A How can Magnetism be emplained with the nelp of special relativity)

Ans: Before going through the answer, lets, first tack about some prerequistic knowledge.

1. 2nd postulate of special theory of relativity

3. Lorentz Prantformation

2"d postulate: Speed of light is constant, independent of the relative motion of the sowice.

length contraction:

suppose an observer moves with some speed 'VAN. when he calculates me speed of light he win tind it to be

suppose, another observer moving with a very high speed (comparable to the speed of light), when he will measure the speed of to uguly he will bind it to be c

" According to and postulate of S.T.R , we com

say the above thengs"

13ut, there are cartain 6/27 avre consequences regarding tuis. Two se observers in relative motion, may see that the distance bet n points may shrink (length contraction) to the time period betn. diff phyrical events may disate (Time disation) just to make sure the speed of lights remain constant.

we are interested in knowing about the length controllion.

when we meanuel something, the meanurement of the length in the dirn of the relative motion "appears" to have contracted wit the observer in relative motion.

will only happen in the direct relative motion. (*) length contraction

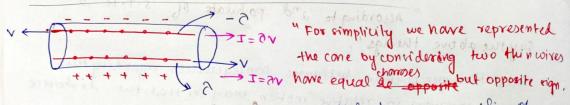
In order to derive the enpression for Longth contraction, we are going to use a set of transformation eqns known as Loventz Transformation eqns. [This are eqns that connect the coordinate of space to time but two observers in relative motion]

In are not included the derivations of Length contraction using Lorentzing transformation, otherwise the presentation will be too long".

tength contracted length (ΔL) = Original length (ΔL_0) · $\sqrt{1-\frac{v^2}{c^2}}$ i.e $\Delta L = \Delta L_0 \sqrt{\frac{1-v^2}{c^2}}$ unhere, $v \rightarrow vel$, of observer observer $c \rightarrow speed$ of light.

End of prerequisites

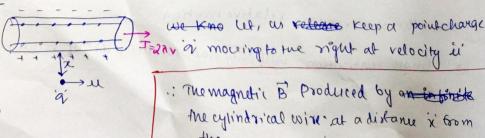
NOW, let us explain the magnetism using speci length conteaction from special theory of Relativity with a simple case!



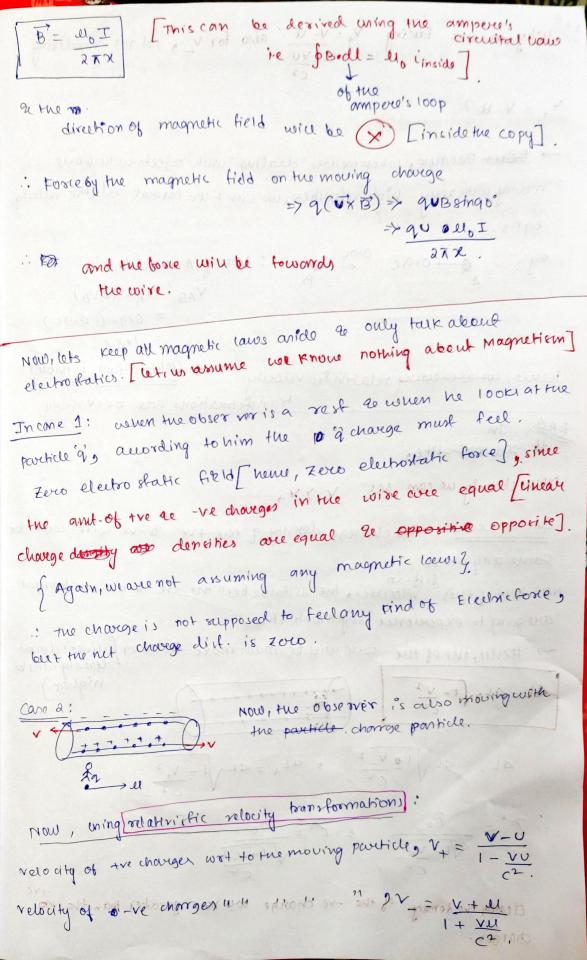
let, the big wire conxists of two distinct wires convisting of two charges moving to the right (with vel. i) the -ve charges moving to the right (with vel. i) the recommendate having linear to the left (with vel. v). The two wises are having linear charge density +2 9e - 2 respectively.

Total coverent = AV+AV = AAV.

Couse 1:



the axis of the wise.



 $t \stackrel{\text{th}}{=} w \text{ by , we were fairing } V_{+} = \frac{v - u}{t - \frac{v \cdot u}{c^{2}}}$ also for v_, not not considering 1 = N-W.) -> Beaux Because, when we are dealing with objects which are maing with very high velocities, we can't use normal relative velocity eg: Nel. of A worl B VAB = VA - (OVB) = 019-(-0,90) = 1.8 C. cat know nothing about magnetism? Gehich isnot possible! hence, we were aning relativistic velocity transformations are necessary # 30 700 1160 erro electro static fittid himi, now, clearly we can see 90 -ve wire wont be Consequences: The charge denvity of the tre change descript as Same any more Now relit q thès velocities, the distance bet the tre & the ve changes are going to experience length contraction. Henry one of the coise will be much more or denser [linear chame higher) AL = ALOVI - V2 $\Delta L = \Delta L \sqrt{\frac{10 V_{2}}{C^{2}}}$, $\Delta L_{1} = \Delta L \sqrt{1 - \frac{V_{1}^{2}}{C^{2}}}$ $\frac{1}{\sqrt{1-\frac{v_{+}^{2}}{v_{+}^{2}}}}, \frac{1}{\sqrt{1-\frac{v_{-}^{2}}{v_{-}^{2}}}}$

charges.

:. The cross section of the colve has non-zero net charge. - As a result, the external charged particle would now experience some net force. to a moving observer there should be force ". But we can show that, the monient we do transormation from moving

charge to the rest observer, It will be enactly me same as

we know about the magnetic field.

Now, Net charge density: met = 12+ - 2- = 20 10 - 20 10-1 $\sqrt{1-\frac{\sqrt{2}}{c^2}} \sqrt{1-\frac{\sqrt{2}}{c^2}}$

on volving we get,

met tobal chourge density with the mouning observer. 2net = -270 vu $\frac{1-v^2}{c^2}\sqrt{1-\frac{u^2}{c^2}}$ Clearly, the net charac dentity is a non-zero term. - are not any direction town

And nother force the wire will enert to the particle win be toward

the wire E due to a line charge density at a distance x' from the wix:

E'= Prot [com, be derived using gaun law.

area x

of the com gaussian surfall. Force experienced by the particle, when the objection is in moving frame

Remembring of free space 2 KEOXC2 VI- 412

Now, using relativistic force transformation:

Force experienced by the change (F')

Foxe enpercined by the charge wellen the observer is ingroud (F)

1-e F'= F VI-UL (2)

 $: F = F^{1} \times \sqrt{1 - u^{2}} = -\frac{92 \times u \times u}{4 \times \sqrt{1 - u^{2}}} \times \sqrt{1 - u^{2}}$

= - 92 valo [:]= 20 v

 $=-9u\left(\frac{u_0T}{\lambda\kappa\eta}\right)$

Clearly, we bound a force acting on the particle of

having a magnitude (-au. uoz and direction toward the wire.

This, enactly looks like the force enested by the magnetic tield on the pareticle.

Hener, aring Relativity we explained the Magnetism