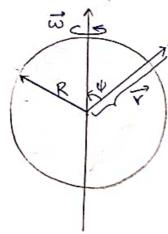
PHYSICS Presentation of question

RETHUVARNA. K.M 20211210

A spherical shell of radius R, corrying a unitorm surtace change or is set spinning at angular velocity w. Find the veltor potential it produced at point r.

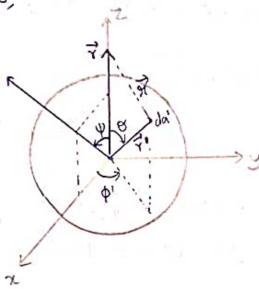


- ans. The surface charge on spherical shell is .
 - >> When the spherical shell is set spinning at angular velocity w the charge present on the surface is also spining along with it, that is it is moving. Due to this surface current density will be produced.
 - => Due to this surface current density, a magnetic field is produced in the sphere. If magnetic field is produced, the magnetic vector potential will also be there. : B = DXA.

here in this question we want to find the value of. magnetic vector potential.

It might seem natural bo set the polaraxis along w, but infact the integration is easier if we let if he on the zaxis, so that w is tilted at an angle 4.

We may as well orient the x-axis so that w lies in xz plane, as shown below,



- Now, there can be a point on zaxis, where the magnetic field due to surface current produced on a small onea element da' can be felt.
- So, the source point (7) from origin) and field point & from brigin) is what shown in the figure. The separation vector is 97.
- The angle between the z-axis and the axis about which the sphere rotates is 4.
- The spherical coordinate of the sorarce point is (r', 0, 0)
 - > so now we want to find the magnetic potential at that point (on z axis) due to the surface current produced on the sphere.
 - This point can be inside or outside sphere.

 If r < R then it is inside and if r > R it is outside sphere.

we know,
$$\overrightarrow{A}(\overrightarrow{r}) = \frac{\mu_0}{4\pi} \int \frac{K(\overrightarrow{r}) d\overrightarrow{a}}{\cancel{r}} - \overrightarrow{0}$$
.

where 97 is the separation vector.

da' - small area element of the source coordinate k(Ti) - surface current density.

$$91 = |\vec{91}| = |\vec{R}^2 + \vec{r}^2 - 2R\vec{r} \cos \theta|$$
 (finding anknown side of a triangle)

- (ri is R here) surface of shell.

value of da in terms of spherical coordinate is da'= R2 sino'do'do' (da'= Rsino'do'do')

we know the formulas

$$\therefore \vec{k} = \sigma(\vec{w} \times \vec{r})$$
put in (1) \(\vec{A}(\vec{r}) = \frac{M_0}{4\pi} \in \frac{\sigma(\vec{w} \times \vec{r})}{\sigma} \sigma^2

w is in nz plane, it can be resolved as, - wsmy in + weary &

$$\nabla = \vec{w} \times \vec{r} = \begin{bmatrix} \hat{n} & \hat{y} & \hat{z} \\ \hat{w} & \hat{y} & \hat{z} \end{bmatrix}$$

$$\begin{cases}
\nabla = \vec{w} \times \vec{r} = \begin{bmatrix} \hat{n} & \hat{y} & \hat{z} \\ \hat{w} & \hat{w} & \hat{w} & \hat{w} \\ \hat{k} & \hat{w} & \hat{w} & \hat{w} \\ \hat{k} & \hat{v} & \hat{w} & \hat{w} \\ \end{pmatrix}$$

$$\begin{cases}
R \sin \theta \cos \theta & R \sin \theta \sin \theta \\ \hat{v} & R \cos \theta
\end{cases}$$

$$\overrightarrow{W} \times \overrightarrow{F}' = \begin{vmatrix} \overrightarrow{A} & \overrightarrow{G} & \overrightarrow{G} \\ wsin \Psi & O \\ Rsin O'sind' & Rios \Phi' \end{vmatrix}$$

