

## Table 3.5 The Maxwell relations

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From  $U$ : 
$$\left(\frac{\partial T}{\partial V}\right)_S = - \left(\frac{\partial p}{\partial S}\right)_V$$

From  $H$ : 
$$\left(\frac{\partial T}{\partial p}\right)_S = \left(\frac{\partial V}{\partial S}\right)_p$$

From ~~F~~  
~~A~~: 
$$\left(\frac{\partial p}{\partial T}\right)_V = \left(\frac{\partial S}{\partial V}\right)_T$$

From  $G$ : 
$$\left(\frac{\partial V}{\partial T}\right)_p = - \left(\frac{\partial S}{\partial p}\right)_T$$

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## TdS equations

$$dS = \left(\frac{\partial S}{\partial T}\right)_V dT + \left(\frac{\partial S}{\partial V}\right)_T dV$$

$$TdS = T\left(\frac{\partial S}{\partial T}\right)_V dT + T\left(\frac{\partial S}{\partial V}\right)_T dV$$

Consider reversible isochoric change

$$TdS = (dTQ)_{V=\text{const}} \quad T\left(\frac{\partial S}{\partial T}\right)_V = C_V$$

$$TdS = C_V dT + T\left(\frac{\partial P}{\partial T}\right)_V dV$$

1st TdS eqn.

; Maxwell's reln

$$\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial P}{\partial T}\right)_V$$

$$dS = \left(\frac{\partial S}{\partial T}\right)_P dT + \left(\frac{\partial S}{\partial P}\right)_T dP$$

$$TdS = T \underbrace{\left(\frac{\partial S}{\partial T}\right)_P}_{dT} + T \underbrace{\left(\frac{\partial S}{\partial P}\right)_T}_{dP} \cdot$$

$$\boxed{TdS = C_P dT - T \left(\frac{\partial V}{\partial T}\right)_P dP}$$

$$\left\{ \left(\frac{\partial S}{\partial P}\right)_T = - \left(\frac{\partial V}{\partial T}\right)_P \right.$$

$\hookrightarrow$  2nd TdS eqn.

## Internal Energy eqns

$$dU = TdS - PdV$$

$$\frac{dU}{dP} = T \frac{dS}{dP} - P \frac{dV}{dP}$$

Holding T const

$$\left(\frac{\partial U}{\partial P}\right)_T = T \left(\frac{\partial S}{\partial P}\right)_T - P \left(\frac{\partial V}{\partial P}\right)_T$$

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Maxwell

$$\left(\frac{\partial S}{\partial P}\right)_T = - \left(\frac{\partial V}{\partial T}\right)_P$$

$$\boxed{\left(\frac{\partial U}{\partial P}\right)_T = - T \left(\frac{\partial V}{\partial T}\right)_P - P \left(\frac{\partial V}{\partial P}\right)_T}$$

→ 2<sup>nd</sup> internal energy identity

$$dU = TdS - PdV$$

$$\frac{dU}{dV} = T \frac{\partial S}{\partial V} - P$$

$$\left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial S}{\partial V}\right)_T - P$$

$$\boxed{\left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial P}{\partial V}\right)_T - P}$$

$$\left\{ \left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial P}{\partial T}\right)_V \right.$$

Maxwell

## Heat Capacity Equation

$$TdS = C_P dT - T \left( \frac{\partial V}{\partial T} \right)_P dP \quad \text{--- (1)}$$

$$TdS = C_V dT + T \left( \frac{\partial P}{\partial T} \right)_V dV \quad \text{--- (2)}$$

Equating (1) & (2)

$$C_P dT - T \left( \frac{\partial V}{\partial T} \right)_P dP = C_V dT + T \left( \frac{\partial P}{\partial T} \right)_V dV$$

$$dT = \frac{T \left( \frac{\partial P}{\partial T} \right)_V dV}{C_P - C_V} + \frac{T \left( \frac{\partial V}{\partial T} \right)_P dP}{C_P - C_V} \quad \text{--- (3)}$$

But

$$dT = \left( \frac{\partial T}{\partial V} \right)_P dT + \left( \frac{\partial T}{\partial P} \right)_V dT \quad \text{--- (4)}$$

$$\Rightarrow \left( \frac{\partial T}{\partial V} \right)_P = \frac{T \left( \frac{\partial P}{\partial T} \right)_V}{C_P - C_V} ; \quad \left( \frac{\partial T}{\partial P} \right)_V = \frac{T \left( \frac{\partial V}{\partial T} \right)_P}{C_P - C_V} .$$

$$C_P - C_V = T \left( \frac{\partial P}{\partial T} \right)_V \left( \frac{\partial V}{\partial T} \right)_P$$

