

Table 3.5 The Maxwell relations

From U :
$$\left(\frac{\partial T}{\partial V}\right)_S = -\left(\frac{\partial p}{\partial S}\right)_V$$

From H :
$$\left(\frac{\partial T}{\partial p}\right)_S = \left(\frac{\partial V}{\partial S}\right)_p$$

From ~~A~~^F:
$$\left(\frac{\partial p}{\partial T}\right)_V = \left(\frac{\partial S}{\partial V}\right)_T$$

From G :
$$\left(\frac{\partial V}{\partial T}\right)_p = -\left(\frac{\partial S}{\partial p}\right)_T$$

Tds equations

$$dS = \left(\frac{\partial S}{\partial T} \right)_V dT + \left(\frac{\partial S}{\partial V} \right)_T dV$$

$$TdS = T \left(\frac{\partial S}{\partial T} \right)_V dT + T \left(\frac{\partial S}{\partial V} \right)_T dV$$

Consider reversible isochoric change

$$TdS = (dQ)_{V=\text{const}}$$

$$T \left(\frac{\partial S}{\partial T} \right)_V = C_V$$

$$TdS = C_V dT + T \left(\frac{\partial P}{\partial T} \right)_V dV$$

↓
1st TdS eqn.

∴ Maxwell's reln

$$\left(\frac{\partial S}{\partial V} \right)_T = \left(\frac{\partial P}{\partial T} \right)_V$$

$$dS = \left(\frac{\partial S}{\partial T} \right)_P dT + \left(\frac{\partial S}{\partial P} \right)_T dP$$

$$TdS = T \left(\frac{\partial S}{\partial T} \right)_P dT + T \left(\frac{\partial S}{\partial P} \right)_T dP$$

$$TdS = C_p dT - T \left(\frac{\partial V}{\partial T} \right)_P dP$$

$$\left\{ \left(\frac{\partial S}{\partial P} \right)_T = - \left(\frac{\partial V}{\partial T} \right)_P \right.$$

↳ 2nd TdS eqn.

Internal Energy eqns

$$dU = TdS - PdV$$

$$\frac{dU}{dP} = T \frac{dS}{dP} - P \frac{dV}{dP}$$

Holding T const

$$\left(\frac{\partial U}{\partial P} \right)_T = T \left(\frac{\partial S}{\partial P} \right)_T - P \left(\frac{\partial V}{\partial P} \right)_T$$

Maxwell
↓
 $\left(\frac{\partial S}{\partial P} \right)_T = - \left(\frac{\partial V}{\partial T} \right)_P$

$$\boxed{\left(\frac{\partial U}{\partial P} \right)_T = -T \left(\frac{\partial V}{\partial T} \right)_P - P \left(\frac{\partial V}{\partial P} \right)_T}$$

→ 2nd internal energy identity

$$dU = T dS - P dV$$

$$\frac{dU}{dV} = T \frac{dS}{dV} - P$$

$$\left(\frac{\partial U}{\partial V} \right)_T = T \left(\frac{\partial S}{\partial V} \right)_T - P$$

$$\boxed{\left(\frac{\partial U}{\partial V} \right)_T = T \left(\frac{\partial P}{\partial T} \right)_V - P}$$

$$\left\{ \left(\frac{\partial S}{\partial V} \right)_T = \left(\frac{\partial P}{\partial T} \right)_V \right.$$

Maxwell

Heat Capacity Equation

$$Tds = C_p dT - T \left(\frac{\partial V}{\partial T} \right)_P dP \quad \text{--- (1)}$$

$$Tds = C_v dT + T \left(\frac{\partial P}{\partial T} \right)_V dV \quad \text{--- (2)}$$

Equating (1) & (2)

$$C_p dT - T \left(\frac{\partial V}{\partial T} \right)_P dP = C_v dT + T \left(\frac{\partial P}{\partial T} \right)_V dV$$

$$dT = \frac{T \left(\frac{\partial P}{\partial T} \right)_V dV}{C_p - C_v} + \frac{T \left(\frac{\partial V}{\partial T} \right)_P dP}{C_p - C_v} \quad \text{--- (3)}$$

But

$$dT = \left(\frac{\partial T}{\partial V} \right)_P dV + \left(\frac{\partial T}{\partial P} \right)_V dP \quad \text{--- (4)}$$

$$\Rightarrow \left(\frac{\partial T}{\partial V} \right)_P = \frac{T \left(\frac{\partial P}{\partial T} \right)_V}{C_p - C_v} \quad ; \quad \left(\frac{\partial T}{\partial P} \right)_V = \frac{T \left(\frac{\partial V}{\partial T} \right)_P}{C_p - C_v} .$$

$$C_p - C_v = T \left(\frac{\partial P}{\partial T} \right)_V \left(\frac{\partial V}{\partial T} \right)_P$$

