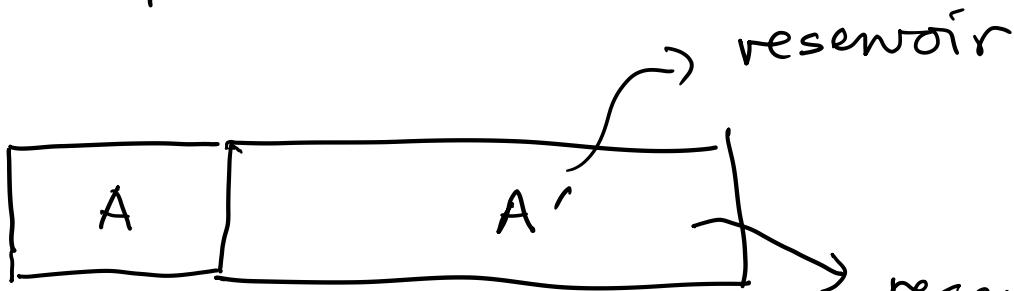


Recap



$A^{(o)} = A + A' \Rightarrow$ isolated
reservoir at a given temp T

canonical ensemble

$$P_r = \frac{e^{-\beta E_r}}{\sum_r e^{-\beta E_r}}$$

Boltzmann distribution

$$\beta = \frac{1}{kT}$$

P_r : prob. of finding A in one particular microstate r

$P(E)$: prob. of finding A in an energy range E to $E + \delta E$

$$P(E) = \sum_r P_r$$

$r \rightarrow$ over all states in given energy range

$$E < E_r < E + \delta E$$

$\Omega(E)$: # of states in this energy range
 \rightarrow all have equal prob. $\sim e^{-\beta E}$

$$P(E) = C \Omega(E) e^{-\beta E} \Rightarrow \text{valid no matter how small } A \text{ is}$$

$\Omega(E) \sim$ rapidly rising fn. of E



$P(E)$ has a peak .

Say y is some parameter characterizing the system

$$\left\{ \begin{array}{l} \bar{y} = \frac{\sum_r y_r e^{-\beta \bar{E}_r}}{\sum_r e^{-\beta \bar{E}_r}} \end{array} \right.$$

Applications of the canonical distribution

Paramagnetism :

No magnetic atoms per unit volume

Placed in an external mag field H .

- Each atom has spin $\frac{1}{2}$ and an intrinsic mag moment μ .
- can point only either parallel ~~or~~ or antiparallel to H .

Q : Substance at temp T , what is the mean mag moment M_H (in the direction of H) of such an atom?

System A : single atom

reservoir A' : rest of the atoms in material

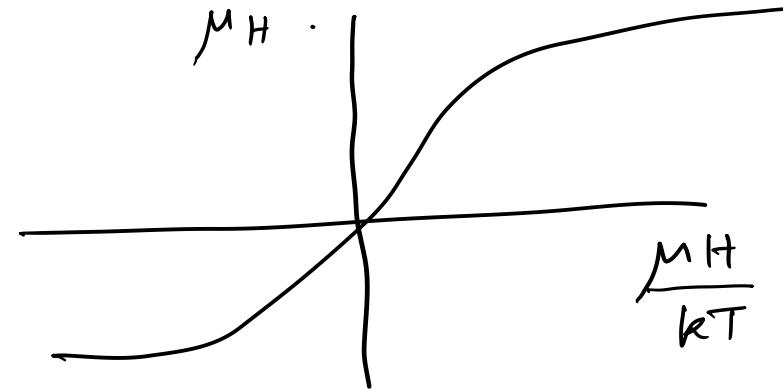
- each atom is in two possible states (+) or (-)
- (+) is || to H , $\mu_H = \mu$, $\epsilon_+ = -\mu H$
- (-) is antill to H , $\mu_H = -\mu$, $\epsilon_- = \mu H$

$$P_+ = C e^{-\beta \epsilon_+} = C e^{\beta \mu H}$$

$$P_- = C e^{-\beta \epsilon_-} = C e^{-\beta \mu H}$$

$$\bar{\mu}_H = ?$$

$$\bar{\mu}_H = \frac{P_+ \mu + P_- (-\mu)}{P_+ + P_-}$$



$$= \mu \frac{e^{\beta \mu_H} - e^{-\beta \mu_H}}{e^{\beta \mu_H} + e^{-\beta \mu_H}}$$

$$\boxed{\bar{\mu}_H = \mu \tanh \frac{\mu_H}{kT}}$$

$\rightarrow T \rightarrow \infty$

$$\bar{\mu}_H = 0$$

$\rightarrow T \rightarrow 0$

$$\bar{\mu}_H = \mu$$

Magnetization \bar{M}_o = mean magnetic moment per unit volume in the direction of H

$$\boxed{\bar{M}_o = N_o \bar{\mu}_H}$$

$$= N_o \mu \tanh \frac{\mu H}{kT}$$

$$\frac{\mu H}{kT} \ll 1$$

$$\bar{M}_o \simeq N_o \mu \frac{\mu H}{kT}$$

$$\simeq \frac{N_o \mu^2 H}{kT}$$

$$\boxed{\bar{M}_o \simeq \chi H}$$

susceptibility .