Respectively with the contraction of the present of the contraction of the contractio

where wit it are resolved into components

$$V = RW \left[-(\omega \cos \psi \sin \theta' \cos \phi') \pi^{2} + (\cos \psi \sin \theta' \cos \phi') - \sin \psi \cos \theta') \hat{g} + (\cos \psi \sin \theta' \sin \phi') \hat{g} + (\cos \psi \sin \theta' \sin \phi') \hat{g} + (\sin \psi \sin \theta') \hat{g} + (\cos \psi \sin \theta' \cos \theta') \hat{g} + (\cos \psi \sin \theta') \hat{g} + (\cos \theta') \hat{g}$$

take
$$V = R^2 + Y^2 - 2RYU$$

$$AV = -2RYUU$$

when $U = -1 \implies V = (R+Y)^2$

$$U = +1 \implies V = (R-Y)^2$$

... the integral become:
$$(R-Y)^2 \frac{V^2 - R^2 - Y^2}{\sqrt{V}(-2RY)} \frac{dV}{-2RY}$$

$$= \frac{1}{4R^2Y^2} \frac{(R-Y)^2}{(R+Y)^2} \frac{(R^2+Y^2)}{\sqrt{V}(-2RY)} \frac{dV}{\sqrt{V}}$$

on evaluation this becomes
$$\frac{1}{4R^2Y^2} \left[\frac{2V\sqrt{V}}{3} \frac{(R-Y)^2}{(R+Y)^2} - (R^2+Y^2) \frac{2V\sqrt{(R+Y)^2}}{(R+Y)^2} \right]$$

$$= \frac{1}{3R^2Y^2} \left[(R^2 + Y^2 + RY) | R - Y| - (R^2 + Y^2 - RY) (R+Y) \right]$$
Thirde the sphere for $R > Y$

$$= \frac{1}{-3R^2Y^2} \left[(R^3 - YR^2 + RY^2 - Y^3 + YR^2 - RY^2) - (R^3 + YR^4 + RY^2 + Y^3 - YR^2 - RY^2) \right]$$

$$= \frac{1}{2R^2Y^2} \left[(R^3 - Y^3) - (R^3 + Y^3) \right]$$

$$= \frac{1}{2R^2Y^2} \left[(R^3 - Y^3) - (R^3 + Y^3) \right]$$

 $\therefore \int \frac{u du}{\sqrt{R^2 + v^2 - 2Rvu}} = \frac{2\gamma^3}{3R^2 v^2} = \frac{2r}{3R^2}$ This ide the sphere for RTY.

smilory for T>R can be done.

Now, 3. becomes.

outside the sphere For RET RKY

$$\vec{A} = - \frac{MoR^3 - wsin}{z} \Psi \left(\frac{2R}{3r^2} \right) \hat{g}$$

$$\vec{A} = \mu_0 P^3 = (-\omega r, \sin \psi \vec{g}) \left(\frac{2r}{3r^3}\right)$$
 and $\vec{w} \times \vec{r} = -\omega r \sin \psi \vec{g}$

Inside the sphere for R>Y

(i) $\overrightarrow{A}(\overrightarrow{F}) = \begin{cases} \frac{M_0 R \in (\overrightarrow{U} \times \overrightarrow{F})}{3}, & \text{for points abstitle the sphere} \\ \frac{3}{M_0 R \cdot \cancel{F}} (\overrightarrow{U} \times \overrightarrow{F}), & \text{for points butside the sphere} \end{cases}$

In newborral spherical coordinates, r has angle o making with z-axis. Let w is along zaxis, angle between wand v is O. Then. Dx7 = -wy sino \$

$$\overrightarrow{A} = \underbrace{M_0 R^4 \in \overrightarrow{W} \times \overrightarrow{7}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 3}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}{3 \times 2}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}_{3 \times 3}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}_{3 \times 3}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}_{3 \times 3}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}_{3 \times 3}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}_{3 \times 3}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}_{3 \times 3}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}_{3 \times 3}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-w \times sno \beta)}_{3 \times 3}}_{3 \times 3} = \underbrace{\frac{M_0 R^4 \in (-$$

and inside the sphere.

$$\overrightarrow{A}(\gamma_0, \phi) = -\frac{M_0 W R \sigma}{3} r \sin \theta$$

curiously, the field inside this spherical shell is unitom.

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$$\vec{B} = \vec{\nabla} \times \vec{A} = 2 \underline{MoRW \in (100 \hat{\tau} - smo \hat{\theta})}$$