But  $\left( \frac{\partial P}{\partial T} \right)_V = - \left( \frac{\partial V}{\partial T} \right)_P \left( \frac{\partial P}{\partial V} \right)_T \rightarrow$  chain rule

$$C_P - C_V = -T \left( \frac{\partial V}{\partial T} \right)_P^2 \left( \frac{\partial P}{\partial V} \right)_T$$

$$C_p - C_v = -T \left( \frac{\partial V}{\partial T} \right)_P^2 \left( \frac{\partial P}{\partial V} \right)_T$$

•  $\left( \frac{\partial P}{\partial V} \right)_T < 0$  for all known substances.

$$\left( \frac{\partial V}{\partial T} \right)^2 > 0 \quad \therefore \boxed{C_p - C_v > 0}$$

•  $T \rightarrow 0 \quad C_p \rightarrow C_v \quad \text{at absolute zero} \quad C_p = C_v$

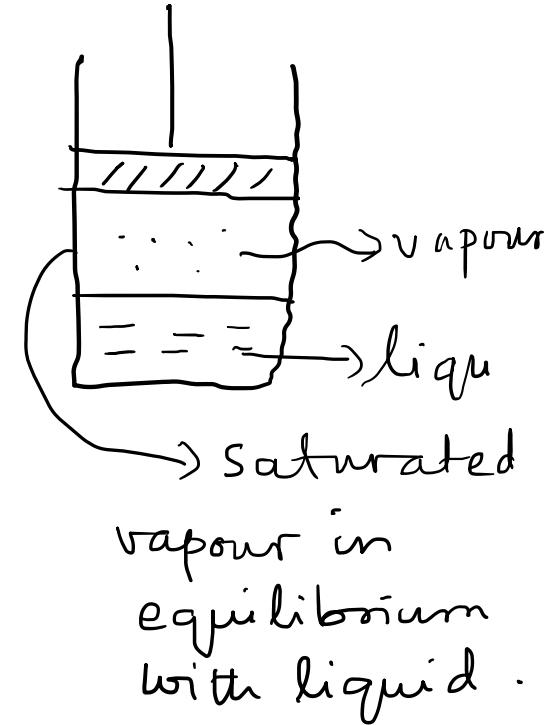
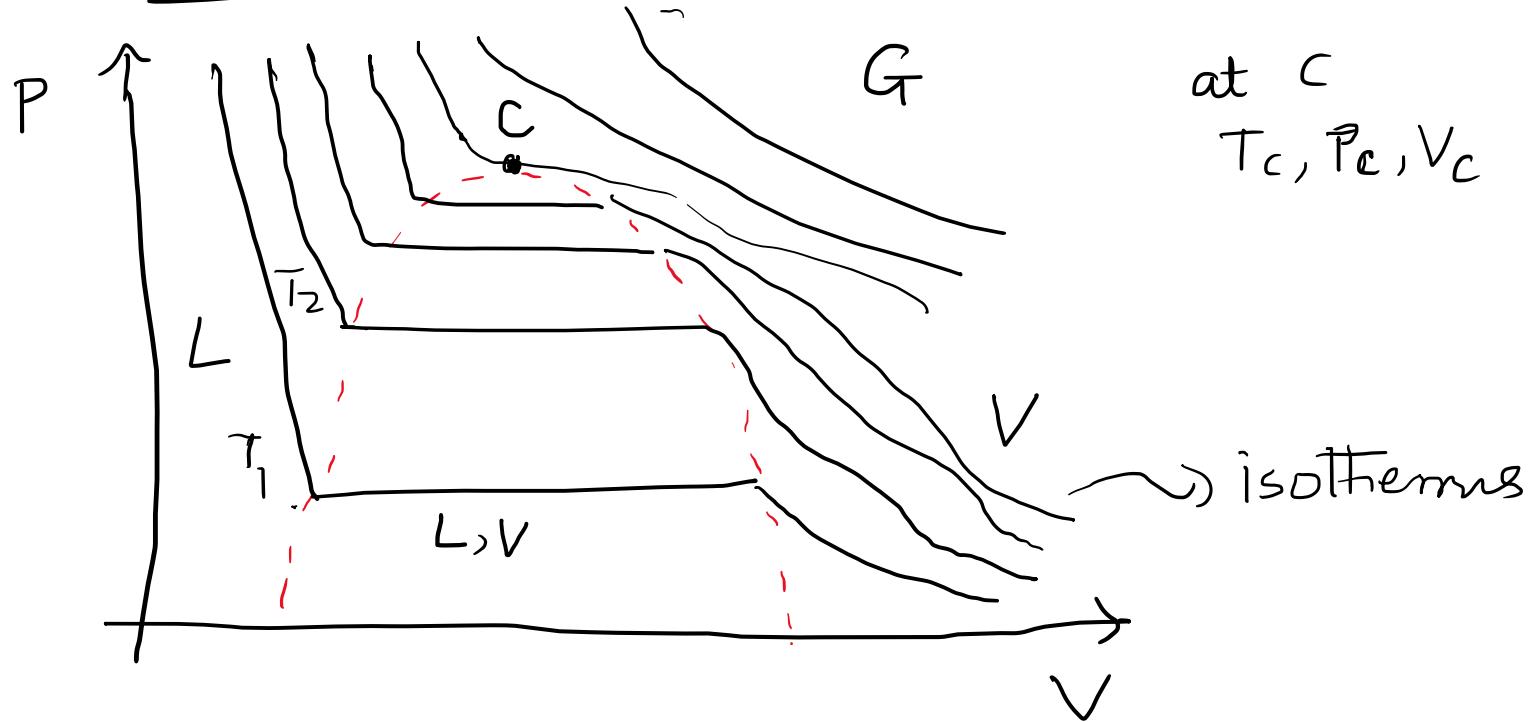
•  $C_p = C_v$  when  $\left( \frac{\partial V}{\partial T} \right)_P = 0$  at  $4^\circ C$   $P_w \max$

