

# Physics LATAM - Riemannian and Complex Geometry — Homework 3 Solutions

Course: Physics LATAM - Riemannian and Complex Geometry

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Date: October 29, 2025

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## Exercise 3.1 — Equivalence of two definitions of vector bundle

Problem. Show the two standard definitions of a rank- $r$  real vector bundle are equivalent: (A) a total space with local trivializations  $\phi_i : \pi^{-1}(U) \rightarrow U \times \mathbb{R}^r$ ; and (B) an open cover with smooth transition maps  $g_{\alpha\beta} : U_\alpha \cap U_\beta \rightarrow GL(r, \mathbb{R})$  satisfying the cocycle conditions.

Solution (construction in both directions).

Direction (A)  $\rightarrow$  (B).

1. Let  $(U, \phi)$  and  $(V, \psi)$  be two trivializations. On the overlap  $U \cap V$  both give vector-space isomorphisms of each fibre with  $\mathbb{R}^r$ . Define  $g_{UV} : U \cap V \rightarrow GL(r, \mathbb{R})$  by the relation

$$g_{UV}(x) = (p_2 \circ \psi) \circ (p_2 \circ \phi)^{-1} \in GL(r, \mathbb{R}),$$

where  $p_2 : U \times \mathbb{R}^r \rightarrow \mathbb{R}^r$  is projection to the second factor. Concretely, if  $\phi(v) = (x, v_\phi)$  and  $\psi(v) = (x, v_\psi)$  for  $v \in E_x$ , then  $v_\psi = g_{UV}(x)v_\phi$ .

2. The functions  $g_{UV}$  are smooth because they are compositions of smooth maps coming from the trivializations. They satisfy the cocycle conditions:  $g_{UU} = \text{id}$  and on triple overlaps  $g_{UV}g_{VW}g_{WU} = \text{id}$  by straightforward composition.

Direction (B)  $\rightarrow$  (A).

1. Start from the disjoint union  $\bigsqcup_{\alpha} (U_{\alpha} \times \mathbb{R}^r)$ . Impose the equivalence relation:  $(x, v, \alpha) \sim (x, g_{\alpha\beta}(x)v, \beta)$  for  $x \in U_{\alpha} \cap U_{\beta}$ . Denote the quotient by  $E$  and write  $\pi : E \rightarrow M$  for the map induced by  $(x, \cdot, \alpha) \mapsto x$ .
2. The cocycle condition ensures this relation is transitive, so  $E$  is well-defined. Local charts  $\phi_{\alpha} : \pi^{-1}(U_{\alpha}) \rightarrow U_{\alpha} \times \mathbb{R}^r$  are obtained by mapping the equivalence class of  $(x, v, \alpha)$  to  $(x, v)$ . The transition functions between these canonical charts are exactly the  $g_{\alpha\beta}$ , so  $E$  is a smooth manifold with the required local trivializations and fibrewise vector-space structure.
3. One checks the local charts give a smooth atlas (compatibility of charts follows from smoothness of  $g_{\alpha\beta}$ ). Thus the two definitions agree.

### Exercise 3.2 — Line bundle triviality from a nowhere-vanishing section

**Problem.** Show that a real line bundle  $L \rightarrow M$  which admits a smooth section  $s$  with  $s(x) \neq 0$  for all  $x$  is trivial.

**Solution.**

1. Define a map  $\Phi : M \times \mathbb{R} \rightarrow L$  by

$$\Phi(x, t) = t s(x).$$

This map is smooth and fibre-preserving:  $\Phi(x, t) \in L_x$  for all  $(x, t)$ . It is linear in the second factor and so a bundle map.

