

Double slit

$$I = I_0 \underbrace{\frac{\sin^2 \beta}{\beta^2}}_{\text{diffraction}} \underbrace{\cos^2 \gamma}_{\text{interference}}$$

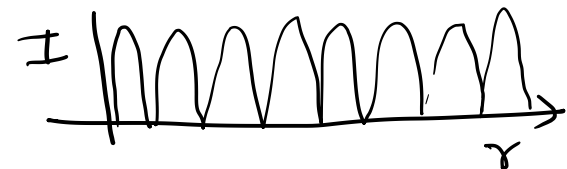
$$\beta = \frac{\pi b \sin \theta}{\lambda}$$
$$\gamma = \frac{\pi d \sin \theta}{\lambda}$$

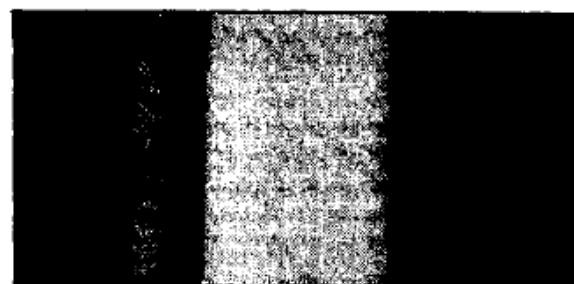
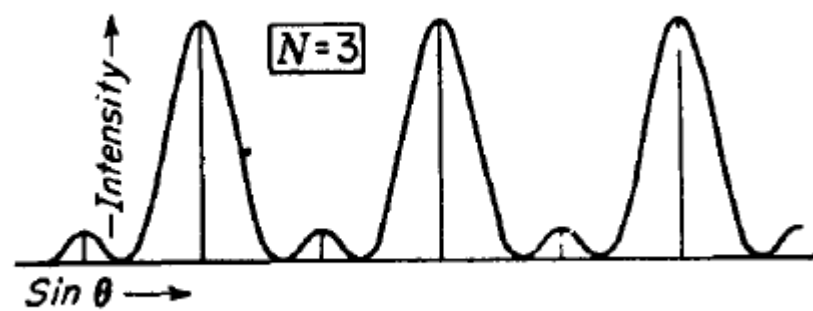
Effect of increasing # of slits.

 → grating : large # of parallel equidistant slits.

→ # of slits increasing → narrowing of maxima, sharper lines

→ appearance of weak secondary maxima in between principal maxima. Recall in double slit

