

# Feynman's Paradox

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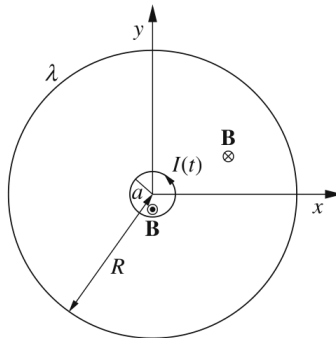
## 1 Introduction

Feynman's disk paradox was a thought experiment that said if a free-floating disk of charges surrounded a solenoid with current flowing through it, the disk would rotate when the current was shut off due to a transfer of angular momentum from the crossed electric and magnetic fields..

## Problem

A non-conducting ring of radius  $R$  is at rest on the  $xy$  plane, with its center at the origin of the coordinate system. The ring has mass  $m$ , negligible thickness, and an electric charge  $Q$  distributed uniformly on it, so that the ring has a linear charge density  $\lambda = Q/(2\pi a)$ . The ring is free to rotate around its axis without friction. A superconducting circular ring of radius  $a \ll R$ , coaxial to the charged ring and carrying an electric current  $I_0$ , also lies on the  $xy$  plane, as shown. At time  $t = 0$  the superconducting loop is heated above its critical temperature, and switches to normal conductivity. Consequently, its current decays to zero according to a law  $I = I(t)$ .

- Neglecting self-induction effects, evaluate the angular velocity  $\omega(t)$  of the charged ring as a function of the current  $I(t)$  in the smaller ring. Evaluate the final angular velocity  $\omega_f$ , and the final angular momentum  $L_f$ , of the charged ring.
- Evaluate the magnetic field at the ring center,  $B_c$ , generated by the rotation of the ring.
- Discuss how the results of a) are modified by taking the "self-inductance"  $L$  of the charged ring into account.



### S-6.6 Feynman's “Paradox”

a) The mutual inductance  $M$  between the charged ring and the superconducting ring is, assuming  $a \ll R$  (see Problem 6.12),

$$M = 4\pi k_m b_m \frac{\pi a^2}{2R}. \quad (\text{S-6.54})$$

Thus, when a current  $I(t)$  is circulating in the smaller ring of radius  $a$ , the magnetic flux through the charged ring is

$$\Phi_I = MI(t) = 4\pi k_m b_m \frac{\pi a^2}{2R} I(t). \quad (\text{S-6.55})$$

If  $\Phi_I$  is time-dependent, it gives origin to an induced electric field  $\mathbf{E}_I$ , whose line-integral around the charged ring is

$$\oint \mathbf{E}_I \cdot d\boldsymbol{\ell} = -b_m \frac{d\Phi_I}{dt} = -4\pi k_m b_m^2 \frac{\pi a^2}{2R} \partial_t I(t). \quad (\text{S-6.56})$$

Due to the symmetry of our problem, field  $\mathbf{E}_I$  is azimuthal on the  $xy$  plane, and independent of  $\phi$ . Its magnitude on the charged ring is thus

$$E_I = \frac{1}{2\pi R} \oint \mathbf{E}_I \cdot d\boldsymbol{\ell} = -k_m b_m^2 \frac{\pi a^2}{R^2} \partial_t I(t), \quad (\text{S-6.57})$$

and the force exerted on an infinitesimal element  $d\boldsymbol{\ell}$  of the charged ring is

$$d\mathbf{f} = \mathbf{E}_I \lambda d\boldsymbol{\ell} = -\hat{\phi} k_m b_m^2 \frac{\pi a^2}{R^2} \lambda d\boldsymbol{\ell} \partial_t I(t), \quad (\text{S-6.58})$$

corresponding to a torque  $d\boldsymbol{\tau}$  about the center of the ring

$$d\boldsymbol{\tau} = \mathbf{r} \times d\mathbf{f} = -\hat{\mathbf{z}} k_m b_m^2 \frac{\pi a^2}{R} \lambda d\boldsymbol{\ell} \partial_t I(t). \quad (\text{S-6.59})$$

The total torque on the charged ring is thus

$$\boldsymbol{\tau} = \int d\boldsymbol{\tau} = -\hat{\mathbf{z}} k_m b_m^2 \frac{\pi a^2}{R} \lambda 2\pi R \partial_t I(t) = -\hat{\mathbf{z}} k_m b_m^2 \frac{\pi a^2}{R} Q \partial_t I(t), \quad (\text{S-6.60})$$

where  $Q = 2\pi R \lambda$  is the total charge of the ring. The equation of motion for the charged ring is thus

$$mR^2 \frac{d\omega}{dt} = \boldsymbol{\tau} = -k_m b_m^2 \frac{\pi a^2}{R} Q \partial_t I(t), \quad (\text{S-6.61})$$

where  $mR^2$  is the moment of inertia of the ring. The solution for  $\omega(t)$  is

$$\omega(t) = -k_m b_m^2 \frac{\pi a^2}{mR^3} Q \int_0^t \partial_{t'} I(t') dt' = k_m b_m^2 \frac{\pi a^2}{mR^3} Q [I_0 - I(t)], \quad (\text{S-6.62})$$

and the final angular velocity is

$$\omega_f = k_m b_m^2 \frac{\pi a^2}{mR^3} Q I_0 = \begin{cases} \frac{\mu_0 a^2 Q}{4mR^3} I_0, & \text{SI,} \\ \frac{\pi a^2 Q}{c^3 mR^3} I_0, & \text{Gaussian,} \end{cases} \quad (\text{S-6.63})$$

corresponding to a final angular momentum

$$L_f = mR^2 \omega_f = k_m b_m^2 \frac{\pi a^2 Q}{R} I_0 = \begin{cases} \frac{\mu_0 a^2 Q}{R} I_0, & \text{SI,} \\ \frac{\pi a^2 Q}{c^3 R} I_0, & \text{Gaussian,} \end{cases} \quad (\text{S-6.64})$$

independent of the mass  $m$  of the ring.