2. Show  $\Phi$  is a vector-bundle isomorphism. Surjectivity: for any  $v \in L_x$  the fibre is one-dimensional, so there exists a scalar  $t$  with  $v = ts(x)$ ; hence  $v = \Phi(x, t)$ . Injectivity: if  $\Phi(x, t) = 0$  then  $ts(x) = 0$  and since  $s(x) \neq 0$  we get  $t = 0$ .
3. The inverse map is explicitly  $\Phi^{-1}(v) = (\pi(v), \lambda(v))$ , where  $\lambda(v)$  is the unique real number with  $v = \lambda(v)s(\pi(v))$ ; the dependence on  $v$  is smooth because in any local trivialization we represent  $s$  by a nonvanishing function and dividing by it is a smooth operation. Thus  $\Phi$  is a diffeomorphism and linear on each fibre. Therefore  $L \cong M \times \mathbb{R}$  is trivial.

**Conclusion.** A nowhere-vanishing global section provides a global frame for a line bundle, so the bundle is trivial.

### Exercise 3.3 — Poincaré disk isometric to hyperbolic upper half-plane

Problem. Show the unit disk model (Poincaré disk) with metric

$$ds^2 = \frac{4|dz|^2}{(1-|z|^2)^2}, \quad |z| < 1,$$

is isometric to the upper half-plane model  $\mathbb{H} = \{w \in \mathbb{C} : \operatorname{Im} w > 0\}$  with metric

$$ds^2 = \frac{|dw|^2}{(\operatorname{Im} w)^2}.$$

Solution (explicit Möbius map and pullback check).

1. Consider the Möbius transformation

$$w = \phi(z) = i \frac{1+z}{1-z}, \quad |z| < 1.$$

This map sends the unit disk to the upper half-plane. Two elementary algebraic identities are key:

- (i) the derivative

$$\phi'(z) = \frac{2i}{(1-z)^2}, \quad |\phi'(z)| = \frac{2}{|1-z|^2};$$

- (ii) the imaginary part of  $w$ :

$$\operatorname{Im} \phi(z) = \frac{1-|z|^2}{|1-z|^2}.$$

2. Pull back the upper-half-plane metric via  $\phi$ . One gets

$$\begin{aligned} \phi^* \left( \frac{|dw|^2}{(\operatorname{Im} w)^2} \right) &= \frac{|\phi'(z)|^2 |dz|^2}{(\operatorname{Im} \phi(z))^2} \\ &= \frac{(2/|1-z|^2)^2 |dz|^2}{((1-|z|^2)/|1-z|^2)^2} \\ &= \frac{4|dz|^2}{(1-|z|^2)^2}, \end{aligned}$$

which is exactly the Poincaré disk metric. Therefore  $\phi$  is an isometry between the two models.

### Exercise 3.4 — Christoffel symbols for the round sphere metric

Problem. Let  $S^n \subset \mathbb{R}^{n+1}$  be the unit sphere with the metric induced from the Euclidean metric on  $\mathbb{R}^{n+1}$ . Compute the Christoffel symbols of the Levi-Civita connection in stereographic coordinates (the standard local coordinates on  $S^n \setminus \{N\}$ ).

Solution (use conformal metric from stereographic projection).

1. Stereographic projection from the North pole gives coordinates  $u = (u^1, \dots, u^n) \in \mathbb{R}^n$  on the sphere minus the pole. The pullback metric is conformally flat with conformal factor

$$g_{ij}(u) = \lambda(u)^2 \delta_{ij}, \quad \lambda(u) = \frac{2}{1 + |u|^2},$$

where  $|u|^2 = \sum_k (u^k)^2$ . (This is standard; derive by the projection formula or recall the textbook expression.)

2. For a conformal metric  $g_{ij} = \lambda^2 \delta_{ij}$  the Christoffel symbols (with respect to the flat coordinate basis) are

$$\Gamma_{ij}^k = \delta_{jk} \partial_i (\ln \lambda) + \delta_{ik} \partial_j (\ln \lambda) - \delta_{ij} \partial_k (\ln \lambda).$$

This formula follows from the general expression

$$\Gamma_{ij}^k = \frac{1}{2} g^{k\ell} (\partial_i g_{j\ell} + \partial_j g_{i\ell} - \partial_\ell g_{ij}),$$

substituting  $g_{ij} = \lambda^2 \delta_{ij}$  and simplifying.

