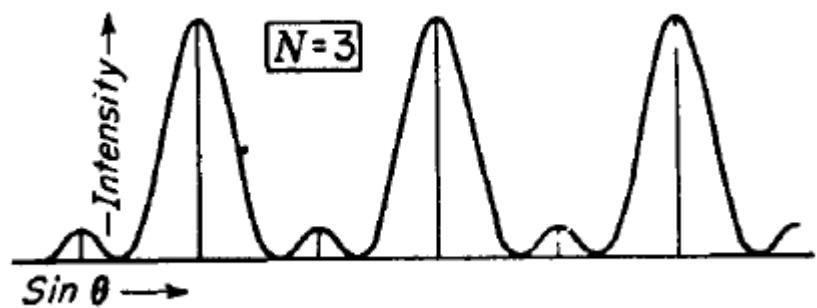


→ grating : large # of parallel equidistant slits.

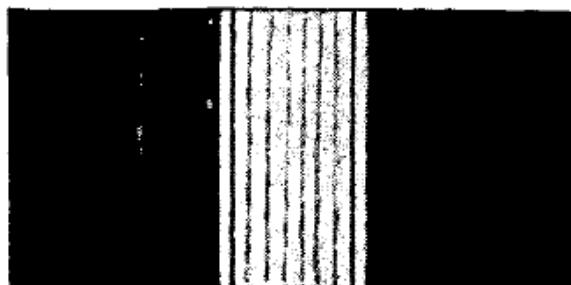
→ # of slits increasing → narrowing of maxima, sharper lines

→ appearance of weak secondary maxima in between principal maxima. Recall in double slit