**b)** The rotating charged ring is equivalent to a circular loop carrying a current  $I_{\text{rot}} = Q\omega/2\pi$ . Thus, after the current in the small ring is switched off, there is still a magnetic field due to the rotation of the charged ring. The final magnetic field at the center of the rings is

$$\begin{aligned}\mathbf{B}_c &= \hat{\mathbf{z}} \frac{k_m}{2} \frac{I_{\text{rot}}}{R} = \hat{\mathbf{z}} \frac{k_m}{4\pi} \frac{Q\omega}{R} \\ &= \hat{\mathbf{z}} \frac{k_m^2 b_m^2 a^2 Q^2}{4mR^4} I_0 = \begin{cases} \hat{\mathbf{z}} \frac{\mu_0^2 a^2 Q^2}{64\pi^2 mR^4} I_0, & \text{SI,} \\ \hat{\mathbf{z}} \frac{a^2 Q^2}{4c^4 mR^4} I_0, & \text{Gaussian,} \end{cases}\end{aligned}\quad (\text{S-6.65})$$

parallel to the initial field  $\mathbf{B}_0 = \hat{\mathbf{z}} k_m I_0/(2a)$ , in agreement with Lenz's law. We further have

$$\pi a^2 B_c = M I_{\text{rot}}, \quad (\text{S-6.66})$$

where  $M$  is the mutual inductance of the rings (S-6.54).

**c)** As seen above at point **b)**, the rotating charged ring generates a magnetic field all over the space. This field modifies the magnetic flux through the rotating ring itself, giving origin to self-induction. Let  $\mathcal{L}$  be the "self-inductance" of the rotating ring. The magnetic flux generated by the rotating ring through itself is

$$\Phi_{\text{rot}} = \frac{1}{b_m} \mathcal{L} I_{\text{rot}} = \frac{1}{b_m} \mathcal{L} \frac{Q\omega}{2\pi}. \quad (\text{S-6.67})$$

Correspondingly, (S-6.56) for the line integral of the electric field around the charged ring is modified as follows:

$$\oint \mathbf{E}_I \cdot d\boldsymbol{\ell} = -b_m \left( \frac{d\Phi_I}{dt} + \frac{d\Phi_{\text{rot}}}{dt} \right) = -\frac{4\pi^2 k_m b_m^2 a^2}{2R} \partial_t I - \mathcal{L} \frac{Q}{2\pi} \frac{d\omega}{dt}. \quad (\text{S-6.68})$$

The torque on the ring becomes

$$\tau = -\hat{\mathbf{z}} \left( \frac{k_m b_m^2 \pi a^2 Q}{R} \partial_t I + \mathcal{L} \frac{Q^2 a^2}{2\pi} \frac{d\omega}{dt} \right), \quad (\text{S-6.69})$$

and the equation of motion (S-6.61) becomes

$$mR^2 \frac{d\omega}{dt} = -\frac{k_m b_m^2 \pi a^2 Q}{R} \partial_t I - \mathcal{L} \frac{Q^2 a^2}{2\pi} \frac{d\omega}{dt},$$

or

$$\left( mR^2 + \mathcal{L} \frac{Q^2 a^2}{2\pi} \right) \frac{d\omega}{dt} = -\frac{k_m b_m^2 \pi a^2 Q}{R} \partial_t I, \quad (\text{S-6.70})$$

which is equivalent to (S-6.61) if we replace the mass of the charged ring by an effective value

$$m_{\text{eff}} = m + \mathcal{L} \frac{Q^2 a^2}{2\pi R^2}. \quad (\text{S-6.71})$$

Thus we obtain for the dependence of  $\omega$  on  $I(t)$

$$\omega(t) = k_m b_m^2 \frac{\pi a^2 Q}{m_{\text{eff}} R^3} [I_0 - I(t)], \quad (\text{S-6.72})$$

and for its final value

$$\omega_f = k_m b_m^2 \frac{\pi a^2 Q}{m_{\text{eff}} R^3} I_0, \quad (\text{S-6.73})$$

corresponding to a final angular momentum

$$\begin{aligned} L_f &= mR^2 \omega_f = k_m b_m^2 \frac{\pi a^2 Q}{R + \mathcal{L} a^2 Q^2 / (2\pi m R)} I_0 \\ &= \begin{cases} \frac{\mu_0}{4\pi} \frac{\pi a^2 Q}{R + \mathcal{L} a^2 Q^2 / (2\pi m R)} I_0, & \text{SI,} \\ \frac{1}{c^3} \frac{\pi a^2 Q}{R + \mathcal{L} a^2 Q^2 / (2\pi m R)} I_0, & \text{Gaussian.} \end{cases} \end{aligned} \quad (\text{S-6.74})$$

The final magnetic flux through the charged ring is

$$\Phi_f = \frac{1}{b_m} \mathcal{L} \frac{Q \omega_f}{2\pi} = k_m b_m \frac{\mathcal{L} a^2 Q^2}{2mR^3 + \mathcal{L} Q^2 a^2 R / \pi} I_0, \quad (\text{S-6.75})$$

and the approximations of point a) are valid only if

$$\Phi_f \ll \Phi_0 = 4\pi k_m b_m \frac{\pi a^2}{2R} I_0, \quad \text{or} \quad \frac{\mathcal{L} Q^2}{4\pi^2 m R^2 + 2\pi \mathcal{L} Q^2 a^2} \ll 1. \quad (\text{S-6.76})$$