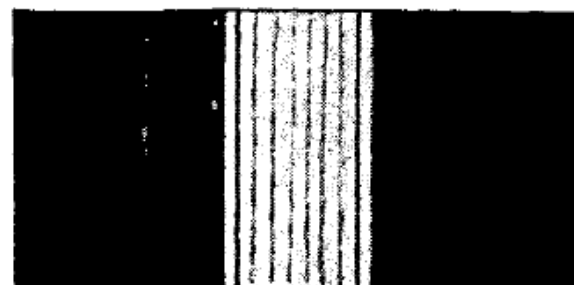




(a) 1 slit



(d) 5 slits



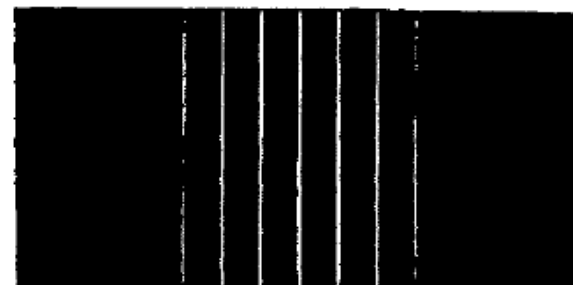
(b) 2 slits



(e) 6 slits



(c) 3 slits



(f) 20 slits

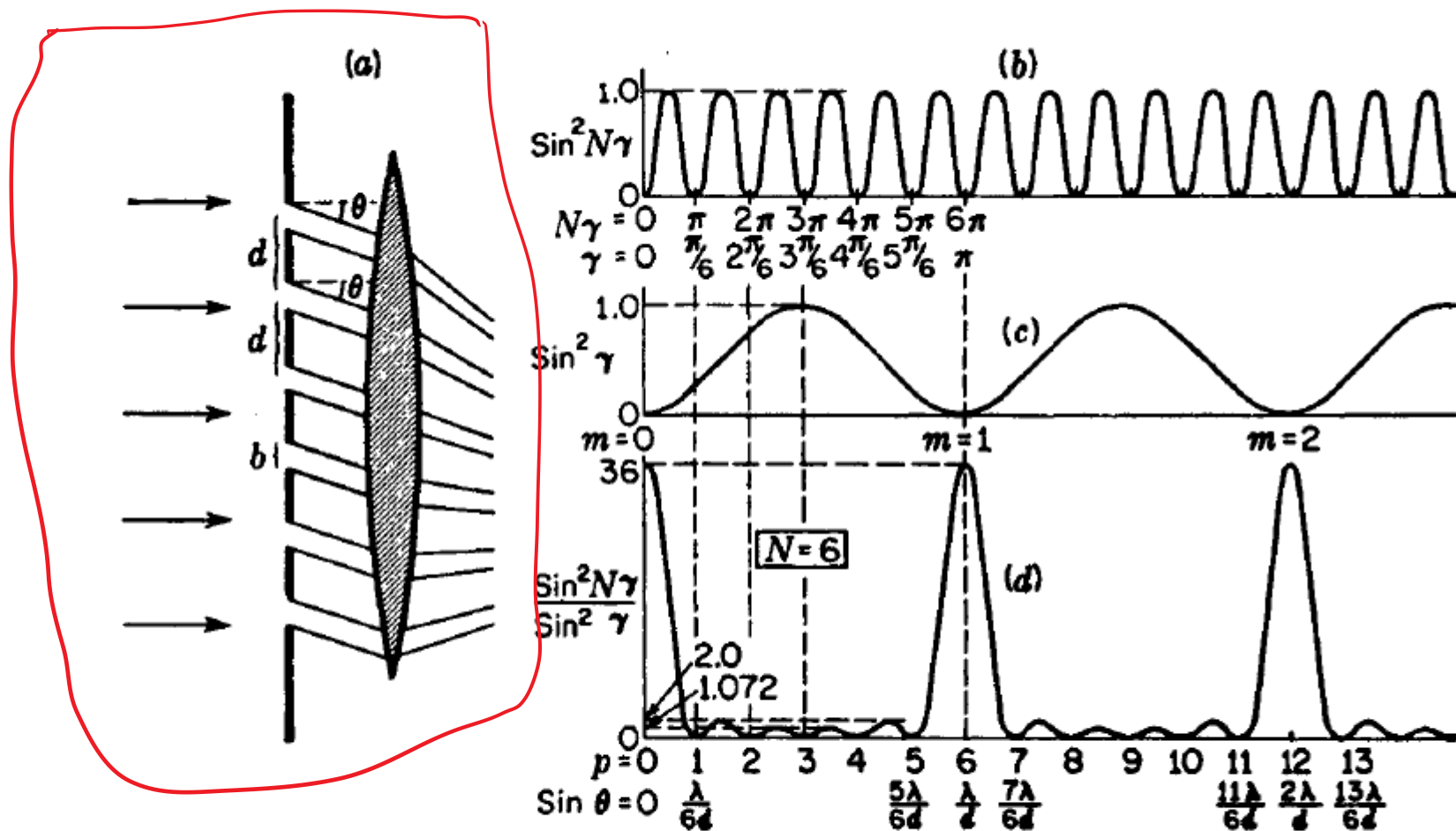


FIGURE 17C

Fraunhofer diffraction by a grating of six very narrow slits and details of the intensity pattern.