(a) 1 slit



(b) 2 slits



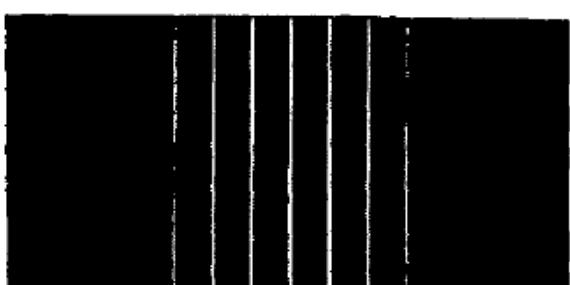
(c) 3 slits



(d) 5 slits



(e) 6 slits



(f) 20 slits

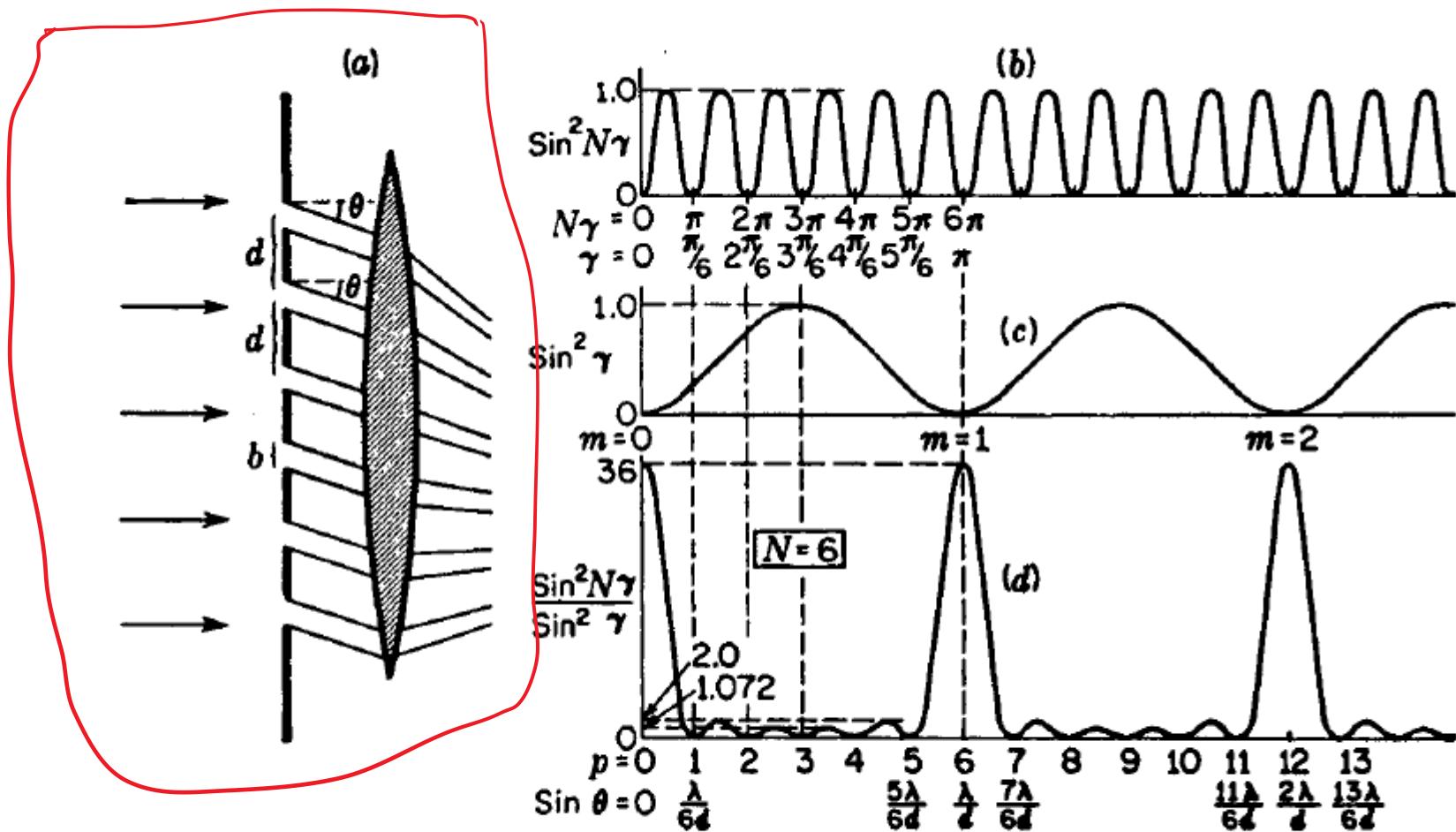


FIGURE 17C

Fraunhofer diffraction by a grating of six very narrow slits and details of the intensity pattern.

Intensity distribution from N slit grating :

Complex amplitude method

Net amplitude (complex) at P

$$\underbrace{A e^{i\theta}}_{\tilde{A}} = a \left( 1 + e^{i\delta} + e^{2i\delta} + \dots + e^{i(N-1)\delta} \right)$$
$$= a \frac{1 - e^{iN\delta}}{1 - e^{i\delta}}$$

$$I = |\tilde{A}|^2 = a^2 \frac{(1 - e^{iN\delta})(1 - e^{-iN\delta})}{(1 - e^{i\delta})(1 - e^{-i\delta})} = a^2 \frac{1 - \cos N\delta}{1 - \cos \delta}$$

