

Haushal Patel

20221112

Batch 3

Q. Suppose the actual force of interaction between two point charges is

$$\vec{F} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2} \left(1 + \frac{r}{\lambda}\right) e^{-r/\lambda} \hat{r}$$

where λ is a constant and is of the order of the half of the radius of the universe so that the force reduces to Coulomb's law for ' r ' small enough. Formulate the new laws of electrostatics.

Solⁿ:

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \left(1 + \frac{r}{\lambda}\right) e^{-r/\lambda} \hat{r}$$

For a system of charges, by superposition principle, we have

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \sum_{i=1}^n \frac{q_i}{r_i^2} \left(1 + \frac{r_i}{\lambda}\right) e^{-r_i/\lambda} \hat{r}_i$$

For a continuous charge distribution,

$$\boxed{\vec{E} = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}') \hat{r}}{r^2} \left(1 + \frac{r}{\lambda}\right) e^{-r/\lambda} d\tau}$$

As the electric field is radial and symmetric,

$$\vec{\nabla} \times \vec{E} = 0$$

$\Rightarrow \int_a^b \vec{E} \cdot d\vec{r}$ is independent of path and therefore it has a scalar potential such that,

$$\boxed{\vec{E} = -\vec{\nabla} V}$$

$$V(\vec{r}) - V(\infty) = - \int_{\infty}^{\vec{r}} \vec{E} \cdot d\vec{r}$$

as $\int \vec{E} \cdot d\vec{r}$ is independent of path, we choose the path ~~at~~ along radial line $d\vec{r} = dr \hat{r}$

$$\Rightarrow V(\vec{r}) - 0 = - \int_{\infty}^{\vec{r}} \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \left(1 + \frac{r}{\lambda}\right) e^{-r/\lambda} dr$$

$$\text{let } \frac{-r}{\lambda} = t \Rightarrow dt = -\frac{dr}{\lambda}$$

$$\Rightarrow V(\vec{r}) = \frac{-q}{4\pi\epsilon_0} \int_{-\infty}^{\frac{-r}{\lambda}} \frac{1}{(-\lambda t)^2} (1-t) e^t (-\lambda dt)$$

$$= \frac{1}{\lambda} \frac{q}{4\pi\epsilon_0} \int_{-\infty}^{\frac{-r}{\lambda}} \left(\frac{1}{t^2} - \frac{1}{t}\right) e^t dt$$

$$= \frac{1}{\lambda} \frac{q}{4\pi\epsilon_0} \left(\frac{e^t}{t} \right) \Big|_{-\infty}^{\frac{-r}{\lambda}}$$

$$\boxed{V(\vec{r}) = \frac{q}{4\pi\epsilon_0} \frac{e^{-r/\lambda}}{r}}$$

Now, Let us calculate the equivalent Gauss's law in this universe for a single charge at origin,

$$\begin{aligned} \oint_S \vec{E} \cdot d\vec{a} &= \iint \frac{q}{4\pi\epsilon_0} \frac{1}{r^2} \left(1 + \frac{r}{\lambda}\right) e^{-r/\lambda} \times r^2 \sin\theta d\theta d\phi \\ &= \frac{q}{4\pi\epsilon_0} \left(1 + \frac{r}{\lambda}\right) e^{-r/\lambda} \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi \\ &= \frac{q}{\epsilon_0} \left(1 + \frac{r}{\lambda}\right) e^{-r/\lambda} \end{aligned}$$

which is not constant and thus the "usual" Gauss' law does not hold. But if we calculate volume integral of potential

$$\begin{aligned} \int_V V d\tau &= \iiint \frac{q}{4\pi\epsilon_0} \frac{e^{-r/\lambda}}{r} \times r^2 \sin\theta dr d\theta d\phi \\ &= \frac{q}{4\pi\epsilon_0} \int r e^{-r/\lambda} dr \int \sin\theta d\theta \int d\phi \\ &= \frac{q}{\epsilon_0} \int r e^{-r/\lambda} dr \end{aligned}$$

By simple substitution $t = -\frac{\rho}{\lambda}$ and then applying integration-by-parts we see that,

$$\int_V V d\tau = (-\lambda^2) \frac{q}{\epsilon_0} \left[\left(1 + \frac{\rho}{\lambda}\right) e^{-\rho/\lambda} - 1 \right]$$

$$\Rightarrow -\frac{1}{\lambda^2} \int_V V d\tau = \underbrace{\frac{q}{\epsilon_0} \left(1 + \frac{\rho}{\lambda}\right) e^{-\rho/\lambda}}_{\oint_S \vec{E} \cdot d\vec{a}} - \frac{q}{\epsilon_0}$$

$$\Rightarrow \boxed{\oint_S \vec{E} \cdot d\vec{a} + \frac{1}{\lambda^2} \int_V V d\tau = \frac{q}{\epsilon_0}}$$

We obtained this for a single charge at origin but, in general, it applies to any distribution of charge and any surface and not only a sphere.

Now by divergence theorem,

$$\int_V (\vec{\nabla} \cdot \vec{E}) d\tau + \frac{1}{\lambda^2} \int_V V d\tau = \int \frac{\rho}{\epsilon_0} d\tau$$

$$\Rightarrow \int_V \left(\vec{\nabla} \cdot \vec{E} + \frac{V}{\lambda^2} \right) d\tau = \int \left(\frac{\rho}{\epsilon_0} \right) d\tau$$

Since this is valid for any volume, the integrands must be same,

$$\boxed{\vec{\nabla} \cdot \vec{E} + \frac{V}{\lambda^2} = \frac{\rho}{\epsilon_0}} \quad \text{--- (1)}$$

Suppose there is some conductor with free charges as well as some extra charge residing on it.

As \vec{E} is zero inside a conductor,

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\text{Also, } \vec{E} = -\vec{\nabla} V \Rightarrow V = \text{constant}$$

$$\text{As } \vec{\nabla} \cdot \vec{E} = 0, \quad \frac{V}{\lambda^2} = \frac{\rho}{\epsilon_0} \text{ and therefore}$$

ρ = charge density ~~on~~ inside the conductor should also be constant.