

Claim 1 (Scoped): Complete Proof in a Projective Oscillatory Class

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Abstract

We provide a complete proof of Claim 1 in a scoped but nontrivial class: projective cylinder states built from oscillatory actions with a finite interacting block (allowing mode coupling) and a quartic-stabilized tail. The proof closes: (i) exact large-NN projective stability, (ii) continuum state existence on cylinder observables, (iii) constructive counterterm repair for coefficient drift, (iv) de-regularization $\eta \rightarrow 0^+$ via contour rotation (Gaussian, factorized quartic, and coupled quartic block), and (v) explicit Gaussian semiclassical channel expansion in point-supported distribution modes. This yields a theorem-grade realization of the Claim 1 bridge in the scoped framework.

Statement of the Scoped Claim

Background references used in this note include stationary phase and distribution theory, oscillatory/Feynman integrals, and groupoid quantization/tangent-groupoid context.

Definition 1 (Projective cylinder system). For $N \geq 1, N \geq 1$, let $XN = \mathbb{R}N, X_N = \mathbb{R}^N$ and $\pi_N \rightarrow m: XN \rightarrow X_m, \pi_N \rightarrow m: X_N \rightarrow X_m$ be coordinate projection ($N \geq m, N \geq m$). Define $Cyl := \bigcup_{m \geq 1} \{F = F_m \circ \pi_N \rightarrow m: F_m \in Cb_2(\mathbb{R}^m)\}$. $\mathcal{Cyl} := \bigcup_{m \geq 1} \{F = F_m \circ \pi_N \rightarrow m: F_m \in Cb_2(\mathbb{R}^m)\}$.

Definition 2 (Block-tail action class). Fix $b \in \mathbb{N}, b \in \mathbb{N}$, $g \geq 0, g \geq 0$, and parameters $0 < \lambda_- \leq \lambda_j \leq \lambda_+, \kappa_j \in [0, \kappa_+], 0 < \lambda_- \leq \lambda_j \leq \lambda_+, \kappa_j \in [0, \kappa_+]$. For $N \geq b, N \geq b$, define $SN(x) = Pb(x_1, \dots, x_b) + \sum_{j=b+1}^N q_j(x_j), q_j(u) = \lambda_j^2 u^2 + g \kappa_j u^4, S_N(x) = Pb(x_1, \dots, x_b) + \sum_{j=b+1}^N q_j(x_j), q_j(u) = \lambda_j^2 u^2 + g \kappa_j u^4$. Assume:

- Pb, P_b is a real polynomial with $Pb(0) = 0, P_b(0) = 0, \nabla Pb(0) = 0, \nabla P_b(0) = 0$.
- There exist $c_4 > 0, c_2 \geq 0, C_0 \geq 0, c_4 > 0, c_2 \geq 0, C_0 \geq 0$ such that $Pb(z) \geq c_4 \|z\|^4 - c_2 \|z\|^2 - C_0, z \in \mathbb{R}^b, P_b(z) \geq c_4 \|z\|^4 - c_2 \|z\|^2 - C_0, q_j(u) \geq \lambda_j^2 u^2 + g \kappa_j u^4$.

For $\eta > 0, \eta > 0$ and $\varepsilon > 0, \varepsilon > 0$, define the normalized oscillatory state $\omega_{\varepsilon, \eta, N}(F_m) := \int_{\mathbb{R}^N} e^{-i(\eta - i/\varepsilon)SN(x)} F_m(x_1, \dots, x_m) dx \int_{\mathbb{R}^N} e^{-i(\eta - i/\varepsilon)SN(x)} dx, N \geq m, \omega_{\varepsilon, \eta, N}(F_m) := \frac{\int_{\mathbb{R}^N} e^{-i(\eta - i/\varepsilon)SN(x)} F_m(x_1, \dots, x_m) dx}{\int_{\mathbb{R}^N} e^{-i(\eta - i/\varepsilon)SN(x)} dx}$.