Intensity distribution from N slit grating:

Complex amplitude method

Net amplitude (complex) at P

$$\begin{aligned}\underbrace{A e^{i\theta}}_{\tilde{A}} &= a \left(1 + e^{i\delta} + e^{2i\delta} + \dots + e^{i(N-1)\delta} \right) \\ &= a \frac{1 - e^{iN\delta}}{1 - e^{i\delta}}\end{aligned}$$

$$I = |\tilde{A}|^2 = \frac{a^2 (1 - e^{iN\delta})(1 - e^{-iN\delta})}{(1 - e^{i\delta})(1 - e^{-i\delta})} = a^2 \frac{1 - \cos N\delta}{1 - \cos \delta}$$

$$I = \frac{a^2 \sin^2 N \frac{\delta}{2}}{\sin^2 \delta/2}$$

$$\gamma = \frac{\delta}{2} = \frac{\pi d \sin \theta}{\lambda}$$

\swarrow a = amplitude from single slit.

$$I = I_0 \frac{\sin^2 \beta}{\beta^2} \frac{\sin^2 N \gamma}{\sin^2 \gamma}$$

$$\beta = \frac{\pi b \sin \theta}{\lambda}$$

for $N = 2$.

$$I = I_0 \frac{\sin^2 \beta}{\beta^2} \frac{\sin^2 2\gamma}{\sin^2 \gamma} = 4 I_0 \frac{\sin^2 \beta}{\beta^2} \cos^2 \gamma.$$

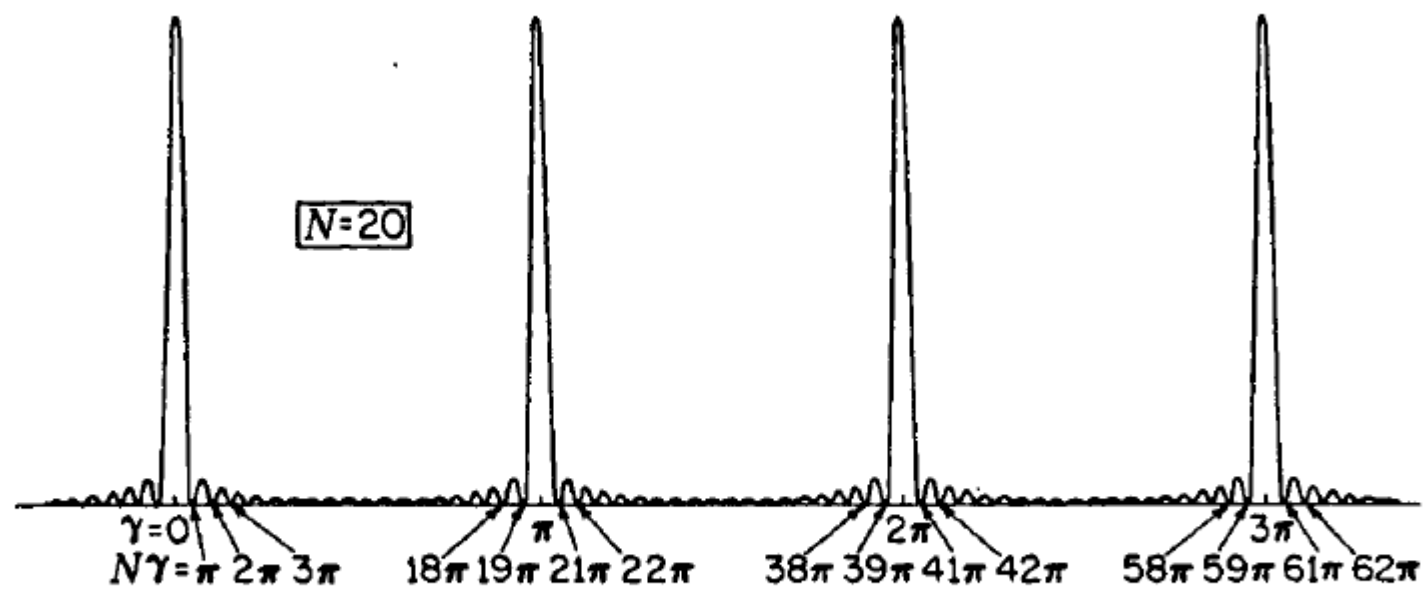


FIGURE 17D
Intensity pattern for 20 narrow slits.

