

Applications of the canonical distribution

Ex 2. Molecule in an ideal gas

Monatomic gas at temp T, vol V .

molecule : system

reservoir : rest of the gas

$$E = \frac{1}{2}mv^2 = \frac{\vec{p}^2}{2m}$$

molecule has position between \vec{r} & $\vec{r} + d\vec{r}$
momentum between \vec{p} & $\vec{p} + d\vec{p}$

volume of phase space corresponding to this

$$d^3\vec{r} d^3\vec{p} = dx dy dz dp_x dp_y dp_z$$

$$P(\vec{r}, \vec{p}) d^3\vec{r} d^3\vec{p} \propto \left(\frac{d^3\vec{r} d^3\vec{p}}{h_0^3} \right) e^{-\frac{\beta p^2}{2m}}$$

prob. that molecule
be found in this
range

Prob. density .

$P(\vec{p}) d^3 \vec{p} \rightarrow$ Prob that molecule has momentum between \vec{p} & $\vec{p} + d\vec{p}$ (irrespective of location)

$P'(\vec{v}) d^3 \vec{v} \rightarrow \dots$ velocity between \vec{v} & $\vec{v} + d\vec{v}$

$P'(\vec{v}) d^3 \vec{v} \Rightarrow$ integrate over all values of position

$$\iint \underbrace{d^3 \vec{r}}_{=V}$$

$$P'(\vec{v}) d^3 \vec{v} = C e^{-\frac{\beta m v^2}{2}} d^3 \vec{v}$$

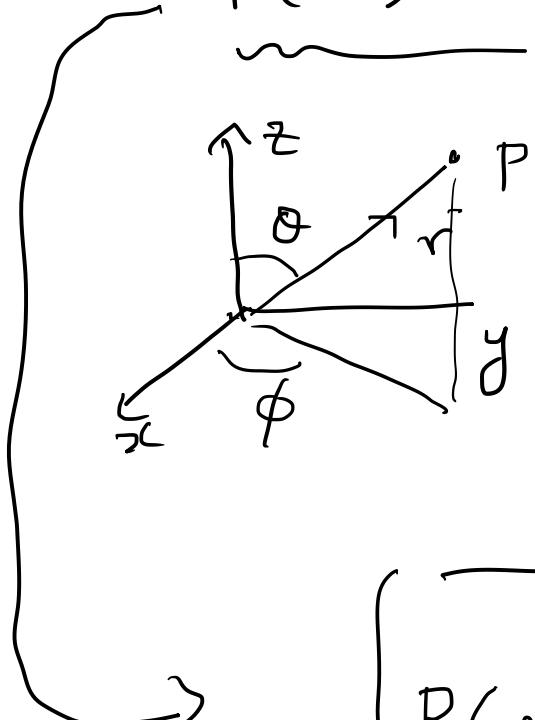
Maxwell's
velocity distribution

v factor

Speed distribution \rightarrow irrespective of direction

molecule has speed between v & $v + dv$

$$P(v)dv = C \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} e^{-\beta mv^2/2} v^2 dv \sin\theta d\theta d\phi$$