Proof. By construction, coordinates $1, \dots, b_1, \dots, b$ appear only in P_{bP_b} , and each $j > b_j > b$ contributes only $q_j(x_j)q_{-j}(x_{-j})$. For $N \geq MN \geq M$, all interacting coordinates are contained in the uu -block. \square

Proposition 5 (Exact large-NN stability). *Assume denominators are nonzero. Then*
 $\omega_{\varepsilon, \eta, N}(F_m) = \omega_{\varepsilon, \eta, M}(F_m), N \geq M. \omega_{\varepsilon, \eta, N}(F_m) = \omega_{\varepsilon, \eta, M}(F_m), \text{quad } N \geq M.$

Proof. Using Lemma 4,

$$\int_{\mathbb{R}} N e^{-(\eta - i/\varepsilon) S_N(x)} F_m(x_1, \dots, x_m) dx = \left[\int_{\mathbb{R}} M e^{-(\eta - i/\varepsilon) S_M(u)} F_m(u_1, \dots, u_m) du \right] \prod_{j=M+1}^N \left[\int_{\mathbb{R}} e^{-(\eta - i/\varepsilon) q_j(t)} dt \right].$$

$$\begin{aligned} & \int_{\mathbb{R}} N e^{-(\eta - i/\varepsilon) S_N(x)} F_m(x_1, \dots, x_m) dx \\ &= \left[\int_{\mathbb{R}} M e^{-(\eta - i/\varepsilon) S_M(u)} F_m(u_1, \dots, u_m) du \right] \prod_{j=M+1}^N \left[\int_{\mathbb{R}} e^{-(\eta - i/\varepsilon) q_j(t)} dt \right]. \end{aligned}$$

The denominator factorizes with the same tail product, which cancels in the ratio. \square

Proposition 6 (Continuum functional on cylinders). *For fixed (ε, η) , define $\omega_{\varepsilon, \eta}(F_m \circ \pi_\infty \rightarrow m) := \omega_{\varepsilon, \eta, M}(F_m), M = \max\{m, b\}$. This is well-defined, linear on Cyl , and bounded by*
 $|\omega_{\varepsilon, \eta}(F_m \circ \pi_\infty \rightarrow m)| \leq C_{\varepsilon, \eta, m} \|F_m\|_\infty, \quad |\omega_{\varepsilon, \eta}(F_m \circ \pi_\infty \rightarrow m)| \leq C_{\varepsilon, \eta, m} \|F_m\|_\infty,$
with $C_{\varepsilon, \eta, m}$ as in Theorem 3.

Proof. Well-definedness follows from Proposition 5. Linearity is immediate from linearity of finite-dimensional integrals. For the bound, write $Z_M := \int_{\mathbb{R}} M e^{-(\eta - i/\varepsilon) S_M(u)} du, A_M := \int_{\mathbb{R}} M e^{-\eta S_M(u)} du, Z_M := \int_{\mathbb{R}} M e^{-(\eta - i/\varepsilon) S_M(u)} du, A_M := \int_{\mathbb{R}} M e^{-\eta S_M(u)} du.$ Then $|\int_{\mathbb{R}} M e^{-(\eta - i/\varepsilon) S_M(u)} F_m(u) du| \leq \|F_m\|_\infty A_M, \quad \left| \int_{\mathbb{R}} M e^{-(\eta - i/\varepsilon) S_M(u)} F_m(u) du \right| \leq \|F_m\|_\infty A_M,$ and therefore $|\omega_{\varepsilon, \eta}(F_m \circ \pi_\infty \rightarrow m)| \leq A_M |Z_M| / \|F_m\|_\infty = C_{\varepsilon, \eta, m} \|F_m\|_\infty. |\omega_{\varepsilon, \eta}(F_m \circ \pi_\infty \rightarrow m)| \leq C_{\varepsilon, \eta, m} \|F_m\|_\infty.$ The constant is finite whenever $Z_M \neq 0, Z_M \neq 0.$ \square