$$I = \frac{a^2 \sin^2 N \frac{\delta}{2}}{\sin^2 \frac{\delta}{2}}$$

$\gamma = \frac{\delta}{2} = \frac{\pi d \sin \theta}{\lambda}$

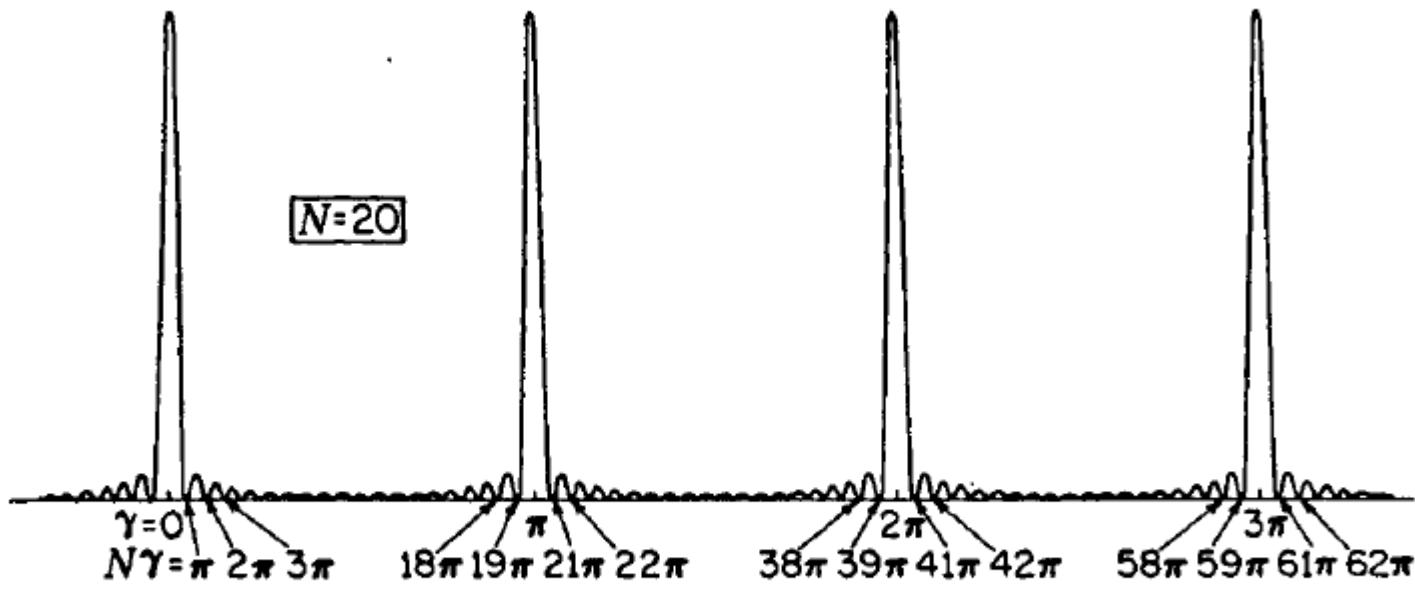
$a$  = amplitude from single slit.

$$\boxed{I = I_0 \frac{\sin^2 \beta}{\beta^2} \frac{\sin^2 N \gamma}{\sin^2 \gamma}}$$

$$\beta = \frac{\pi b \sin \theta}{\lambda}$$

for  $N = 2$ .

$$I = I_0 \frac{\sin^2 \beta}{\beta^2} \frac{\sin^2 2 \gamma}{\sin^2 \gamma} = 4 I_0 \frac{\sin^2 \beta}{\beta^2} \cos^2 \gamma$$



**FIGURE 17D**  
Intensity pattern for 20 narrow slits.

## Analyze intensity distribution

### Principal Maxima

$\frac{\sin^2 N\gamma}{\sin^2 \gamma}$  m > new factor represents interference term for N slits.

→ maximum f value is  $N^2$ .

$$\lim_{\gamma \rightarrow m\pi} \frac{\sin N\gamma}{\sin \gamma} = \pm N$$

These maxima correspond to those in the double slit since for above values of  $\gamma$

$$\rightarrow \boxed{ds \sin \theta = m\lambda}$$

Relative intensities of diff orders of m is governed by diffraction envelope relation between  $\beta$  &  $r$  is same, condition for missing orders also same

## Minima & secondary maxima

Minimum of  $\frac{\sin^2 N\gamma}{\sin^2 \gamma}$  : Numerator becomes zero more often than denominator



$$N\gamma = p\pi \rightarrow \text{minimum}$$

except for special cases where  $p = mN \rightarrow$  principal maximum

Cond'n. for minimum  $\gamma = \frac{p\pi}{N}$  excluding  $p = mN$

path diff

$$d\sin\theta = \frac{\lambda}{N}, \frac{2\lambda}{N}, \dots, \underbrace{\frac{(N-1)\lambda}{N}}, \underbrace{\frac{(N+1)\lambda}{N}} \text{ Minima .}$$