}

$$\frac{\mu H}{kT} = y$$

$$\tanh y = \frac{e^y - e^{-y}}{e^y + e^{-y}}$$

$$\simeq \frac{1+y-1+y}{1+y+1-y}$$

$$y \ll 1$$

$$1+y+1-y$$

$$\simeq y \cdot \boxed{y \gg 1}$$

$$\frac{\mu H}{kT} \ll 1 \cdot \boxed{\tanh y = 1}$$

$$\bar{M}_o = \chi H$$

$$\chi = \frac{N_o M^2}{kT}$$

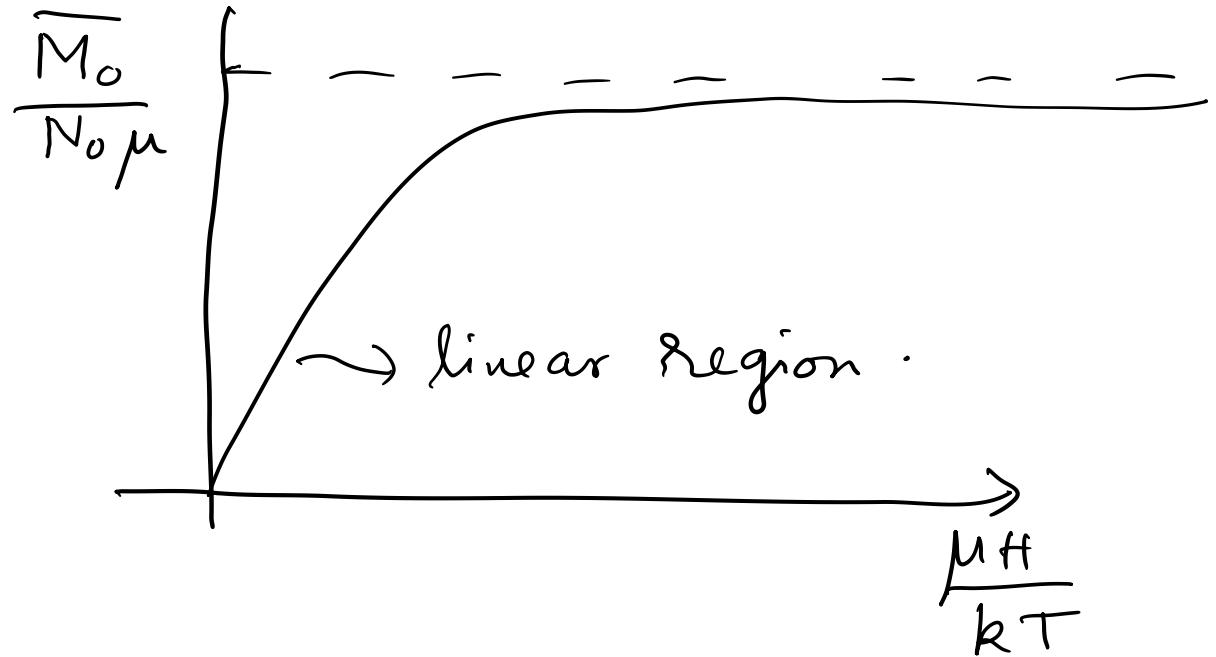
$$x \propto \frac{1}{T}$$

Curies' Law .

On the other hand

$$\bar{M}_o \rightarrow N_o \mu \quad \frac{\mu H}{kT} \gg 1 .$$

→ independent of $H \rightarrow$ equal to the saturation magnetization



non-interacting
mag atoms of
spin half &
mag moment μ .

Example 2 : Molecule in an ideal gas .

molecule \rightarrow system

other
molecules \rightarrow reservoir .

$$E = \frac{1}{2}mv^2 = \frac{p^2}{2m} .$$

$$\begin{aligned} & -E/kT \\ & \text{Boltzmann factor} \sim e^{-\frac{mv^2}{2kT}} . \end{aligned}$$