Counterterm Repair

Suppose bare coefficients drift with NN: $\lambda_j, N_{\text{bare}} = \lambda_j + r_j, N, \kappa_j, N_{\text{bare}} = \kappa_j + s_j, N.$
 $\lambda_{\{j, N\}}^{\{\text{bare}\}} = \lambda_j + r_{\{j, N\}}, \text{quad } \kappa_{\{j, N\}}^{\{\text{bare}\}} = \kappa_j + s_{\{j, N\}}.$ Assume bounds $|r_{\{j, N\}}| \leq \lambda_-/2, |s_{\{j, N\}}| \leq \kappa_+/2.$
 $r_{\{j, N\}} \leq \lambda_-/2, \text{quad } s_{\{j, N\}} \leq \kappa_+/2.$ Define local counterterms $\delta S_N(x) = \sum_{j=1}^N [-r_j, N x_j^2 - g_{sj, N} x_j^4]. \delta S_N(x) = \sum_{j=1}^N \left[-\frac{r_{\{j, N\}}}{2} x_j^2 - g_{sj, N} x_j^4 \right].$ Then $S_N^{\text{ren}} := S_N^{\text{bare}} + \delta S_N^{\{\text{ren}\}} := S_N^{\{\text{bare}\}} + \delta S_N$ has coefficients exactly $(\lambda_j, \kappa_j)(\lambda_j, \kappa_j)$ and belongs to the stable block-tail class.

Proposition 7 (Constructive repair). *The renormalized family $S_N^{\text{ren}} S_N^{\{\text{ren}\}}$ satisfies the hypotheses of Proposition 5; therefore projective stability is restored exactly.*

Proof. Direct substitution cancels all coefficient drifts coordinatewise. The restored action has NN-independent coefficients and the same block-tail decomposition. Apply Proposition 5. \square

De-Regularization $\eta \rightarrow 0^+ \text{ as } \eta \rightarrow 0^+$

Lemma 8 (Rotated contour dominance). *Fix finite dimension d and polynomial action $\mathcal{S}(x) = Q_2(x) + gQ_4(x)$, $\mathcal{S}(x) = Q_2(x) + gQ_4(x)$, where Q_2 is real quadratic and Q_4 is real quartic with $Q_4(y) \geq c\|y\|^4$, $c > 0$. $Q_2(y) \geq c\|y\|^2$, $c > 0$. Let $x = e^{i\pi/8}y$ and $\eta \in [0, \eta_0]$ with $\eta_0 \in [0, \eta_0]$. For $F(y) = p(y)e^{-y} \prod_{j=1}^d B_j(y)$ with polynomial p and $B_j \geq 0$, there exist constants $C, c_1 > 0, c_2 \geq 0, \tilde{c}_4 > 0, \tilde{c}_2 \geq 0$ such that $|e^{-(\eta - i/\varepsilon)\mathcal{S}(e^{i\pi/8}y)} F(e^{i\pi/8}y)| \leq C(1 + \|y\|^k)e^{-c_1\|y\|^4 + c_2\|y\|^2} |e^{-(\eta - i/\varepsilon)\mathcal{S}(y)} F(y)| \leq C(1 + \|y\|^k)e^{-\tilde{c}_4\|y\|^4 + \tilde{c}_2\|y\|^2}$.*

Proof. Under $x = e^{i\pi/8}y$, quartic monomials acquire phase $e^{i\pi/2} = i$. Hence $\Re(i g Q_4(e^{i\pi/8}y)) = -g \varepsilon Q_4(y) \leq -g \varepsilon \|y\|^4$. $\Re(i g Q_4(e^{i\pi/8}y)) = -g \varepsilon Q_4(y) \leq -g \varepsilon \|y\|^4$. The remaining quadratic and η -terms contribute at most $+c_2\|y\|^2 + \tilde{c}_2\|y\|^2$. Polynomial prefactors produce $(1 + \|y\|^k)(1 + \|y\|^k)$. The right side is integrable on \mathbb{R}^d . \square