$$C_p - C_v = -T \frac{\left(\frac{\partial V}{\partial T}\right)_P^2}{\left(\frac{\partial V}{\partial P}\right)_T}$$

$$\boxed{C_p - C_v = \frac{T v \beta^2}{u}}$$

$v$  : specific volume .

## Gibbs potential application



We will focus on flat region

$$m = m_1 + m_2$$

1 subscript liq.  
2 " vapour

$$V = m_1 v_1(T) + m_2 v_2(T) \quad v : \text{specific volume}$$

$$U = m_1 u_1(T) + m_2 u_2(T)$$

$$\left. \begin{array}{l} U = U_1 + U_2 \\ S = S_1 + S_2 \\ V = V_1 + V_2 \end{array} \right\} \quad \begin{array}{l} G = U - TS + PV \\ T, P \text{ const region} \\ G = G_1 + G_2 \end{array}$$

$$\left. \begin{array}{l} G_1 = m_1 g_1 \\ G_2 = m_2 g_2 \end{array} \right\} \quad \begin{array}{l} P, T \text{ const} \\ m_1, m_2 \text{ can vary} \end{array}$$

$$m_1 \rightarrow m + dm$$

$$m_2 \rightarrow m - dm$$

$$(m_1 + dm) g_1 + (m_2 - dm) g_2 = G + dG$$

$$dG = dm(g_1 - g_2) \quad \text{In equilibrium}$$

$$\Rightarrow dG = 0 \Rightarrow \boxed{g_1 = g_2}$$

$$u_1 - Ts_1 + Pv_1 = u_2 - Ts_2 + Pv_2$$

$$(u_2 - u_1) - T(s_2 - s_1) + P(v_2 - v_1) = 0$$

$$\begin{aligned} \frac{d}{dT}(u_2 - u_1) - T \frac{d}{dT}(s_2 - s_1) - (s_2 - s_1) + \frac{dP}{dT}(v_2 - v_1) \\ + P \frac{d}{dT}(v_2 - v_1) = 0 \end{aligned}$$

But 1st Law

$$T \frac{dS}{dT} = \frac{du}{dT} + P \frac{dV}{dT}$$

$$\frac{d}{dT} (u_2 - u_1) - T \frac{d}{dT} (s_2 - s_1) - (s_2 - s_1) + \frac{dP}{dT} (v_2 - v_1) + P \frac{d}{dT} (v_2 - v_1) = 0$$



$$- (s_2 - s_1) + \frac{dP}{dT} (v_2 - v_1) = 0$$

$$\boxed{\left( \frac{dP}{dT} \right) = \frac{\lambda}{T(v_2 - v_1)}}$$

But  $(s_2 - s_1) = \frac{\lambda}{T}$

latent heat

Clausius - Clapeyron  
eqn.