omit values  $0, \frac{N\lambda}{N}, \dots$

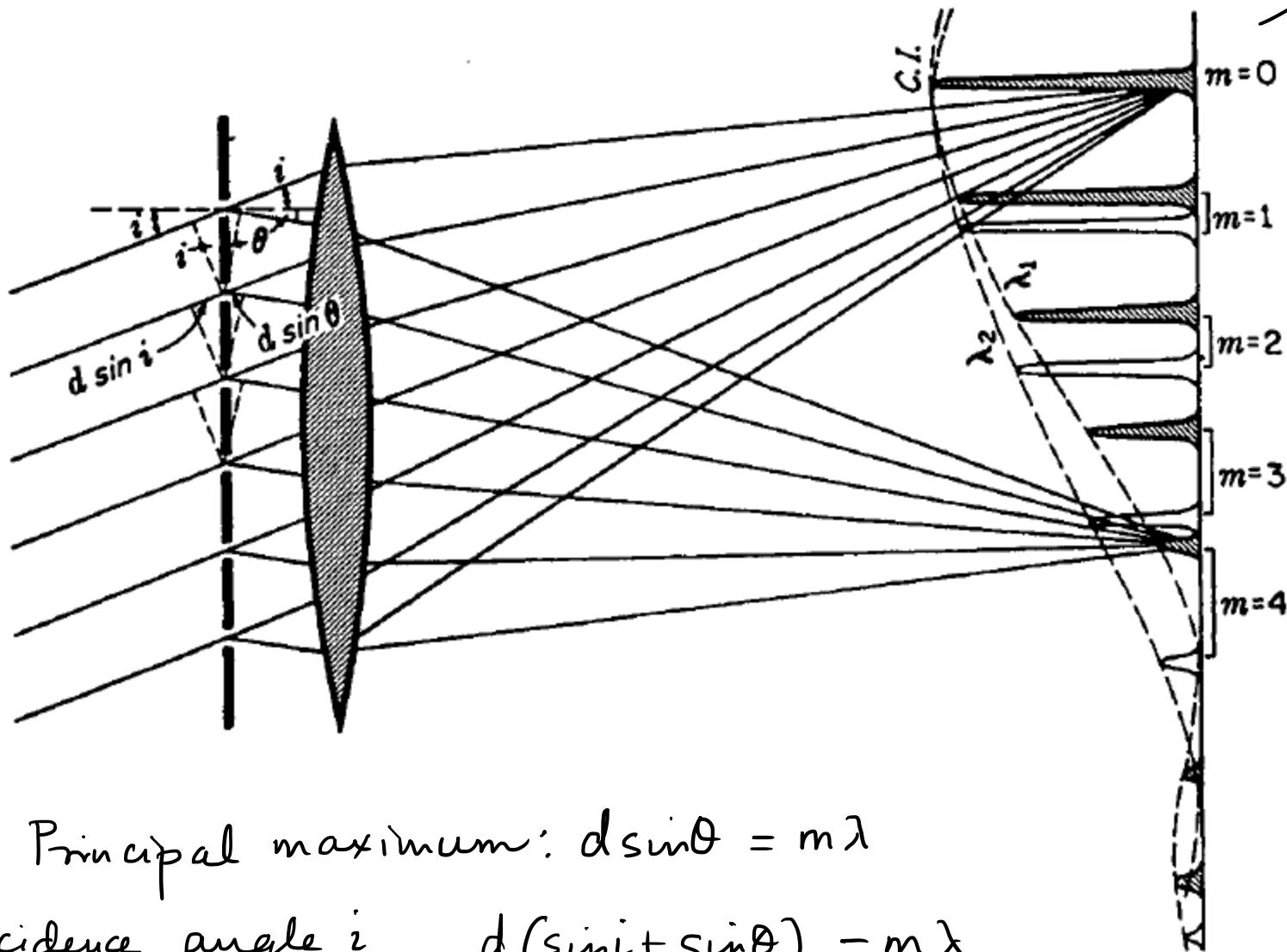
$d\sin\theta = m\lambda \rightarrow$  Principal maxima

Between two principal maxima  $(N-1)$  point of zero intensity

Between two minima intensity rises again  $\rightarrow$  secondary maxima

they are much smaller intensity than principal maxima.