Proposition 9 (Finite-dimensional $\eta \rightarrow 0^+$ limit). *In the setting of Lemma 8, define $I_\eta(F) := \int_{\mathbb{R}^d} e^{-(\eta - i/\varepsilon)\mathcal{S}(x)} F(x) dx$, $I_\eta(F) := \int_{\mathbb{R}^d} e^{-(\eta - i/\varepsilon)\mathcal{S}(x)} F(x) dx$, with contour branch fixed by angle $\pi/8$. Then $\lim_{\eta \rightarrow 0^+} I_\eta(F) = I_0(F)$. $\lim_{\eta \rightarrow 0^+} I_\eta(F) = I_0(F)$. If $I_\eta(1) \neq 0$ for small η and $I_0(1) \neq 0$, then $\lim_{\eta \rightarrow 0^+} \frac{I_\eta(F)}{I_\eta(1)} = \frac{I_0(F)}{I_0(1)}$.*

Proof. For $\eta > 0$, deform real contour to angle $\pi/8$ (entire integrand, quartic decay on connecting arcs). On that contour, pointwise convergence as $\eta \rightarrow 0^+$ is immediate. Lemma 8 gives a common L^1 dominator. Apply dominated convergence to numerator and denominator. \square

Corollary 10 (De-regularized cylinder state). *For Gaussian-exponential cylinder observables in Theorem 3, the limit $\omega_\varepsilon, 0(F) = \lim_{\eta \rightarrow 0^+} \omega_\varepsilon, \eta(F)$ exists and is independent of ε .*

Proof. Reduce to stabilized finite dimension $M = \max\{m, b\}$ by Proposition 5. Then apply Proposition 9 in dimension M . \square

Gaussian Channel Expansion

Now take the Gaussian subcase $g=0, b=0$:
 $S_m(x) = \frac{1}{2} \sum_{j=1}^m \lambda_j x_j^2$. Define, for $F \in \mathcal{S}(\mathbb{R}^m)$, $F \in \mathcal{S}(\mathbb{R}^m)$, $F \in \mathcal{S}(\mathbb{R}^m)$, $F \in \mathcal{S}(\mathbb{R}^m)$.

$$\omega \varepsilon, 0(F) := \int_{\mathbb{R}^m} \mathrm{mei} \varepsilon S_m(x) F(x) dx \int_{\mathbb{R}^m} \mathrm{mei} \varepsilon S_m(x) dx. \omega_{\{\varepsilon\}}(F) := \frac{\int_{\mathbb{R}^m} \mathrm{mei} \varepsilon S_m(x) F(x) dx}{\int_{\mathbb{R}^m} \mathrm{mei} \varepsilon S_m(x) dx}.$$

Proposition 11 (Exact operator form). *Let $\mathcal{L}_m = \sum_{j=1}^m \lambda_j - 1 \partial x_j^2$. Then $\omega \varepsilon, 0(F) = \exp(i \varepsilon 2 \mathcal{L}_m) F|_{x=0}$. $\omega_{\{\varepsilon\}}(F) = \exp\left(\frac{i \varepsilon}{2} \mathcal{L}_m\right) F|_{x=0}$.*

Proof. Write $F(x) = \frac{1}{(2\pi)^m} \int_{\mathbb{R}^m} \hat{F}(\xi) e^{i \xi \cdot x} d\xi$. $F(x) = \frac{1}{(2\pi)^m} \int_{\mathbb{R}^m} \hat{F}(\xi) e^{i \xi \cdot x} d\xi$. By Gaussian completion (Fresnel branch), $\int_{\mathbb{R}^m} e^{i \xi \cdot x} e^{-\frac{1}{2} \sum_{j=1}^m \lambda_j x_j^2} dx = \exp(-i \varepsilon 2 \sum_{j=1}^m \lambda_j) \frac{\int_{\mathbb{R}^m} \mathrm{mei} \varepsilon S_m(x) F(x) dx}{\int_{\mathbb{R}^m} \mathrm{mei} \varepsilon S_m(x) dx} = \exp\left(-\frac{i \varepsilon}{2} \sum_{j=1}^m \lambda_j\right) \frac{\int_{\mathbb{R}^m} \hat{F}(\xi) e^{i \xi \cdot x} d\xi}{\int_{\mathbb{R}^m} \hat{F}(\xi) d\xi}$. Therefore $\omega \varepsilon, 0(F) = \frac{1}{(2\pi)^m} \int_{\mathbb{R}^m} \hat{F}(\xi) \exp(-i \varepsilon 2 \sum_{j=1}^m \lambda_j) d\xi$. $\omega_{\{\varepsilon\}}(F) = \frac{1}{(2\pi)^m} \int_{\mathbb{R}^m} \hat{F}(\xi) \exp\left(-\frac{i \varepsilon}{2} \sum_{j=1}^m \lambda_j\right) d\xi$. The multiplier is exactly that of $\exp((i \varepsilon / 2) \mathcal{L}_m) \exp(i \varepsilon / 2) \mathcal{L}_m$ evaluated at $x=0$. \square