Analyze intensity distribution

Principal Maxima

$\frac{\sin^2 N\gamma}{\sin^2 \gamma} \sim$ new factor represents interference term for N slits.

→ maximum of value is N^2 .

$$\lim_{\gamma \rightarrow m\pi} \frac{\sin N\gamma}{\sin \gamma} = \pm N.$$

these maxima correspond to those in the double slit since for above values of γ

$$\rightarrow \boxed{d \sin \theta = m\lambda}$$

Relative intensities of diff orders of m is governed by diffraction envelope relation between β & γ is same, condition for missing orders also same

Minima & secondary maxima

Minimum of $\frac{\sin^2 N\gamma}{\sin^2 \gamma}$: Numerator becomes zero more often than denominator



$$N\gamma = p\pi \rightarrow \text{minimum}$$

except for special cases where $p = mN \rightarrow$ principal maximum

Cond'n. for minimum $\gamma = \frac{p\pi}{N}$ excluding $p = mN$

path diff

$$d \sin \theta = \frac{\lambda}{N}, \frac{2\lambda}{N}, \dots, \frac{(N-1)\lambda}{N}, \frac{(N+1)\lambda}{N} \text{ Minima.}$$

omit values $0, \frac{N\lambda}{N}, \dots$

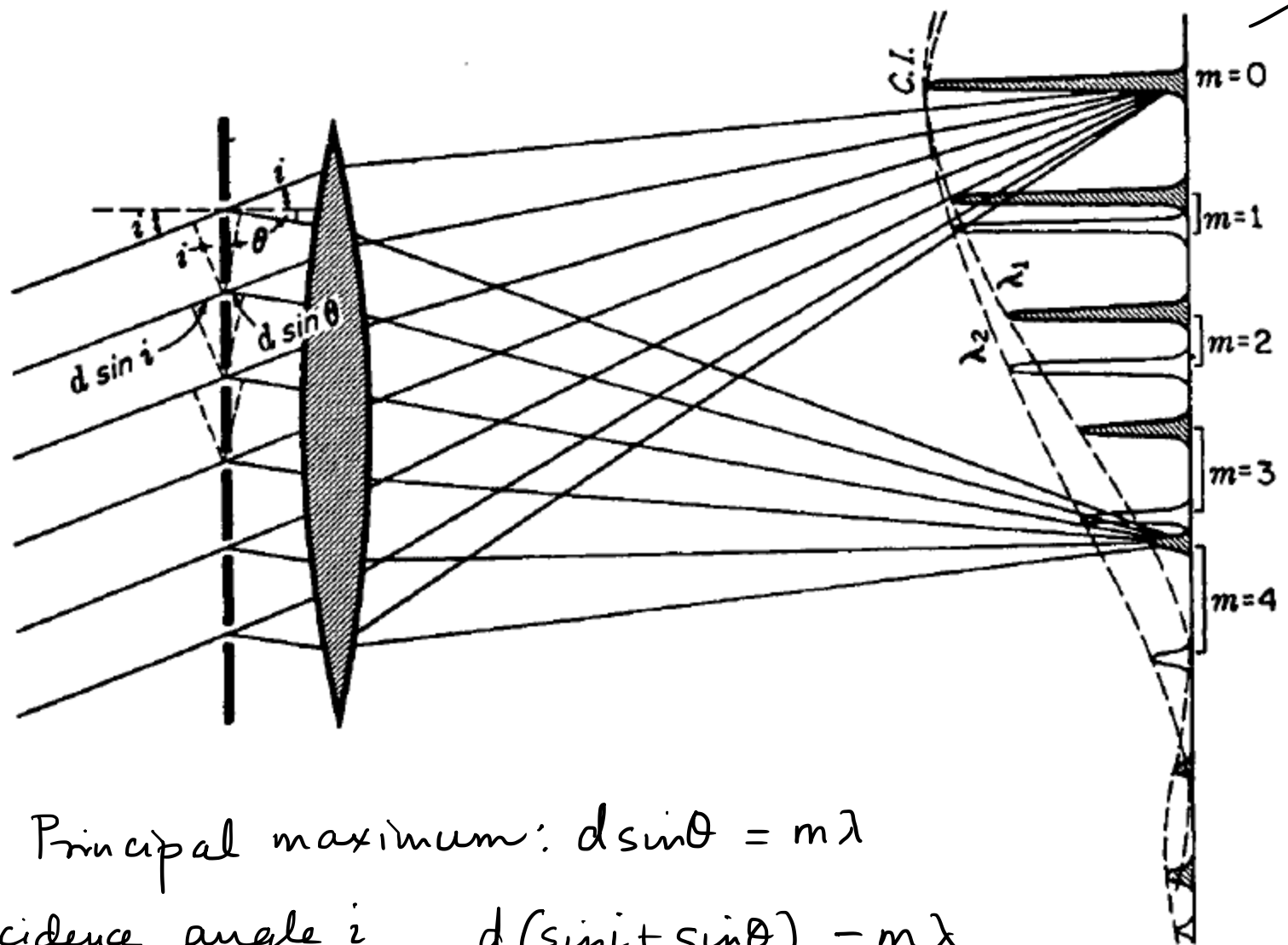
$$\boxed{d \sin \theta = m\lambda} \rightarrow \text{Principal maxima}$$