But $\left(\frac{\partial P}{\partial T} \right)_V = - \left(\frac{\partial V}{\partial T} \right)_P \left(\frac{\partial P}{\partial V} \right)_T \rightarrow$ chain rule

$$C_p - C_v = -T \left(\frac{\partial V}{\partial T} \right)_P^2 \left(\frac{\partial P}{\partial V} \right)_T$$

$$C_p - C_v = -T \left(\frac{\partial V}{\partial T} \right)_P^2 \left(\frac{\partial P}{\partial V} \right)_T$$

- $\left(\frac{\partial P}{\partial V} \right)_T < 0$ for all known substances.

$$\left(\frac{\partial V}{\partial T} \right)^2 > 0 \quad \therefore \boxed{C_p - C_v > 0}$$

- $T \rightarrow 0 \quad C_p \rightarrow C_v$ at absolute zero $C_p = C_v$

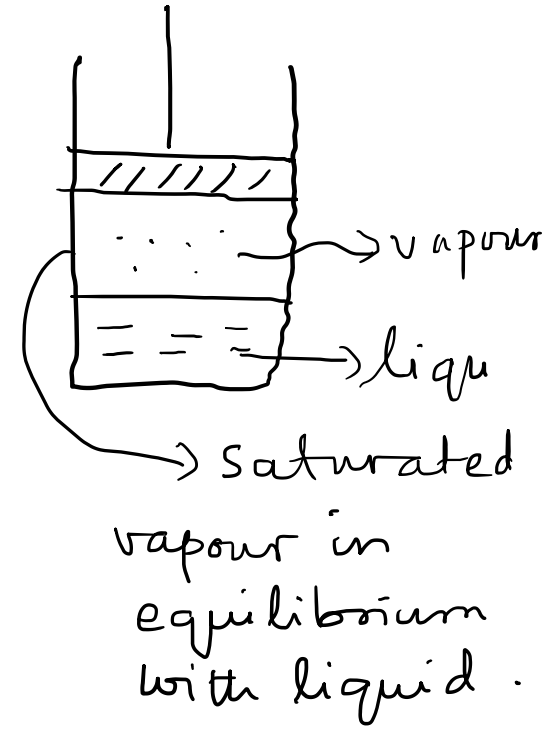
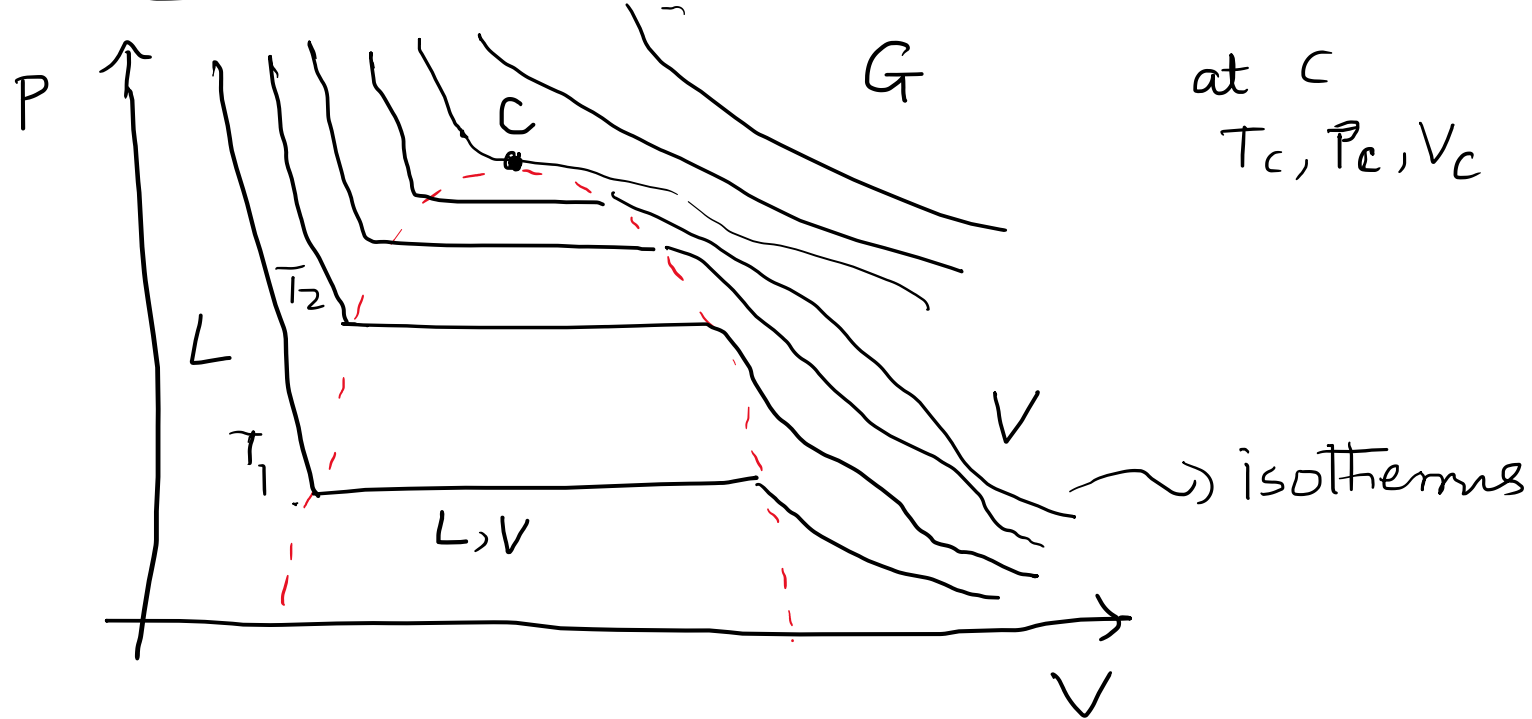
- $C_p = C_v$ when $\left(\frac{\partial V}{\partial T} \right)_P = 0$ at 4°C $\rho_w \text{ max}$