- The secondary maxima are not of equal intensity and they fall off in intensity on either side of principal maxima. Not equally spaced either

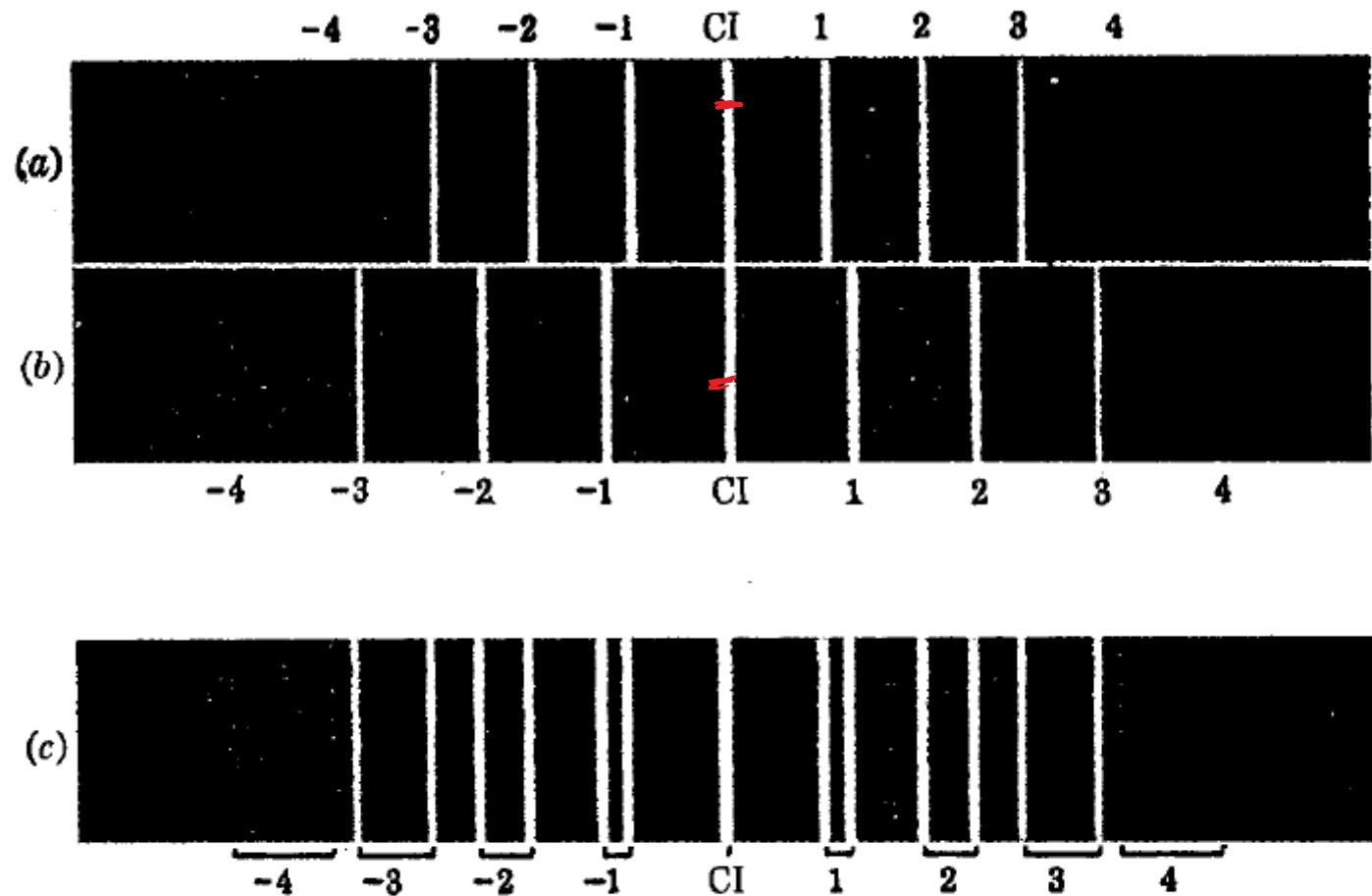


$d = 5b$  .  
5<sup>th</sup> order  
missing .

Principal maximum:  $d \sin \theta = m\lambda$

at incidence angle  $i$        $d(\sin i + \sin \theta) = m\lambda$  .

$m=0$  central image .  
maximum irrespective of  $\lambda$  .



**FIGURE 17F**

Grating spectra of two wavelengths: (a)  $\lambda_1 = 4000 \text{ \AA}$ ; (b)  $\lambda_2 = 5000 \text{ \AA}$ ; (c)  $\lambda_1$  and  $\lambda_2$  together.