Between two principal maxima $(N-1)$ point of zero intensity

Between two minima intensity rises again \rightarrow secondary maxima

they are much smaller intensity than principal maxima.

- The secondary maxima are not of equal intensity and they fall off in intensity on either side of principal maxima. Not equally spaced either



$d = 5b$

5th order missing.

Principal maximum: $d \sin \theta = m \lambda$

at incidence angle i $d (\sin i + \sin \theta) = m \lambda$

$m=0$ central image.
maximum irrespective of λ .

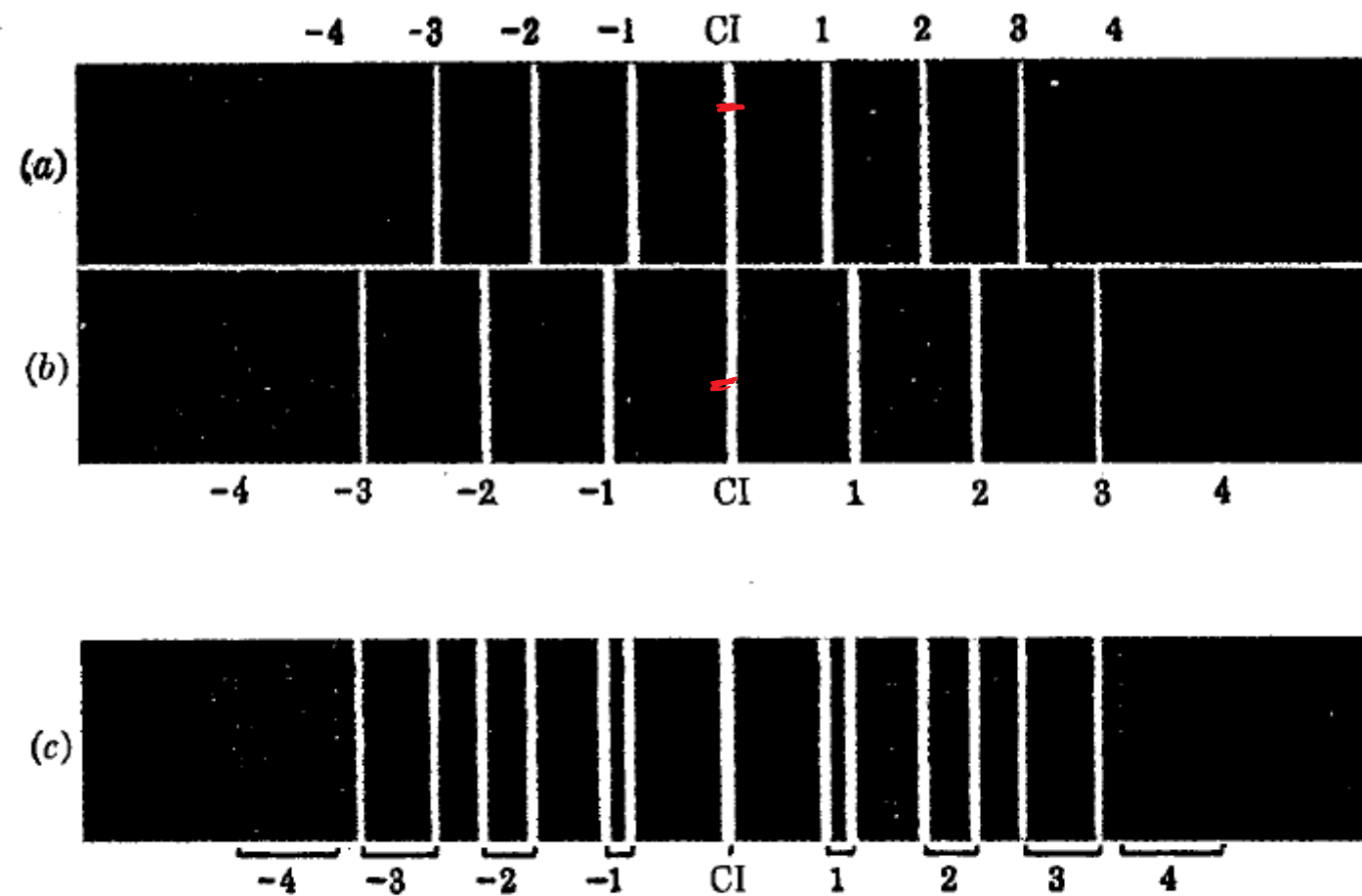
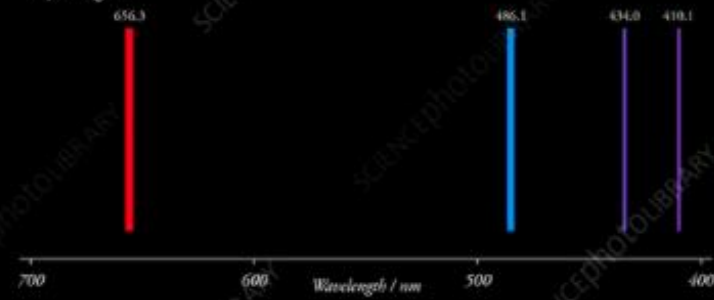


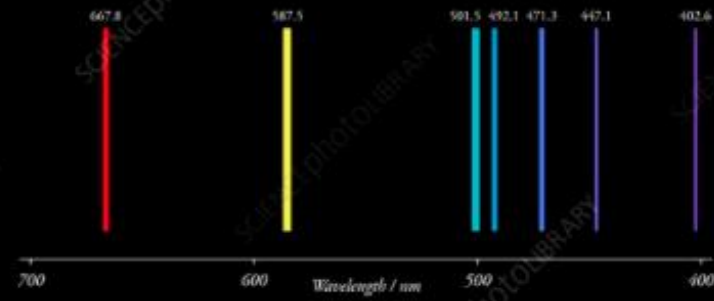
FIGURE 17F

Grating spectra of two wavelengths: (a) $\lambda_1 = 4000 \text{ \AA}$; (b) $\lambda_2 = 5000 \text{ \AA}$; (c) λ_1 and λ_2 together.