$$C_p - C_v = -T \frac{\left(\frac{\partial V}{\partial T}\right)_P^2}{\left(\frac{\partial V}{\partial P}\right)_T}$$

$$C_p - C_v = \frac{T v \beta^2}{\kappa}$$

v : specific volume .

Gibbs potential application



We will focus on flat region

$$m = m_1 + m_2$$

1 subscript liq.
2 " vapour

$$V = m_1 v_1(T) + m_2 v_2(T) \quad v : \text{specific volume} .$$

$$U = m_1 u_1(T) + m_2 u_2(T) .$$

$$\left. \begin{aligned} U &= U_1 + U_2 \\ S &= S_1 + S_2 \\ V &= V_1 + V_2 \end{aligned} \right\} \begin{aligned} G &= U - TS + PV \\ T, P &\text{ const region} \\ G &= G_1 + G_2 . \end{aligned}$$

$$\left. \begin{aligned} G_1 &= m_1 g_1 \\ G_2 &= m_2 g_2 \end{aligned} \right\} \begin{aligned} P, T &\text{ const} \\ m_1, m_2 &\text{ can vary} . \end{aligned}$$

$$\begin{aligned} m_1 &\rightarrow m + dm \\ m_2 &\rightarrow m - dm \end{aligned}$$

$$(m_1 + dm) g_1 + (m_2 - dm) g_2 = G + dG$$

$$dG = dm(g_1 - g_2) \quad \text{In equilibrium}$$

$$\Rightarrow dG = 0 \quad \Rightarrow \boxed{g_1 = g_2}$$


$$u_1 - Ts_1 + Pv_1 = u_2 - Ts_2 + Pv_2$$

$$(u_2 - u_1) - T(s_2 - s_1) + P(v_2 - v_1) = 0$$

$$\frac{d}{dT}(u_2 - u_1) - T \frac{d}{dT}(s_2 - s_1) - (s_2 - s_1) + \frac{dP}{dT}(v_2 - v_1) + P \frac{d}{dT}(v_2 - v_1) = 0$$

But 1st Law

$$T \frac{ds}{dT} = \frac{du}{dT} + P \frac{dv}{dT}$$


$$\frac{d}{dT} (u_2 - u_1) - T \frac{d}{dT} (s_2 - s_1) - (s_2 - s_1) + \frac{dP}{dT} (v_2 - v_1) + P \frac{d}{dT} (v_2 - v_1) = 0$$

↓

$$-(s_2 - s_1) + \frac{dP}{dT} (v_2 - v_1) = 0$$

$$\therefore \begin{cases} \text{But} \\ (s_2 - s_1) = \frac{\lambda}{T} \end{cases}$$

latent heat

$$\boxed{\left(\frac{dP}{dT} \right) = \frac{\lambda}{T(v_2 - v_1)}}$$

→ Clausius-Clapeyron
eqn.

