

Q. An intense energetic electron beam can pass normally through a grounded metal foil. The beam is switched on at $t = 0$ at a current $I = 3 \times 10^6$ amp & a cross-section area $A = 1000 \text{ cm}^2$. After the beam has run for 10^{-8} sec, calculate the electric field at the point P on the output face of the foil & near the beam axis due to space charge of the beam.

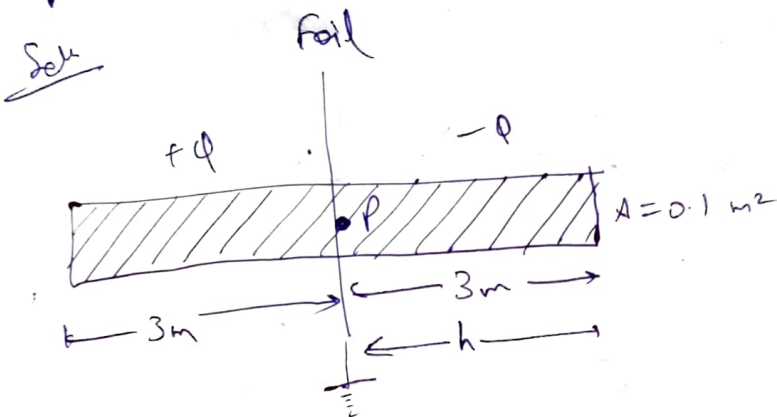


fig. 1

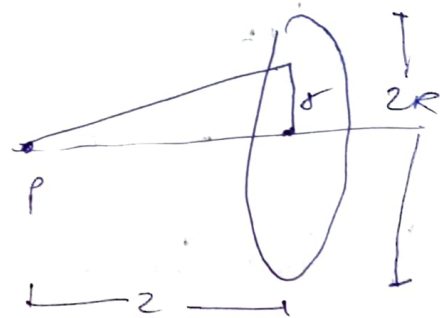


fig. 2

Assuming, electron beam travels c

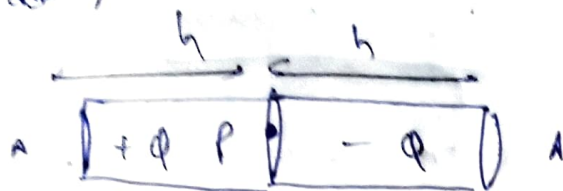
$$h = ct = 3 \times 10^8 \times 10^{-8} = 3\text{m}$$

$$\text{Also, } Q = -It = -3 \times 10^6 \times 10^{-8} = -3 \times 10^{-2} \text{ C}$$

The charge on the left side of the foil does not contribute to the electric field at point P (due to shielding effect). The foil is at net voltage 0 so, the action of the grounded metal foil can be ~~separated~~ replaced by an image charged cylinder. This image cylinder & the real cylinder are symmetrical w.r.t to the metal foil & their charges are opp. in sign in fig 1.

So, the problem is now ~~is~~ \Rightarrow

Electric field at P



We first calculate the electric field at point P on the axis of a uniformly charged disc of surface charge density σ as shown in Fig 2. The V is

$$V_p = \frac{1}{4\pi\epsilon_0} \int_0^R \frac{\sigma \cdot 2\pi r dr}{\sqrt{z^2 + r^2}} = \frac{1}{4\pi\epsilon_0} \int_0^R \frac{\sigma \cdot 2\pi r dr}{\sqrt{z^2 + r^2}}$$

$$= \frac{\pi\sigma}{4\pi\epsilon_0} \left[2\sqrt{z^2 + r^2} \right]_0^R = \frac{\sigma}{2\epsilon_0} \left[\sqrt{z^2 + R^2} - z \right]$$

to the field intensity E.

$$E_p = -\frac{\partial V}{\partial z} = -\frac{\sigma}{2\epsilon_0} \left(\frac{z}{\sqrt{z^2 + R^2}} - 1 \right)$$

from Fig. 1.62, the field at the point P produced by the right charge cylinder is:

$$E_p = \frac{1}{2\epsilon_0} \int_0^h \left(\frac{Q}{h\pi R^2} dz \right) \left(\frac{z}{\sqrt{z^2 + R^2}} - 1 \right)$$

$$= \frac{Q}{2\pi\epsilon_0 h R^2} \int_0^h \left(\frac{\frac{1}{2} dz^2}{\sqrt{z^2 + R^2}} - dz \right)$$

$$= \frac{Q}{2\pi\epsilon_0 h R^2} \int_0^h \frac{\frac{1}{2} dz^2}{\sqrt{z^2 + R^2}} - dz$$

$$= \frac{Q}{2\pi\epsilon_0 h R^2} \left[\sqrt{z^2 + R^2} \Big|_0^h - h \right]$$

$$= \frac{Q}{2\pi\epsilon_0 h R^2} \left[\sqrt{h^2 + R^2} - R - h \right]$$

$$\therefore E_{\text{net at P}} = 2 \times E_0$$

$$= \frac{-Q}{4\pi\epsilon_0 h R^2} \left[R + h - \sqrt{R^2 + h^2} \right]$$

$$= \frac{-3 \times 10^{-2}}{4 \times 8.85 \times 10^{-12} \times 3 \times \frac{0.1}{4}} \times \left[3 + \sqrt{\frac{0.1}{4}} - \sqrt{3^2 + \frac{0.1}{4}} \right]$$

$$= -1.42 \times 10^9 \text{ V/m}$$

The minus sign indicates that the electric field points to the right.

Motivation: The physical importance of this phenomenon lies in the ability to study the properties of material at a very small scale. Electron microscopy uses this phenomenon to produce highly detailed images of the structure of materials. By analyzing the behaviour of electrons as they interact with a material, scientists can gain insight into the material's properties & behaviour.

Other applications include using electron beams to manipulate & modify materials at a small scale, such as in the manufacturing of computer chips.