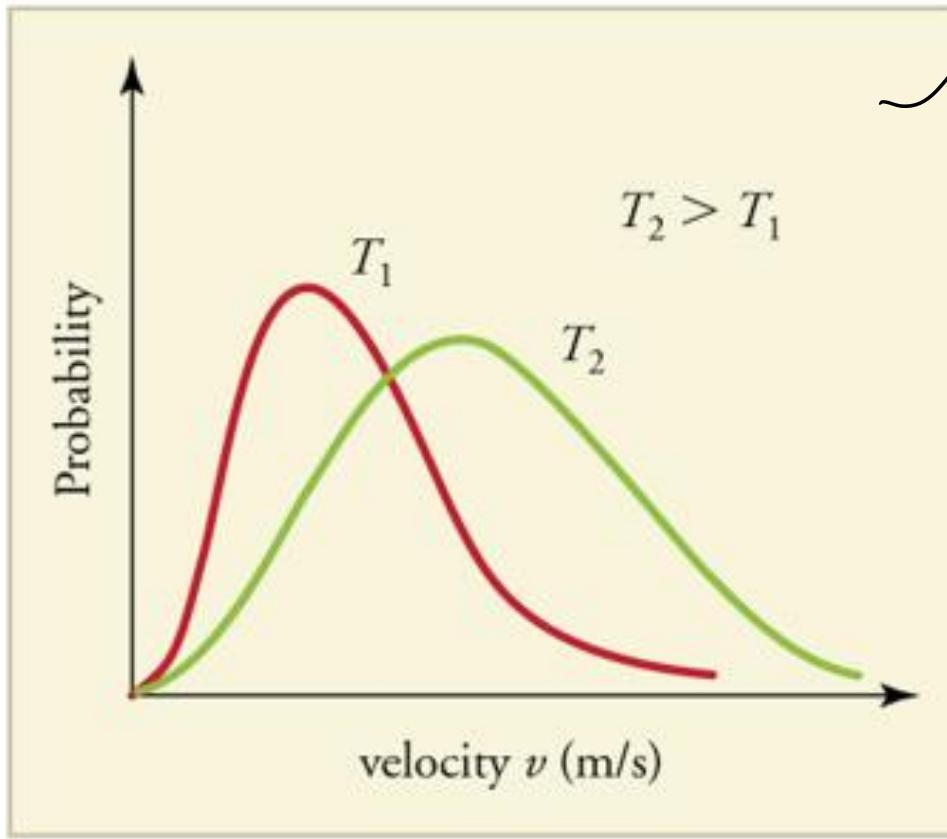
$$d^3 \vec{r} = \underbrace{r^2 dr \sin\theta d\theta d\phi}_{dx dy dz}$$

$$d^3 \vec{v} = v^2 dv \sin\theta d\theta d\phi$$

$$\left. \begin{aligned} \int \sin\theta d\theta d\phi \\ = 4\pi \\ C' = 4\pi C \end{aligned} \right\}$$

$$P(v)dv = C' e^{-\beta mv^2/2} v^2 dv$$

speed dist.



$$e^{-\frac{\beta m v^2}{2}} v^2$$

$$e^{-\frac{m v^2}{2kT}} v^2$$

Example 3 : Thermal excitation of atoms.

Hydrogen atom in the sun, $T = 5800\text{ K}$.

$P(s_2)$: prob that the atom is in $n=2$ state
(first excited state)

$P(s_1)$: prob. is in $n=1$ state
ground state.

$$\frac{P(s_2)}{P(s_1)} = \frac{e^{-E_2/kT}}{e^{-E_1/kT}} = e^{-\frac{(E_2 - E_1)}{kT}}$$

$$\Delta E = 10.2 \text{ eV}$$

$$kT = (8.62 \times 10^{-5} \text{ eV}/\text{k}) (5800 \text{ K}) \\ = 0.5 \text{ eV}$$

$$\frac{P(s_2)}{P(s_1)} = e^{-20.4} \approx 1.4 \times 10^{-9}$$

} 1.4 atoms in
a billion will
be in an excited
state.

$n=1$ \rightarrow only 1 state with E_1

$n=2$ \rightarrow has 4 states with E_2 .

Total # of atoms in $n=2$ state $\sim 1.4 \times 4 \times 10^{-9}$
 ~ 5.6 .

Atoms in atm of sun absorb sunlight on way to earth

→ induces transition to excited states .