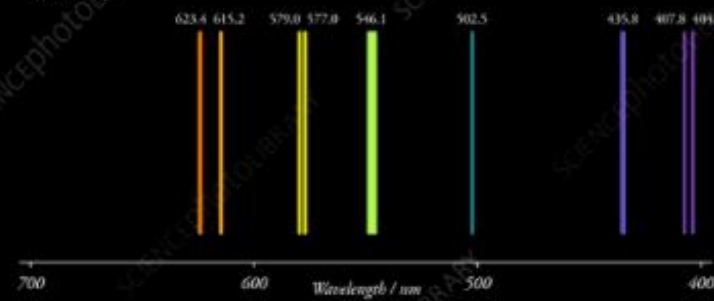
Hydrogen

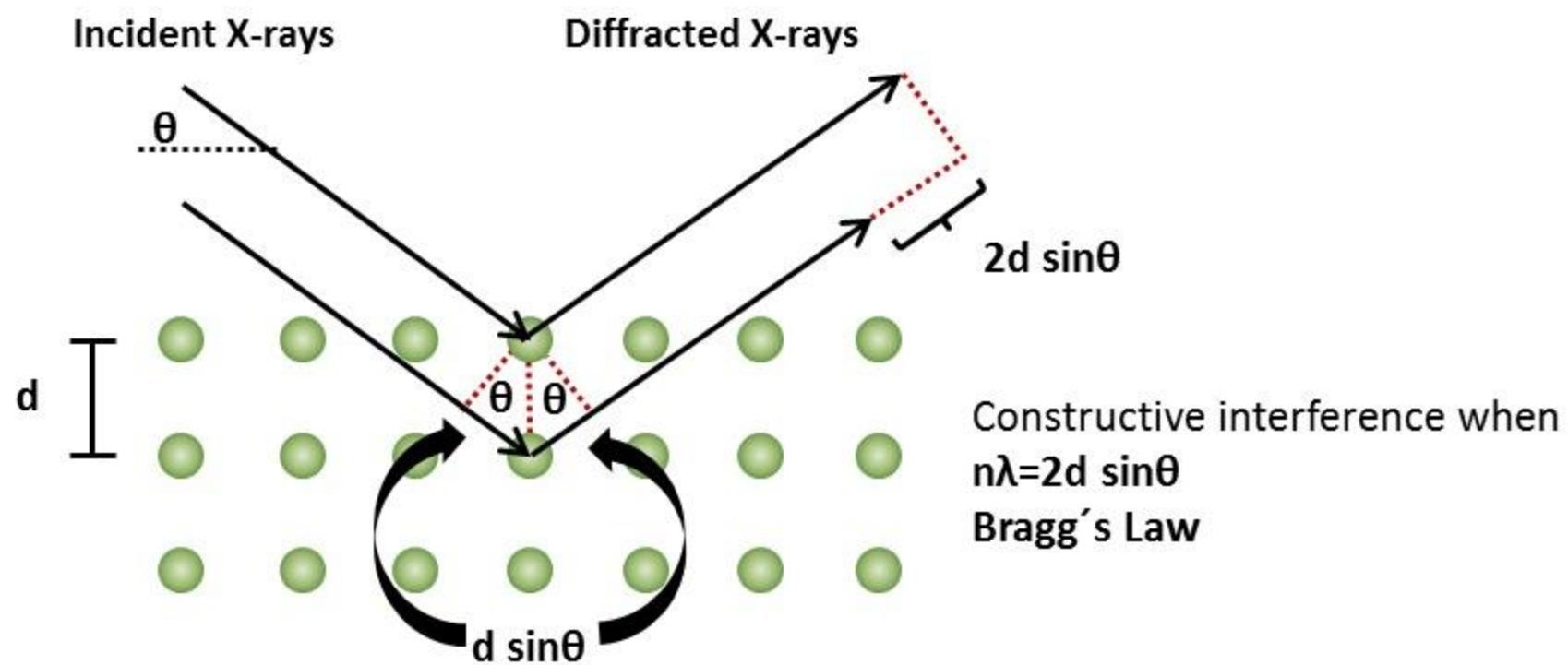


Helium



Mercury





Alumina Powder Diffraction