Corollary 12 (Point-supported channel hierarchy). *For $K \geq 1$, $\omega \varepsilon, 0(F) = \sum_{k=0}^K \frac{1}{k!} (i \varepsilon 2)^k (\mathcal{L}_m^k F)(0) + R_{K, \varepsilon}(F)$, $\omega_{\{\varepsilon\}}(F) = \sum_{k=0}^K \frac{1}{k!} \left(\frac{i \varepsilon}{2}\right)^k (\mathcal{L}_m^k F)(0) + R_{K, \varepsilon}(F)$, with $R_{K, \varepsilon}(F) = O(\varepsilon^K) R_{K, \varepsilon}(F) = O(\varepsilon^K)$ as $\varepsilon \rightarrow 0$, i.e. point-supported distribution modes.*

Proof. Expand the exponential operator in power series and use Schwartz regularity. \square

Static Extremum Localization and the Variational-Delta Ladder

Proposition 13 (Static Morse localization). *Let $f \in C^\infty(\mathbb{R}^d) \cap C^\infty(\mathbb{R}^d)$ with unique nondegenerate critical point x^* : $\nabla f(x^*) = 0$, $\det \nabla^2 f(x^*) \neq 0$. $\nabla f(x^*) = 0$, $\det \nabla^2 f(x^*) \neq 0$. For $O \in C_c^\infty(\mathbb{R}^d)$ in $C_c^\infty(\mathbb{R}^d)$, $A_\varepsilon(O) := \varepsilon^{-d/2} \int_{\mathbb{R}^d} e^{i \xi \cdot x} O(x) dx$, $A_{\varepsilon}(O) := \varepsilon^{-d/2} \int_{\mathbb{R}^d} \mathrm{mei} \varepsilon S_m(x) f(x) O(x) dx$ satisfies $|A_\varepsilon(O)|^2 \rightarrow (2\pi)^d |O(x^*)|^2 |\det \nabla^2 f(x^*)|$. $|A_{\varepsilon}(O)|^2 \rightarrow (2\pi)^d \frac{|O(x^*)|^2}{|\det \nabla^2 f(x^*)|}$. Equivalently, $|A_\varepsilon(O)|^2 \rightarrow (2\pi)^d \frac{|\delta(\nabla f)|}{|O|^2}$. $|A_{\varepsilon}(O)|^2 \rightarrow (2\pi)^d \frac{|\delta(\nabla f)|}{|O|^2}$.*

Proof. Standard stationary phase at a single Morse critical point. \square

Corollary 14 (Finite-dimensional QM/QFT truncations). *For any finite-dimensional discretization S_N of a variational action (time-sliced QM or lattice QFT), the same stationary-phase mechanism localizes on $\nabla S_N = 0$, $\nabla S_N = 0$, providing the finite-dimensional realization of $\delta(\delta S)$ as an extremum selector.*

Conclusion

Theorem 3 gives a complete proof of Claim 1 in the scoped projective class:

1. exact cylinder-limit closure in NN,
2. constructive renormalization/counterterm repair,
3. regulator removal $\eta \rightarrow 0^+$ including a coupled quartic block,
4. explicit semiclassical point-supported channel expansion in the Gaussian core.

Open frontier (outside this theorem): genuinely growing mode-coupled large-NN interactions with uniform bounds beyond fixed interacting blocks.

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