1st excited 656 nm , 486 nm , 434 nm → Balmer series

\downarrow $n=2$. → Solar spectrum these appear as dark lines

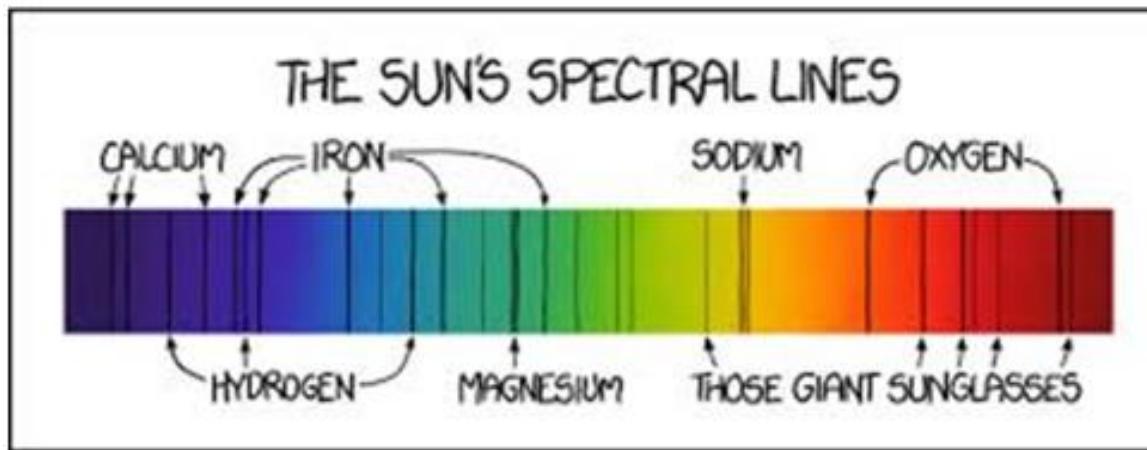
Absorption bands . also have (Fe, Mg, Na, Ca) etc .

Weird thing : all the other wavelengths that are absorbed by atoms start out in states $< 3 \text{ eV}$ from ground state .

→ These lines ought to be more prominent .

- Fact is Balmer lines are quite prominent among others.

\Rightarrow H is much more abundant in the sun
than other atoms !!



Calculation of mean values in a canonical ensemble

$$P_r = \frac{e^{-\beta E_r}}{\sum_r e^{-\beta E_r}}$$

$$\text{Mean energy} = \bar{E} = \frac{\sum_r E_r e^{-\beta E_r}}{\sum_r e^{-\beta E_r}}$$

Sums are over all accessible states

Can be expressed in a much more compact form using the partition function

$$Z = \sum_r e^{-\beta E_r}.$$

$$\sum_r E_r e^{-\beta E_r} = - \frac{\partial}{\partial \beta} \sum_r e^{-\beta E_r} = - \frac{\partial Z}{\partial \beta}.$$

$$\boxed{\bar{E} = - \frac{1}{Z} \frac{\partial Z}{\partial \beta}}$$

$$\boxed{\bar{E} = - \frac{1}{\partial \beta} \ln Z}$$

Dispersion

$$\overline{(\Delta E)^2} = \overline{E^2} - (\bar{E})^2.$$

$$\overline{E^2} = \frac{\sum E_r^2 e^{-\beta E_r}}{\sum e^{-\beta E_r}}.$$

$$= \frac{1}{Z} \frac{\partial^2 Z}{\partial \beta^2}.$$

$$\bar{E}^2 = \frac{\partial}{\partial \beta} \left(\frac{1}{Z} \frac{\partial Z}{\partial \beta} \right) + \frac{1}{Z^2} \left(\frac{\partial Z}{\partial \beta} \right)^2$$

$$\overline{E^2} = \frac{\partial}{\partial \beta} \left(\frac{1}{Z} \frac{\partial Z}{\partial \beta} \right) + \frac{1}{Z^2} \left(\frac{\partial Z}{\partial \beta} \right)^2$$

$$= - \frac{\partial \overline{E}}{\partial \beta} + (\overline{E})^2 .$$

$$\overline{(\Delta E)^2} = \overline{E^2} - (\overline{E})^2$$

$$\boxed{(\Delta E)^2 = - \frac{\partial \overline{E}}{\partial \beta} = - \frac{\partial^2}{\partial \beta^2} \ln Z} \quad > 0$$

$-\frac{\partial \overline{E}}{\partial \beta} > 0$
 $\frac{\partial \overline{E}}{\partial T} > 0$
 $\frac{\partial \overline{E}}{\partial P} > 0$