3. Compute the derivatives of  $\ln \lambda$  for the stereographic factor. We have

$$\ln \lambda = \ln 2 - \ln(1 + |u|^2), \quad \partial_i \ln \lambda = -\frac{2u^i}{1 + |u|^2}.$$

4. Substitute into the conformal formula to obtain an explicit expression:

$$\Gamma_{ij}^k(u) = -\frac{2}{1 + |u|^2} (\delta_{jk} u^i + \delta_{ik} u^j - \delta_{ij} u^k).$$

This is the Christoffel symbol in these coordinates. It is symmetric in  $i, j$  as required, and one may check it agrees with the geometric fact that geodesics are great circles when transformed back to the sphere.

### Exercise 3.5 — Christoffel symbols for the hyperbolic upper-half metric

Problem. For  $H^n(\mathbb{R}) = \{(x^1, \dots, x^{n-1}, y) : y > 0\}$  with metric

$$g = \frac{R^2}{y^2} \sum_{i=1}^{n-1} dx^i \otimes dx^i + \frac{R^2}{y^2} dy \otimes dy = \frac{R^2}{y^2} \delta_{ab} dx^a dx^b,$$

compute the Christoffel symbols.

Solution (conformal-flat calculation).

1. This metric is conformal to the Euclidean metric with conformal factor  $\lambda = R/y$ . So we may use the same conformal formula as in the previous exercise:

$$\Gamma_{ij}^k = \delta_{jk} \partial_i (\ln \lambda) + \delta_{ik} \partial_j (\ln \lambda) - \delta_{ij} \partial_k (\ln \lambda).$$

2. Compute derivatives of  $\ln \lambda$ :

$$\ln \lambda = \ln R - \ln y, \quad \partial_{x^i} \ln \lambda = 0 \quad (i = 1, \dots, n-1), \quad \partial_y \ln \lambda = -\frac{1}{y}.$$

3. Using indices  $a, b, c$  ranging over  $(x^1, \dots, x^{n-1}, y)$ , the only nonzero partial derivatives above involve  $y$ . Plugging into the formula gives the nonzero Christoffel components (writing the last coordinate index as  $n$  to denote  $y$ ):

- For  $k$  corresponding to an  $x$ -direction ( $k \leq n-1$ ):

$$\Gamma_{in}^k = \Gamma_{ni}^k = -\frac{1}{y} \delta_{ik}, \quad (1 \leq i, k \leq n-1).$$

- For  $k = n$  (the  $y$ -direction):

$$\Gamma_{ij}^n = \frac{1}{y} \delta_{ij}, \quad (1 \leq i, j \leq n-1), \quad \Gamma_{nn}^n = -\frac{1}{y}.$$

4. Collecting the nonzero components succinctly (where indices  $i, j, k \leq n-1$  run over the  $x$ -coordinates and index  $n$  denotes  $y$ ):

$$\begin{aligned}\Gamma_{in}^k &= \Gamma_{ni}^k = -\frac{1}{y} \delta_{ik}, & (1 \leq i, k \leq n-1), \\ \Gamma_{ij}^n &= \frac{1}{y} \delta_{ij}, & (1 \leq i, j \leq n-1), \\ \Gamma_{nn}^n &= -\frac{1}{y}.\end{aligned}$$

For the  $n = 2$  case (the familiar upper-half plane), this reduces to the components given in the worked example above.

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### Exercise 3.6 — Metric compatibility equivalent to parallel metric

**Problem.** Let  $g$  be a Riemannian metric and  $\nabla$  a connection. Prove that  $\nabla$  is compatible with  $g$  (i.e. preserves inner products under parallel transport) iff  $\nabla g = 0$  (the metric is covariantly constant).

**Solution.**

1. Recall the coordinate formula for the covariant derivative of a  $(0, 2)$ -tensor  $g$  is

$$(\nabla_X g)(Y, Z) = X(g(Y, Z)) - g(\nabla_X Y, Z) - g(Y, \nabla_X Z).$$

2. If  $\nabla$  is compatible with  $g$ , parallel transport along any curve preserves the inner product. Differentiating the inner product of two parallel vector fields along the curve gives zero; writing that derivative in the form above yields  $(\nabla_X g)(Y, Z) = 0$  for all vector fields  $X, Y, Z$ . Hence  $\nabla g = 0$ .
3. Conversely, if  $\nabla g = 0$ , let  $Y(t), Z(t)$  be parallel vector fields along a curve  $\gamma(t)$  (so  $\nabla_{\dot{\gamma}} Y = 0$  and similarly for  $Z$ ). Then

$$\frac{d}{dt} g(Y, Z) = (\nabla_{\dot{\gamma}} g)(Y, Z) + g(\nabla_{\dot{\gamma}} Y, Z) + g(Y, \nabla_{\dot{\gamma}} Z) = 0 + 0 + 0 = 0.$$

Hence parallel transport preserves inner products; this is exactly metric compatibility.

Therefore the two formulations are equivalent.

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### Exercise 3.7 — Induced connection on tensor product preserves induced metric

Problem. Let  $(M, g)$  be Riemannian and consider the induced metric on  $T^*M \otimes T^*M$  defined by  $g(T, S) := T_{ij}S_{kl}g^{ik}g^{jl}$ . If a connection  $\nabla$  on  $T^*M$  is compatible with  $g$ , show the induced connection on  $T^*M \otimes T^*M$  is compatible with the induced metric.

Solution.

1. The induced connection on the tensor product is the Leibniz (product) connection: for  $T, S$  tensor fields,

$$\nabla_X(T \otimes S) = (\nabla_X T) \otimes S + T \otimes (\nabla_X S).$$

2. Compute the derivative of the inner product:

$$\begin{aligned} X(g(T, S)) &= X(g^{ik}g^{jl}T_{ij}S_{kl}) \\ &= (\nabla_X g^{ik})g^{jl}T_{ij}S_{kl} + g^{ik}(\nabla_X g^{jl})T_{ij}S_{kl} \\ &\quad + g^{ik}g^{jl}(\nabla_X T_{ij})S_{kl} + g^{ik}g^{jl}T_{ij}(\nabla_X S_{kl}). \end{aligned}$$

3. Because  $\nabla$  is compatible with  $g$  on  $T^*M$ , we have  $\nabla_X g^{ik} = 0$  (equivalently  $\nabla_X g_{ik} = 0$  and raising indices commutes with  $\nabla$ ). Therefore the first two terms vanish and we get

$$X(g(T, S)) = g(\nabla_X T, S) + g(T, \nabla_X S),$$

which is exactly the condition that the induced connection is metric-compatible on  $T^*M \otimes T^*M$ .

This completes the proof.

### Exercise 3.8 — Trace and covariant derivative identity

Problem. Let  $(M, g)$  be a Riemannian manifold with Levi-Civita connection  $\nabla$ . For a  $(2, 0)$ -tensor  $T$  prove that for any vector field  $X$ ,

$$X(\text{tr}_g T) = g^{ij}(\nabla_X T)_{ij},$$

where  $\text{tr}_g T = g^{ij}T_{ij}$  is the trace using the inverse metric  $g^{ij}$ .

Solution.

1. Compute the derivative of the trace using the product rule:

$$\begin{aligned} X(\operatorname{tr}_g T) &= X(g^{ij}T_{ij}) \\ &= (\nabla_X g^{ij})T_{ij} + g^{ij}X(T_{ij}). \end{aligned}$$

2. Express  $X(T_{ij})$  in terms of covariant derivatives:

$$(\nabla_X T)_{ij} = X(T_{ij}) - T_{kj}\Gamma_i^k(X) - T_{ik}\Gamma_j^k(X),$$

where the last two terms are the Christoffel contributions. Contracting this with  $g^{ij}$  and using  $\nabla g = 0$  (so  $\nabla_X g^{ij} = 0$ ) removes the connection correction terms, leaving exactly

$$X(\operatorname{tr}_g T) = g^{ij}(\nabla_X T)_{ij}.$$

3. This identity is a direct and useful consequence of metric compatibility of the Levi-Civita connection.